



Low-frequency Carbon Radio Recombination Lines. II. The Diffuse Interstellar Medium

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Abstract

In the second paper of the series, we have modeled low-frequency carbon radio recombination lines (CRRLs) from the interstellar medium. Anticipating the Low Frequency Array survey of Galactic CRRLs, we focus our study on the physical conditions of the diffuse, cold neutral medium. We have used the improved departure coefficients computed in the first paper of the series to calculate line-to-continuum ratios. The results show that the line width and integrated optical depths of CRRLs are sensitive probes of the electron density, gas temperature, and emission measure of the cloud. Furthermore, the ratio of CRRL to the [C II] at the 158 μm line is a strong function of the temperature and density of diffuse clouds. Guided by our calculations, we analyze CRRL observations and illustrate their use with data from the literature.

Key words: atomic processes – line: formation – methods: numerical – radiative transfer – radio lines: general – radio lines: ISM

1. Introduction

The interstellar medium (ISM) plays a central role in the evolution of galaxies. The formation of new stars slowly consumes the ISM, locking it up for millions to billions of years while stars, as they age, return much of their mass increase in metallicity back to the ISM. Stars also inject radiative and kinetic energy into the ISM, and this controls the physical characteristics (density, temperature, and pressure) as well as the dynamics of the gas, as revealed in observed spectra. This interplay of stars and surrounding gas leads to the presence of distinct phases (e.g., Field et al. 1969; McKee & Ostriker 1977). Diffuse atomic clouds (the cold neutral medium, CNM) have densities of about 50 cm^{-3} and temperatures of about 80 K, where atomic hydrogen is largely neutral but carbon is singly ionized by photons with energies between 11.2 and 13.6 eV. The warmer (~ 8000 K) and more tenuous ($\sim 0.5\text{ cm}^{-3}$) intercloud phase (the warm neutral medium, WNM, and warm ionized medium, WIM) is heated and ionized by far ultraviolet (FUV) and extreme ultraviolet (EUV) photons escaping from H II regions (Wolfire et al. 2003). While these phases are often considered to be in thermal equilibrium and in pressure balance, the observed large turbulent width and presence of gas at thermally unstable, intermediate temperatures may indicate that kinetic energy input is important. Thermally unstable gas could indicate that the gas does not have sufficient time to cool between subsequent passages of a shock or after intermittent dissipation of turbulence (e.g., Kim et al. 2011). In addition, the ISM also hosts molecular clouds, where hydrogen is in the form of H_2 and self-gravity plays an important role. All of these phases are directly tied to key questions on the origin and evolution of the ISM, including the energetics of the CNM, WNM, and the WIM; the evolutionary relationship of atomic and molecular gas; the relationship of these ISM phases with newly formed stars; and the conversion of their radiative and kinetic power into thermal and turbulent energy of the ISM (e.g., Elmegreen

& Scalo 2004; Scalo & Elmegreen 2004; Cox 2005; McKee & Ostriker 2007).

The diffuse ISM has been long studied using, in particular, the 21 cm hyperfine transition of neutral atomic hydrogen (e.g., Kulkarni & Heiles 1987, p. 87; Heiles & Troland 2003a). These observations have revealed the prevalence of a two-phase structure in the ISM of cold clouds embedded in a warm intercloud medium. However, it has been notoriously difficult to determine the physical characteristics (density, temperature) of these structures in the ISM because H I by itself does not provide a good probe. Optical and UV observations of atomic lines can provide the physical conditions but are by necessity limited to pinpoint experiments toward bright background sources. However, with the opening up of the low-frequency radio sky with modern interferometers such as the Low Frequency Array for Radioastronomy (LOFAR; van Haarlem et al. 2013), the Murchison Wide Field Array (Tingay et al. 2013), the Long Wavelength Array (Ellingson et al. 2013), and, in the future, the Square Kilometer Array (SKA), systematic surveys of low-frequency ($\nu \lesssim 300$ MHz) carbon radio recombination lines (CRRLs) have come into reach, and these surveys can be expected to quantitatively measure the conditions in the emitting gas (Oonk et al. 2015; Oonk et al. 2017).

Carbon has a lower ionization potential (11.2 eV) than hydrogen and can be ionized by radiation fields in regions where hydrogen is largely neutral. Recombination of carbon ions with electrons to high Rydberg states will lead to CRRLs in the submillimeter to decameter range. CRRLs have been observed in the ISM of our Galaxy toward two types of clouds: diffuse clouds (e.g., Konovalenko & Sodin 1981; Erickson et al. 1995; Roshi et al. 2002; Stepkin et al. 2007; Oonk et al. 2014) and photodissociation regions (PDRs), the boundaries of H II regions and their parent molecular clouds (e.g., Natta et al. 1994; Wyrowski et al. 1997; Quireza et al. 2006). Recently, Morabito et al. (2014) discovered extragalactic CRRLs associated with the nucleus of the nearby starburst galaxy M82. Theoretical models for CRRLs were first developed by

Watson et al. (1980) and Walmsley & Watson (1982), including the effects of dielectronic recombination³ with the simultaneous excitation of the $^2P_{3/2}$ fine structure level and later extended by Ponomarev & Sorochenko (1992) and by Payne et al. (1994). However, these studies were hampered by the limited computer resources available at that time.

In the coming years, we will use LOFAR to carry out a full northern hemisphere survey of CRRL-emitting clouds in the Milky Way. This will allow us to study the thermal balance, chemical enrichment, and ionization rate of the CNM from degree scales down to scales corresponding to individual clouds and filaments in our Galaxy. Furthermore, following the first detection of low-frequency CRRLs in an extragalactic source (M82; Morabito et al. 2014) we will also use LOFAR to perform the first flux-limited survey of CRRLs in extragalactic sources. Given the renewed observational interest in CRRLs, a new theoretical effort seems warranted. In the first paper of this series (Salgado et al. 2017, hereafter Paper I), we studied the level population of hydrogenic atoms including the effects of dielectronic recombination in carbon atoms. The level population of atoms, however, is not the only process that influences the strength of an observed line, as radiative transfer effects can alter the strength or depth of an observed line. In this paper, we use the results of Paper I to develop CRRLs as a tool to derive the physical conditions in the emitting gas. In this, we will focus on cold, diffuse clouds as these are expected to dominate the low-frequency CRRL sky. The paper is organized as follows. In Section 2.1 we review radiative transfer theory in the context of radio recombination lines. We review the line-broadening mechanisms of CRRLs in Section 2.3. In Section 3, we present the results of our models and compare them with observations from the literature and provide guidelines to analyze such observations. Finally, in Section 6, we summarize our results and provide the conclusions of our work.

2. Theory

2.1. Radiative Transfer of CRRLs

The physical conditions of the diffuse ISM (temperatures of $T_e \approx 100$ K and electron densities $N_e \approx 10^{-2} \text{ cm}^{-3}$) favor an increase in level population at high quantum levels via dielectronic recombination (Paper I). Moreover, the presence of an external radiation field can also alter the level population of carbon atoms. In addition, while low-frequency CRRLs are observed in absorption against a background continuum (e.g., Kantharia & Anantharamaiah 2001; Morabito et al. 2014; Oonk et al. 2014), high-frequency recombination lines are observed in emission. Therefore, radiative transfer effects must be analyzed in order to derive meaningful physical parameters from observations.

We begin our analysis by revisiting the radiative transfer problem in the context of CRRLs. At a given frequency, the observed emission has two components, corresponding to the line transition itself and the underlying continuum emission. In Appendix B, we summarize the standard general solution to the one-dimensional radiative transfer equation of a line in a homogeneous medium. Here, we show the result for a cloud at

a constant temperature T_e ⁴:

$$\frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} = \frac{\eta B_\nu(T_e)(1 - e^{-\tau_\nu^{\text{total}}}) + I_0(\nu)e^{-\tau_\nu^{\text{total}}}}{B_\nu(T_e)(1 - e^{-\tau_\nu^{\text{c}}}) + I_0(\nu)e^{-\tau_\nu^{\text{c}}}} - 1, \quad (1)$$

where I_ν^{line} is the intensity of a line at a frequency ν , I_ν^{cont} is the intensity of the continuum, η is a correction factor to the Planck function due to non-LTE effects (as defined in Strelitski et al. 1996; Gordon & Sorochenko 2009, see Appendix B), $B_\nu(T_e)$ is the Planck function, τ_ν^{total} is the sum of the line and continuum optical depth (τ_ν^{l} and τ_ν^{c} , respectively), and $I_0(\nu)$ is the intensity of a background continuum source at the frequency of the line.⁵

In the presence of a strong background radiation field, as is the case for low-frequency lines in the diffuse ISM ($I_0 \gg \eta B_\nu(T_e)$, see below), the background term (I_0) dominates, and the first term in the numerator and denominator on the right-hand side of Equation (1) can be ignored, and this equation simplifies to

$$\frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} = e^{-\tau_\nu^{\text{l}}} - 1, \quad (2)$$

independent of the background source. Assuming that the line is optically thin ($|\tau_\nu^{\text{l}}| \ll 1$), Equation (2) is approximated by (e.g., Kantharia & Anantharamaiah 2001)

$$\frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} = -\tau_\nu^{\text{l}}. \quad (3)$$

Note that, due to the minus sign on the right-hand side of Equation (3), when τ_ν^{l} is positive the line is observed in absorption against the background source.

From the definition of τ_ν^{l} (see Appendix B) and explicitly considering the normalized line profile, $\phi(\nu)$, (with $\int \phi(\nu) d\nu = 1$), we obtain

$$\frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} = -\kappa_\nu^{\text{l}} \phi(\nu) L. \quad (4)$$

Introducing the departure coefficients from LTE, b_n , and the correction factor for stimulated emission or absorption, β_n (Brocklehurst & Seaton 1972; Gordon & Sorochenko 2009), we can write

$$\begin{aligned} \frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} &= -\kappa_\nu^{\text{l}}(\text{LTE}) \phi(\nu) b_n \beta_{nn'} L, \\ \frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} &\approx -1.069 \times 10^7 \Delta n M(\Delta n) \frac{b_n \beta_{nn'}}{T_e^{5/2}} e^{x_n} \text{EM}_{C+} \phi(\nu), \end{aligned} \quad (5)$$

assuming $h\nu \ll kT_e$ and $\Delta n/n \ll 1$; in Equation (5) we have inserted the value for the κ_ν^{l} absorption coefficient (Appendix B). Here, $\text{EM}_{C+} = N_e N_{C+} L$ is the emission measure in units of $\text{cm}^{-6} \text{ pc}$, N_e is the electron density, N_{C+} is the carbon ion density, and L is the path length of the cloud in pc; $\Delta n = n' - n$ is the difference between the levels involved in the transition, and the factor $M(\Delta n)$ ⁶ comes from the approximation to the oscillator

⁴ Throughout this article we assume a filling factor of unity.

⁵ In Appendix A, we provide a comprehensive list of the symbols used in this article.

⁶ Some values for $M(\Delta n) = 0.1908, 0.02633, 0.008106, 0.003492, 0.001812$, for $\Delta n = 1, 2, 3, 4, 5$, respectively (Menzel 1968).

³ As in Salgado et al. (2015), following common usage in the astronomical literature, we refer to this process as dielectronic recombination rather than the more appropriate dielectronic capture.

strength of the transition, as given by Menzel (1968; see Appendix C). The $b_n\beta_{nn'}$ factor relates the line emission or absorption to the level population of the emitting atoms and has been calculated following the method described in Paper I; $\chi_n = hcZ^2Ry/n^2kT_e$, as defined in Appendix B.

At the line center, the line-to-continuum ratio depends on the broadening of the line (see Section 2.3). However, we can remove the dependence on the line profile by integrating the line over frequency:

$$\int \frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} d\nu = -1.069 \times 10^7 \Delta n M(\Delta n) \times \frac{b_n\beta_{nn'}}{T_e^{5/2}} e^{\chi_n} \text{EM}_{\text{C}^+} \text{ Hz.} \quad (6)$$

Note that by setting $\Delta n = 1$ (i.e., for $\text{Cn}\alpha$ lines⁷) in Equation (6), we recover Equation (70) in Shaver (1975) and Equation (5) in Payne et al. (1994).

For high densities, $b_n\beta_n$ approaches unity at high n levels, and the integrated line-to-continuum ratio changes little with n for a given T_e and EM_{C^+} . When the β_n factor in Equation (6) is positive (negative), the line is in absorption (emission). The strong dependence on electron temperature of the integrated line-to-continuum ratio ($\propto T_e^{-2.5}$) favors the detection of low-temperature clouds. An increase of a factor of 2 (3) in the temperature reduces the integrated line-to-continuum ratio by a factor of about 6 (15), all other terms being equal.

