

Rapid directional changes associated with a 6.5 kyr-long Blake geomagnetic excursion at the Blake-Bahama Outer Ridge.

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Abstract

Geomagnetic excursions are recognised as intrinsic features of the Earth's magnetic field. High-resolution records of field behaviour, captured in marine sedimentary cores, present an opportunity to determine the temporal and geometric character of the field during geomagnetic excursions and provide constraints on the mechanisms producing field variability. We present here the highest resolution record yet published of the Blake geomagnetic excursion (~ 125 ka) measured in three cores from Ocean Drilling Program (ODP) Site 1062 on the Blake-Bahama Outer Ridge. The Blake excursion has a controversial structure and timing but these cores have a sufficiently high sedimentation rate (~ 10 cm ka⁻¹) to allow detailed reconstruction of the field behaviour at this site during the excursion. Palaeomagnetic measurements of the cores reveal rapid transitions (< 500 years) between the

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contemporary stable normal polarity and a completely reversed state of long duration which spans a stratigraphic interval of 0.7 m. We determine the duration of the reversed state during the Blake excursion using oxygen isotope stratigraphy, combined with ^{230}Th excess measurements to assess variations in the sedimentation rates through the sections of interest. This provides an age and duration for the Blake excursion with greater accuracy and with constrained uncertainty. We date the directional excursion as falling between 129 and 122 ka with a duration for the deviation of 6.5 ± 1.3 kyr. The long duration of this interval and the fully reversed field suggest the existence of a pseudo-stable, reversed dipole field component during the excursion and challenge the idea that excursions are always of short duration.

Keywords:

paleomagnetism, geomagnetism, geomagnetic excursion, Blake, ^{230}Th , ODP (Ocean Drilling Program)

1. Introduction

Palaeomagnetic measurements have revealed that since the last full reversal the Earth's magnetic field has, for brief intervals, deviated from the behaviour expected during 'normal' secular variation (Laj and Channell, 2007). These significant deviations in direction and intensity of the Earth's field have become known as geomagnetic excursions. While geomagnetic excursions have long been recognised in palaeomagnetic records (e.g. Channell, 2006; Lund et al., 1998; Laj et al., 2006) the mechanism that causes them and their relationship to reversals remain enigmatic (Merrill and McFadden, 1999; Roberts, 2008). The durations of excursions are of particular interest

11 because they may relate to the stability of magnetic field states and provide
12 information about processes during directional changes in the field.

13 For instance, Gubbins (1999), expanding on work by Hollerbach (1993),
14 proposed a hypothesis to explain why the field rapidly returned to its original
15 polarity following an excursion, rather than remain in the opposite state, as in
16 a reversal. The hypothesis emphasized the influence of interactions between
17 the inner and outer core on the geomagnetic field, whereby the inner core
18 imposes a magnetic inertia on the system. Gubbins (1999) suggested that
19 an excursion, in contrast to a full reversal, occurs when the magnetic field in
20 the liquid outer core does not persist in a reversed state for sufficient time for
21 the solid inner core's field to also achieve reversal. The implication is that
22 the duration of excursions should be relatively short (~ 3 kyr); otherwise
23 the inner core would stabilise the reversed polarity interval, resulting in a
24 long-duration (10s-100s kyr) chron. By understanding excursions and their
25 relation to full reversals, we can test this hypothesis, and more generally,
26 gain insight into processes in the Earth's core and the causes of variation in
27 the Earth's field.

28 Detailed analyses of recent geomagnetic excursions have tried to deter-
29 mine typical durations for excursions. Two of the most recent have been
30 studied in particular detail. The Laschamp excursion (LE) (~ 41 ka), first
31 identified in the Laschamp and Olby lava flows in central France (Bonhommet
32 and Babkin, 1967), has since been found in globally distributed marine core
33 records and is now generally recognised as being of relatively short duration
34 (~ 2 kyr) (Channell, 2006; Laj et al., 2000, 2006).

35 Initial studies suggested that the duration of the earlier, Iceland Basin

36 excursion (IBE) (~ 185 ka) might also be relatively short (~ 2 kyr) (Laj et al.,
37 2006). However, using a novel geochemical approach to capture variations in
38 sedimentation rate in cores from the North Atlantic, Knudsen et al. (2007)
39 demonstrated that the IBE had a longer duration of up to 7 kyr, apparently
40 challenging the idea that excursions would have a short duration.

41 The difference in duration for these two excursions and the limited number
42 of excursions studied at sufficiently high resolution, leaves unclear whether
43 excursions have a characteristic duration. The rate at which the
44 Earth's field changes direction, and thus the time taken for the field to re-
45 verse, is an additional constraint on the nature of fluid flow in the outer core
46 that generates the geomagnetic field.

47 A virtual geomagnetic pole (VGP) does not necessarily represent a true
48 magnetic pole, rather it is calculated on the assumption that the magnetic
49 field direction at a location is the result of a dipolar field. The rate at
50 which the VGP 'moves' should therefore be understood as a representation
51 of the rate of change of the field direction at the site location rather than the
52 movement of any 'real' pole.

53 Clement (2004) presented an analysis of records of the last four geomag-
54 netic reversals that estimated the mean time taken for the VGP to reverse.
55 the duration of the transition zones was found to be 5 ± 1 kyr, correspond-
56 ing to a rate of latitudinal change of $18^\circ \text{ kyr}^{-1}$ (where the transition zone is
57 defined as the period within which the VGP was between 45°N and 45°S).

58 However, the rate of change of field direction appears to be significantly
59 faster during excursions. By virtue of the short duration of the Laschamp
60 excursion, during which VGPs from some records reach latitudes beyond

61 45°S (Laj et al., 2006), the sum of the two transition intervals must be less
62 than 2 kyr. This implies rates of VGP change of at least 90° kyr⁻¹; a rate
63 significantly faster than for reversals. The records of the IBE presented by
64 Knudsen et al. (2007), suggest minimum transition rates of ~30° kyr⁻¹, where
65 transitions take 3 kyr, from a core on the Gardar Drift but far more rapid
66 transition rates of ~200° kyr⁻¹ from a core on the Bermuda Rise. While geo-
67 metric effects could potentially account for some of the discrepancy between
68 the excursions and full reversal transition rates, the difference is sufficiently
69 great to suggest that polarity transitions may be more rapid during excur-
70 sional events than during full reversals.

71 **2. The Blake Excursion**

72 The Blake excursion was first documented by Smith and Foster (1969)
73 in seven deep-sea sediment cores from the Blake Outer Ridge in the west-
74 ern North Atlantic. It has since become commonly accepted as a global
75 geomagnetic feature occurring between the Laschamp and Iceland Basin ex-
76 cursions (Laj and Channell, 2007). However, whilst the Blake excursion
77 is represented in many cores and palaeomagnetic records it remains one of
78 the most ambiguous of the recent geomagnetic excursions. What is clear
79 though is that, like other polarity transitions, the Blake excursion is asso-
80 ciated with an interval of significantly reduced geomagnetic field intensity.
81 Global relative palaeointensity stacks reveal a long duration interval of re-
82 duced palaeointensity lasting (~25 kyr), demonstrating that the excursion is
83 an intrinsic feature of the global field (Frank et al., 1997; Valet et al., 2005;
84 Channell et al., 2009; Ziegler et al., 2011).

85 Previous studies of the timing and duration of the Blake event are com-
86 piled in Table 1. Although the field intensity is well-constrained, estimates
87 for the duration and timing of any directional deviation during the Blake ex-
88 cursion vary and the uncertainties are unconstrained. Radiometric dating of
89 lavas suggest reversed directions at 123 ± 7 ka (Zhu et al., 2000) and a reef car-
90 bonate record places the beginning of a reversed event at 132.8 ± 1.3 ka (Ménabréaz
91 et al., 2010) (Table 1). Tric et al. (1991) provide one of the most detailed ma-
92 rine records of a ‘Blake’ event (Fig. 1a), identifying a double event in a core
93 from the Nile River Fan in the Eastern Mediterranean (following work by Tu-
94 cholka et al. (1987)). Using oxygen isotope stratigraphy and tephrochronol-
95 ogy, Tric et al. (1991) found the event within marine isotope stage (MIS) 5e
96 and calculated an upper bound for the excursion’s duration of 4.5 kyr.

97 However, even within the Mediterranean, cores suggest durations rang-
98 ing from 3.5–8.6 kyr (Tucholka et al., 1987). This discrepancy is most likely
99 the result of three factors: the potential for unrecognised variations in sedi-
100 mentation rates between widely spaced age-constrained boundaries and the
101 combined effects of limited resolution and significant magnetic remanence
102 lock-in depths at low-sedimentation rates (< 10 cm kyr⁻¹) (Roberts and Win-
103 kelhofer, 2004). Palaeomagnetic records from cores on the Yermack Plateau in
104 the Arctic suggested a relatively long duration of up to ~ 10 kyr (Nowaczyk
105 et al., 1994). However, these Arctic sediments have since been identified as
106 potentially subject to a chemical remanent magnetization (CRM) that can
107 result in self-reversal (Channell and Xuan, 2009) and thus these records are
108 potentially unreliable.

109 Some of the previous records that suggest the most complex field be-

110 haviour during the Blake excursion, such as those of Zhu et al. (1994) and
111 Fang et al. (1997), have been obtained from sections from the western Loess
112 Plateau of China (Fig. 1b and c). In both of these records, the Loess ap-
113 pears to record multiple reversed periods with an estimated total duration of
114 ~ 5 kyr. Yet dating loess records is not trivial; age models for loess rely on
115 thermoluminescence dating and correlating palaeosol horizons and magnetic
116 parameters with global isotope curves (e.g. Kukla et al., 1988; Evans and
117 Heller, 2001; Fang et al., 1997). The limited resolution of age models means
118 that average sedimentation rates are used over relatively long intervals.

119 As noted by Parés et al. (2004), on examination the two Loess records are
120 significantly different. Most importantly, the reversed directions appear in
121 different stratigraphic horizons in the sections. Zhu et al. (1998) and Parés
122 et al. (2004) demonstrated that there are ambiguities in the faithfulness of
123 loess palaeomagnetic records of the Blake excursion after failing to find the
124 event in many sections in the Chinese Loess Plateau. Parés et al. (2004)
125 attribute the absence of the event to delayed magnetization acquisition and
126 a previously unrecognised CRM. In the face of these difficulties and incon-
127 sistencies, it is hard to be confident in either of the two loess records of the
128 Blake and the field behaviour they describe.

