

On the depletion and accretion time-scales of cold gas in local early-type galaxies

Timothy A. Davis^{1,2★} and Martin Bureau³

¹*Centre for Astrophysics Research, University of Hertfordshire, Hatfield, Herts AL1 9AB, UK*

²*School of Physics & Astronomy, Cardiff University, Queens Buildings, The Parade, Cardiff CF24 3AA, UK*

³*Sub-department of Astrophysics, University of Oxford, Denys Wilkinson Building, Keble Road, Oxford OX1 3RH, UK*

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ABSTRACT

We consider what can be learnt about the processes of gas accretion and depletion from the kinematic misalignment between the cold/warm gas and stars in local early-type galaxies. Using simple analytic arguments and a toy model of the processes involved, we show that the lack of objects with counter-rotating gas reservoirs strongly constrains the relaxation, depletion and accretion time-scales of gas in early-type galaxies. Standard values of the accretion rate, star-formation efficiency and relaxation rate are not simultaneously consistent with the observed distribution of kinematic misalignments. To reproduce that distribution, both fast gas depletion ($t_{\text{dep}} \lesssim 10^8$ yr; e.g. more efficient star formation) and fast gas destruction (e.g. by active galactic nucleus feedback) can be invoked, but both also require a high rate of gas-rich mergers (>1 Gyr^{−1}). Alternatively, the relaxation of misaligned material could happen over very long time-scales ($\simeq 100$ dynamical times or ≈ 1 –5 Gyr). We explore the various physical processes that could lead to fast gas depletion and/or slow gas relaxation, and discuss the prospects of using kinematic misalignments to probe gas-rich accretion processes in the era of large integral-field spectroscopic surveys.

Key words: ISM: evolution – ISM: molecules – galaxies: elliptical and lenticular, cD – galaxies: evolution – galaxies: ISM – galaxies: kinematics and dynamics.

1 INTRODUCTION

In all its forms, gas is one of the most important drivers of galaxy evolution. The availability of cold gas (atomic and molecular) regulates star formation, and thus the evolution and fate of galaxies, from the beginning of the Universe till today. Many factors affect the supply and removal of gas. Some processes allow gas to flow on to galaxies and form stars (e.g. cold accretion, mergers, hot halo cooling and stellar mass-loss; e.g. Keres et al. 2005; Parriott & Bregman 2008; Bauermeister, Blitz & Ma 2010), while others compete to remove, destroy, and/or prevent gas accretion (e.g. stellar and active galactic nucleus (AGN) feedback, outflows, the development of hot haloes and virial shocks; e.g. Silk & Rees 1998; Birnboim & Dekel 2003; Croton et al. 2006). The balance reached by these processes determines if a galaxy forms stars throughout its life, or if it becomes a gas-starved ‘red and dead’ object. Understanding these processes is thus vital to further develop theories of galaxy evolution.

One laboratory where we are able to study these processes at work is gas-rich early-type galaxies (ETGs; ellipticals and lenticulars).

ETGs typically have red optical colours, forming a ‘red sequence’ in optical colour-magnitude diagrams (e.g. Baldry et al. 2004). The majority of the stars in current-day ETGs were already in place at $z = 2$ (e.g. Bower, Lucey & Ellis 1992; Thomas et al. 2005), suggesting that some of the physical processes listed above have kept them relatively gas-free ever since.

Nevertheless, a small proportion of ETGs do have cold gas and associated star formation, consistent with a recent regeneration of their gas reservoir (Young et al. 2014). Such low-level activity has been found to contribute upto 30 per cent of the mass of some ETGs since $z = 1$ (e.g. Kaviraj 2014).

A great deal of work has in fact been done in recent years to quantify the amount of gas present in the ETG population. Studies of statistically complete samples have shown that ≈ 25 per cent of ETGs have $> 10^7 M_{\odot}$ of cold molecular gas (e.g. Combes, Young & Bureau 2007; Welch, Sage & Young 2010; Young et al. 2011), while ≈ 40 per cent have sizeable atomic gas reservoirs ($> 10^7 M_{\odot}$) (e.g. Morganti et al. 2006; Sage & Welch 2006; Serra et al. 2012). Dust is also present in a large fraction of ETGs (e.g. Smith et al. 2012), and the presence of gas has been inferred from low-level residual star formation (e.g. Knapp et al. 1989; Yi et al. 2005; Kaviraj et al. 2007; Temi, Brighenti & Mathews 2009; Salim & Rich 2010; Shapiro et al. 2010).

*E-mail: DavisT@cardiff.ac.uk

A quenched object has only two paths by which to rebuild its cold ISM: an internal one (cooling of material lost from stars during their evolution) and an external one (wet mergers, cold accretion, and/or cooling of hot/shocked accreted gas). Stellar mass-loss must of course always be present, but the cold gas detection rate smaller than unity suggests that the majority of stellar ejecta does not cool but rather remains warm or hot (e.g. settling at the virial temperature of the galaxy; e.g. Parriott & Bregman 2008). Various observational studies have suggested that the external channel dominates in field environments. The evidence for this comes from the large fraction of gas-rich ETGs that are morphologically disturbed in deep imaging (e.g. van Dokkum 2005; Duc et al. 2015), the presence of young kinematically decoupled cores in IFU surveys (e.g. McDermid et al. 2006), the orders of magnitude surplus of interstellar medium (ISM) compared to expectations from stellar mass-loss (e.g. Knapp et al. 1989; Merluzzi 1998; Kaviraj et al. 2012; Rowlands et al. 2012), the large spread of gas-to-dust ratios within their cold ISM (Smith et al. 2012; Davis et al. 2015) and the kinematic decoupling of the cold and ionized gas from the stars (Sarzi et al. 2006; Davis et al. 2011b; Katkov, Sil'chenko & Afanasiev 2014).