From Equation (6), we note that for $\text{Cn}\alpha$ lines

$$\int \frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} d\nu = -20.4 b_n\beta_{nn'} \left(\frac{T_e}{100 \text{ K}} \right)^{-2.5} \text{EM}_{\text{C}^+} \text{ Hz,} \quad (7)$$

$$= -0.2 b_n\beta_{nn'} \left(\frac{T_e}{100 \text{ K}} \right)^{-2.5} \left(\frac{N_e}{0.1 \text{ cm}^{-3}} \right)^2 \left(\frac{L}{\text{pc}} \right) \text{ Hz,} \quad (8)$$

assuming that electrons are produced by singly ionized carbon ($N_e = N_{\text{C}^+}$) and for a high n level ($n \gg \sqrt{(1.6 \times 10^5/T_e)}$). The typical optical depths that can be observed with current instruments are $\sim 10^{-3}$. As we already mentioned, for high n $b_n\beta_n \simeq 1$. Hence, clouds ($L \simeq 5$ pc) with electron densities greater than 10^{-2} cm^{-3} (hydrogen densities $> 50 \text{ cm}^{-3}$) are readily observable.

2.2. The Far-infrared Fine Structure Line of C+

The fine structure transition $^2P_{1/2}-^2P_{3/2}$ of carbon ions is one of the main coolants in diffuse neutral clouds at $158 \mu\text{m}$. Moreover, the $[\text{C II}]$ $158 \mu\text{m}$ line is directly linked to the level population of carbon atoms at low temperatures through the dielectronic recombination process. In Section 3, we show how observations of this line combined with CRRLs can be used as powerful probes of the temperature of diffuse neutral clouds. Here, we give a description of an emission model of the line. The intensity of the $[\text{C II}]$ $158 \mu\text{m}$ line in the optically thin limit

is given by (e.g., Sorochenko & Tsivilev 2000)

$$I_{158} = \frac{h\nu}{4\pi} A_{3/2,1/2} N_{3/2}^+ L = \frac{h\nu}{4\pi} \frac{A_{3/2,1/2} 2 \exp(-92/T_e) R}{1 + 2 \exp(-92/T_e) R} N_{\text{C}^+} L, \quad (9)$$

with ν the frequency of the $^2P_{1/2}-^2P_{3/2}$ transition, $A_{3/2,1/2} = 2.4 \times 10^{-6} \text{ s}^{-1}$ is the spontaneous transition rate, $N_{3/2}^+$ the number density of carbon ions in the $3/2$ state, L the length along the line of sight of the observed cloud, and N_{C^+} the density of carbon ions; R is defined in Ponomarev & Sorochenko (1992) and Payne et al. (1994) (see Paper I):

$$R = \frac{N_e \gamma_e + N_H \gamma_H}{N_e \gamma_e + N_H \gamma_H + A_{3/2,1/2}}, \quad (10)$$

where γ_e and γ_H are the de-excitation rates due to electrons and hydrogen atoms, respectively. The rates involved are detailed in Paper I. We assume that collisions with electrons and hydrogen atoms dominate over molecular hydrogen and neglect collisions with H_2 , as in Paper I. This is a good approximation for diffuse clouds with column densities up to $\sim 10^{21} \text{ cm}^{-2}$. For larger column densities, the H/H_2 transition will have to be modeled in order to evaluate R .

The optical depth of the C^+ fine structure line for the transition $^2P_{1/2}-^2P_{3/2}$ is given by Crawford et al. (1985) and Sorochenko & Tsivilev (2000):

$$\tau_{158} = \frac{c^2}{8\pi\nu^2} \frac{A_{3/2,1/2}}{1.06\Delta\nu} 2\alpha_{1/2}\beta_{158} N_{\text{C}^+} L, \quad (11)$$

where $\Delta\nu$ is the full width at half maximum (FWHM) of the line (assumed to be Gaussian); the $\alpha_{1/2}(T_e)$ and $\beta_{158}(T_e)$ coefficients depend on the electron temperature of the cloud and are defined by Sorochenko & Tsivilev (2000):

$$\alpha_{1/2}(T_e) = \frac{1}{1 + 2 \exp(-92/T_e) R}, \quad (12)$$

$$\beta_{158}(T_e) = 1 - \exp(-92/T_e) R. \quad (13)$$

Adopting a line width of 2 km s^{-1} , at low electron temperatures and densities, the far-infrared $[\text{C II}]$ line is optically thin for hydrogen column densities less than about $1.2 \times 10^{21} \text{ cm}^{-2}$. For a cloud size of 5 pc , this corresponds to hydrogen densities of $\sim 10^2 \text{ cm}^{-3}$ and electron densities of $\simeq 10^{-2} \text{ cm}^{-3}$ if carbon is the dominant ion.

2.3. Line Profile of Recombination Lines

The observed profile of a line depends on the physical conditions of the cloud, as an increase in electron density and temperature or the presence of a radiation field can broaden the line, and this is particularly important for high n . Therefore, in order to determine the detectability of a line, the profile must be considered. Conversely, the observed line width of recombination lines provides additional information on the physical properties of the cloud.

The line profile is given by the convolution of a Gaussian and a Lorentzian profile and is known as a Voigt profile (Shaver 1975; Gordon & Sorochenko 2009). Consider a cloud of gas of carbon ions at a temperature T_e . Random thermal motions of the atoms in the gas produce shifts in frequency that

⁷ We will refer to electron transitions in carbon for levels $n+1 \rightarrow n$ as $\text{Cn}\alpha$, $n+2 \rightarrow n$ as $\text{Cn}\beta$, and $n+3 \rightarrow n$ as $\text{Cn}\gamma$ (Gordon & Sorochenko 2009).

are reflected in the line profile as a Gaussian broadening (Doppler broadening). In the most general case, turbulence can increase the width of a line and, as is common in the literature (e.g., Rybicki & Lightman 1986), we describe the turbulence by an rms turbulent velocity. Thus, the Gaussian line profile can be described by

$$\Delta\nu_D = \frac{\nu_0}{c} \sqrt{\frac{2kT_e}{m_C} + \langle v_{\text{rms}} \rangle^2}, \quad (14)$$

where m_C is the mass of the carbon atom and $\langle v_{\text{rms}} \rangle$ is the rms turbulent velocity. The Gaussian width in frequency space is proportional to the frequency of the line transition.

At low frequencies, collisions and radiation broadening dominate the line width. The Lorentzian (FWHM) broadening produced by collisions is given by

$$\Delta\nu_{\text{col}} = \frac{1}{\pi} \sum_{n \neq n'} N_e C_{n'n}, \quad (15)$$

where $C_{n'n}$ is the collision rate for electron-induced transitions from level n' to n , and N_e is the electron density. Note that $C_{n'n}$ depends on temperature (Shaver 1975; Gordon & Soroichenko 2009). In order to estimate the collisional broadening, we fitted the following function at temperatures between 10 and 30,000 K:

$$\sum_{n \neq n'} C_{n'n} = 10^{a(T_e)} n^{\gamma_c(T_e)}, \quad (16)$$

which is valid for levels $n > 100$. Values for $a(T_e)$ and $\gamma_c(T_e)$ as a function of electron temperature are given in Table 3.

In a similar way as for collisional broadening, the interaction of an emitter with a radiation field produces a broadening of the line profile. In Appendix C, we give a detailed expansion for different external radiation fields. Here, we discuss the case of a synchrotron radiation field characterized by a power law with a temperature T_0 at a reference frequency $\nu_0 = 100$ MHz and a spectral index $\alpha_{\text{pl}} = -2.6$ (see Section 3). Under the above considerations, the FWHM for radiation broadening is given by

$$\Delta\nu_{\text{rad}} = 6.096 \times 10^{-17} T_0 n^{5.8} \text{ (s}^{-1}\text{)}. \quad (17)$$

As is the case for collisional broadening, radiation broadening depends only on the level and the strength of the surrounding radiation field. The dependence on n is stronger than that of collisional broadening at low densities, and radiation broadening dominates over collisional broadening. As the density decreases, the level n where radiation broadening dominates decreases. In order to estimate where this occurs, we define t_n as

$$t_n(T_e, T_0, N_e) = \frac{\Delta\nu_{\text{rad}}}{\Delta\nu_{\text{col}}} = \left[\frac{6.096 \times 10^{-17}}{10^{a(T_e)}} \right] \left(\frac{T_0}{N_e} \right) n^{5.8 - \gamma_c(T_e)}. \quad (18)$$

Note that the dependence on electron temperature is contained within the fitting coefficients, a and γ_c . For $T_e = 100$ K, we find $t_n \approx 5.82 \times 10^{-7} \sqrt{n} (T_0/N_e)$. In Figure 1, we show t_n as a function of electron density for $T_0 = 1000$ K. For a given electron temperature and density, the influence of an external radiation field is larger for higher levels since $t_n \propto n^{5.8 - \gamma_c}$ and $\gamma_c < 5.8$ (see appendix). For a given density, the influence of

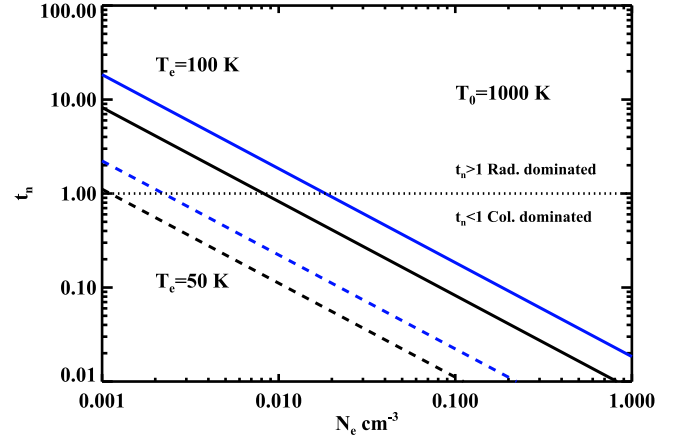


Figure 1. The t_n factor defined in Equation (18) as a function of electron density for quantum n levels between 200 (black line) and 1000 (blue line). The figure is presented for two electron temperatures: $T_e = 50$ K (dashed lines) and $T_e = 100$ K (solid lines). The dotted line marks the boundary for the line widths being in the collision-dominated regime ($t_n < 1$) and the radiative-dominated regime ($t_n > 1$).

the radiation field on the line width is larger at higher electron temperatures. For the typical conditions of the CNM, that is, at $T_e = 100$ K and $N_e = 0.02 \text{ cm}^{-3}$, the value of $t_n \approx 1$ and both the radiation field and electron density affect the line width in similar amounts.

3. Method

In order to study the radiative transfer effects on the lines, we use the method outlined in Paper I to compute the departure coefficients for different electron temperatures and densities and considering an external radiation field. The cosmic microwave radiation field (CMB) and the Galactic synchrotron power-law radiation field spectra are included. We represent the CMB by a 3 K blackbody and the galactic radiation field by a power law $[I_0(\nu) = T_0(\nu/\nu_0)^{\alpha_{\text{pl}}}]$ with $T_0 = 1000$ K at a frequency $\nu_0 = 100$ MHz and $\alpha_{\text{pl}} = -2.6$ (Landecker & Wielebinski 1970; Bennett et al. 2003). In the Galactic plane, the Galactic radiation field can be much larger than 1000 K at 100 MHz (Haslam et al. 1982). At frequencies higher than 1 GHz, corresponding to $\text{Cn}\alpha$ transitions from levels with $n < 200$, the background continuum is dominated by the CMB (Figure 2). At even higher frequencies, the background continuum can be dominated by dust and free-free emission, which are strongly dependent on the local conditions of the cloud and its position in the Galaxy. For simplicity, we focus our study on levels with $n > 200$.

Departure coefficients were computed for $T_e = 20, 50, 100$, and 200 K and electron densities in the range 10^{-2} – 1 cm^{-3} . Once the departure coefficients were obtained, we computed the corresponding optical depths assuming a fixed length along the line of sight of 1 pc from the usually adopted approximated optical depth solution to the radiative transfer problem (Equation (6)). The value of 1 pc corresponds to emission measures in the range of $\text{EM}_{\text{C}+} = 10^{-4}$ to $1 \text{ cm}^{-6} \text{ pc}$. Our calculations assume a homogeneous density distribution in a cloud and should be taken as illustrative since it is well known that inhomogeneities exist in most clouds. The fixed length of 1 pc corresponds to column densities of 10^{18} – 10^{21} cm^{-2} , with the adopted density range. Diffuse clouds show a power-law distribution function in HI column density with a median

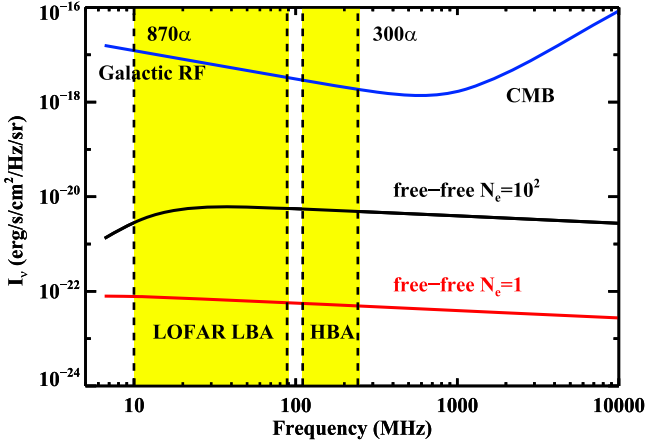


Figure 2. A comparison between the continuum radiation fields. The galactic synchrotron radiation field dominates over the free-free cloud continuum at $T_e = 100$ K. Therefore, the strong background approximation is valid for the low-temperature cases considered in this analysis. The yellow zone marks the range in frequency observable by LOFAR.

column density of $0.76 \times 10^{20} \text{ cm}^{-2}$ (Heiles & Troland 2003b). Reddening studies are weighted to somewhat large clouds, and the standard “Spitzer”-type cloud (Spitzer 1978) corresponds to a column density of $3.6 \times 10^{20} \text{ cm}^{-2}$. Local H I complexes associated with molecular clouds have $N_{\text{H}} \approx 10^{21} \text{ cm}^{-2}$.