129 Excursionial field behaviour during the Blake excursion therefore remains
130 ambiguous. A coherent global picture of the excursion is limited by the
131 quality of previous records and the difficulties in dating marine records. It is
132 difficult to discern whether previously observed anomalous directions within
133 this period of low field intensity correspond to a global directional excursion
134 feature or represent localised fluctuations apparent during a period of reduced

135 dipole intensity.

136 **3. Study Site and Sampling**

137 We present a new palaeomagnetic record of the Blake excursion, at the
138 highest resolution yet achieved, using discrete samples obtained from Ocean
139 Drilling Program (ODP) Site 1062 on the Bahama Outer Ridge (Fig. 2a,
140 $28^{\circ} 14.8'N$, $74^{\circ} 25.0'W$). ODP Leg 172 has been invaluable in demonstrating
141 the existence of multiple excursions during the Brunhes chron (Lund et al.,
142 1998, 2001). Using published shipboard long-core measurements for ODP
143 Leg 172 (Shipboard Scientific Party, 1998a), three cores from the Bahama
144 Outer Ridge were selected for this study (Site 1062, Cores B, D and E). We
145 derive a new high resolution chronology for the Blake excursion at Site 1062
146 using $^{230}\text{Th}_{xs}$ measurements coupled to an oxygen isotope stratigraphy.

147 Discrete cubic samples ($\sim 7 \text{ cm}^3$ sample volume) were taken from the
148 three cores at the ODP repository in Bremen. In all, 78 samples were taken
149 from core 1062B (16.25–25.40 mbsf), 157 from 1062D (10.37–34.08 mbsf) and
150 151 from 1062E (8.50–27.53 mbsf). Within the excursions interval, sampling
151 was undertaken at 2.5 cm intervals so that this record is virtually continuous.
152 Considering the average sedimentation rate ($\sim 10 \text{ cm kyr}^{-1}$) the individual
153 samples represent an average of 200 years.

154 **4. Age Model**

155 *4.1. Composite Depth Scale*

156 Although the three cores at Site 1062 are close to one another, gaps in
157 recovery mean that a composite depth scale is required to compare equiva-

158 lent intervals between them. The procedure employed shipboard to develop a
159 composite depth scale links the cores using constant depth offsets at selected
160 tie-points. (Shipboard Scientific Party, 1998b). However, the limitation of
161 using linear depth offsets is that not all prominent features within the in-
162 dividual cores are aligned. Within the interval studied here, correlation of
163 decimetre-scale features in the shipboard magnetic susceptibility record was
164 found to be poor. Thus, to determine the precise relationship between sam-
165 ples in neighbouring cores, we have developed a more refined correlation
166 across the studied interval. A new cross-correlation of the cores was cre-
167 ated by comparing records of three sedimentological parameters measured
168 shipboard that are independent of magnetic field direction: L^* , a measure
169 of the lightness of the sediment; gamma ray count; and magnetic suscepti-
170 bility. Features in cores 1062B and D are correlated to those in 1062E. All
171 subsequent depths from 1062B, D and E are quoted in ‘revised metres com-
172 posite depth’ (rmcd) which is their equivalent depth in 1062E (mbsf). Note
173 therefore, that for 1062E, ‘rmcd’ is the same as ‘mbsf’.

174 *4.2. Oxygen Isotope Stratigraphy*

175 Age models for the cores exist from shipboard and later work but have
176 been refined by two separate approaches for this study. In the absence of
177 isotope data, Grützner et al. (2002) derived age models for the cores at Site
178 1062 by tuning variations of estimated carbonate content to the orbital pa-
179 rameters. Although variations in carbonate content show a first-order trend
180 associated with glacial-interglacial cycles for some Leg 172 sites (Chaisson
181 et al., 2002), variance in sedimentological parameters may be affected by
182 local, rather than global, factors.

183 To refine the chronology, we use two separate approaches for this study.
184 We have cross-correlated Core 1062E with a benthic $\delta^{18}\text{O}$ record in a neigh-
185 bouring piston core, Knorr 31 GPC9 (28°14.7'N, 74°26.4'W, 4758 m) (Keig-
186 win et al., 1994) and measure a planktic $\delta^{18}\text{O}$ record in Core 1062E itself.

187 We correlate Core 1062E with Knorr 31 GPC9 using the variations in es-
188 timated carbonate content derived from diffuse spectral reflectance (Giosan
189 et al., 2001). GPC9 has a measured carbonate content record and a mea-
190 sured oxygen isotope record from analyses of benthic foraminifera *Cibici-*
191 *doides* spp. (Keigwin et al., 1994) (Figure 2b). On the basis of our correla-
192 tion between the two cores, the GPC9 benthic $\delta^{18}\text{O}$ record is placed on the
193 1062E depth scale in Figure 2c.

194 To further confirm this $\delta^{18}\text{O}$ chronology, we also measured the oxygen iso-
195 topes of planktonic foraminifera from Core 1062E. Following magnetic analy-
196 sis, each sample was submerged in deionised water and mechanically agitated
197 until suspension of the sediment was achieved. The sediment suspension was
198 then sieved and the 63 μm fraction retained. This fraction was subsequently
199 dried in an oven at 50°C. Specimens of the planktonic *Globorotalia inflata*
200 were the most consistently present species throughout the studied core inter-
201 val and were therefore chosen for analysis. Benthic foraminifera were rare.
202 Due to the small volume of sediment per sample ($\sim 7 \text{ cm}^3$), only a limited
203 number of samples contained sufficient specimens for analysis. Oxygen iso-
204 tope analyses were conducted using a Kiel IV attached to a Delta V isotope
205 ratio mass spectrometer at the University of Oxford. All measurements are
206 given in standard delta notation in per mil (‰) relative to the Vienna PeeDee
207 Belemnite (V-PDB) carbonate standard. Isotopic values were related to the

208 V-PDB standard by repeated measurements of National Bureau of Standards
209 NBS-19 and NBS-18 (with $\delta^{18}\text{O}$ values of -2.2‰ and -23.2‰ respectively).
210 Analytical precision for replicate standard measurements ($n=13$) of $\delta^{18}\text{O}$ in
211 NBS-19, made concurrently with 1062E analyses, was $\pm 0.09\text{‰}$ (1σ).

212 In total, 53 samples from 1062E were analysed. Three well-preserved
213 specimens of *G. inflata* were used for each stable isotopic measurement and
214 10 samples were repeated to check the reproducibility of the results and assess
215 the potential for seasonal bias. Two measurements, at 12.73 and 14.78 rmcd,
216 were identified as anomalous and are probably the result of diagenetic alter-
217 ation of fractured shells. These measurements were disregarded from further
218 analysis. The results of our $\delta^{18}\text{O}$ analysis are shown in Figure 2c.

219 4.3. Chronology

220 Biostratigraphy conducted during the ODP expedition identified the last
221 appearance of *Globorotalia tumida flexuosa* (68 ka) to be at a minimum depth
222 of 13.23 rmcd in core 1062E (Shipboard Scientific Party, 1998a). Using this
223 biostratigraphic datum as a guide, we correlate the GPC9 and 1062E oxygen
224 isotope curves to the global record.

225 We derive two age-constrained boundaries (Fig. 2c.) from the oxygen
226 isotope records to construct our age model. We place our younger age
227 constraint, the transition from marine isotope stage (MIS) 5 to MIS 4 at
228 13.64 ± 0.35 rmcd. This value, with a conservative estimate for the uncer-
229 tainty, encompasses the transition to more positive values in both the 1062E
230 and GPC9 $\delta^{18}\text{O}$ records. We assign this depth an age of 72.0 ± 1.5 ka (Cut-
231 ler et al., 2003; Lisiecki and Raymo, 2005). For our older age constraint,
232 the change to more negative isotopic values at 20.20 ± 0.40 rmcd is corre-

233 lated with the transition from MIS 6 to MIS 5 and is assigned an age of
234 135 ± 1.5 ka (Cheng et al., 2009; Kawamura et al., 2007; Henderson et al.,
235 2001; Thomas et al., 2009; Drysdale et al., 2009). The assigned ages and
236 uncertainties represent the best estimates for these climatic transitions using
237 published ages determined by uranium-series dating of carbonate material
238 such as corals and stalagmites.

239 Although broadly consistent with the age-models determined on-board
240 the ODP cruise ship (Shipboard Scientific Party, 1998a) and later by Grützner
241 et al. (2002) (using variations in magnetic susceptibility and carbonate con-
242 tent respectively), the $\delta^{18}\text{O}$ stratigraphy reveals that the transition between
243 MIS 5 and MIS 4 occurs a little deeper in the core (13.64 ± 0.35 rmcd) than
244 suggested by the carbonate content record (12.96 rmcd) (Fig. 2d.). The $\delta^{18}\text{O}$
245 transition between MIS 6 and 5 (20.20 ± 0.40 rmcd) is consistent with that
246 suggested by the carbonate record (20.58 rmcd).

247 **5. Palaeomagnetic Methods & Results**

248 All samples were subjected to stepwise alternating-field (AF) demagne-
249 tization at applied peak fields of 3, 6, 9, 12, 15, 20, 25, 30, 35, 40, 50, 60,
250 70, 80, 90 and 100 mT. All measurement and demagnetization steps were
251 performed using a 2G Enterprises DC-SQUID cryogenic magnetometer with
252 an in-line, triaxial, alternating field (AF) demagnetizer in a shielded room at
253 the University of Oxford.

254 Typical orthogonal demagnetization diagrams of discrete samples from
255 1062B, D and E are shown in Figure 3. The vast majority of samples dis-
256 played a steeply inclined remanent magnetization component which was typ-

257 ically removed from each sample by peaks fields between 9 and 25 mT. This
258 is most likely an isothermal remanent magnetization (IRM) induced by ex-
259 posure of the sediment to the magnetized steel core barrel during drilling.
260 This is a common feature of samples from oceanic boreholes and was first
261 noticed in Deep-Sea Drilling Project (DSDP) cores (Ade-Hall and Johnson,
262 1976).

263 Principal component analysis (PCA) was used to determine the charac-
264 teristic remanent magnetization (ChRM). Following removal of the drilling-
265 induced overprint, the majority of samples (80%) display a single, high-
266 stability ChRM. Typically, less than 10% of the total remanence remained
267 after application of the 80 mT peak field. These samples all had ChRMs with
268 maximum angular deviation (MAD) values less than 5° . The high-stability
269 component of magnetization determined suggests that the sediment retains
270 an excellent record of past changes in field direction.