In this paper, we consider how gas accretion and depletion affect this last piece of observational evidence: the kinematic decoupling of the cold and ionized gas from the stars. In Section 2, we discuss how the observed distribution of angles between the projected angular momenta of the gas and stars can be used to estimate the gas-rich merger rate of ETGs and the time taken for gas to be depleted. In Section 3, we use a simple model to show that there is tension between these estimates and other observations. We discuss how to reconcile the two in Section 4 and conclude briefly in Section 5.

2 ACCRETION AND DEPLETION TIMESCALES

2.1 Origin of the gas

Davis et al. (2011b) studied the kinematic misalignment (the difference between the projected angular momenta) of the molecular gas and stars in the complete, volume-limited ATLAS^{3D} survey of ETGs (Cappellari et al. 2011). Following Sarzi et al. (2006), they assumed that gas that has cooled from stellar mass-loss must always be kinematically aligned with those same stars. On the other hand, material that has been accreted on to a galaxy from external sources can have any angular momentum. Davis et al. (2011b) found that roughly two-thirds of the galaxies are kinematically aligned and are thus *consistent* with an internal gas origin (i.e. stellar mass-loss and/or gas remaining gas from the galaxy formation). One-third of the galaxies are kinematically misaligned and *must* therefore have acquired their gas externally (i.e. cold accretion and/or minor mergers).

Most importantly, as shown in Fig. 1, this behaviour varies drastically as a function of environment. Indeed, there is essentially *no* kinematically misaligned galaxy in the Virgo cluster (the only cluster surveyed by ATLAS^{3D}), implying that external accretion of cold gas is shut off in dense environments. Conversely, at least *half* of the galaxies in the field have acquired their molecular gas externally. The same applies to ionized gas probed through optical emission lines, that can be detected in some cases in objects with ionized gas masses as low as $\approx 10^3 M_\odot$. This results in a higher detection rate overall (61 per cent for ionized gas versus 22 per cent for molecular gas; Sarzi et al. 2006; Young et al. 2011) and thus better number statistics.

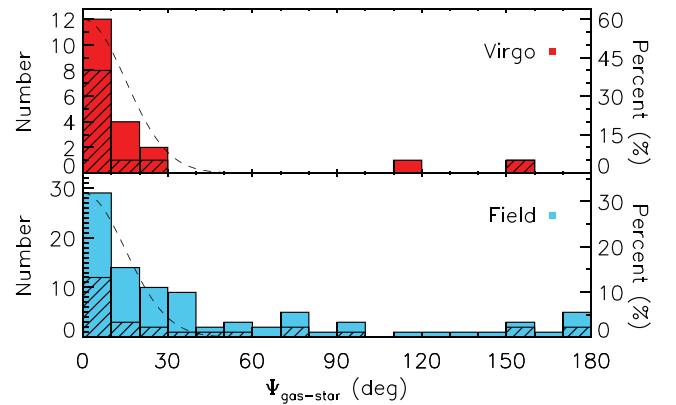


Figure 1. Distribution of kinematic misalignment between gas and stars in the ETGs of the ATLAS^{3D} sample. Ionized gas is shown in solid colours and molecular gas as the shaded histogram. Top: Virgo cluster galaxies. Bottom panel: field and low-density environment galaxies. Adapted from Davis et al. (2011b).

2.2 Depletion time-scale

Although its importance was not recognized at the time, a key feature of the histograms presented in Fig. 1 is the absence of a peak at a kinematic misalignment of 180° for field galaxies. A kinematic misalignment generally traces a morphological misalignment (Krajinović et al. 2011), and structurally-misaligned gas in a fast-rotating ETG (one with a clear disc structure) should be forced back to the equatorial plane over time. This is primarily because the disc gravity provides a substantial torque towards the plane, and because gas dissipation prevents long lived precession of the gas reservoir (Tohline, Simonson & Caldwell 1982). To first order this relaxation should occur symmetrically, with only the initial accretion angle determining if the gas relaxes to be exactly co- or counter-rotating, although stellar mass-loss (corotating by definition) should create a drag on any misaligned material (including that at 180°). The peak of galaxies with kinematically aligned gas (at $|\psi_{\text{gas-stars}}| = 0^\circ$) is thus likely caused by a combination of this relaxation process for externally accreted gas, drag from stellar mass-loss, and gas with an internal origin (i.e. cooled stellar mass-loss itself). By contrast, the bin at 180° should contain only relaxed externally-accreted gas.

Assuming we are not observing local ETGs at a special epoch in their evolution, the absence of a peak at a kinematic misalignment angle $|\psi_{\text{gas-stars}}| = 180^\circ$ in Fig. 1 thus implies that the externally-accreted gas is depleted on a time-scale t_{dep} (equal to the gas survival time-scale t_{surv}) shorter than the torquing time-scale t_{torque} , irrespective of the depletion process.

Previous studies of the relaxation of gas discs in elliptical galaxies were conducted to understand the dynamics of gaseous polar rings. Tohline et al. (1982) and Lake & Norman (1983) considered relaxation in both axisymmetric and triaxial haloes, and found that in either case the alignment process typically takes a few dynamical times (t_{dyn}). Lake & Norman (1983) found that

$$t_{\text{torque}} \approx \frac{t_{\text{dyn}}}{\epsilon}, \quad (1)$$

where ϵ is the ellipticity of the potential. For a typical bulge in a lenticular galaxy $\epsilon \approx 0.2$ (Mendez-Abreu et al. 2008), thus $t_{\text{torque}} \approx 5 t_{\text{dyn}}$.