4. Results

4.1. Line Widths

We begin our discussion with the results for the line widths. We show the line widths for our diffuse cloud models in Figure 3. At high frequencies (low n), the Gaussian (Doppler) core of the line dominates the line profile in frequency space, and the line width increases with frequency. At low frequencies (high n), on the other hand, the Lorentzian profile dominates—because of either collisional or radiation broadening—and the line width decreases with increasing frequency. In order to guide the discussion, we have included observed line widths for $\text{Cn}\alpha$ transitions for Cas A (Payne et al. 1994; Kantharia et al. 1998), Cyg A (Oonk et al. 2014), and M82 (Morabito et al. 2014).

When the Doppler core dominates, CRRL observations provide both an upper limit on the gas temperature and an upper limit on the turbulent velocity of the diffuse ISM (compare Equations (17) and (16)). For typical parameters of the turbulent ISM (1 km s^{-1}), turbulence dominates over thermal velocities when $T_e \lesssim 700$ K.

Radiation broadening and collisional broadening show a very similar dependence on n , and it is difficult to disentangle these two processes from CRRL observations. For the Galactic radiation field (i.e., a synchrotron spectrum with $T_0 = 1000$ K at 100 MHz), the two processes contribute equally to the line width at a density $N_e \approx 0.03 \text{ cm}^{-3}$ (Figure 3). Low-frequency observations can, thus, provide an upper limit on the density and radiation field. As illustrated in Figure 3, the transition from a Doppler to a Lorentzian broadened line is quite rapid (in frequency space), but the actual value of n where it occurs depends on the physical conditions of the cloud (i.e., T_e , N_e , T_0 , and $\langle v_{\text{rms}} \rangle$).

In Figure 3 we can see that the RRLs from Cas A and Cyg A fall in a region of the diagram corresponding to densities lower

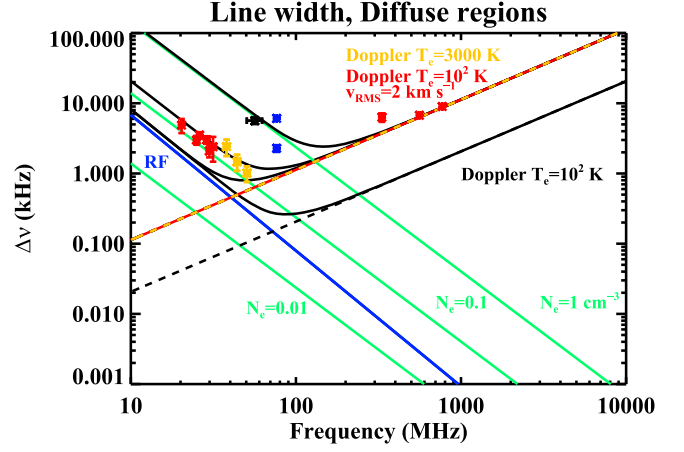


Figure 3. A comparison between broadening produced by the Galactic radiation field (blue line), collisional broadening at $N_e = 1, 0.1$, and 0.01 cm^{-3} (green lines), and thermal (Doppler) broadening at 100 K (black dashed line). The red and yellow curves correspond to a turbulent Doppler parameter $\langle v_{\text{rms}} \rangle^2 = 2 \text{ km s}^{-1}$ and $T_e = 300$ K, respectively. We include data for Cas A (Payne et al. 1994; Kantharia et al. 1998) as red points, Cyg A (Oonk et al. 2014) as yellow points, regions for the inner galaxy (Erickson et al. 1995) as blue points, and data for M82 (Morabito et al. 2014) as a black point.

than about 0.1 cm^{-3} , and the detection in M82 corresponds to higher densities, to a much stronger radiation field, or to the blending of multiple broad components. From observations at high frequencies, it is known that the lines observed toward Cas A are the result of three components at different velocities in the Perseus and Orion arms. Therefore, the physical parameters obtained from line widths should be taken as upper limits.

When the line profile is dominated by the Doppler core, the ratio of the β to α line width is unity. However, radiation or collisional broadening affects the $\text{Cn}\alpha$ and $\text{Cn}\beta$ lines differently because at the same frequency $\text{Cn}\alpha$ and $\text{Cn}\beta$ lines originate from different n levels. In Figure 4, we show the ratio $\Delta\nu(\beta)/\Delta\nu(\alpha)$. We notice that, when radiation broadening dominates the line width, this ratio goes to a constant value, independent of the background temperature. From the radiation broadening formula (Equation (17)), we see that $\Delta\nu(\beta)/\Delta\nu(\alpha) = (n_\beta/n_\alpha)^{-3\alpha_{\text{pl}}-2}$ and, for a power law $\alpha_{\text{pl}} = -2.6$, the ratio approaches $\Delta\nu(\beta)/\Delta\nu(\alpha) = 3.8$ as n increases. At high electron densities, collisional processes dominate the broadening of the lines. From Equation (16), the $\Delta\nu(\beta)/\Delta\nu(\alpha)$ ratio tends to a constant value of $(n_\beta/n_\alpha)^{\gamma_c} = 1.26^{\gamma_c}$. There is a temperature dependence in the exponent γ_c , and, for electron temperatures less than 1000 K, we find that $\Delta\nu(\beta)/\Delta\nu(\alpha) \approx 3.1\text{--}3.6$ (see Table 3), similar to the radiation-broadening case.

4.2. Integrated Line-to-continuum Ratio

As discussed in Section 2.1, the line-to-continuum ratio of CRRLs is often solved approximately, using Equation (6). In this subsection, we will discuss when this approximation is justified. In this, we have to recognize that, under the conditions of the diffuse ISM, recombining carbon atoms are not in LTE (Paper I). Indeed, electrons can recombine to high levels due to dielectronic recombination, thus increasing the population in comparison to the LTE values. This increase in the level population leads to an increase in the values of the

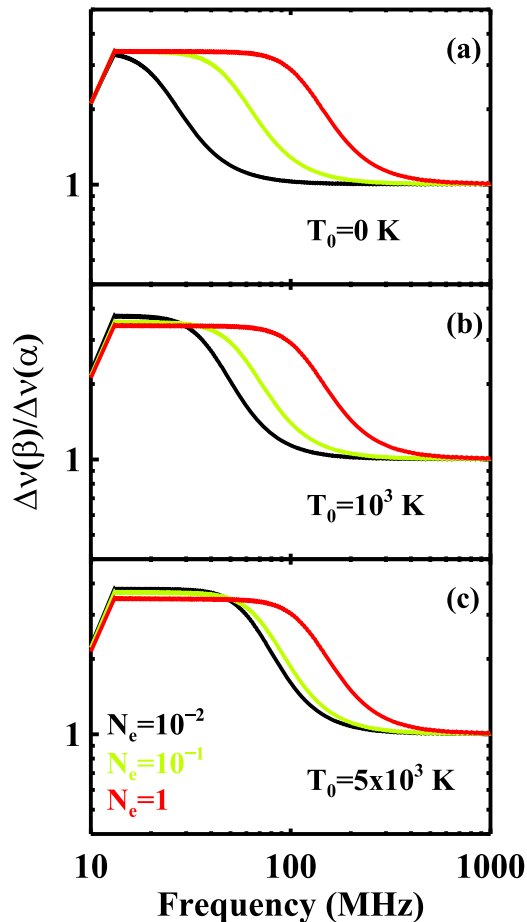


Figure 4. The α and β line width transitions for diffuse regions as a function of frequency for different power-law radiation fields: (a) without an external radiation field; (b) a power-law radiation field with $T_0 = 1000$ K; and (c) same as (b) for $T_0 = 5000$ K. The line widths correspond to electron densities of $N_e = 1$, 0.1 , and 0.01 cm^{-3} (red, green, and black lines).

$b_n \beta_n$ coefficients in Equation (6) and, consequently, to an increase in the optical depth of the lines.

In Figure 5 we show the integrated line-to-continuum ratio as a function of level n for $T_e = 100$ K. We compare the values obtained using the approximated expression given in Equation (6) (red lines) and by solving the radiative transfer equation (Equation (1), black lines). The agreement between the two approaches is good for levels $n \gtrsim 250$, since at these high levels the approximations that lead to Equation (6) are valid. For levels lower than $n \approx 250$, differences appear. In particular, at low electron densities ($N_e \approx 0.01$ cm^{-3}), the results using Equation (6) show lines in absorption, while the results derived from solving the radiative transfer equation predict lines in emission. The difference between the two approaches can be understood in terms of the excitation temperature (see appendix). As can be seen in Figure 6, the red zones correspond to low n levels, where the excitation temperature is higher than the background continuum temperature, and the lines appear in emission (despite the β_n being positive). At higher n values, $\beta_n < 0$ (yellow zones), and the excitation temperature is negative, reflecting an inversion in the level population, so lines appear in emission. While there is an inversion of the level population, the line optical depths are too low ($\tau_l \sim 10^{-3}$) to produce a maser (compare with Equation (7)). At even higher levels (blue zones in Figure 6),

the excitation temperature is less than the background continuum temperature, and the lines are in absorption. As the electron density increases, dielectronic recombination is less efficient, and the levels for which β_n is negative shift to lower n values, resembling the values for hydrogenic level population (Hummer & Storey 1987, Paper I). Furthermore, for high quantum numbers and high densities, $\beta_n = 1$, and the excitation temperature is equal to the electron temperature of the gas.

From this analysis, we conclude that Equation (6) is valid for high ($n \gtrsim 250$) quantum numbers, and the ratio of two lines depends only on the temperature and electron density of the cloud through the departure coefficients. In Figure 7, we demonstrate this by showing the integrated line-to-continuum ratio of $Cn\alpha$ as a function of quantum number normalized to the level 500 (similar results can be obtained using other n levels). The normalized ratio becomes smaller for high densities because b_n/β_n values change little with n as the levels are closer to equilibrium. As the electron density decreases, dielectronic recombination is more efficient in overpopulating intermediate levels (Paper I), producing large changes in the values of the ratios.

4.3. CRRLs as Diagnostic Tools for the Physical Conditions of the ISM

4.3.1. Line Ratios

We have already discussed the use of the line width to constrain the properties of the emitting/absorbing gas. As Figure 7 illustrates, line ratios are very sensitive to the physical conditions in the gas. Moreover, the use of line ratios “cancels out” the dependence on the emission measure. Here, we demonstrate the use of line ratios involving widely different n values as diagnostic tools in “ratio versus ratio” plots. As an example, we show three line ratios in Figure 8, normalized to $n = 500$. The lines are chosen to sample the full frequency range of LOFAR and the different regimes (collisional, radiative) characteristic for CRRLs. The $n = 300, 400, 500$ lines are a particularly good probe of electron density for regions with temperatures less than about 100 K. The use of the $n = 500$ level does not affect our results, and other levels (e.g., $n = 600$ or 800) may be used for computing the ratios. We note that in a limited but relevant electron density range ($N_e \sim 1\text{--}5 \times 10^{-2}$ cm^{-3}), these lines can be good tracers of temperature. At higher densities, the departure coefficients approach unity and the ratios tend to group in a small region of the plot, and the use of the ratios as probes of temperature requires measurements with high signal-to-noise ratio to derive physical conditions from the observations.

4.3.2. The Transition from Absorption to Emission

In Paper I, we discussed the use of the level where lines transition from emission to absorption (n_t) as a constraint on the density of a cloud (Figure 9). The limited observations in the Galactic plane (Erickson et al. 1995; Kantharia & Anantharamiah 2001) indicate that $400 > n_t > 350$ and n_t depends on both temperature and density. The transition level can be used to estimate the electron density for electron densities lower than about 10^{-1} cm^{-3} . For increasing electron density, it becomes more difficult to constrain this quantity from the transition level alone.