271 The remaining samples exhibited more complex behaviour; a few samples
272 displayed a randomly oriented low intensity component of remanence that
273 was easily removed by the lowest peak field, 3 mT. This was interpreted as a
274 viscous overprint that was probably acquired during storage. A further group
275 of samples also exhibited a component above 80-90 mT which appeared to
276 have no preferred direction, accounting for 1-5% of the total remanence. At
277 this level, the low signal to noise ratio obscures natural remanence directions.

278 Samples from intervals during the transitions between polarities often
279 displayed the most complex behaviour accompanied by significantly reduced
280 NRM intensities. In these intervals, the ChRMs were more ambiguous: 17
281 MAD values exceeded 10° but none of the samples used in later analysis had

282 MAD values greater than 15° . Only seven samples failed to yield directions
283 fulfilling the above criteria and were subsequently rejected. In general, their
284 ChRMs were not in disagreement with the rest of the record but their MAD
285 values were greater than 15° .

286 *5.1. Palaeomagnetic Directions*

287 All core sections were originally extracted unoriented. Furthermore, the
288 ChRM directions showed progressive trends in declination along core. This
289 was taken as an indication that the cores had been twisted during the drilling
290 process. To examine changes in latitude in the virtual geomagnetic pole we
291 assume that the records conform to a geocentric axial dipole model (GAD)
292 whereby the Earth's field approximates a dipole aligned along the axis of
293 rotation. We oriented all core sections (~ 1.5 m length) such that the mean
294 declination outside the excursions interval is oriented towards the geographic
295 North Pole (0°) (Knudsen et al., 2006).

296 The resulting three records of palaeomagnetic direction (Fig. 4) are re-
297 markably similar and all exhibit a period of excursions behaviour spanning
298 0.7 m. The average non-excursions inclinations of 46° in 1062B and 43° in
299 both 1062D and 1062E compare well with the expected inclination of 44°
300 for the sites' latitude predicted by the TK03.GAD model (Tauxe and Kent,
301 2004).

302 In all three records, the inclination changes rapidly at ~ 19 rncd marking
303 the initiation of the directional excursion (Fig. 4b). A brief peak in positive
304 inclinations is immediately followed by an abrupt change to negative incli-
305 nations, reaching a maximum of -55 to -65° . Simultaneously, the declina-
306 tions swing to 180° (Fig. 4c). This initial transition occurs within 500 years.

307 Within the excursions interval, 1062B exhibits shallower inclinations than
308 the other two cores. Virtual geomagnetic poles (VGPs) were calculated for
309 each sample using the direction vectors of the respective ChRMs (Fig. 4d).
310 The record shows a rapid transition to fully reversed directions before an
311 equally rapid return to northern latitudes. The limited number of transition
312 VGPs precludes discussion of VGP paths. During non-excursions intervals
313 the variation in the VGP latitude is within that expected of secular variation.

314 In general, the directions derived from the discrete samples are consistent
315 with earlier continuous split-core measurements (Shipboard Scientific Party,
316 1998a) although a second, younger, anomalous feature at ~ 17 rncd apparent
317 in the 1062E split-core record is not replicated in the discrete samples mea-
318 surements (Fig. 4a). The magnetizations of the samples within this interval
319 are indistinguishable from those outside it.

320 *5.2. Core Properties*

321 *5.2.1. Isothermal Remanent Magnetization (IRM) Acquisition*

322 To confirm that the geomagnetic field behaviour has been faithfully recorded,
323 it is important to constrain the effects of the mineralogical properties of the
324 sediment cores. Stepwise application of isothermal remanant magnetization
325 may be used to determine the coercivity spectrum of a sample and thereby
326 its constituent magnetic grain mineralogy. After complete demagnetization,
327 a series of 32 representative samples, 16 each from 1062D and 1062E, were
328 subjected to progressive forward-field IRM acquisition using a Molspin Pulse
329 Magnetizer at the University of Oxford. The samples' magnetizations were
330 measured out-of-field between each step using the 2G magnetometer. Sub-
331 sequently, the samples were subjected to an equivalent series of fields in the

332 opposite direction (back-field).

333 The IRM acquisition curves showed that in some cases saturation isother-
334 mal remanent magnetization (SIRM) was not achieved by application of an
335 800 mT field (90-98% saturation). To determine the SIRM, the acquisition
336 curves were modelled using IRMUNMIX (Heslop et al., 2002) assuming a two
337 component mixture of one lower and one higher coercivity component. The
338 model can be used to calculate an estimate for the SIRM. The remanence
339 coercivity (H_{CR}), the ‘reverse’ field required to reduce a ‘forward’ saturation
340 remanence to zero, varies between 30 and 40 mT (Fig. 5a). The S-Ratio = -
341 $IRM_{0.3T}/SIRM$ was subsequently calculated for all samples. The majority of
342 samples had an S-Ratio greater than 0.85, and no samples measured had an
343 S-Ratio less than 0.7 (Fig. 5b). This indicates that low-coercivity remanence
344 carriers, in this case probably magnetite grains, consistently dominate the
345 magnetic signal. The lack of variation in the H_{CR} values and S-ratios with
346 depth suggests that there has been little change in magnetic grain prove-
347 nance and mineral composition during the studied time interval. Within the
348 reversed direction interval the S-Ratio value is greater than 0.9 indicating
349 that the contribution of any higher coercivity component is minor.

350 *5.2.2. Magnetic Susceptibility*

351 The magnetic susceptibility of the sediments is dependent upon the con-
352 centration, grain-size and mineralogy of the remanence carriers. The IRM
353 acquisition results indicate that the average magnetic mineralogy in the sam-
354 ples does not vary substantially with depth. Therefore variation within the
355 susceptibility record can be used to estimate changes in the size and con-
356 centration of the magnetic grains. The low-field magnetic susceptibility of

357 representative samples was measured using a Geofyzika KLY-2 KappaBridge
358 at Oxford (Fig. 5c). The magnetic susceptibility of the samples shows little
359 variation with depth and is particularly consistent across the excursions interval.
360 Therefore, variation in the intensity of the cores' ChRM is unlikely to
361 be the result of changes in the concentration or grain-size of the remanence
362 carriers.

363 5.2.3. Grain Size

364 The ratio of anhysteretic remanence magnetic susceptibility (χ_{ARM}/χ) to
365 low field magnetic susceptibility may be used to investigate first-order varia-
366 tions in grain-size (Banerjee et al., 1981). The magnitude of the ARM is rel-
367 atively more sensitive to the finer grain-size fraction while low-field magnetic
368 susceptibility (χ) is particularly sensitive to the coarser grain fraction (King
369 et al., 1982). Thus an increase in the proportion of fine to coarse grain
370 magnetite is reflected in shift towards higher χ_{ARM}/χ ratios. An ARM was
371 induced in roughly half of the samples with a 100 mT peak AF field and a
372 0.1 mT bias field using a DTECH-2000 AF unit which can induce ARMs at
373 Imperial College London. The ARM acquired was subsequently demagne-
374 tized at 30 mT.

375 We observe that the excursions interval is characterized by high values of
376 χ_{ARM}/χ (due to high ARM intensities) representing a higher proportion of
377 fine magnetic grains (Fig. 6a). This feature was also recognized in the same
378 interval by Schwartz et al. (1996) in piston cores from the Blake Outer Ridge.
379 Further study of the piston cores by Schwartz et al. (1997) demonstrated that
380 high values of χ_{ARM}/χ and ARM associated with high S-ratios are most
381 likely indicators of bacterial magnetite content. In such cases, high values

382 of χ_{ARM}/χ were indicative of the presence of a very fine grained (single
383 domain) magnetite component with high ARM intensities rather than an
384 overall reduction in the size of the clastic magnetite component.

385 *5.3. Relative Palaeointensity Proxies*

386 Sedimentary records have the potential to record variations in the relative
387 palaeointensity of the Earth's field (e.g. Tauxe, 1993). The magnitude of the
388 ChRM is dependent upon a combination of: the magnetic mineralogy of a
389 sample; the concentration of magnetic grains; and the strength of the original
390 field. Although absolute palaeointensity can not be determined, the relative
391 palaeointensity (RPI) variation may be obtained by accounting for the effects
392 of magnetic grain mineralogy and concentration.

393 We normalize the NRM using the ARM and IRM measurements (Fig. 6c
394 and d). After samples had been subjected to an applied magnetization, they
395 were subsequently demagnetized at 30 mT (Fig. 6b) and the intensities of
396 the remanent NRM, IRMs and ARMs after demagnetization at 30 mT were
397 used to calculate the NRM/ARM and NRM/IRM palaeointensity proxies.

398 The significant variation in χ_{ARM}/χ suggests that we should approach
399 our interpretation of the relative palaeointensity proxies with a degree of cau-
400 tion (Tauxe, 1993). However, the relative palaeointensity proxies do exhibit
401 a high degree of consistency, both between the individual cores and between
402 the two methods employed to determine relative palaeointensity. Further-
403 more, the corroboration of the record by three independent cores is evidence
404 that each has, to a reasonable extent, faithfully recorded the intensity of the
405 local palaeomagnetic field.

406 In general, the relative palaeointensity proxies show a steady decrease in

407 field intensity from a maximum of 3 or 4 times the average intensity in our
408 record. The field eventually collapses to a significant minimum at 19 rncd,
409 coinciding with directional excursion. As the field direction returns to normal
410 polarity, the intensity also begins to recover, but this return occurs over an
411 interval of several metres above the directional excursion. Although these
412 general trends are in agreement with palaeointensity stacks (e.g. Ziegler et al.,
413 2011), the potential for sedimentary influences must limit our confidence in
414 any quantification of the relative intensity changes.

415 **6. Constraining Sedimentation Rate**

416 To determine an age and duration for the event in our cores we must first
417 constrain the sedimentation rate. Alone, our oxygen isotope stratigraphy
418 cannot resolve millennial-scale fluctuations in sedimentation rate between
419 the two age-constrained stratigraphic bounds. However, unlike records of
420 the Blake excursion from other marine cores, we do not have to assume an
421 average sedimentation rate within a broad interval in the stratigraphy. In-
422 stead, using the excess ^{230}Th approach (Francois et al., 2004; Henderson and
423 Anderson, 2003) we can reconstruct sedimentation rate variations relative to
424 the average sedimentation rate between the age-constrained bounds. Follow-
425 ing the approach of Knudsen et al. (2007), we determine an accurate age
426 and duration for the Blake excursion in our record and also provide precise
427 uncertainties. Thus, the chronology we employ to estimate the age and du-
428 ration of the excursion does not suffer from the limited resolution that has
429 hindered many earlier attempts to estimate the ages and durations of short
430 duration events in the past.