As shown in Davis et al. (2013), the CO discs of the ATLAS^{3D} ETGs have typical radii $R \approx 1$ kpc, with rotational velocities at the disc edges of $V_{\text{rot}} \approx 200 \text{ km s}^{-1}$. This therefore yields the following

constraint on the accreted molecular gas depletion time-scale:

$$\begin{aligned}
 t_{\text{dep}} = t_{\text{surv}} &\leq t_{\text{torque}} \\
 &\lesssim \frac{t_{\text{dyn}}}{\epsilon} \\
 &\lesssim \frac{2\pi R / V_{\text{rot}}}{\epsilon} \\
 &\lesssim \frac{3 \times 10^7 \text{ yr}}{0.2} \\
 &\lesssim 1.5 \times 10^8 \text{ yr}.
 \end{aligned} \tag{2}$$

The estimate $t_{\text{dep}} \lesssim 10^8 \text{ yr}$ is surprisingly short. Indeed, if the accreted molecular gas were to form stars at the same rate as observed in later-type disc galaxies (i.e. spiral galaxies), it should be consumed on a time-scale of $1\text{--}2 \times 10^9 \text{ yr}$ (e.g. Kennicutt 1998; Bigiel et al. 2011). Furthermore, the star-formation efficiency (SFE) of normal ETGs is *lower* than that of spiral galaxies (e.g. Kawata, Cen & Ho 2007; Saintonge et al. 2012; Martig et al. 2013; Davis et al. 2014), so if anything the depletion time-scale of molecular gas through star formation should be even longer in ETGs. According to the arguments above, the accreted gas observed in local ETGs must therefore be depleted at least 10 times faster than the star-formation consumption time-scale. We discuss this puzzling result in more depth in Section 4.

2.3 Accretion time-scale

Combining the accreted molecular gas depletion (or survival) time-scale with our detection rate ($N_{\text{gas}}/N_{\text{tot}}$) yields a simple estimate of the accretion time-scale t_{acc} , that is the time interval between successive accretion events: $t_{\text{acc}} \approx t_{\text{dep}} / (N_{\text{gas}}/N_{\text{tot}})$.

The overall molecular gas detection rate of 22 per cent (Young et al. 2011) is an overestimate of the true value, as it includes kinematically aligned galaxies to which the current argument does not apply. Only kinematically misaligned gas should be counted. About one-third of all the detected galaxies are kinematically misaligned, but this is an equally misleading (under)estimate, since we know that accretion does not take place in dense environments (Davis et al. 2011b). Indeed, the detection rate of kinematically-misaligned gas in the Virgo cluster is essentially nil, effectively yielding an infinite gas accretion time-scale. The right detection rate to use is therefore that of kinematically-misaligned molecular gas in the field (and low-density environments such as groups). The overall molecular gas detection rate in the field is 29 per cent (Young et al. 2011), and roughly half of this gas is kinematically misaligned (Davis et al. 2011b), yielding a relevant detection rate of ≈ 14 per cent. The molecular gas accretion time-scale is therefore

$$\begin{aligned}
 t_{\text{acc}} &= \frac{t_{\text{dep}}}{(N_{\text{gas}}/N_{\text{tot}})|_{\text{field}}} \\
 &\lesssim \frac{t_{\text{dyn}} / \epsilon}{(N_{\text{gas}}/N_{\text{tot}})|_{\text{field}}} \\
 &\lesssim \frac{1.5 \times 10^8 \text{ yr}}{0.14} \\
 &\lesssim 10^9 \text{ yr}.
 \end{aligned} \tag{3}$$

Our estimate therefore suggests that local ETGs in the field accrete molecular gas from an external source at least once every gigayear ($R_{\text{acc}} \gtrsim 1 \text{ Gyr}^{-1}$). We discuss this estimate further, and compare it with those of other studies, in Section 4.

Interestingly, the overall detection rate of diffuse ionized gas emission (H β and/or [O III] emission lines) in the ATLAS^{3D} sample is 61 per cent (Sarzi et al., in preparation). As Davis et al. (2011b) showed that the ionized and molecular gas are always kinematically aligned and therefore must share a common origin, we can in principle replace the molecular gas detection rate used in equation (3) by that for the ionized gas. The ionized gas detection rate for field galaxies in ATLAS^{3D} is 68 per cent, with roughly half kinematically misaligned, so the relevant detection rate is 34 per cent, roughly 2.5 times higher than that used above for the molecular gas. This in turn yields an upper limit on the ionized gas accretion time-scale for local field ETGs that is 2.5 times smaller, only $\lesssim 4 \times 10^8 \text{ yr}$.

The relation between the gas phases in objects without detectable cold gas is, however, unclear. In many objects the ionized gas is likely in an interface layer (ionized by old stars, AGN or low-level star formation) around an atomic/molecular gas disc that has a low enough mass to fall below our sensitivity limit. This interpretation is supported by the similarity in the misalignment histograms (Fig. 1). It is possible that this emission could be more nebulous, however, arising from the envelopes around evolved stars, and thus not be in a disc at all in some objects (in this case, however, the ionized gas emission would have to corotate with the stars). Given the above uncertainty the estimates of Tohline et al. (1982) and Lake & Norman (1983) may not apply (or, equivalently, the scaling factor may be different), and the ionized gas relaxation time may be much longer than that for molecular gas. This would then yield larger estimates for both t_{dep} and t_{acc} . Given these uncertainties, in the rest of this paper we favour the more reliable molecular gas estimates to put other time-scales in context.

3 ACCRETION MODELS

3.1 Toy model

To validate and explore further the simple calculations above, we created a toy model. This model includes in a simple framework the processes of gas accretion, relaxation and depletion. These are characterized by three free parameters, respectively, the accretion rate (R_{acc} , the number of gas accretion events per unit time), the relaxation time-scale (t_{relax} , in units of t_{dyn}), and the gas depletion time (t_{dep}).

We initialised the model at time $t = 0$ with an arbitrary number of gas-free, ‘red and dead’ ETGs. These have realistic luminosities drawn uniformly from a Schechter function representation of the ATLAS^{3D} K -band galaxy luminosity function (Cappellari et al. 2011). From this we calculate the maximum circular velocity at which the gas can rotate by assuming the galaxies follow the K -band CO Tully–Fisher relation of Davis et al. (2011a). Molecular gas reaches beyond the turnover of the rotation curve in ≈ 70 per cent of ATLAS^{3D} ETGs (Davis et al. 2011a, 2013), so this should provide a good estimate of the rotation velocity of the gas at the edges of the gas discs for the bulk of the population.