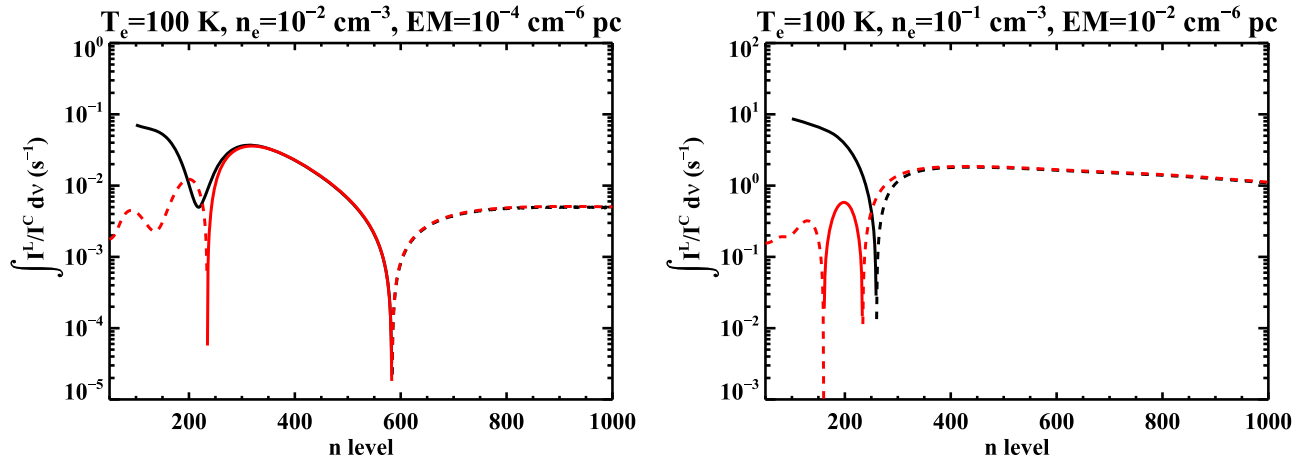


Figure 5. The line-to-continuum ratio of CRRLs as a function of principal quantum number for $T_e = 100$ K and $N_e = 0.01$ and 0.1 cm^{-3} (left and right panels, respectively). The values were computed from the radiative transfer solution (Equation (1)) and the galactic radiation field as a background. Black lines correspond to the result of solving the equation of radiative transfer, while red lines correspond to the approximation expression given in Equation (6). At levels larger than $n \gtrsim 250$, the differences between the approximation (dashed) and the radiative transfer solution (solid) are minor.

4.3.3. Line Ratios as a Function of Δn

Combining observations of $Cn\alpha$ lines with $Cn\beta$ and $Cn\gamma$ lines can provide further constraints on the physical parameters of the cloud. In Figure 10 we show the α -to- β ratio of the integrated line-to-continuum ratio as a function of frequency. Recall that $Cn\alpha$ and $Cn\beta$ lines observed at almost the same frequency probe very different n levels ($n_\alpha = 1.26n_\beta$). Figure 10 shows that both electron density and temperature are involved. At high n levels, the $b_n\beta_n$ values are approximately unity and the α -to- β ratio approaches $M(1)/2M(2) \approx 0.1908/0.0526 = 3.627$ (Equation (6)), making the ratio less useful in constraining temperature and electron density. However, even at high n , this ratio does remain useful for investigating the radiation field incident upon the CRRL-emitting gas.

4.3.4. The CRRL/[C II] Ratio

The [C II] $158 \mu\text{m}$ line is the dominant cooling line of diffuse clouds and acts as a thermostat regulating the temperature (Hollenbach & Tielens 1999). In realistic models of the ISM of galaxies (e.g., Wolfire et al. 1995), the photoelectric effect on polycyclic aromatic hydrocarbon molecules and very small grains heats the gas, and the cooling by the [C II] $158 \mu\text{m}$ line adjusts to satisfy the energy balance. As the heating is a complicated function of the physical conditions (Bakes & Tielens 1994), models become very involved. Here, we sidestep this issue, and we calculate the [C II] $158 \mu\text{m}$ intensity as a function of N_e and T_e for a uniform cloud. The intensity scales with the column density of carbon ions, N_{C^+} , and temperature. In contrast, the CRRLs scale with the emission measure divided by $T_e^{5/2}$ (compare with Equation (6)). Hence, the ratio of the CRRL to the $158 \mu\text{m}$ line shows a strong dependence on temperature (and electron density), but for a constant density this ratio does not depend on column density. In Figure 11 we show the CRRL/[C II] ratio as a function of density for different temperatures. For the physical conditions relevant for diffuse clouds, the CRRL/[C II] ratio is a powerful diagnostic tool. Moreover, as we demonstrate below, low-frequency CRRLs are not expected to be observable at the typical temperatures and densities of classical H II regions. We recognize that [C II] at $158 \mu\text{m}$ can be produced by the WIM. Nevertheless, we expect the

contribution from the WIM to the [C II] line to be $\sim 4\%$ in the general ISM (Pineda et al. 2013). However, COBE observations of the [N II] $205 \mu\text{m}$ line from the Milky Way (Bennett et al. 1994) have demonstrated that [C II] $158 \mu\text{m}$ emission from the WIM may be more important along some sight lines (Heiles 1994).

5. On the Observed Behavior of CRRLs

5.1. General Considerations

CRRLs have been observed toward two types of regions: high-density PDRs and diffuse clouds (Gordon & Sorochenko 2009). In general, low-frequency CRRLs are observed in absorption with values for the integrated line-to-continuum ratio in the range of $1\text{--}5 \text{ Hz}^8$ and a peak line-to-continuum ratio of $\sim 10^{-4}$ to 10^{-3} (Erickson et al. 1995; Kantharia et al. 1998; Roshi et al. 2002; Oonk et al. 2014).

In order to observe CRRLs, carbon atoms must be singly ionized. In H II regions, carbon is found in higher stages of ionization, and the gas is dense and warm. Hence, recombination lines of the type we study here are not expected to be strong. In PDRs of high density, carbon atoms transition from ionized to neutral and into molecular (CO) around a visual extinction $A_V \approx 4$ mag, depending on the density and UV field. Assuming $A_V = N_H/1.9 \times 10^{21} \text{ mag cm}^{-2}$, we can estimate the maximum column density of carbon that can be expected for such a transition region. Assuming that carbon is fully ionized and a carbon abundance of 1.6×10^{-4} , we obtain a column density of carbon of $1.2 \times 10^{18} \text{ cm}^{-2}$.

As mentioned in Section 2.1, CRRLs produced in clouds with high temperatures are faint due to the strong dependence of the line-to-continuum ratio on temperature. Therefore, regions of low temperature are favored to be observed using low-frequency recombination lines. These two considerations (low T_e and N_e) set a range of electron density and temperature for which CRRLs are easier to detect. Specifically, consider a medium with two phases in pressure equilibrium. From

⁸ We quote the integrated line-to-continuum ratio in units of Hz as opposed to km s^{-1} .

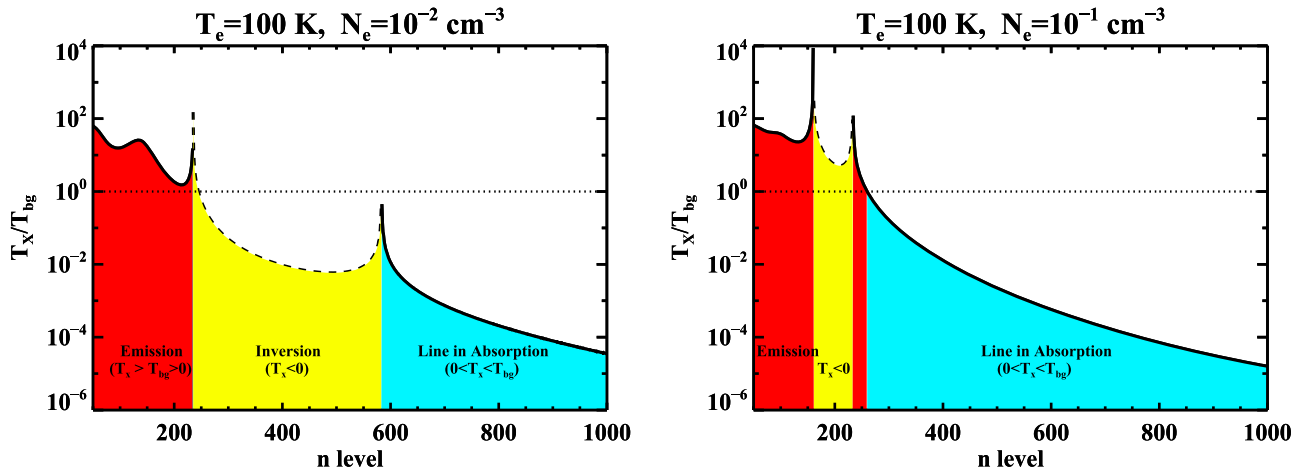


Figure 6. Ratio of the excitation to background temperature (T_X/T_{bg}). Lines are in emission in the red zone since $T_X > T_{bg}$ and in the yellow zone due to an inversion on the level population and $T_X < 0$. Lines appear in absorption in the light blue zone since the background temperature is (much) larger than the excitation temperature.

Equation (4), the optical depth ratio scales then with

$$\frac{\tau_1}{\tau_2} \propto \frac{N_{e,1}^2 T_{e,2}^{5/2} (b_n \beta_n)_1 L_1}{T_{e,1}^{5/2} N_{e,2}^2 (b_n \beta_n)_2 L_2} \propto \left(\frac{T_2}{T_1}\right)^{9/2} \frac{(b_n \beta_n)_1 L_1}{(b_n \beta_n)_2 L_2}. \quad (19)$$

For parameters relevant for the CNM and WNM ($T_{e,1} = 80$ K, $T_{e,2} = 8000$ K, respectively; Tielens 2005), we have then $\tau_1/\tau_2 \sim 10^9 (b_n \beta_n)_1 / (b_n \beta_n)_2 L_1/L_2$. Clearly, CRRLs will overwhelmingly originate in cold, diffuse clouds. Therefore, unlike 21 cm HI observations, the analysis of CRRL observations is not hampered by confusion of CNM and WNM components.

The fact that low-frequency recombination lines are observed in absorption sets a lower limit on the density for the clouds where CRRLs are produced. Our models show that at electron densities lower than 10^{-2} cm^{-3} and for temperatures lower than 200 K, low-frequency CRRLs are in emission.

5.2. Illustrative Examples

In this section we illustrate the power of our models to derive physical parameters from observations of CRRLs. We selected observations toward Cas A because, to our knowledge, the clouds toward Cas A are the best studied using CRRLs. We then expand this illustration by using observations of two regions observed toward the Galactic center from Erickson et al. (1995).

5.2.1. Cas A

We begin our analysis with CRRLs detected toward Cas A from the literature (e.g., Payne et al. 1994; Kantharia et al. 1998; Stepkin et al. 2007). In Figure 12, we summarize the constraints from the integrated line α -to- β ratio as a blue zone using the Stepkin et al. (2007) data. The transition from emission to absorption ($350 < n_t < 400$) is shown as the green zone. The 600–500 ratio vs. 270–500 ratio is included as the red zone.⁹ Finally, the yellow zone is the intersection of all of the above-mentioned zones.

The line width does not provide much of an additional constraint. For Cas A, with an observed line width of 6.7 kHz

at $\nu = 560$ MHz (Kantharia et al. 1998), the implied gas temperature would be $T_e = 3000$ K, and actually we expect that the line is dominated by turbulence with $\langle v_{rms} \rangle \approx 2 \text{ km s}^{-1}$ (Figure 3). Likewise, the Cas A observations from Payne et al. (1994) and Kantharia et al. (1998) are of little additional use as we arrive at $N_e \lesssim 0.1 \text{ cm}^{-3}$ and $T_0 \lesssim 2000$ K.

Perusing Figure 12, we realize that the α -to- β line ratio does not provide strong constraints due to the frequency at which the lines were observed as all the models converge to the high-density limit (Figure 10). The transition level from emission to absorption (n_t) restricts the allowed models to an area in the N_e versus T_e plane. However, at low temperatures ($T \lesssim 50$ K), the constraining power of n_t is limited. The “ratio vs. ratio” plots can be quite useful in constraining both the electron density and the temperature of the line-producing cloud, as we have illustrated here.

The results of our models show that the properties of the cloud are well restricted in density ($N_e = 2\text{--}3 \times 10^{-2} \text{ cm}^{-3}$)—corresponding to H densities of $\sim 100\text{--}200 \text{ cm}^{-3}$ —but somewhat less in temperature ($T_e = 80\text{--}200$ K). We emphasize, though, that these results are ill-defined averages because the CRRLs toward Cas A are known to be produced in multiple velocity components that are blended together. In addition, a preliminary analysis of the LOFAR data indicates variations in CRRL optical depth on angular scales significantly smaller than the beam sizes used in the observational data from the previous literature studies quoted here. Nevertheless, this example illustrates the power of CRRL observations to measure the physical conditions in diffuse interstellar clouds.

5.2.2. Galactic Center Regions

As a second example, we analyze observations of clouds detected toward regions in the galactic plane (Erickson et al. 1995). In view of the scarceness, low spatial resolution, and limited frequency coverage of the data available in the literature, our results should be taken with care and considered illustrative. We chose two regions with good signal-to-noise measurements ($S/N > 10$). In Table 1, we show the line parameters for C441 α and C555 β lines from Erickson et al. (1995) with a beam size of 4° .