431 ^{230}Th is produced in the water column by the decay of dissolved ^{234}U .
432 Unlike ^{234}U , ^{230}Th is highly insoluble. As a result, it is rapidly adsorbed
433 onto the surfaces of sinking particles and transported to the underlying sed-
434 iment: a process known as ‘scavenging’ (Bacon and Anderson, 1982). This
435 component of the ^{230}Th in the sediment is termed ‘excess’ thorium ($^{230}\text{Th}_{xs}$)
436 as it is not supported by any corresponding uranium component.

437 The flux of $^{230}\text{Th}_{xs}$ to the seafloor is dependent only upon the depth of
438 the overlying water column and the concentration of ^{234}U in seawater, both
439 of which are unlikely to have varied considerably over the last several hun-
440 dred thousand years (Henderson, 2002). Modelling the variability of the flux
441 of $^{230}\text{Th}_{xs}$ into the sediment via scavenging using an ocean general circula-
442 tion model (Henderson et al., 1999) indicates that the flux at the location
443 of Site 1062 is likely to have been relatively constant and equivalent to its
444 production rate in the water column. The concentration of the $^{230}\text{Th}_{xs}$ in the
445 sediment is therefore dependent upon the sedimentation rate. By measur-
446 ing the concentration of $^{230}\text{Th}_{xs}$ in the sediment we can determine relative
447 variations in total sediment flux. We assume a simple reciprocal relation-
448 ship whereby a doubling of the sedimentation rate would halve the measured
449 $^{230}\text{Th}_{xs}$ concentration.

450 The relatively high average sedimentation rate between the age-constrained
451 bounds ($\sim 10 \text{ cm kyr}^{-1}$) suggests that there is some lateral input of mate-
452 rial to the site. Our method assumes that the degree of syndepositional
453 lateral redistribution or ‘focussing’ of sediment remains constant at the site
454 throughout this interval. The Blake Outer Ridge and surrounding sediments
455 owe their existence to the influx and deposition of terrigenous sediments car-

456 ried from the north by the deep, geostrophic Western Boundary Undercurrent
 457 (WBUC) (Heezen et al., 1966). The Gulf Stream flows rapidly northwards
 458 along the western margin of the Blake Plateau separating the ridge from
 459 the potential sediment source of the eastern margin of the continent (Heezen
 460 et al., 1966). ODP Site 1062, south of the Blake Outer Ridge and east of the
 461 Blake Plateau, is therefore in a relatively ‘sheltered’ location. The lack of
 462 variation in magnetic grain mineral composition and concentration through-
 463 out our studied interval (Fig. 5) suggests that the sedimentary environment
 464 remained relatively constant during our sampled interval.

465 Scavenging is not, however, the only source of ^{230}Th into the sediment.
 466 To determine the contribution of $^{230}\text{Th}_{xs}$ to the total ^{230}Th content of the
 467 sediment, contributions from ^{230}Th supported by detrital U ($^{230}\text{Th}_{det}$) and
 468 ^{230}Th ingrown from authigenic U ($^{230}\text{Th}_{auth}$) must be subtracted from the
 469 total measured activity of ^{230}Th , such that:

$$^{230}\text{Th}_{xs} = ^{230}\text{Th}_{meas} - ^{230}\text{Th}_{det} - ^{230}\text{Th}_{auth} \quad (1)$$

470 We conducted ^{230}Th measurements on 18 samples spanning the MIS 5
 471 interval of core 1062E. For each sample measured, 0.1 g of dry sediment was
 472 dissolved following the method outlined in Thomas et al. (2007). For two of
 473 the samples, a further 0.1 g of sediment was dissolved independently to test
 474 repeatability. U and Th were then chemically separated from the solution
 475 via anion chromatography (Robinson et al., 2004). The samples were spiked
 476 with a mixed ^{229}Th and ^{236}U spike and concentrations of the isotopes ^{234}U
 477 and ^{238}U were measured statically and ^{230}Th and ^{232}Th were subsequently
 478 measured dynamically, on a Nu Instruments ICP-MS at the Department

479 of Earth Sciences, University of Oxford (Mason and Henderson, 2010) (see
480 Supplementary Material).

481 We calculate the detrital component of the total measured ^{238}U and use
482 this to determine the detrital input of ^{230}Th . $^{238}\text{U}_{det}$ is calculated by as-
483 suming that all of the ^{232}Th in the sample is of detrital origin and using an
484 appropriate value for $(^{238}\text{U}/^{232}\text{Th})_{det}$.

$$^{230}\text{Th}_{det} = \left(\frac{^{230}\text{Th}}{^{238}\text{U}}\right)_{det} \times ^{238}\text{U}_{det} = \left(\frac{^{230}\text{Th}}{^{238}\text{U}}\right)_{det} \times \left[\left(\frac{^{238}\text{U}}{^{232}\text{Th}}\right)_{det} \times ^{232}\text{Th}\right] \quad (2)$$

485 To estimate $(^{238}\text{U}/^{232}\text{Th})_{det}$, we identified the samples unlikely to have a
486 thorium contribution from the decay of authigenic uranium. The distribution
487 of measured $^{238}\text{U}/^{232}\text{Th}$ ratios from our samples includes one group with
488 consistently low values averaging 0.53 ± 0.09 (Fig. 7). The remaining samples
489 have a range of values greater than 0.7. We interpret the first group to
490 represent those samples with limited or no authigenic uranium component,
491 the higher uranium group are therefore identified as those that have had
492 varying contributions of authigenic uranium in addition to the detrital input.
493 We therefore use a value of 0.53 ± 0.09 for $(^{238}\text{U}/^{232}\text{Th})_{det}$, consistent with
494 the range typical for Atlantic sediments ($=0.6 \pm 0.1$ Henderson and Anderson,
495 2003).

496 If ^{230}Th were in secular equilibrium with ^{238}U in marine detrital material,
497 $(^{230}\text{Th}/^{238}\text{U})_{det}$ would equal 1, however, $(^{230}\text{Th}/^{238}\text{U})_{det}$ may be less than 1
498 due to alpha recoil from smaller particles (Osmond and Ivanovich, 1992).
499 We estimate the deviation of the system from equilibrium by averaging the
500 measured $^{234}\text{U}/^{238}\text{U}$ ratios from samples thought to be free of authigenic
501 uranium ($= 0.98 \pm 0.02$). The $(^{230}\text{Th}/^{238}\text{U})_{det}$ ratio is likely to be lower

502 than this due to further alpha recoil during the decay of ^{234}U to ^{230}Th . We
 503 therefore estimate $(^{230}\text{Th}/^{238}\text{U})_{det}$ by doubling the deviation of the measured
 504 $^{234}\text{U}/^{238}\text{U}$ average from secular equilibrium to give a $(^{230}\text{Th}/^{238}\text{U})_{det}$ ratio of
 505 0.96 ± 0.04 .

506 Any remaining ^{238}U not accounted for by lithogenic material is assumed
 507 to be of authigenic origin. The component of ^{230}Th produced by the decay
 508 of authigenic ^{238}U in each sample ($^{230}\text{Th}_{auth}$) is given by (Henderson and
 509 Anderson, 2003):

$$^{230}\text{Th}_{auth} = (^{238}\text{U}_{tot} - ^{238}\text{U}_{det}) \times \left[(1 - e^{-\lambda_{230}t}) + \frac{\lambda_{230}}{\lambda_{230} - \lambda_{234}} \left(\frac{^{234}\text{U}}{^{238}\text{U}}_{init} - 1 \right) (e^{-\lambda_{234}t} - e^{-\lambda_{230}t}) \right] \quad (3)$$

510 Where λ is the decay constant in yr^{-1} , t is the age (derived from the $\delta^{18}\text{O}$
 511 stratigraphy), and $(^{234}\text{U}/^{238}\text{U})_{init}$ is the initial ^{234}U to ^{238}U ratio in seawater
 512 ($=1.146$) (e.g. Robinson et al., 2004). The current $^{230}\text{Th}_{xs}$ activity may then
 513 be determined using Equation 1. The initial $^{230}\text{Th}_{xs}$ concentration was then
 514 calculated by correcting for its own decay:

$$^{230}\text{Th}_{xs}^{\circ} = ^{230}\text{Th}_{xs} e^{\lambda_{230}t} \quad (4)$$

515 The sedimentation rate, $F_{(n)}$ (in cm kyr^{-1}) for each sample interval was
 516 then calculated using the following expression:

$$F_n = F_a \times \frac{\overline{\text{Th}}}{\text{Th}_n} \quad (5)$$

517 Where Th_n is the concentration of $^{230}\text{Th}_{xs}$ in sample n ; F_a is the average
 518 sedimentation rate in cm kyr^{-1} between the two boundaries of known age

519 and $\overline{\text{Th}}$ is the weighted average concentration of $^{230}\text{Th}_{xs}$ between the bound-
520 aries. This assumes a simple reciprocal relationship whereby the variation in
521 the initial concentration of initial $^{230}\text{Th}_{xs}$ about the weighted average initial
522 $^{230}\text{Th}_{xs}$ is proportional to the variation of the sedimentation rate about the
523 average between the boundaries.

524 The average sedimentation rate between the boundaries (72–135 ka) is
525 10.4 cm kyr^{-1} . Our thorium data indicate that the standard deviation of
526 the rate is only 1.6 cm kyr^{-1} and therefore that the sedimentation rate is
527 rather constant during MIS 5, despite climatic changes. The constancy of
528 sedimentation rate allows us to confidently calculate a duration for the Blake
529 excursion from our new high-resolution record and to assess the uncertainty
530 on this duration. Using the dry bulk density measured on-board ship of
531 0.74 g cm^{-3} (Shipboard Scientific Party, 1998a) and employing the equations
532 given by Francois et al. (2004), we calculate a focusing factor (Ψ) of 4.2. Using
533 a third age constraint at the end of MIS 5e, we investigated to what extent the
534 focussing remained constant throughout MIS 5 (available as supplementary
535 material). Our estimates suggest that the focussing factor may vary by up
536 to $\sim 5\%$. This source of uncertainty is included in our uncertainty for the
537 interval duration.