After each timestep we randomly select galaxies that will accrete gas, such that the average number of mergers per galaxy per unit time is equal to R_{acc} . In this simple model, R_{acc} is independent of galaxy mass and morphology, which may not be true in reality. We discuss this assumption further in Section 4.3.

The amount of gas accreted is chosen randomly from the H_2 mass function of Young et al. (2011), that was fitted with a Gaussian at the high-mass end, and was assumed constant below the completeness limit down to an H_2 mass of $10^6 M_{\odot}$ (see Fig. 2), as predicted

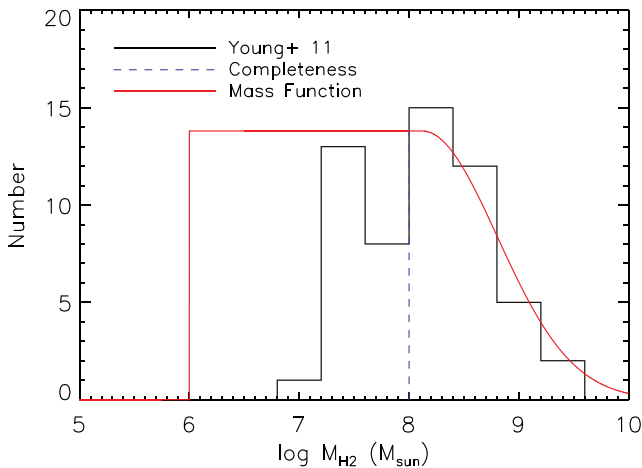


Figure 2. Observed and modelled H_2 mass functions. The measured distribution of H_2 masses from Young et al. (2011) is shown in black, with the approximate completeness limit indicated by a dashed line. The (arbitrarily normalized) H_2 mass function used in the toy models is shown in red (see Section 3.1).

from semi-analytic models (e.g. Lagos et al. 2014). Removing the low-mass cutoff does not affect our results. We here assume that after coalescence each one of our mergers results in the creation of a disc of H_2 in the centre of the remnant, with a mass that is drawn from the observed distribution. For this simple model, we condense the complex time-dependent processes happening in mergers, and assume that this final state is obtained instantaneously. We discuss these assumptions further in Section 4.3.

We determine the radial extent of the gas, and hence the dynamical time at its outer edge, by drawing uniformly from a lognormal fit to the observed distribution of molecular disc radii in Davis et al. (2013).

The angle at which gas is accreted on to galaxies is also chosen uniformly, sampling the full range 0 – 180° . If gas already exists in a galaxy, the interaction between the two reservoirs is roughly included by mass-weighting the resulting position angle of the gas (thus lying between the newly accreted and old material). Our assumption about the exact form of this weighting does not strongly affect our results (see Section 4.3).

The relaxation of the gas is tracked at each timestep. The gas relaxes at a rate that has a cosine dependence on the current misalignment angle (i.e. the rate increases as the disc approaches the galaxy plane, where the torque is higher), following the formalism of Tohline et al. (1982). The total time taken to relax is calculated such that the gas disc relaxes into the plane from a misalignment of 60° in a time t_{relax} . We note that adopting a simpler form, where the gas relaxes linearly (i.e. by a full $\pi/2$ radians in a time t_{relax}) does not affect the results of this paper. More complex prescriptions for the gas relaxation are beyond the scope of our analysis, but we expect any reasonable form to not unduly alter the results derived from these two cases.

Finally, we destroy a fraction of the molecular gas at each time step. In the first set of simulations discussed below, we choose to include only the effect of star formation on the molecular gas, using up ≈ 10 percent of the gas per dynamical time, equivalent to the standard star-formation law derived from local star-forming galaxies (Kennicutt 1998). In Section 3.2.3, we also include a very simple model of gas destruction (such as that caused by AGN feedback). An additional free parameter (t_{AGN}) then sets the duty cycle over

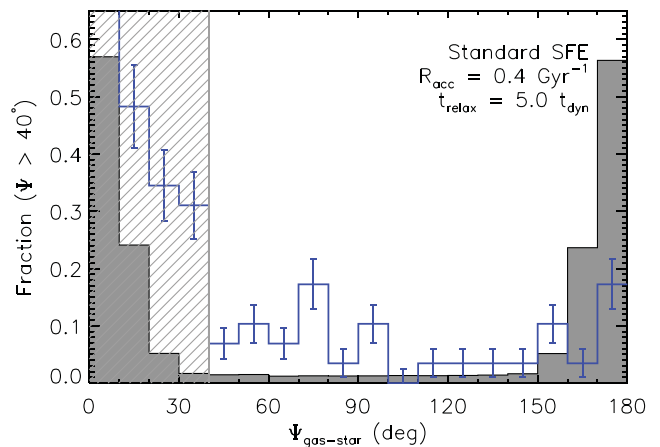


Figure 3. Kinematic misalignment histogram of our reference model, showing the steady state kinematic misalignment distribution (grey histogram) of a model where the gas is depleted only by star formation with a standard efficiency, and where the gas relaxes in five dynamical times. The accretion rate required to match the observed detection rate is then 0.2 Gyr^{-1} . The parameters used are listed in the figure legend. The ATLAS^{3D} data are shown as a blue histogram, and should not be compared with the model in the shaded region below 40° , where other processes that are not modelled here may dominate (e.g. stellar mass-loss), thus increasing the number of detected objects. Both histograms are normalized so the area with $\Psi > 40^\circ$ is equal to unity.

which gas is destroyed. All the gas is instantaneously blown out of the model galaxy a time t_{AGN} after it has been accreted. Clearly, this highly simple model is not meant to capture in detail the highly complex AGN feedback process, but it at least allows us to explore in a simple manner its effect. Any process that removes gas in one discrete event can be modelled by this formalism.