In Figure 13, we summarize the constraints imposed by the integrated α -to- β line ratio as a blue zone, the transition from

⁹ We use the $n = 270$ data from Kantharia et al. (1998) and estimate the data at $n = 600$ from Payne et al. (1994) and analogous plots as in Figure 8.

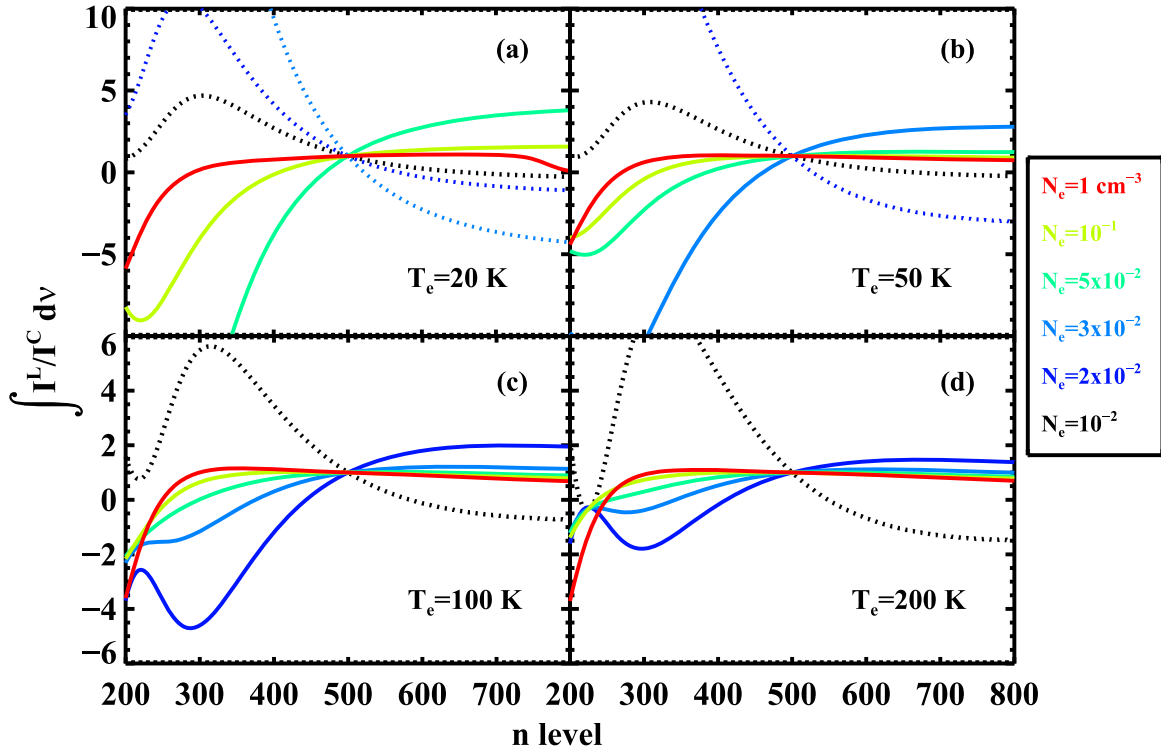


Figure 7. Integrated line-to-continuum ratio normalized to the value at the level $n = 500$ for $T_e = 20, 50, 100,$ and 200 K. Dotted lines indicate that the C500 α line is in emission. The values have been computed considering radiative transfer effects (Equation (1)).

emission to absorption (n_t level) as a green zone (which we estimate to be $350 < n_t < 400$), and the integrated line-to-continuum to the $I(158 \mu\text{m})$ ratio as the orange zone.

From the line widths toward the lines of sight in Table 1, an upper limit to the density can be estimated by assuming pure collisional broadening, as shown in Section 2.3. The upper limits on density are $N_e \leq 1.5 \text{ cm}^{-3}$ for G000.0+0 and $N_e \leq 0.5 \text{ cm}^{-3}$ for G002.0-2. The constraint is even more strict when considering that part of the broadening must be produced by the Galactic radiation field. Assuming no collisional broadening, the upper limits on the background temperatures for the regions are $T_0 \leq 4 \times 10^4 \text{ K}$ for G000.0+0 and $T_0 \leq 1.5 \times 10^4 \text{ K}$ for G002.0-2. These are strict upper limits as the observations from Erickson et al. (1995) were performed with large beams, and the observed lines are likely produced by several “clouds” in the beam.

We estimate the value for $I(158 \mu\text{m})$ to be $8\text{--}12 \times 10^{-5} \text{ erg s cm}^{-2} \text{ sr}^{-1}$ from COBE data (Bennett et al. 1994). Since the data from Erickson et al. (1995) is for the C441 α line, we created a diagnostic plot similar to that in Figure 11 for the level 441. We obtain for G000.0+0 a value for T_e between 20 and 60 K, and N_e between 4×10^{-2} and $1 \times 10^{-1} \text{ cm}^{-3}$. For G002.0-2, we obtain $T_e = 20\text{--}80 \text{ K}$ and $N_e = 4 \times 10^{-2}\text{--}1 \times 10^{-1} \text{ cm}^{-3}$. With these values and using Equation (7), we determine lengths of 2–19 pc for G000.0+0 and 1–9 pc for G002.0-2. Assuming that the electrons are provided by carbon ionization and adopting a carbon gas-phase abundance of 1.6×10^{-4} , we derive thermal pressures between 5000 and 37,500 K cm^{-3} for G000.0+0 and between 5000 and 50,000 K cm^{-3} for G002.0-2. Strictly speaking, the values from COBE include emission produced in the neutral and warm components along the lines of sight. Since CRRLs are

expected to be produced predominantly in cold clouds, the here-determined ratio between the CRRL and the [C II] line can be underestimated. However, Pineda et al. (2013) showed that the contribution from ionized gas to the [C II] line is $\sim 4\%$ toward the inner Galaxy. It is clear from Figure 13 that the n_t level (green zone in the plots) and the integrated α -to- β line ratio provide similar constraints in the N_e vs. T_e plane. By far, the strongest constraint comes from n_t , since the errors in the measurements do not provide strong limits on the α -to- β line ratio. As the error bars are rather large, the derived constraints—given above—are not very precise. Nevertheless, the inherent power of CRRLs for quantitative studies of diffuse clouds in the ISM is quite apparent.

5.3. Discussion

As the examples of Cas A and G000.0+0 and G002.0-2 show, a large amount of relevant physical information on the properties of the clouds can be obtained from CRRL measurements, despite the scarceness of the data used here. The α -to- β line ratios can provide powerful constraints as long as the frequency observed is higher than 30 MHz. As illustrated by our Cas A example, the CRRL ratio plots can be extremely useful in constraining the electron density and temperature, and lines with a large separation in terms of quantum number are expected to be the most useful ratios. As illustrated in Figure 8, ratios between levels around 300 and 500 can provide direct constraints or indirect constraints by using, in addition, the n_t value. An advantage of using ratios is that they only depend on the local conditions and beam-filling factors are of little concern.

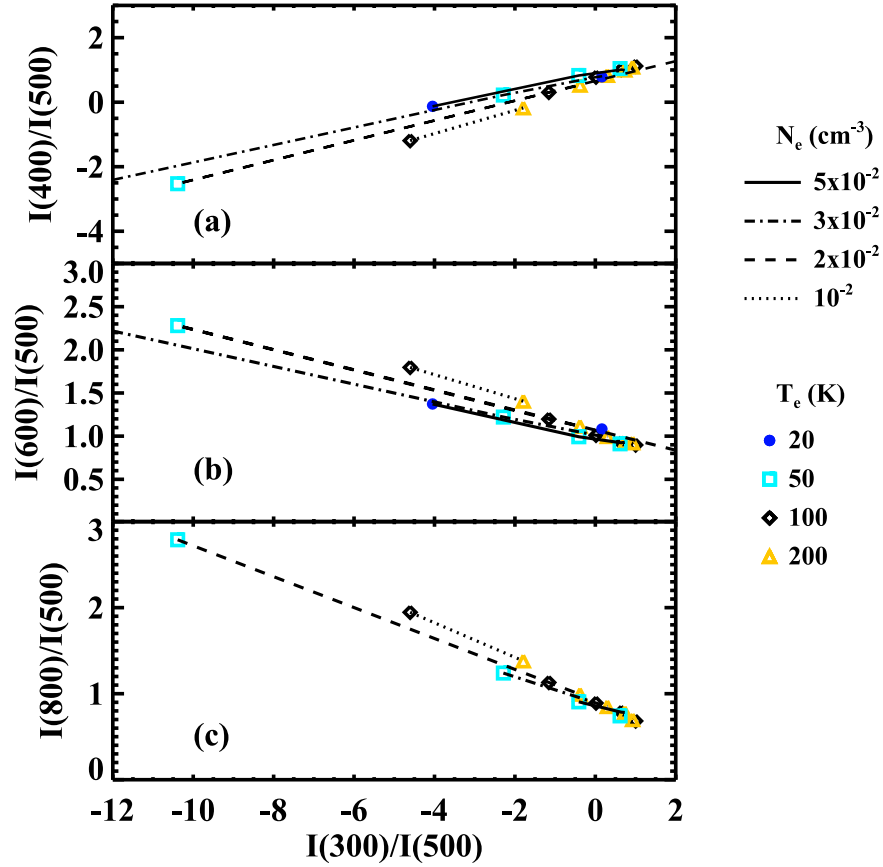


Figure 8. Example ratio diagnostic plots for different electron temperatures and densities. Cyan points are for $T_e = 50$ K, black points for $T_e = 100$ K, and orange points for $T_e = 200$ K. Different densities are joined by dotted lines ($N_e = 10^{-2} \text{ cm}^{-3}$), dashed lines ($N_e = 2 \times 10^{-2} \text{ cm}^{-3}$), dash-dotted lines ($N_e = 3 \times 10^{-2} \text{ cm}^{-3}$), and continuous lines ($N_e = 5 \times 10^{-2} \text{ cm}^{-3}$). (a) Ratio of the integrated line to continuum for levels 400 and 500 vs. 300–500 ratio. (b) Ratio of the integrated line to continuum for levels 600 and 500 vs. 300–500 ratio. (c) Ratio of the integrated line to continuum for levels 800 and 500 vs. 300–500 ratio.

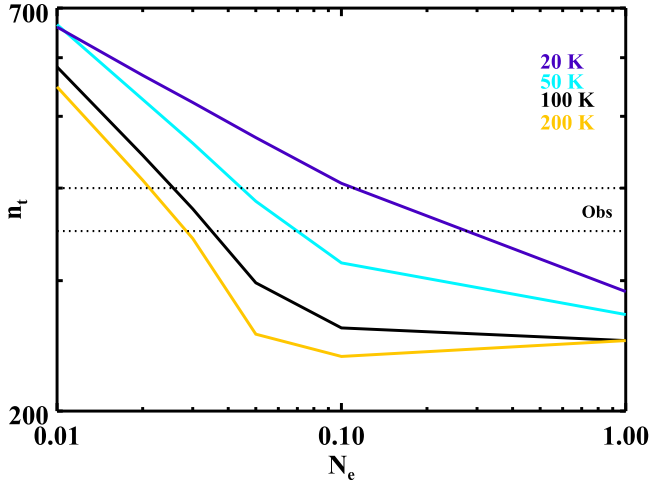


Figure 9. Level where lines transition from emission to absorption (n_t) as a function of electron density (N_e) for $T_e = 50$, 100, and 200 K. The horizontal dashed lines mark the limits as suggested by observations of CRRLs in the Galaxy.

Although we consider our examples illustrative, the determined values for T_e and n_e are within the values expected from theory (e.g., Wolfire et al. 2003; Kim et al. 2011) and H I 21 cm observations (Heiles & Troland 2003b). Moreover, the derived thermal pressures agree well with those derived from CI UV lines in the local ISM (Jenkins & Shaya 1979; Jenkins et al. 1983; Jenkins & Tripp 2011).

6. Summary and Conclusions

In this paper we have analyzed CRRL observations. Anticipating the LOFAR CRRL survey, we focus our study in the low-frequency regime, corresponding to transitions between lines with high principal quantum number. We have studied the radiative transfer of recombination lines and the line-broadening mechanisms in the most general form.

Our results show that line widths provide constraints on the physical properties of the gas. At high frequencies, the observed line widths provide limits on the gas temperature and on the turbulent velocity of the cloud. At low frequencies, the observed line widths provide constraints on the electron density of the intervening cloud and on the radiation field that the cloud is embedded in. Using the departure coefficients obtained in Paper I, we analyzed the behavior of the lines under the physical conditions of the diffuse ISM. Integrated optical depths provide constraints on the electron density, the electron temperature, and the emission measure or size of the cloud. The use of CRRLs together with [C II] at $158 \mu\text{m}$ can constrain the temperature.

As an illustration of the use of our models, we have analyzed existing data in low-frequency CRRLs toward Cas A and the inner galaxy to derive physical parameters of the absorbing/emitting clouds (Payne et al. 1994; Erickson et al. 1995; Stepkin et al. 2007).