538 **7. Duration and Age of the Blake Excursion at Site 1062**

539 For the sake of comparison with previous studies of excursions, we use a
540 conventional co-latitude VGP cut-off whereby the duration of an excursion
541 is the period of time the VGP spends more than 45° from the geographic
542 north pole (assumed to be equivalent to the long-term time-averaged VGP

543 position) (Merrill and McFadden, 1994). Following this definition, the Blake
544 excursion spans 0.7 m (the interval 18.8–19.5 rmc). If we instead examine
545 the deviation of VGPs away from ‘normal’ secular variation using the adap-
546 tive approach of Vandamme (1994), we obtain a VGP cut-off value from our
547 data of 27° co-latitude. Due to the rapid change between reversal states,
548 the use of this alternative cut-off value defines a very similar excursion in-
549 terval. Assuming a constant sedimentation rate, between the boundaries,
550 of 10.4 cm kyr⁻¹ the duration of the directional excursions interval would
551 be 6.7 kyr. Using the refined, thorium-corrected sedimentation rates, the
552 duration of the excursions interval is 6.5±1.3 kyr, with the event occurring
553 between 128.5±4.0 ka and 122.0±3.9 ka (note that while the ages have a
554 larger uncertainty, they co-vary, such that the duration has a relatively low
555 uncertainty). To precisely constrain the uncertainties on the duration and
556 age of the excursion, we use a Monte-Carlo method in which the input iso-
557 topic data are assigned normal distribution probability density functions and
558 the age and depths of the age-constrained boundaries are varied randomly
559 between maxima and minima. The uncertainty on the duration estimate is
560 principally controlled by potential variation in focussing and the uncertain-
561 ties assigned to the age and depths of the boundaries.

562 Our chronology for the Site 1062 cores is in good agreement with the most
563 precise radiometric age for the start of the excursion (132.8±1.3 ka) (Ménabréaz
564 et al., 2010). However, the stratigraphic relationship between MIS 5e and
565 the excursions event provides an easier way to compare our record with pre-
566 viously published marine core and Loess records because previous records
567 have assumed different chronologies for the oxygen isotope record. In our

568 Site 1062 record, the Blake excursion is almost exactly coeval with MIS 5e
569 and the changes in inclination span the entire MIS 5e interval. Mediter-
570 ranean cores at a similar latitude show that the local directional excursion
571 occurred towards the end of the MIS 5e (Tric et al., 1991; Tucholka et al.,
572 1987), and Arctic records similarly show the excursion coinciding with MIS
573 5e (Nowaczyk et al., 1994).

574 If ‘lock-in’ has occurred at depth below the sediment surface, the age of
575 the measured magnetic signal would be younger than the sediment in which
576 it was recorded and thus the stratigraphic chronology would overestimate the
577 age of the excursion. By studying a core with a relatively high sedimenta-
578 tion rate, $\sim 10 \text{ cm kyr}^{-1}$, we reduce the likelihood of there being a significant
579 ‘lock-in’ depth (Roberts and Winklhofer, 2004). The sharp transitions be-
580 tween polarity states in the Site 1062 record suggest that ‘lock-in’ of the
581 magnetic signal occurs across a limited depth range throughout the record
582 of the excursion. The sedimentation rate is also relatively constant. Thus,
583 even if there is a ‘lock-in’ depth, it is likely to be constant, such that while
584 the age of the excursion may be less, the calculated duration of the excursion
585 is unlikely to change significantly.

586 **8. Discussion**

587 Our new data provides the highest resolution record yet analysed of the
588 Blake geomagnetic excursion. The agreement between the 3 cores indicates
589 that the Site 1062 sediments are excellent recorders of the palaeomagnetic
590 field, yielding high-quality palaeomagnetic vectors with low MAD values.

591 *8.1. Geometry of the Excursion Field*

592 The cores from Site 1062 record a fully reversed field, consistent with
593 some previously obtained records of the Blake excursion (e.g. Fang et al.,
594 1997; Zhu et al., 1994). However, unlike the records of the Blake excursion
595 from the Mediterranean (Tric et al., 1991) and Chinese loess (Fang et al.,
596 1997; Zhu et al., 1994), the cores from Site 1062 do not record multiple
597 reversed intervals separated by normal polarity intervals.

598 Although cores that record only a single reversed interval have generally
599 been attributed to the filtering of the magnetic signal by low sedimentation
600 rate cores (Tucholka et al., 1987; Tric et al., 1991), the high sedimentation
601 rate (~ 10 cm kyr⁻¹) of the Site 1062 cores, the high resolution of our new
602 record, and the recording of rapid polarity transitions imply that the geo-
603 magnetic signal from Site 1062 is well-resolved. Furthermore, the relatively
604 constant sedimentation rate and long duration of the excursion suggest that
605 a hiatus in the Site 1062 record, which could encompass a normal polar-
606 ity interval during the excursion, is unlikely and was similarly not found in
607 previous age models (Fig. 2d) (Grützner et al., 2002).

608 Our new record therefore calls the global existence of multiple reversed
609 intervals into question and suggests that such behaviour should not be con-
610 sidered as characteristic of the Blake excursion.

611 *8.2. Duration of Geomagnetic Excursions*

612 Our record from Site 1062 indicates that the geomagnetic field deviated
613 from directions that might be expected during ‘normal’ secular variation for
614 6.5 ± 1.3 kyr. This duration is within previous estimates for the duration of
615 the Blake excursion (3.5–8.6 kyr) (Tric et al., 1991; Zhu et al., 1994; Fang

616 et al., 1997) but is robust due to the high sedimentation rate and $^{230}\text{Th}_{xs}$
617 constrained chronology. Although the duration of the Blake excursion’s ex-
618 pression in the directional record may vary between locations, the internal
619 process that caused the excursion must therefore have had a relatively ‘long’
620 minimum duration of ~ 6 kyr to create the signal found in our record.

621 The duration of the Blake excursion is therefore significantly longer than
622 that determined for the Laschamp excursion (~ 2 kyr) (Laj et al., 2000, 2006)
623 at a number of global sites, including the North Atlantic, and is comparable
624 with the duration of the Iceland Basin excursion (6.8 ± 0.5 kyr) (Knudsen
625 et al., 2007) recorded at Site 1063 in the North Atlantic from the same ODP
626 leg. Directional excursions therefore occupy a broad range of durations and
627 can feature a reversed polarity field state for more than ~ 3 kyr without this
628 resulting in a full polarity reversal.

629 *8.3. Rapid Transitions*

630 Unlike previous records of the Blake excursion (Tric et al., 1991; Tucholka
631 et al., 1987), Site 1062 fails to record the transitional field in any detail de-
632 spite a relatively high sedimentation rate (~ 10 cm kyr $^{-1}$). Furthermore, the
633 Site 1062 record of the Blake excursion implies rapid and complete rever-
634 sals in field direction within periods of less than 500 years corresponding to a
635 VGP transition rate of $\sim 180^\circ$ kyr $^{-1}$. The rapid transition rate in this record
636 suggests that our long duration for the Blake excursion is compatible with
637 much shorter duration estimates for other excursions such as the Laschamp
638 whereby the difference in duration of the event arises as a result of the length
639 of time the field remains in a fully reversed state rather than differences in
640 the durations of the transitions. Furthermore, this rate of change is simi-

641 lar to that observed for the IBE transitions in the North Atlantic (Knudsen
642 et al., 2007) and notably faster than that during reversals. During a reversal,
643 the single polarity transition takes ~ 5 kyr, corresponding to a rate of only
644 $18^\circ \text{ kyr}^{-1}$ (Clement, 2004). The rapidity of the transitions observed in high
645 resolution records of excursions is a particularly striking feature, providing
646 clues about the mechanism that generates field change (Fig. 9). Any mech-
647 anism or mechanisms to explain excursions and reversals must therefore be
648 able to account for the significant difference in transition rate between the
649 two.

650 *8.4. Stability of the Excursion Field*

651 Recent simulations of field reversals, using a time-varying observational
652 model of the geomagnetic field for the last 7000 years (Korte, 2005), assume a
653 decreased dipole strength to be the cause of transitional fields (Brown et al.,
654 2007; Valet and Plenier, 2008). Excursions are simulated in these models by
655 setting the strength of the dipole to a reduced value for a set time interval.
656 The axial dipole field strength must be decreased significantly (to around
657 20%) before large excursions are observed in the simulations (Brown
658 et al., 2007). Such a decrease in intensity is observed in our record for the
659 Blake, and similarly during other polarity transitions (Ziegler et al., 2011),
660 whereby the relative palaeointensity decreases over a significantly longer du-
661 ration interval before the directional excursion is recorded, and its subsequent
662 increase in strength occurs after the field direction has returned to its initial
663 polarity. However, it is difficult to imagine how a time-varying non-dipole
664 field could maintain fully reversed directions for such a long duration as that
665 found for the Blake excursion in our record (6 kyr+).

666 Valet and Plenier (2008) simulate excursions in a similar model by de-
667 creasing the axial dipole strength to zero and then increasing it back to its
668 initial state - a situation equivalent to that of a full reversal. This approach
669 explores therefore the possibility that a geomagnetic excursion could repre-
670 sent an ‘aborted’ polarity interval (Valet et al., 2008). However, they could
671 only simulate directional records with VGPs in high southerly latitudes by
672 adding a short interval of a weakly reversed dipole field before the return of
673 the field to its initial state. This suggests that an extended duration interval
674 with fully reversed directions, such as that found in our Blake record, must
675 represent a short period when there is a weak, reversed polarity dipole field.
676 Our record does not, however, document an increase in field intensity within
677 the directional excursion, which implies that the strength of the dipole must
678 still be minimal and that this interval still represents an unstable field state.

679 9. Conclusions

- 680 • Combining a $\delta^{18}\text{O}$ stratigraphy and the ^{230}Th excess method, our new
681 high-resolution palaeomagnetic record of the Blake excursion at ODP
682 Site 1062 indicates that the geomagnetic field deviated from directions
683 expected during ‘normal’ secular variation for 6.5 ± 1.3 kyr between 129
684 and 122 ka. This is within the range of previous estimates of the
685 duration of the Blake excursion (3.5–8.6 kyr) at other sites, but previous
686 estimates have generally relied on average sedimentation rates across
687 broad depth intervals. ‘
- 688 • Our relatively long duration for the Blake excursion at Site 1062 is
689 evidence that excursions do not necessarily have short durations (i.e.

690 < 3 kyr) but instead occupy a continuous range of durations. Further-
691 more, the maintenance of fully reversed directions for a long duration
692 implies the presence of a weakly reversed dipole during the excursion.

693 • However, despite the length of the Blake excursion, the record of the
694 event at Site 1062 is characterized by rapid polarity transitions (< 500
695 years) that are significantly faster than those observed during reversals.