At the end of each timestep, we are thus left with the observable gas mass and kinematic position angle of each galaxy. We convolve the kinematic misalignment angles with a Gaussian kernel to match the observational errors, and set a limit of $10^7 M_\odot$ to define objects that are detectable, this to match the observations of the ATLAS^{3D} survey (see Fig. 1; Young et al. 2011). Each model presented in Section 3.2 includes 10^5 simulated galaxies and was run for 5 Gyr, although the misalignment distribution histograms reach a steady state more quickly.

3.2 Model results

3.2.1 Reference model

In Fig. 3, we present the misalignment histogram resulting from our reference model, based on our best current understanding of the processes involved. This model uses only the standard star-formation law to deplete the gas, a relaxation time of $5 t_{\text{dyn}}$ (as derived from the results of Lake & Norman 1983), and an accretion rate of 0.4 Gyr^{-1} (calculated so the detection rate of kinematically-misaligned objects in the model matches the observed rate). This accretion rate is consistent with observations of the (minor and minor+major) merger rate estimated from other methods (see Section 4.1).

One should not compare the misalignment histogram produced by the model to the data within the shaded region in Fig. 3 ($\Psi < 40^\circ$), as material produced by internal processes not included in our model (e.g. stellar mass-loss) will contribute (and may dominate) here. Outside this region, clear discrepancies are immediately obvious

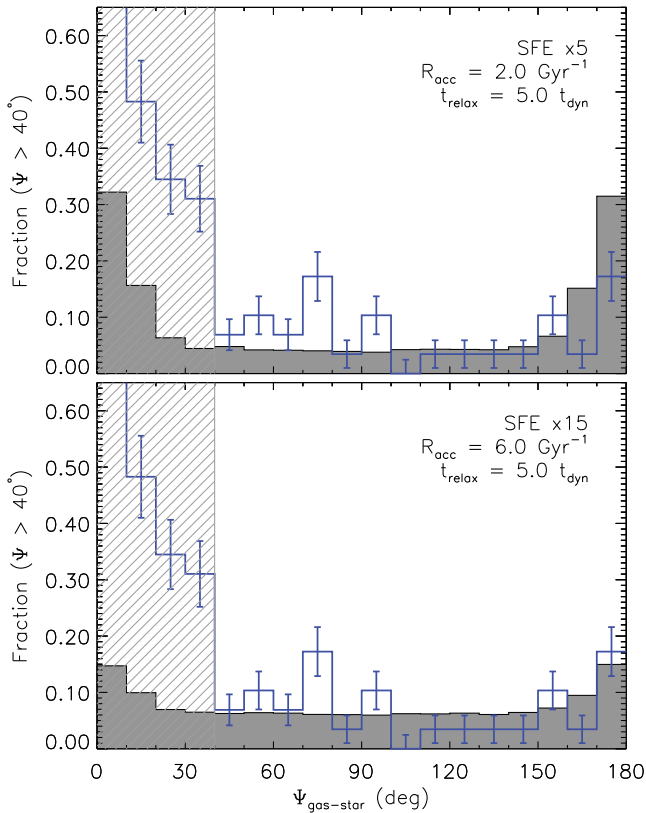


Figure 4. As Fig. 3, but for an enhanced SFE of five (top panel) and 15 (bottom panel) times the standard value. The accretion rates required to match the observed detection rate are then 2 and 6 Gyr^{-1} , respectively.

when comparing the reference model histogram to the observed one (Fig. 3). First, almost all model galaxies (≈ 90 percent) lie in the exactly kinematically aligned and misaligned bins (at 0 and 180° , respectively). Secondly, the bins for kinematically aligned and misaligned gas contain approximately equal numbers of model galaxies. To produce an asymmetric histogram, one therefore clearly requires another process that preferentially produces/preserves gas that is always kinematically aligned with the stars (or, conversely, preferentially removes gas that is kinematically misaligned). The reference model thus highlights graphically the points made in Section 2.2, that the standard relaxation and gas depletion time-scales cannot reproduce the observed misalignment histograms. In the sections below, we thus consider possible solutions to this problem.

3.2.2 Quicker gas depletion

As argued in Section 2.2, a faster gas depletion time may be required to understand the observed gas kinematic misalignment histograms. We show in Fig. 4 a model where the SFE has been increased by a factor of 5 (top panel) and 15 (bottom panel) with respect to its standard value. This increase can be used to model any process that removes gas over time, not just star formation. Clearly, an order of magnitude increase in the depletion rate is required, over and above that expected from normal star formation, before the distribution of misalignments in the model can match that of the observations.

One drawback of this solution is that the accretion rate required to match the observed detection rate is inversely proportional to the gas depletion time, that is shorter depletion times require proportionally

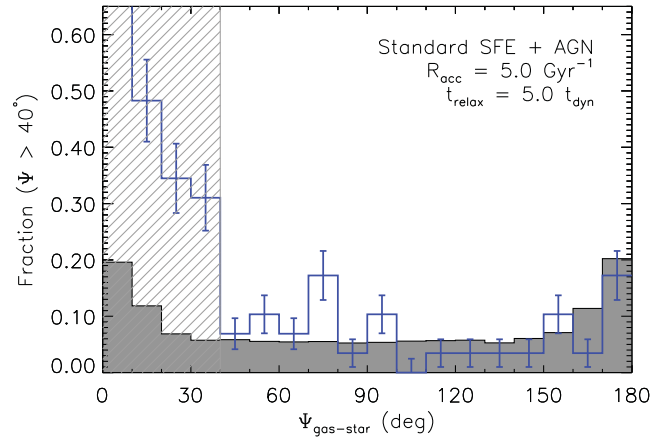


Figure 5. Same as Fig. 3, but with our simple gas destruction prescription enabled (approximating AGN feedback). The assumed AGN duty cycle is 5×10^7 yr. The accretion rate required to match the observed detection rate is then 5 Gyr^{-1} .

more accretion events per unit time to maintain the same detection rate. In these models with increased SFEs, one thus requires a merger rate of $\approx 6 \text{ Gyr}^{-1}$ to match the observed detection rate. We discuss this hypothesis further in Section 4.1.