Our models predict that detailed studies of CRRLs should be possible with currently available instrumentation. By using

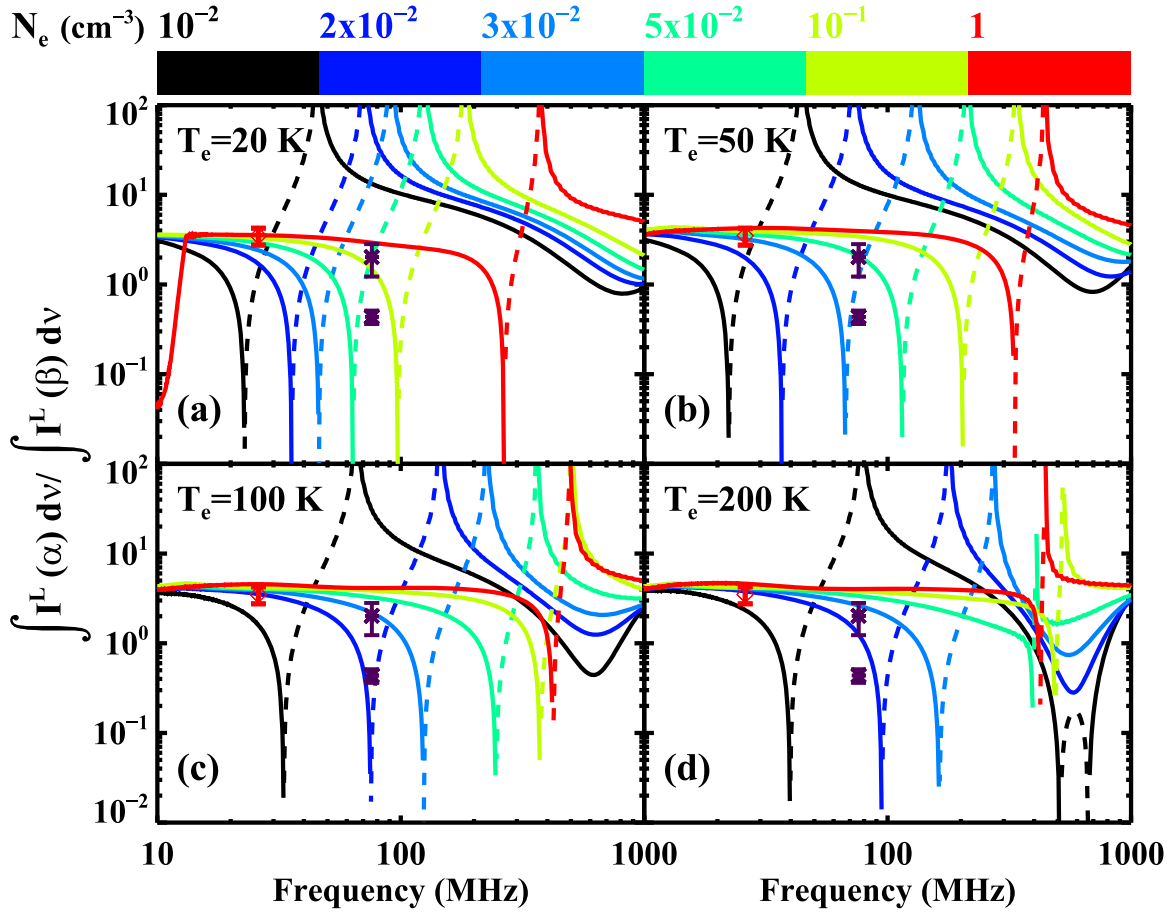


Figure 10. Comparison between the integrated line-to-continuum $I(\alpha)/I(\beta)$ ratio as a function of frequency for different densities (color bar); dashed lines indicate that the ratio is negative, and the colors of the lines are the same as in Figure 7. The values for the ratios approach the LTE value of 3.6 at high n . Large differences can be observed for different densities because lines observed at the same frequency correspond to different levels. We have included the data points for Cas A from Stepkin et al. (2007; red point) and for the inner Galaxy from Erickson et al. (1995; dark blue points).

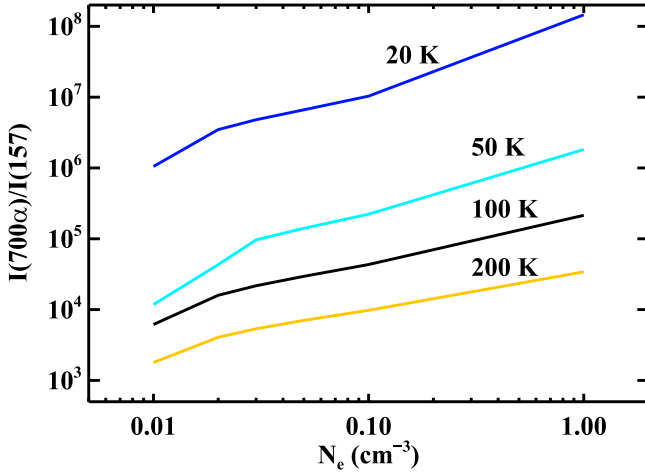


Figure 11. Ratio of the line-to-continuum ratio, for $n = 700$, to the [C II] 158 μm line as a function of density. This example ratio shows how $\text{CRRL}/[\text{C II}]$ can be used as a diagnostic plot to constrain electron density and temperature.

realistic estimates for the properties of the diffuse ISM, we obtain optical depths that are within the capabilities of LOFAR and of the future SKA (Oonk et al. 2015). Given the clumpy nature of the ISM, we encourage observations with high angular resolution. Observations with large beams are biased

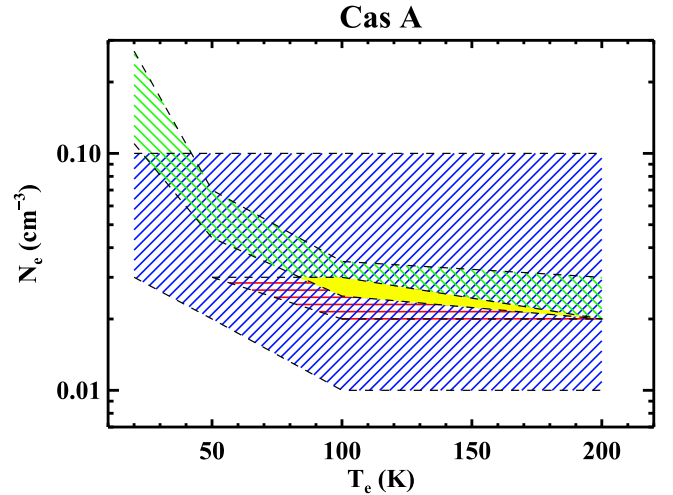


Figure 12. Summary of the constraints for the $Cn\alpha$ and $Cn\beta$ transitions from Stepkin et al. (2007) toward Cas A. The blue zone shows the region allowed by the integrated α -to- β ratio constraints. The green zone is the region allowed from the n_i constraints. The red zone is the region allowed from the 600 to 500 ratio vs. 270 to 500. The yellow zone shows the overlap region from all the constraints. The electron density is well constrained to be $2\text{--}3 \times 10^{-2} \text{ cm}^{-3}$. The temperature is constrained to be within 80 and 200 K.

toward lines of sight with large optical depth and narrow lines, and these happen to be clouds of low density for a given temperature. High spectral resolution is also encouraged in

Table 1
Selected Values for $Cn\alpha$ and $Cn\beta$ Lines for Regions Observed by Erickson et al. (1995)

Name	$\tau(441\alpha)$ $\times 10^{-3}$	$\Delta\nu(441\alpha)$ km s^{-1}	$\tau(555\beta)$ $\times 10^{-3}$	$\Delta\nu(555\beta)$ km s^{-1}
G000.0+0	0.73 ± 0.03	24 ± 1	0.35 ± 0.03	24 ± 2
G002.0-2	0.97 ± 0.08	9 ± 1	0.75 ± 0.04	25 ± 2

order to distinguish multiple components along the lines of sight. Once the temperature and the density have been determined, the observed intensities yield the C+ column density, which can be combined with the H I column density from 21 cm observations to determine the gas-phase carbon abundance.

The main conclusions of our work are as follows:

- (1) CRRLs provide a powerful probe of the physical conditions of diffuse interstellar clouds.
- (2) Meaningful constraints on gas properties can be derived from combining information on the location of the transition from emission to absorption, α -to- β ratios, and α -line ratios spread in frequency. Further limits are provided by the low-frequency line width.
- (3) Comparison of CRRLs with the [C II] 158 μm line measured by COBE (Bennett et al. 1994), BICE (Nakagawa et al. 1998), and Herschel (GOT C+; Pineda et al. 2013), in addition to new observations with the German Receiver for Astronomy at Terahertz Frequencies (Heyminck et al. 2012) on board SOFIA, will provide important constraints primarily on the temperature, but will also aid in further constraining the density and size of diffuse clouds.

Appendix A List of Symbols

A detailed list of the symbols used in this article is presented in Table 2.

Appendix B Radiative Transfer

The radiative transfer equation for a line in the plane-parallel approximation is given by

$$\frac{dI_\nu}{dx}(x) = -k_\nu(x)I_\nu(x) + j_\nu(x) \quad (20)$$

$$k_\nu = k_\nu^l + k_\nu^c \quad (21)$$

$$j_\nu = j_\nu^l + j_\nu^c, \quad (22)$$

where k_ν^l is the line absorption coefficient, k_ν^c is the continuum absorption coefficient, j_ν^l is the line emission coefficient, j_ν^c is the continuum emission coefficient, and $I_\nu(x)$ is the specific intensity of a nebula at a frequency ν as a function of depth in the cloud x . The line absorption and emission coefficients are given by

$$j_\nu^l = \frac{h\nu}{4\pi} A_{n'n} N_{n'} \phi(\nu), \quad (23)$$

$$k_\nu^l = \frac{h\nu}{4\pi} (N_n B_{nn'} - N_{n'} B_{n'n}) \phi(\nu), \quad (24)$$

where $N_{n'}$ is the level population of a given upper level, and N_n is the level population of the lower level; ν is the frequency of the transition, and $A_{n'n}$, $B_{n'n}$ ($B_{nn'}$) are the Einstein coefficients for spontaneous and stimulated emission (absorption), related to each other by

$$A_{n'n} = \frac{2h\nu^3}{c^2} B_{n'n}, \quad (25)$$

$$B_{nn'} = \frac{\omega_{n'}}{\omega_n} B_{n'n}. \quad (26)$$

The factor $\phi(\nu)$ in Equation (23) is the normalized line profile ($\int \phi(\nu) d\nu = 1$). The effects on the emission are analyzed in Section 2.3. Here, we assume that j_ν is evaluated at the line center where the frequency of the transition is ν_0 , and we omit the $\phi(\nu_0)$ factor. Note that due to the normalization, $\phi(\nu_0) < 1$. Under thermodynamic equilibrium, the level population of a level n ($N_n(\text{LTE})$) is given by

$$N_n(\text{LTE}) = N_e N_{\text{ion}} \left(\frac{h^2}{2\pi m_e k T_e} \right)^{1.5} \frac{\omega_n}{2\omega(i)} e^{\chi_n},$$

$$\chi_n = \frac{hcZ^2 R_y}{n^2 k T_e}, \quad (27)$$

where N_e is the electron density in the nebula, T_e is the electron temperature, N_{ion} is the ion density, m_e is the electron mass, h is the Planck constant, k is the Boltzmann constant, c is the speed of light, R_y is the Rydberg constant, and ω_n is the statistical weight of the level n ($\omega_n = 2n^2$, for hydrogen). In the ISM, levels can be out of local thermodynamic equilibrium (Paper I). The level population can then be described by the departure coefficients $b_n = N_n/N_n(\text{LTE})$, that is, the ratio of the level population of a given level to its LTE value. From the definitions of j_ν^l and k_ν^l , we can write the emission and absorption coefficients in terms of the departure coefficients:

$$j_\nu^l = j_\nu^l(\text{LTE}) b_n, \quad (28)$$

$$k_\nu^l = \frac{h\nu}{4\pi} (b_n N_n(\text{LTE}) B_{nn'} - b_{n'} N_{n'}(\text{LTE}) B_{n'n}), \quad (29)$$

$$= k_\nu^l(\text{LTE}) b_n \frac{1 - \frac{b_{n'}}{b_n} e^{-h\nu/kT_e}}{1 - e^{-h\nu/kT_e}}. \quad (30)$$

The correction factor for stimulated emission/absorption, $\beta_{nn'}$, is

$$\beta_{nn'} = \frac{1 - \frac{b_{n'}}{b_n} e^{-h\nu/kT_e}}{1 - e^{-h\nu/kT_e}}. \quad (31)$$

Deviations from equilibrium can also be described in terms of the excitation temperature (T_X) of a transition, defined as

$$\frac{N_{n'}/\omega_{n'}}{N_n/\omega_n} = \exp\left(\frac{-h\nu}{kT_X}\right). \quad (32)$$