696 • The ^{230}Th -excess method shows great promise in constraining the du-
697 ration of short (kyr) events for other records where sedimentation rates
698 may potentially vary more significantly than at Site 1062.

699 • To understand further the true nature of the global geomagnetic field
700 during the Blake excursion, high quality records are required from a
701 number of locations from across the globe particularly from the south-
702 ern hemisphere (Roberts, 2008). Our new record provides the first
703 steps toward a global dataset of high-quality palaeomagnetic data with
704 precise chronological control for the Blake excursion.

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711 NE/G001391/1.

Location	Medium*	Site	Lat.	Long.	Dating	Sed. Rate (cm kyr ⁻¹)	Duration (kyr)	Age / MIS (ka)	Principal Reference
New Mexico	B	Laguna Basalt	35.1	252.3	K/Ar	-	-	128±33	1
Hawaii	B	Hawaii	19.7	155.1	Ar/Ar	-	-	132±32	2
NE China	B	Jilin Province	42.1	128.0	Ar/Ar	-	-	123±7	3
Tahiti	RC	IODP310 M0005	17.8	-149.6	U/Th	-	-	133±1/5e	4
Xining, W China	L	Xining	36.6	101.8	TL*	9.8	5.3	5d	5
W. Plateau, China	L	Lanzhou	36.0	103.8	TL*	2	5.5	5e	6
Blake Ridge, Caribbean	MS	RC10-49	16.6	-79.5	d180	3.4	5	-	7
Blake Ridge, Caribbean	MS	A179-4	16.6	-74.8	d180	2.4	7	5e	7
E. Mediterranean	MS	MD84641	33.0	31.4	d180	2.3	3.5	-	8
W. Mediterranean	MS	KET8004	39.7	12.4	d180	6.4	8.6	5e/5d	8
W. Mediterranean	MS	KET8022	40.6	10.3	d180	5.5	6.4	Spans 5e	8
E. Mediterranean	MS	MD84640	33.1	-32.9	d180	8	4	end of 5e	8
E. Mediterranean	MS	MD84627	32.2	32.3	d180	10.5	4.5	end of 5e	8,9
W. Mediterranean	MS	DED8707	39.7	-12.4	d180	10.5	4	5e	9
Yermack Plateau, Arctic	MS	PS 1533-3 SL	82.0	15.2	d180	2	10	Start 5e	10
Yermack Plateau, Arctic	MS	PS 2212-3 KAL	82.1	15.9	d180	6	10	Start 5e	10
Amazon Fan	MS	942A/946A	5.7	-49.1	d180	100	-	5e	11

Table 1: Age and duration estimates for the Blake Excursion. *B: Basalt, L: Loess, MS: Marine sediment, TL: Thermoluminescence dating and correlation using magnetic susceptibility. RC: Reef carbonate. References: (1) Champion et al. (1988), (2) Holt and Kirschvink (1996), (3) Zhu et al. (2000), (4) Ménabréaz et al. (2010) - note that date is the start of excursion, (5) Zhu et al. (1994), (6) Fang et al. (1997), (7) Smith and Foster (1969), (8) Tucholka et al. (1987), (9) Tric et al. (1991), (10) Nowaczyk et al. (1994), (11) Cisowski and Hall (1997).

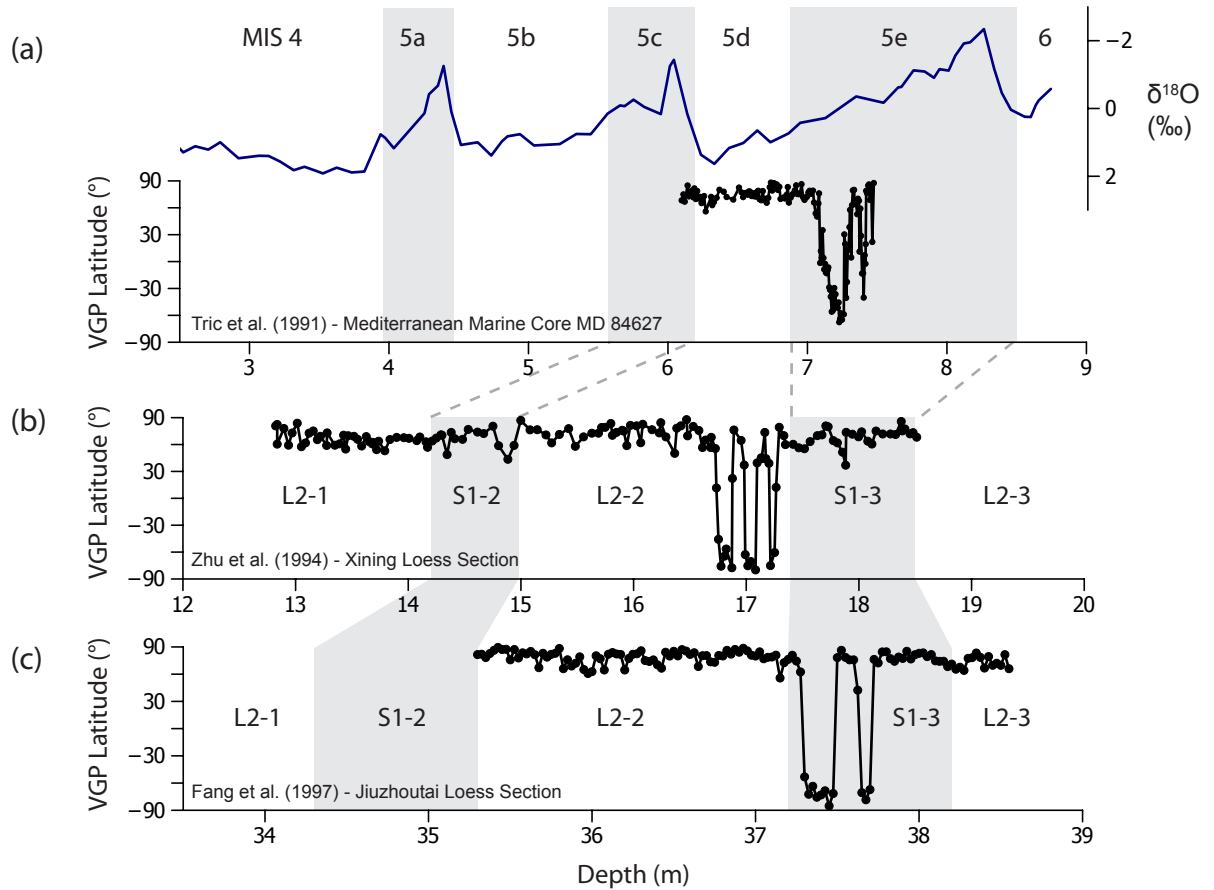


Figure 1: Three of the highest resolution records of the Blake excursion from the literature. (a) $\delta^{18}\text{O}$ stratigraphy for Mediterranean core MD 84627 with Marine Isotope Stages (MIS) identified, and below, the VGP latitude curve from the same core (Tric et al., 1991). (b) VGP latitude record from Xining Loess section, with palaeosol (grey bars) and loess stratigraphy correlated to MIS (dashed lines) (Zhu et al., 1994). (c) VGP latitude record for Jiuzhoutai Loess section and section stratigraphy (Fang et al., 1997).

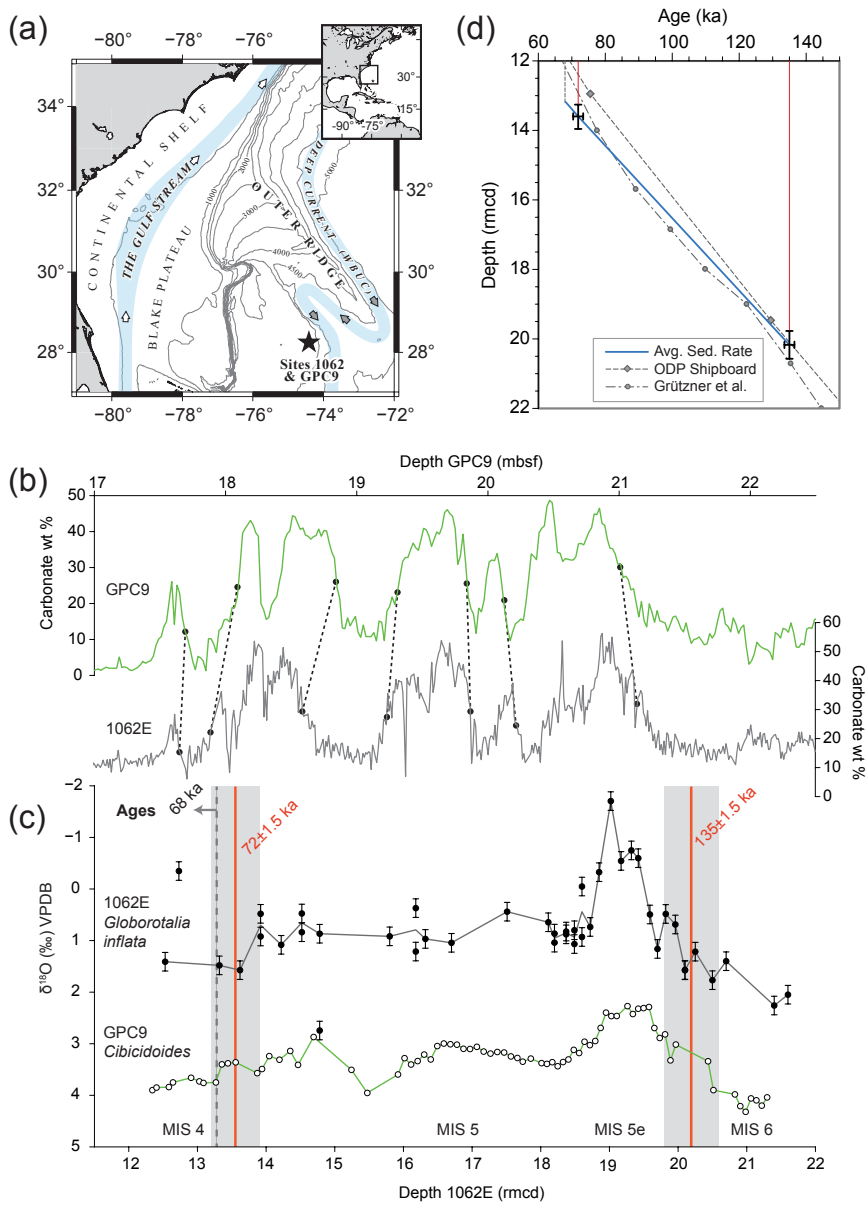


Figure 2: (a) Location of Site 1062 and GPC9 (WBUC - Western Boundary Undercurrent). Bathymetry data taken from ‘The GEBCO One Minute Grid, version 2.0, <http://www.gebco.net>’. (b) Carbonate records from GPC9 (Keigwin et al., 1994) and 1062E (Giosan et al., 2001). Correlated features used as tie-points marked by grey dashed lines. 1062E Depth in ‘rmcd’ (see text) equivalent to metres below seafloor (mbsf). (c) 1062E and GPC9 $\delta^{18}\text{O}$ records on Core 1062E depth scale (rmcd). $\delta^{18}\text{O}$ measured with respect to the Vienna PeeDee Belemnite carbonate standard in delta notation (‰). Age constraints in the $\delta^{18}\text{O}$ stratigraphy for the MIS 6/5 & MIS 5a/4 transitions are marked in red with assumed depth uncertainties shown by the grey bars. Last appearance of *Globorotalia tumida flexuosa* (68 ka) marked with grey arrow. (d) Age-depth plot of core 1062E. Age constraints used in this study are marked in red. Solid blue line assumes average sedimentation between selected age constraints, dashed grey lines show previous age models (Shipboard Scientific Party, 1998a; Grützner et al., 2002).