3.2.3 Gas destruction

In Section 2.2, we discussed the need to decrease the survival time of gas in ETGs. One mechanism that has become popular in recent years to quench ETGs is AGN feedback (e.g. Silk & Rees 1998; Fabian 2012). In this paradigm, the central black hole injects a significant amount of energy into the ISM, removing or destroying it. As described in Section 3.1, we can include a simple feedback prescription in our model, that blows all of the gas out of the galaxies a time t_{AGN} after the gas has been accreted.

We show the result of turning on this simple AGN feedback prescription in Fig. 5. We use the same standard SFE and relaxation time-scale as the reference model (Section 3.2.1), and additionally set $t_{\text{AGN}} = 10^8$ yr. This is at the upper end of the range 10^6 – 10^8 yr that has been suggested in the literature for the typical AGN duty cycle (e.g. Soker et al. 2001; Mazzotta et al. 2002; Hopkins et al. 2006). We again adjust the accretion rate to match the observed detection rate of molecular gas. The AGN feedback prescription does lower the fraction of galaxies within the co- and counter-rotating bins sufficiently to match the observed predictions. However, it also requires a very high accretion rate (approximately 10 times higher than required in the reference model) to match the observed CO detection rate of kinematically-misaligned objects.

3.2.4 Slower relaxation

An alternative explanation for the large fraction of model galaxies whose gas reservoirs have yet to relax is that this process is slower than usually assumed. Fig. 6 shows model runs assuming a standard SFE for depleting the gas and a standard accretion rate of 0.4 Gyr^{-1} (so the detection rate in the model matches the observed rate by construction, as in Section 3.2.1). The relaxation time is however increased to 20 and $80 t_{\text{dyn}}$, respectively, in the top and bottom panels. These models show that the relaxation time would have to be much longer than normally assumed (by around an order

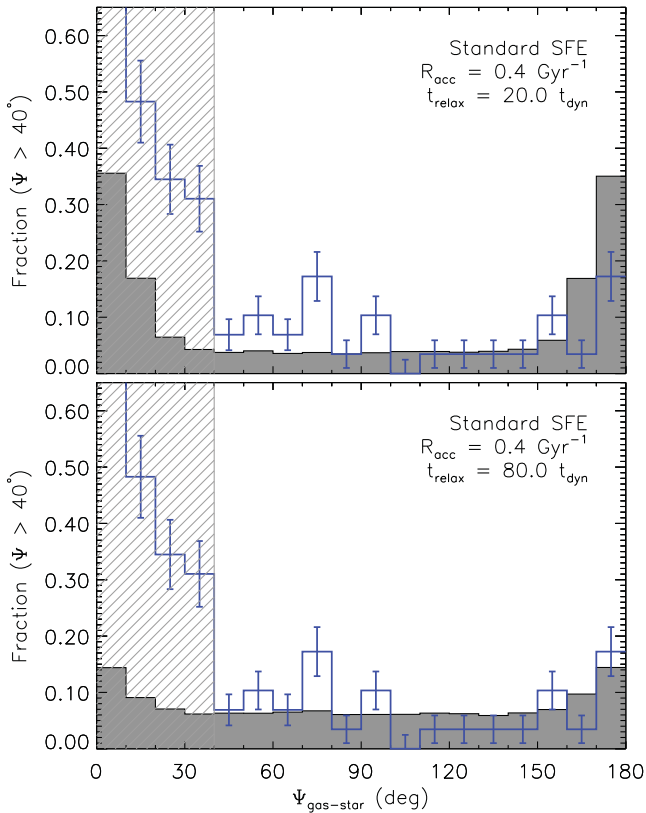


Figure 6. Same as Fig. 3, but for a slower gas relaxation of 20 (top panel) and 150 (bottom panel) dynamical times.

of magnitude) to reproduce the observed lack of counter-rotating objects. This solution has the benefit that it removes the need for an extremely large accretion rate, but it requires additional processes to slow gas relaxation. We discuss this further in Section 4.2.

4 DISCUSSION

In the sections above, we discussed the distribution of kinematic misalignments between the gas and stars of local ETGs, and the way in which it encodes information about the depletion and accretion time-scales of the gas. We showed that the standard gas relaxation and depletion time-scales from the literature cannot reproduce the observed misalignment histogram. We discuss below possible solutions to this problem.

4.1 Gas depletion and destruction

In Section 3.2.2, we showed that one can reproduce the observed kinematic misalignment histograms from ATLAS^{3D} by increasing the rate at which gas is consumed. According to this argument, the accreted gas observed in local ETGs would have to be depleted at least 10 times faster than the standard star-formation consumption time-scale. Extreme starbursts have been observed to have such small depletion times (e.g. Mihos & Hernquist 1994), but studies of ETGs (e.g. Saintonge et al. 2012; Martig et al. 2013; Davis et al. 2014) show that their SFE is actually slightly lower than in spiral galaxies (and therefore very much lower than in starburst objects). Fast starbursts thus cannot be invoked to explain the short depletion times required here. Other potential physical mechanisms

that could remove gas include dissociation of H₂ by X-rays from the ETGs' hot gas haloes, friction and/or shocks against kinematically-aligned gas from stellar mass-loss of the bulk (corotating) galaxy stellar populations, and AGN feedback. Whatever the depletion mechanism is, it needs to be very rapid.

The depletion time-scale of $\lesssim 10^8$ yr we estimated in Section 2.2 is the same as the typical time-scale over which central supermassive black holes are active (e.g. Soker et al. 2001; Mazzotta et al. 2002; Hopkins et al. 2006). The energy input from an AGN is observed to be expelling the gas from at least some ETGs (e.g. NGC1266; Alatalo et al. 2011; Davis et al. 2012), hence we tried to include them in our toy models in a simple way (see Section 3.2.3). We then find that a duty cycle of $\lesssim 10^8$ yr is sufficient to bring the model kinematic misalignments in line with the observed data.