It is easy to see that T_X is related to $\beta_{nn'}$ by

$$\beta_{nn'} = \frac{1 - e^{-h\nu/kT_X}}{1 - e^{-h\nu/kT_e}}. \quad (33)$$

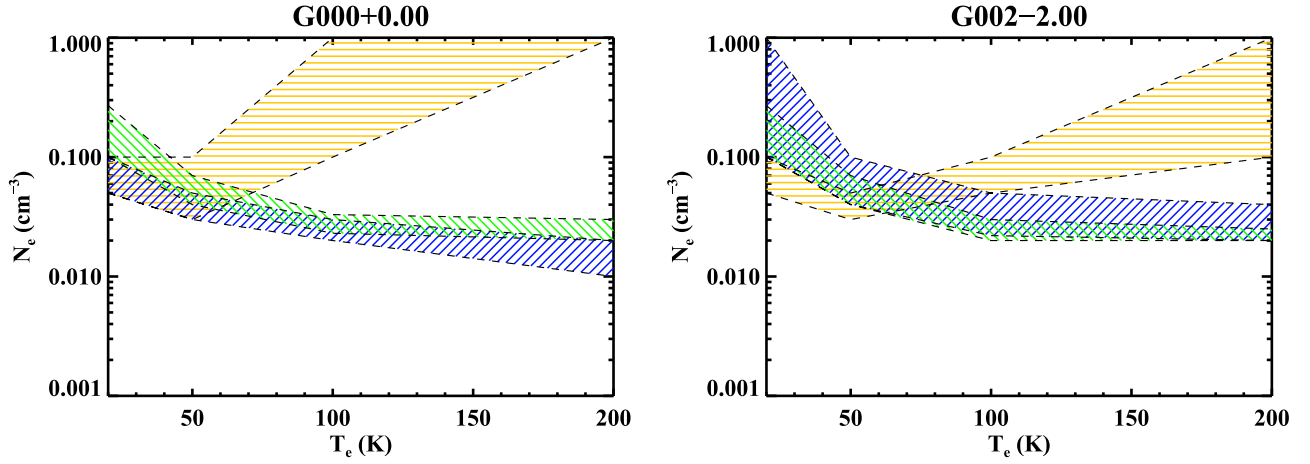


Figure 13. Same as Figure 12 for regions toward the Galactic center (data from Erickson et al. 1995). The α -to- β ratio constraints are shown as a blue region. The constraints derived from n_t are shown as a green zone. In addition, we have added the constraint from the CRRL to [C II] 158 μm ratio as the orange shaded zone.

Clearly, under LTE conditions, the excitation temperature approaches the value of the electron temperature, $T_X = T_e$. The description of the level population in terms of T_X is useful to explain the behavior of the lines, as we show in Section 3. For a homogeneous cloud, the radiative transfer equation can be solved. At a given frequency, the observed flux has contributions from both the line and the continuum, which can be written as

$$I_\nu^{\text{total}} = \frac{j_\nu^c + j_\nu^l}{k_\nu^c + k_\nu^l} [1 - e^{-(\tau_\nu^c + \tau_\nu^l)}] + I_0(\nu) e^{-(\tau_\nu^c + \tau_\nu^l)}, \quad (34)$$

$$I_\nu^c = \frac{j_\nu^c}{k_\nu^c} (1 - e^{-\tau_\nu^c}) + I_0(\nu) e^{-\tau_\nu^c}, \quad (35)$$

where a background continuum source $I_0(\nu)$ has been introduced. The coefficients $\tau_\nu^x = \int k_\nu^x(s) ds$ are the optical depth for x of either the continuum or the line. Assuming homogeneity, $\tau_\nu^x = k_\nu^x L$, where L is the length along the line of sight of the cloud, we can separate the contribution from the line itself since it is given by

$$I_\nu^{\text{line}} = I_\nu^{\text{total}} - I_\nu^c \quad (36)$$

$$\begin{aligned} I_\nu^{\text{total}} - I_\nu^{\text{continuum}} &= \frac{j_\nu^c + j_\nu^l}{k_\nu^c + k_\nu^l} (1 - e^{-\tau_\nu^{\text{total}}}) + I_0(\nu) e^{-\tau_\nu^{\text{total}}} \\ &\quad - \frac{j_\nu^c}{k_\nu^c} (1 - e^{-\tau_\nu^c}) - I_0(\nu) e^{-\tau_\nu^c} \end{aligned} \quad (37)$$

$$\tau_\nu^{\text{total}} = \tau_\nu^c + \tau_\nu^l. \quad (38)$$

We can write the line contribution in terms of the source function (S_ν) by using Kirchoff's law [$j_\nu = \kappa_\nu B_\nu(T_e)$, with $B_\nu(T_e)$ the Planck function]:

$$\begin{aligned} S_\nu &= \frac{j_\nu^c + j_\nu^l}{k_\nu^c + k_\nu^l} \\ &= \left[\frac{k_\nu^c + b_n k_\nu^l(\text{LTE})}{k_\nu^c + b_n \beta_{nn'} k_\nu^l(\text{LTE})} \right] B_\nu(T_e). \end{aligned} \quad (39)$$

We identify a correction factor to the Planck function for departures from LTE:

$$\eta = \frac{k_\nu^c + b_n k_\nu^l(\text{LTE})}{k_\nu^c + b_n \beta_{nn'} k_\nu^l(\text{LTE})}, \quad (40)$$

as in, for example, Strelitski et al. (1996) and Gordon & Sorooshenko (2009).

With these definitions, we can write

$$\begin{aligned} I_\nu^{\text{line}} &= \eta B_\nu(T_e) (1 - e^{-\tau_\nu^{\text{total}}}) - B_\nu(T_e) (1 - e^{-\tau_\nu^c}) + \\ &\quad + I_0(\nu) e^{-\tau_\nu^c} (e^{-\tau_\nu^l} - 1), \end{aligned} \quad (41)$$

and the intensity of a line relative to the continuum is

$$\frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} = \frac{\eta B_\nu(T_e) (1 - e^{-\tau_\nu^{\text{total}}}) + I_0(\nu) e^{-\tau_\nu^{\text{total}}} - B_\nu(T_e) (1 - e^{-\tau_\nu^c}) - I_0(\nu) e^{-\tau_\nu^c}}{B_\nu(T_e) (1 - e^{-\tau_\nu^c}) + I_0(\nu) e^{-\tau_\nu^c}} - 1. \quad (42)$$

In the absence of a background radiation field ($I_0 = 0$), this reduces to

$$\frac{I_\nu^{\text{line}}}{I_\nu^{\text{cont}}} = \frac{\eta (1 - e^{-\tau_\nu^{\text{total}}})}{(1 - e^{-\tau_\nu^c})} - 1. \quad (43)$$

In Section 2.3, we showed that under the conditions of the diffuse ISM the line profile is expected to be Lorentzian in shape, and, at the line center, $\phi(\nu_0) = 2/\pi \Delta\nu_L$ with $\Delta\nu_L$ the FWHM of the line. This sets a range of physical parameters for which the approximation $|\tau_\nu^l| \ll 1$ is valid.

B.1. Doppler and Lorentzian Broadening

Doppler broadening occurs due to turbulent motions in the gas, and thermal motions are given by a Gaussian distribution with a Doppler width (Rybicki & Lightman 1986):

$$\Delta\nu_D = \frac{\nu_0}{c} \sqrt{\frac{2kT}{m_{\text{atom}}}} + \langle v_{\text{rms}} \rangle^2, \quad (44)$$

where m_{atom} is the mass of the atom and $\langle v_{\text{rms}} \rangle$ is the rms turbulent velocity. The line profile as a function of frequency is given by the expression

$$\phi_\nu^G(\nu) = \frac{1}{\Delta\nu_D \sqrt{\pi}} \exp \left(- \left(\frac{\nu - \nu_0}{\Delta\nu_D} \right)^2 \right). \quad (45)$$

Table 2
List of Symbols

Symbol	Description
$A_{n'n}$	Einstein coefficient for spontaneous transitions
$A_{3/2,1/2}$	Spontaneous transition rate of the carbon fine structure line $^2P_{3/2}-^2P_{1/2}$
$a(T_e)$	Fitting coefficient for collisional broadening
$B_{nn'}$	Einstein coefficient for stimulated transition
b_n	Departure coefficient for level n
$B_\nu(T)$	Planck function at frequency ν for a temperature T
$C_{n'n}$	Rates for energy-changing collisions between level n' and n
$Cn\alpha$	Carbon recombination line with $\Delta n = 1$
c	Speed of light
EM_{C+}	Emission measure of carbon ions
h	Planck constant
$I_0(\nu)$	Intensity of the background continuum
I_ν^{line}	Intensity of the line
I_ν^{cont}	Intensity of the continuum
I_{158}	Intensity of the fine structure line of carbon at 158 μm
j_ν^l	Line emission coefficient
j_ν^c	Continuum emission coefficient
k	Boltzmann constant
k_ν^l	Line absorption coefficient
k_ν^c	Continuum absorption coefficient
L	Path length of the cloud
$M(\Delta n)$	Approximation factor for the oscillator strength, as given by Menzel (1968)
m_C	Mass of a carbon atom
$N_{3/2}^+$	Level population of carbon ions in the $^2P_{3/2}$ core
N_{C+}	Number density of carbon ions
N_e	Electron density
\mathcal{N}_{C+}	Carbon column density
\mathcal{N}_H	Hydrogen column density
n	Lower principal quantum number
n'	Upper principal quantum number
n_i	Level where the observed lines transition from emission to absorption
R	Ratio between the fine structure ($^2P_{3/2}-^2P_{1/2}$) level population and the fine structure level population in LTE
Ry	Rydberg constant
T_0	Temperature of power law background spectrum at frequency ν_0
t_n	Ratio of radiation to collisional broadening
T_X	Excitation temperature
T_e	Electron temperature
$\langle v_{\text{rms}} \rangle$	rms turbulent velocity
$\alpha_{1/2}$	Fraction of carbon ions in the $^2P_{1/2}$ level
α_{pl}	Exponent of the power-law background spectrum
$\beta_{nn'}$	Correction factor for stimulated emission
β_{158}	Correction for simulated emission to the [C II] fine structure line $^2P_{3/2}-^2P_{1/2}$
$\gamma_c(T_e)$	Fitting coefficient for collisional broadening
γ_e	De-excitation rate for carbon ions in the $^2P_{3/2}$ core due to collisions with electrons
γ_H	De-excitation rate for carbon ions in the $^2P_{3/2}$ core due to collisions with hydrogen atoms
Δn	$n'-n$, difference between the upper and lower principal quantum number
$\Delta\nu_D$	Doppler width
$\Delta\nu_{\text{rad}}$	Radiation broadening
$\Delta\nu_{\text{col}}$	Collisional broadening
ν	Frequency of a transition
η	Correction factor to the Planck function due to non-LTE level population
τ_{158}	Optical depth for the [C II] fine structure line $^2P_{3/2}-^2P_{1/2}$

Table 2
(Continued)

Symbol	Description
τ_ν^l	Optical depth of the line
τ_ν^c	Optical depth of the continuum
τ_ν^{total}	Sum of τ_ν^l and τ_ν^c
$\phi(\nu)$	Line profile
ν_0	Reference frequency for the power-law background spectrum
χ_n	hcZ^2Ry/n^2kT_e

With this definition, the FWHM is $\Delta\nu_D(\text{FWHM}) = 2\sqrt{\ln(2)}\Delta\nu_D$. Note that Doppler broadening is dominated by turbulence for $T_e < 60.5 (m_{\text{atom}}/m_H)(\langle v_{\text{rms}} \rangle/\text{km s}^{-2})^2 \text{ K}$, where m_H is the mass of a proton.

The Lorentzian width of a line produced by a transition from a level n' to n is related to the net transition out of the levels (Shaver 1975; Rybicki & Lightman 1986):

$$\Gamma_{n'n} = \Gamma_{n'} + \Gamma_n, \quad (46)$$

$$\Gamma_{n'} = \sum_{n < n'} A_{n'n} + \sum_{n \neq n'} N_e C_{n'n} + \sum_{n \neq n'} B_{n'n} I_\nu, \quad (47)$$

$$= \Gamma_{\text{natural}} + \Gamma_{\text{collisions}} + \Gamma_{\text{radiation}} \quad (48)$$

with an analogous formula for Γ_n . Here, we have to consider collisions with electrons and transitions induced by an external radiation field. This produces a Lorentzian line profile:

$$\phi_\nu^L(\nu) = \frac{\gamma}{\pi(\nu - \nu_0)^2 + \gamma^2}, \quad (49)$$

and the FWHM is $\Delta\nu_L(\text{FWHM}) = 2\gamma$. The width γ of a line transition between levels n and n' is given by $\gamma = (\Gamma_n + \Gamma_{n'})/4\pi$. For transitions between lines with $n \approx n'$, we have $\Gamma_n \approx \Gamma_{n'}$ and $\gamma \approx \Gamma_n/2\pi$.