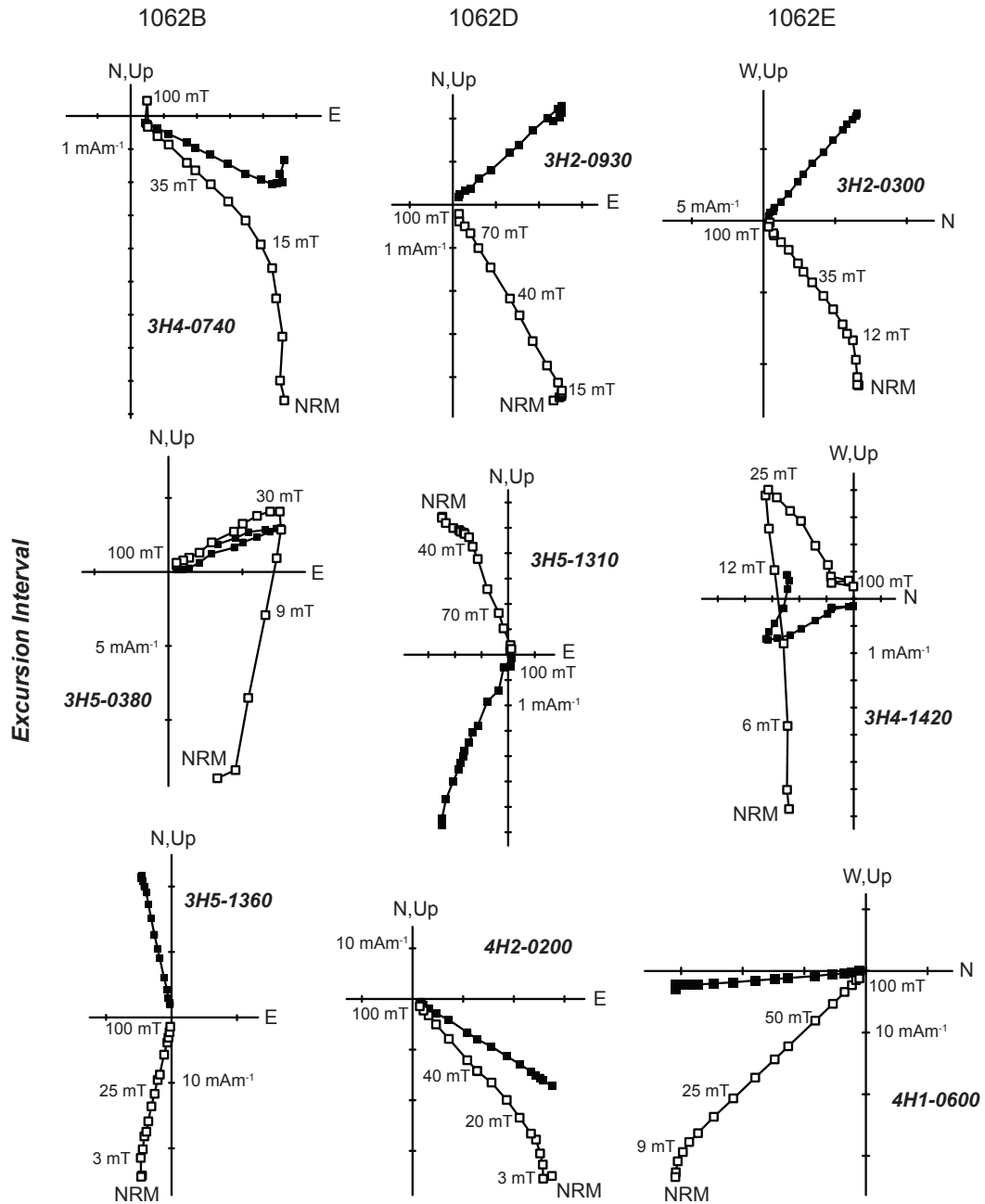


Figure 3: Typical orthogonal demagnetization plots for samples from Cores 1062B, D and E. The top and bottom rows are representative of samples younger and older than the excursions interval (middle row) respectively. Sample declinations are unoriented. Magnetization of sample in milli-Amps per metre (mAm^{-1}). Applied alternating-field (AF) demagnetization field strength in milli-Tesla (mT). Open (solid) squares represent the projection of the magnetic vector onto the vertical (horizontal) plane.

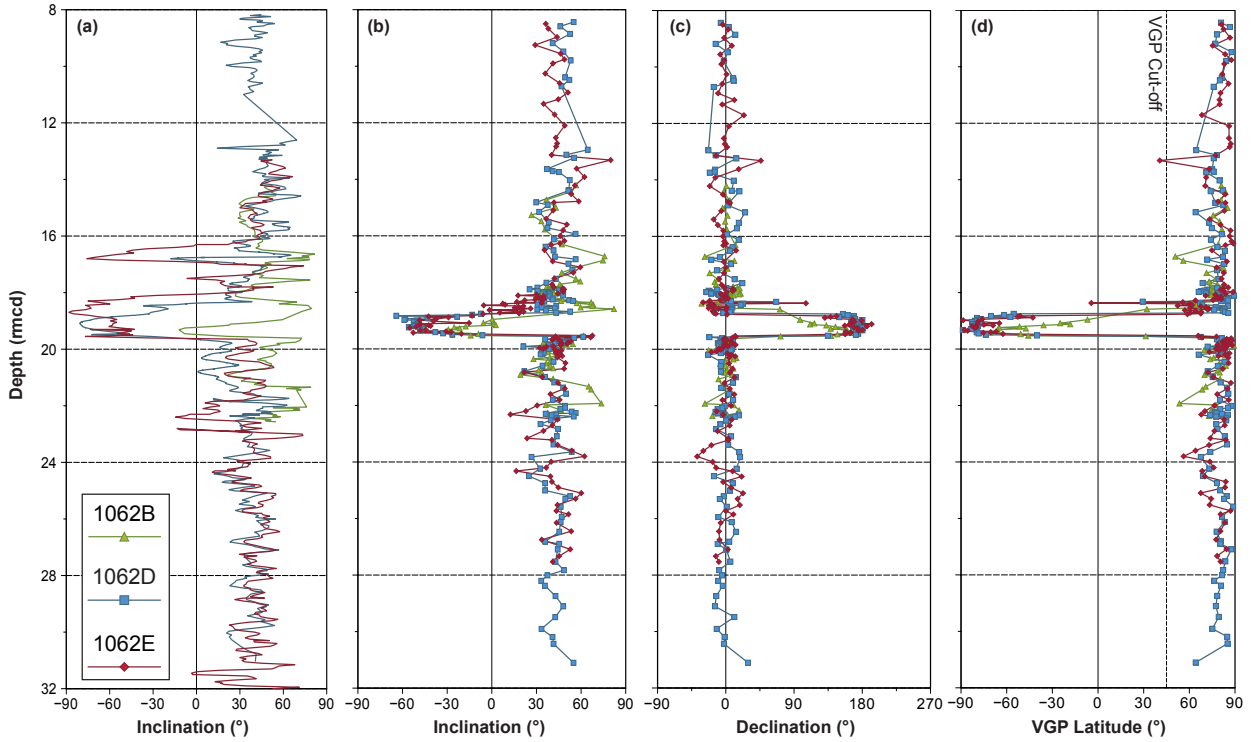


Figure 4: (a) Inclinations of the ChRMs of samples from cores 1062B,D and E against depth (mcd) measured from split-cores on-board ship (Shipboard Scientific Party, 1998a), and (b) from discrete samples (this study). (c) Corrected declinations of discrete samples. (d) Calculated latitude of the virtual geomagnetic poles (VGPs) for all samples. VGP cut-off co-latitude used to define limits of directional excursion marked by dashed line at 45° .