The drawback of both fast depletion and/or destruction of gas is that the accretion rate required to match the observed detection rate is inversely proportional to the gas depletion time. With such small gas depletion/destruction time-scales, a merger rate of ≈ 5 Gyr⁻¹ is required to maintain the detection rate as observed. These time-scales can also be used to constrain the likely mechanism(s) supplying the gas. They rule out major mergers (that are too rare), but as the accretion time-scale becomes ever smaller, the case for nearly continuous accretion becomes stronger, and accretion of diffuse gas (rather than galaxies) along filaments may need to be considered. However, the large fraction of gas-rich ETGs that are morphologically disturbed in deep imaging (e.g. van Dokkum 2005; Duc et al. 2015) suggests that stars are also accreted along with the gas, leading to the idea that minor mergers may dominate over cold mode accretion (see e.g. Kaviraj et al. 2009, 2011).

Interestingly, there is evidence of two separate recent minor mergers within such a short time interval in at least one ATLAS^{3D} ETG galaxy, NGC4150. Indeed, Kaviraj et al. (in preparation) have shown that although this galaxy has a young ($\approx 10^8$ yr) counter-rotating stellar core, presumably formed in a gas-rich minor merger, the molecular gas currently present is corotating with the bulk of the stars, suggesting an even more recent accretion event.

In the current Λ CDM cosmology, the universe is (primarily) hierarchical in nature, making galaxy mergers an essential aspect of galaxy formation and evolution. The gas-rich merger rate at $z = 0$ is thus an interesting quantity to determine from the gas misalignment data, and to compare with that measured from other observations and simulations (e.g. Lacey & Cole 1993; Maller et al. 2006; Guo & White 2008; Genel et al. 2009; Lotz et al. 2011; Bluck et al. 2012; Newman et al. 2012; Casteels et al. 2014; Rodriguez-Gomez et al. 2015).

The rate of gas-rich mergers derived here is somewhat higher than previous estimates of the incidence of minor mergers. For instance, Lotz et al. (2011) presented a range of estimates, and compared it extensively with other samples. They estimated a total merger rate of 0.2 to 0.5 Gyr⁻¹ with 75 percent of these mergers being minor (mass ratios of 1:4 to 1:10). Theoretically, studies such as that of Hopkins et al. (2010) find a major+minor merger rate of ≈ 0.2 Gyr⁻¹ for objects of $10^{11} M_{\odot}$, but predict that this merger rate scales positively with galaxy mass. The more massive galaxies in our sample (with stellar masses above $10^{12} M_{\odot}$) are predicted to have merger rates of up to 0.9 Gyr⁻¹, still below our estimate. To make things worse, the estimates discussed above count all mergers, not only those that are gas rich.

We thus conclude that although accretion rates as fast as those require here may be possible for individual objects, it is unlikely that the average accretion rate in the universe is as high as our models require.

4.2 Slow gas relaxation

Another way to reconcile the observed kinematic misalignments of ETGs with the theoretical predictions is to adopt a relaxation (i.e. torquing) time-scale much longer than a few dynamical times. This explanation is appealing, as it does not require high gas accretion rates, but the possible physical processes slowing the relaxation down are currently unconstrained. One possibility is that the accreted misaligned gas simply precesses. This would be possible if the accreted molecular gas were moving purely ballistically, as discrete molecular clouds can do, but it seems unlikely for the complex multiphase ISM characteristic of galaxies, and especially for the warmer more diffuse ionized gas.

If gas is accreted on to ETGs over long time-scales, then the continued addition of misaligned material could slow the relaxation process. This could happen, for instance, if a gas-rich merger throws out long tidal tails of gas, that fall back slowly over time (e.g. Barnes 2002). Atomic gas tails are observed in a few objects (e.g. Morganti et al. 2006; Serra et al. 2012), but not in the entire population of misaligned ETGs. Low-density tails, or accretion of debris in hotter gas phases, could however be missed. van de Voort et al. (2015) looked at this process in a hydrodynamic simulation of a massive early-type galaxy. They found that continued accretion did indeed slow relaxation, as the angular momentum change induced by accreted gas dominates over that induced by stellar torques. The total time taken for the disc to settle in these simulations was $80\text{--}100\,t_{\text{dyn}}$, very similar to the time-scale we estimate in this work.

Another option is that the hot haloes around galaxies have an effect. Lagos et al. (2014); Lagos et al. (2015) found that the hot haloes around ETGs may be misaligned (either because of mergers or continuous accretion from the cosmic web), and that this could be the cause of the observed kinematic misalignment distribution. These models make various assumptions that still need to be verified through hydrodynamic simulations, and they did not include any relaxation process, but it is clear that continued misaligned cooling could extend relaxation times. In addition, if the hot ISM of ETGs has a net angular momentum and can exert pressure on the relaxing accreted gas, then gas relaxation would not be a symmetric process, which would help explain the observed asymmetric kinematic misalignment distribution. These solutions, however, require the hot haloes of galaxies to have significant angular momenta. It will not be possible to observationally verify whether the hot haloes of galaxies have any significant rotation, and/or if it correlates with the rotation of the stellar body, until the X-ray satellite *ASTRO-H* launches (see e.g. Bianconi, Ettori & Nipoti 2013).

4.3 Model limitations and associated uncertainties

The toy model we present in this paper makes various assumptions, that could potentially affect the results we derive, and are thus discussed below.

We know from previous studies that the gas is not sufficiently extended to reach the flat part of the rotation curve in ≈ 30 per cent of ETGs (Davis et al. 2011a, 2013). In these cases our model will overestimate the gas rotation velocity, and thus underestimate the dynamical time, relaxing the gas too quickly. However, this always happens in gas reservoirs with small radial extents, and thus the shortest dynamical times. Even an underestimation of the dynamical time by a factor of 2 in these objects (the largest difference between the CO velocity and the circular velocity in objects with small gas reservoirs; Davis et al. 2011a) would not help reconcile the reference model and the observations.