In the most general case, the line profile is given by the Voigt profile, that is, the convolution of the Gaussian and the Lorentzian profile (e.g., Gordon & Sorochenko 2009):

$$\phi_\nu^V(\nu) = \int_{-\infty}^{\infty} \phi_\nu^L(\nu) \phi_\nu^G(\nu) d\nu. \quad (50)$$

This can be written in terms of the Voigt function $[H(a, u)]$ by using the proper normalization:

$$\phi_\nu^V(\nu) = \frac{1}{\Delta\nu_D \sqrt{\pi}} H(a, u) \quad (51)$$

$$H(a, u) = \frac{a}{\pi} \int_{-\infty}^{\infty} \frac{e^{-y^2} dy}{a^2 + (y - u)^2}, \quad (52)$$

with $a = \gamma/\Delta\nu_D$ and $u = (\nu - \nu_0)/\Delta\nu_D$. The FWHM of the Voigt profile can be approximated by

$$\Delta\nu_V(\text{FWHM}) = 0.5346\Delta\nu_L(\text{FWHM}) + \sqrt{0.2166\Delta\nu_L(\text{FWHM})^2 + \Delta\nu_D(\text{FWHM})^2}. \quad (53)$$

B.2. Collisional/Stark Broadening

Collisions with electrons produce line broadening:

$$\Delta\nu_{\text{col}} = \frac{2}{\pi} \sum_{n \neq n'} N_e C_{n'n}, \quad (54)$$

Table 3
Coefficients for Equation (16)

T_e (K)	a	γ_c
10	-10.97	5.482
20	-10.67	5.435
30	-10.49	5.407
40	-10.37	5.386
50	-10.27	5.369
60	-10.19	5.354
70	-10.12	5.341
80	-10.06	5.329
90	-10.01	5.318
100	-9.961	5.308
200	-9.620	5.228
300	-9.400	5.170
400	-9.234	5.122
500	-9.085	5.077
600	-8.969	5.041
700	-8.869	5.009
800	-8.780	4.980
900	-8.701	4.953
1000	-8.630	4.929
2000	-8.272	4.806
3000	-8.009	4.706
4000	-7.834	4.636
5000	-7.708	4.583
6000	-7.613	4.542
7000	-7.538	4.509
8000	-7.477	4.482
9000	-7.427	4.458
10000	-7.386	4.439
20000	-7.181	4.329
30000	-7.113	4.281

where $C_{n'n}$ is the collision rate for electron-induced transitions from level n' to n , and N_e is the electron density. For levels $n > 100$, we fitted the following function to depopulating collisions:

$$\sum_{n \neq n'} N_e C_{n'n} = N_e 10^a \times n^{\gamma_c}. \quad (55)$$

Values for a and γ_c are given in Table 3. The values used here agree with those from Shaver at low temperatures, but at temperatures larger than about 1000 K they can differ by factors larger than about 4. The values presented here agree well with those of Griem (1967) at large temperatures. At low frequencies, collisional broadening is large and dominates over Doppler broadening in the absence of a background radiation field. As can be seen from the dependence on the electron density, clouds with higher densities have broader lines than those with lower densities at a given level n .

Appendix C Radiation Broadening

The depopulation of a given level n' due to stimulated transitions is given by

$$\Gamma_{n'}^{\text{radiation}} = \sum_{n \neq n'} B_{n'n} I_\nu, \quad (56)$$

where $B_{n'n}$ is the Einstein B coefficient for stimulated transitions from level n' to n , and I_ν is an external radiation field.

We can write the Einstein $B_{n'n}$ coefficients in terms of the Einstein $A_{n'n}$ coefficients:

$$B_{n+\Delta n,n} I_\nu = \frac{c^2}{2h\nu^3} A_{n+\Delta n,n} I_\nu, \quad (57)$$

(e.g., Shaver 1975 and Gordon & Soroichenko 2009) where we have used the $n' = n + \Delta n$. Assuming a power-law-like radiation field, with temperature $T_R = T_0(\nu/\nu_0)^{\alpha_{\text{pl}}}$ we can write

$$B_{n+\Delta n,n} I_\nu = \frac{kT_0}{h\nu_0^{\alpha_{\text{pl}}}} A_{n+\Delta n,n} \nu^{\alpha_{\text{pl}}-1}. \quad (58)$$

The Einstein A coefficient can be written in terms of the oscillator strength, $f_{n,n+\Delta n}$ (e.g., Shaver 1975):

$$A_{n+\Delta n,n} = \frac{8\pi^2 e^2 \nu^2}{m_e c^3} \left(\frac{n}{n + \Delta n} \right)^2 f_{n,n+\Delta n}, \quad (59)$$

Replacing in Equation (58) leads to

$$\begin{aligned} B_{n+\Delta n,n} I_\nu &= \frac{kT_0}{h\nu_0^{\alpha_{\text{pl}}}} \frac{8\pi^2 e^2 \nu^2}{m_e c^3} \left(\frac{n}{n + \Delta n} \right)^2 f_{n,n+\Delta n} \nu^{\alpha_{\text{pl}}-1}, \\ &= \frac{8\pi^2 e^2}{m_e c^3} \left(\frac{kT_0}{h\nu_0^{\alpha_{\text{pl}}}} \right) \left(\frac{n}{n + \Delta n} \right)^2 f_{n,n+\Delta n} \nu^{\alpha_{\text{pl}}+1}. \end{aligned} \quad (60)$$

Menzel (1968) gives a simple approximation for computing the oscillator strength:

$$\frac{f_{n+\Delta n,n}}{n} \approx M(\Delta n) \left(1 + \frac{3}{2} \frac{\Delta n}{n} \right), \quad (61)$$

with $M(\Delta n) = 4/3 J_{\Delta n}(\Delta n) J'_{\Delta n}(\Delta n) / \Delta n^2$, where $J_{\Delta n}(\Delta n)$ is the Bessel function of order Δn . The $M(\Delta n)$ can be approximated by $M(\Delta n) \approx 0.1908/\Delta n^3$ to less than 16% in accuracy for $\Delta n = 5$, and to better than 3% accuracy by changing the exponent from 3 to 2.9. The values for $M(\Delta n) = 0.1908, 0.02633, 0.008106, 0.003492, 0.001812$ are for $\Delta n = 1, 2, 3, 4, 5$.

The frequency of a line in the hydrogenic approximation is given by

$$\nu_{n+\Delta n,n} = R_{\text{yc}} Z^2 \left(\frac{1}{n^2} - \frac{1}{(n + \Delta n)^2} \right), \quad (62)$$

$$\approx 2R_{\text{yc}} Z^2 \frac{\Delta n}{n^3} \left(1 - \frac{3}{2} \frac{\Delta n}{n} \right), \quad (63)$$

(e.g., Shaver 1975; Gordon & Soroichenko 2009). Replacing ν in Equation (60) and for $Z = 1$, we obtain

$$\begin{aligned} B_{n+\Delta n,n} I_\nu &= \frac{8\pi^2 e^2}{m_e c^3} \left(\frac{kT_0}{h\nu_0^{\alpha_{\text{pl}}}} \right) \left(\frac{n}{n + \Delta n} \right)^2 M(\Delta n) \\ &\times \left(1 + \frac{3}{2} \frac{\Delta n}{n} \right) n \left[2R_{\text{yc}} \frac{\Delta n}{n^3} \left(1 - \frac{3}{2} \frac{\Delta n}{n} \right) \right]^{\alpha_{\text{pl}}+1}, \\ &\approx \frac{8\pi^2 e^2 (2R_{\text{yc}})^{\alpha_{\text{pl}}+1} kT_0}{m_e c^3 h\nu_0^{\alpha_{\text{pl}}}} M(\Delta n) n \left(\frac{\Delta n}{n^3} \right)^{\alpha_{\text{pl}}+1}, \end{aligned}$$

for $\Delta n/n \ll 1$. Rearranging the expression, we arrive at

$$\begin{aligned} B_{n+\Delta n,n} I_\nu &= \frac{8\pi^2 e^2 (2Ryc)^{\alpha_{pl}+1} kT_0}{m_e c^3 h \nu_0^{\alpha_{pl}}} M(\Delta n) n^{-3\alpha_{pl}-2} \Delta n^{\alpha_{pl}+1}, \\ &= \frac{8\pi^2 e^2 (2Ryc)^{\alpha_{pl}+1} kT_0}{m_e c^3 h \nu_0^{\alpha_{pl}}} 0.1908 n^{-3\alpha_{pl}-2} \Delta n^{\alpha_{pl}+2}, \\ &= 2.137 \times 10^4 \left(\frac{6.578 \times 10^{15}}{\nu_0} \right)^{\alpha_{pl}+1} \\ &\quad \times kT_0 \nu_0 n^{-3\alpha_{pl}-2} \Delta n^{\alpha_{pl}+2}. \end{aligned} \quad (64)$$

Evaluating Equation (64) for $T_0 = 22.6 \times 10^3$ K, $\nu_0 = 30$ MHz, $\alpha_{pl} = -2.55$ at $n = 100$, and $\Delta n = 1$, we recover formula 2.177 of Gordon & Sorochenko (2009). Assuming $\alpha_{pl} = -2.6$ at a reference frequency of 100 MHz, Equation (64) is

$$B_{n+\Delta n,n} I_\nu = 0.662 kT_0 n^{5.8} \Delta n^{-4.6} (s^{-1}). \quad (65)$$

Other relevant values are for an optically thick thermal source $\alpha_{pl} = 0$ and an optically thin thermal source $\alpha_{pl} = -2.1$.

The broadening due to a radiation field in terms of the FWHM is

$$\Delta\nu_{\text{rad}}(\text{FWHM}) = \frac{2}{\pi} \sum_{\Delta n} B_{n+\Delta n,n} I_\nu, \quad (66)$$

$$\approx 5.819 \times 10^{-17} T_0 n^{5.8} (1 + 2^{-4.6} + 3^{-4.6}), \quad (67)$$

$$= 6.096 \times 10^{-17} T_0 n^{5.8} (s^{-1}). \quad (68)$$

C.1. The Far-infrared Fine Structure Line of C+

The beam-averaged optical depth of the fine structure line of carbon ions for the transition $^2P_{1/2} - ^2P_{3/2}$ is given by Crawford et al. (1985) and Sorochenko & Tsivilev (2000):

$$\tau_{158} = \frac{c^2}{8\pi\nu^2} \frac{A_{3/2,1/2}}{1.06\Delta\nu} 2\alpha_{1/2}\beta_{158}N_{C^+}L, \quad (69)$$

where $A_{3/2,1/2} = 2.4 \times 10^{-6} s^{-1}$, ν is the frequency of the $^2P_{1/2} - ^2P_{3/2}$ transition, and $\Delta\nu$ is the FWHM of the line. The $\alpha_{1/2}$ and β_{158} defined by Sorochenko & Tsivilev (2000) are

$$\alpha_{1/2} = \frac{1}{1 + 2\exp(-92/T_e)R}, \quad (70)$$

$$\beta_{158} = 1 - \exp(-92/T_e)R. \quad (71)$$

The definition of R is (Ponomarev & Sorochenko 1992; Payne et al. 1994; see also Paper I)

$$R = \frac{N_e\gamma_e + N_H\gamma_H}{N_e\gamma_e + N_H\gamma_H + A_{3/2,1/2}}, \quad (72)$$

where γ_e and γ_H are the de-excitation rates due to electrons and hydrogen atoms, respectively. For consistency we used the same values as in Paper I and neglected collisions with H_2 .

For the physical conditions considered here, we note that the far-infrared [C II] line is optically thin for hydrogen column densities of $\sim 1.2 \times 10^{21} \text{ cm}^{-2}$. This corresponds to hydrogen densities of about 400 cm^{-3} and electron densities of $6 \times 10^{-2} \text{ cm}^{-3}$, assuming a length of the cloud of 1 pc and width of 2 km s^{-1} . The intensity of the [C II] 158 μm line in the

optically thin limit is given by

$$\begin{aligned} I_{158} &= \frac{h\nu}{4\pi} A_{3/2,1/2} N_{3/2}^+ \times L, \\ &= \frac{h\nu}{4\pi} \frac{A_{3/2,1/2} 2\exp(-92/T_e)R}{1 + 2\exp(-92/T_e)R} \mathcal{N}_{C^+}, \end{aligned} \quad (73)$$

with $N_{3/2}^+$ the number density of carbon ions in the 3/2 state, L the path length through the cloud along the line of sight, and \mathcal{N}_{C^+} the column density of carbon ions. Considering radiative transfer effects, the intensity of the line is given by

$$I_{158} = \frac{2h\nu_0}{\lambda^2} \frac{1.06\Delta\nu(\text{FWHM})}{e^{92/T_{158}} - 1}, \quad (74)$$

with T_{158} defined as

$$T_{158} = \frac{92}{\ln[(e^{92/T_e} e^{\tau_{158}}/R - 1)/(e^{\tau_{158}} - 1)]}. \quad (75)$$

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