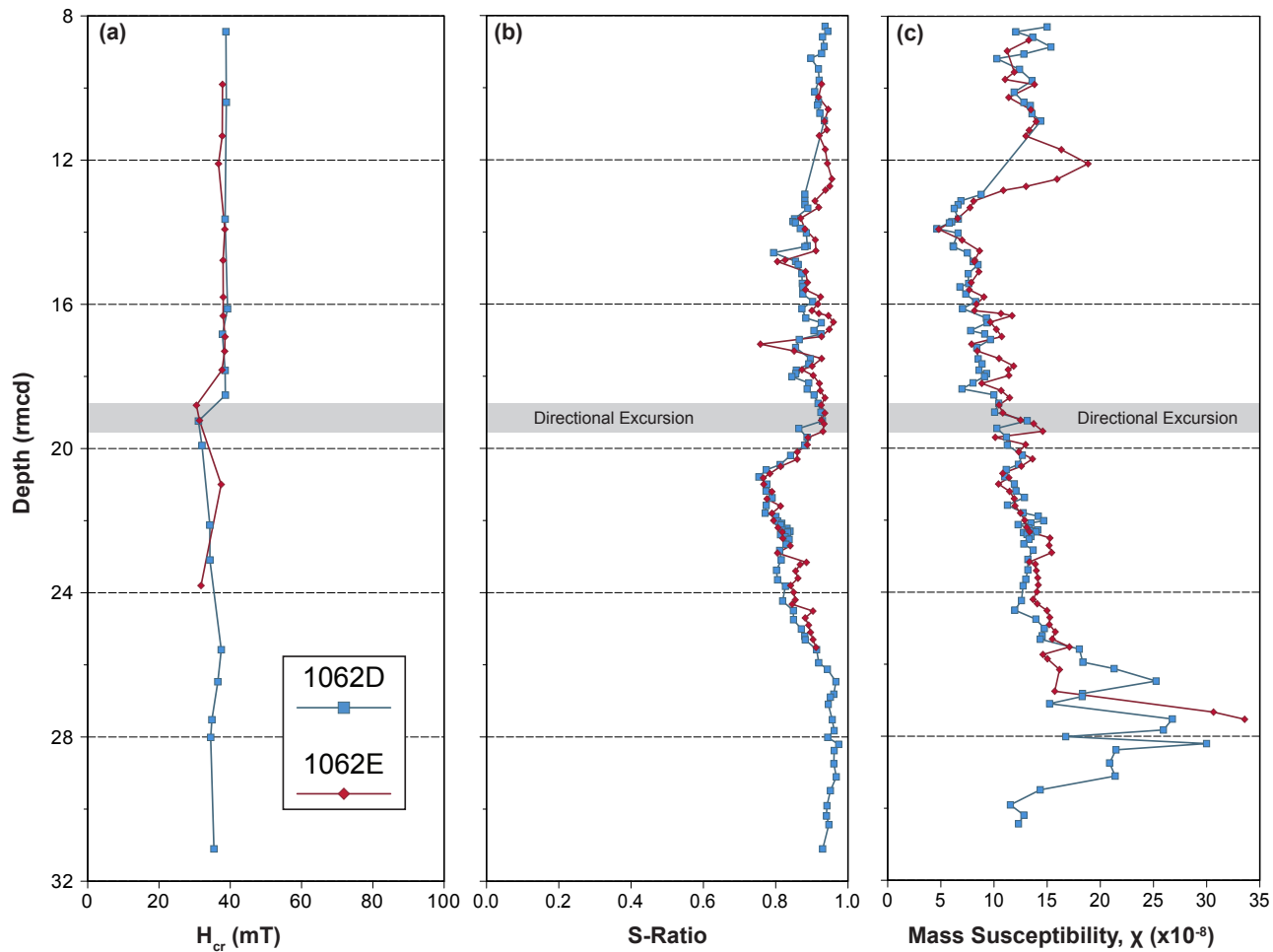


Figure 5: (a) Remanence coercivity (H_{cr} of measured samples in milli-Tesla (mT) from cores 1062D and E against depth (rmcd). (b) S-Ratio (dimensionless - $S\text{-Ratio} = -IRM_{0.3T}/IRM_{0.8T}$). (c) Mass specific magnetic susceptibility (χ). ($10^{-8} \text{ m}^3 \text{ kg}^{-1}$). Directional excursions interval highlighted by grey-shaded bar defined as in Figure 4.

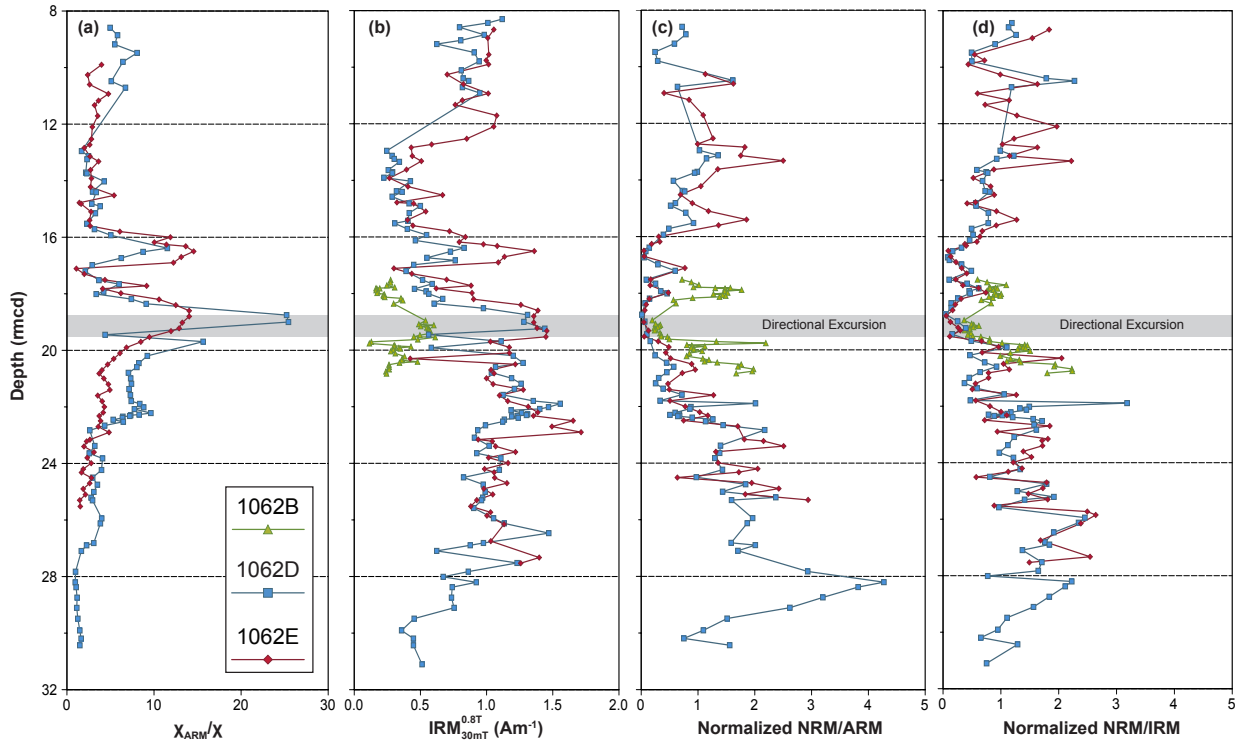


Figure 6: (a) Anhyseretic magnetic susceptibility divided by the measured magnetic susceptibility (χ_{ARM}/χ) for samples from cores 1062D and E. (b) Intensity of applied isothermal remanent magnetization (at 0.8 T) following AF demagnetization at a peak-field of 30 mT. Relative palaeointensity records (c) and (d) for 1062B, D and E derived by dividing the natural remanent magnetizations by the measured anhysteretic and isothermal magnetizations respectively (normalized such that the average values for the records are equal to 1). Directional excursions interval highlighted by grey-shaded bar defined as in Figure 4.

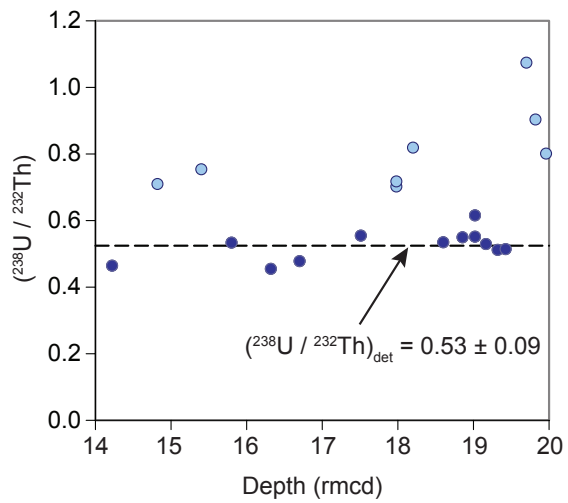


Figure 7: The distribution of measured $^{238}\text{U}/^{232}\text{Th}$ ratios from our samples. One group of samples has consistent low values averaging 0.53 ± 0.09 (dark blue circles), interpreted as the value for $(^{238}\text{U}/^{232}\text{Th})_{det}$. The remaining, higher uranium group (light blue circles) are identified as those that have had varying contributions of authigenic uranium in addition to the detrital ^{238}U .

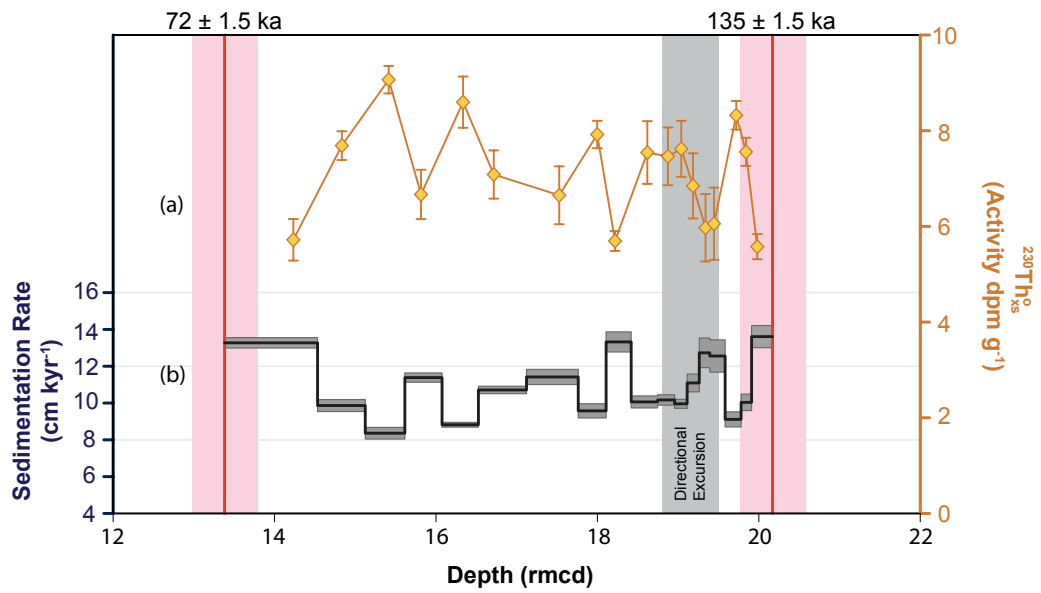


Figure 8: (a) Calculated $^{230}\text{Th}_{xs}^{\circ}$ (orange diamonds) in decays per minute per gram (dpm g^{-1}) with 2σ uncertainty in core 1062E (rmcd). (b) $^{230}\text{Th}_{xs}^{\circ}$ -normalized sedimentation rate (solid line) between the two age constraints marked in red ($72 \pm 1.5 \text{ ka}$ and $135 \pm 1.5 \text{ ka}$ in cm kyr^{-1}) with depth uncertainty envelopes in pink. Grey envelope indicates 2σ uncertainty on sedimentation rate.



Figure 9: VGP latitude record of the Blake Excursion from ODP Site 1062E. Transition zones (VGP within $\pm 45^\circ$) highlighted in grey.

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