The second assumption we make is that R_{acc} , and the amount of gas accreted in each episode, is independent of galaxy mass and morphology. While this may be a reasonable first order approximation, it is known that ETGs follow a morphology-density relation (e.g. Dressler 1980), and interactions preferably happen in group scale environments. We minimize the effect of any such bias here by only comparing the model to observed galaxies in field environments. In addition, the typical merger mass-ratio must vary as a function of galaxy mass, which will affect the typical amount of gas brought in by each encounter. However, as the typical gas fraction in observed ETGs does not seem to depend on galaxy mass or environment (Young et al. 2011) it is likely that this second effect is small, and would not change our results significantly.

Our approach assumes that all of the gas brought in by the merger is molecular. We do not model additional H I gas, that may be brought in by the merger and settle in a larger scale disc (e.g. Serra et al. 2012), and could convert to H₂ over longer time-scales. If significant radial mass/angular momentum exchange occurs, these larger scale H I discs could slow relaxation similarly to the atomic gas tidal tails discussed above.

When gas is already present in a galaxy that experiences a new accretion event, we mass weight the misalignment angle of the new and old gas components to determine the resultant kinematic misalignment. In a real gas-rich merger complex processes take place that we are unable to capture in such a simple model, but we do not expect this simplification to have a major effect. Within our model, when this situation occurs the median amount of accreted gas is an order of magnitude larger than the gas already present, and thus the correction typically changes the accreted angle by only a small amount (a median change of $\approx 4^\circ$). Ignoring this correction entirely, and instead forcing all the gas to rotate with the newly accreted material does not change our result.

We also assume that the radial extent of the gas in these objects is uncorrelated with its mass (and other galaxy properties). We expect this assumption to not unduly affect our results, as although a strong relation between the gas mass and size of atomic gas discs exists (e.g. Broeils & Rhee 1997), any such relation for molecular gas is weak (Davis et al. 2013). We do, however, caution that when gas is accreted on to a galaxy with an existing cold gas reservoir the interaction between the two discs could systematically affect the final radial extent of the settled gas disc.

In our simple model the SFE is a free parameter, but it does not depend on galaxy properties. Gas-rich mergers in the real universe seem to have increased SFEs, that result in faster gas depletion. These events are short-lived, however, and observed ETGs are found to have low SFEs (e.g. Saintonge et al. 2011; Davis et al. 2014). Evidence also exists that gas minor mergers with early-type galaxies can further suppress the SFE (Davis et al. 2015). We thus do not expect merger induced starbursts to be important in these sources.

Overall the limitations discussed above are mostly expected to be second-order effects. Hydrodynamic simulations are required to confirm this, and they will be the topic of a future work.

5 CONCLUSIONS

In this paper, we considered what can be learnt about the processes of gas accretion and depletion from the observed distribution of kinematic misalignment angles between the cold/ionized gas and the stars.

We first presented simple arguments showing that the misalignment distribution encodes information on the relaxation, depletion and accretion time-scales of gas in ETGs. Specifically, we argued

that the lack of a peak of exactly counter-rotating objects strongly constrains these time-scales. Simple calculations, based only on a rough estimate of the relaxation (i.e. torquing) time-scale, imply a short gas depletion time-scale ($t_{\text{dep}} \lesssim 10^8$ yr) and a high rate of gas-rich mergers ($\gtrsim 1 \text{ Gyr}^{-1}$).

We then presented a toy model of the interplay between these processes, allowing us to better constrain and explore the effects of these time-scales on the kinematic misalignment distribution. Realistic distributions of stellar masses and gas radii allow to derive estimates of the relevant time-scales for the whole ETG population. We thus confirmed our simple calculations, clearly showing that the standard values for the accretion rate, SFE and relaxation time are not simultaneously consistent with the observed distribution of kinematic misalignments.

We then explored the effect of varying these parameters. Both faster gas depletion (via e.g. more efficient star formation) and/or faster gas destruction (via e.g. AGN feedback) can be used to explain the kinematic misalignment distribution, but they then require high rates of gas-rich mergers ($\approx 5 \text{ Gyr}^{-1}$). Although some objects have evidence of multiple mergers within such a time-scale, as a population accretion rates this high are unlikely.

An alternative explanation, which does not require high accretion rates, is that the misaligned gas relaxation occurs over longer time-scales ($\simeq 100 t_{\text{dyn}}$ or $\approx 1\text{--}5 \text{ Gyr}$) than usually assumed. We suggest that this could come about because of ongoing accretion of tidal debris in mergers, rotating hot gas haloes, or continued accretion from the cosmic web.

One way in which these scenarios could be tested would be to determine the age of the accretion event in other ways, such as age dating the most recent burst of star formation (via stellar population modelling), or with the ratio of gas phase metallicity to stellar metallicity. Models with fast relaxation require the misaligned objects to have accreted gas recently (typically within the last 200 Myr, but up to $\approx 1 \text{ Gyr}$ for exactly counter-rotating gas). Slow relaxation removes such differences, predicting a flat distribution of accretion times as a function of misalignment.

Of course, all of the processes discussed above may take place simultaneously in galaxies and cannot easily be separated. Many combinations of the parameters in our models can match the observed misalignment distribution. Further work, both observationally and numerically, is required to properly understand the gas misalignment in ETGs. Such work is also timely, since if we can understand the typical time-scale over which misalignments are visible, new large integral-field spectroscopic surveys of galaxies (e.g. the Calar Alto Legacy Integral Field spectroscopy Area survey (CALIFA); García-Benito et al. 2015, the Sydney-AAO Multi-object Integral field spectrograph (SAMI) galaxy survey, Bryant et al. 2015; the Mapping Nearby Galaxies at APO (MaNGA) survey, Bundy et al. 2015) will have the power to investigate these effects as a function of galaxy mass, environment, etc., and thus to accurately constrain the (gas-rich) merger rate in the local universe.

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