

## SDSS IV MANGA: DISCOVERY OF AN H $\alpha$ BLOB ASSOCIATED WITH A DRY GALAXY PAIR – EJECTED GAS OR A ‘DARK’ GALAXY CANDIDATE?

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### ABSTRACT

We report the discovery of a mysterious giant H $\alpha$  blob that is  $\sim 8$  kpc away from the main MaNGA target 1-24145, one component of a dry galaxy merger, identified in the first-year SDSS-IV MaNGA data. The size of the H $\alpha$  blob is  $\sim 3$ -4 kpc in radius, and the H $\alpha$  distribution is centrally concentrated. However, there is no optical continuum counterpart in deep broadband images reaching  $\sim 26.9$  mag arcsec<sup>-2</sup> in surface brightness. We estimate that the masses of ionized and cold gases are  $3.3 \times 10^5 M_{\odot}$  and  $< 1.3 \times 10^9 M_{\odot}$ , respectively. The emission-line ratios indicate that the H $\alpha$  blob is photoionized by a combination of massive young stars and AGN. Furthermore, the ionization line ratio decreases from MaNGA 1-24145 to the H $\alpha$  blob, suggesting that the primary ionizing source may come from MaNGA 1-24145, likely a low-activity AGN. Possible explanations of this H $\alpha$  blob include AGN outflow, the gas remnant being tidally or ram-pressure stripped from MaNGA 1-24145, or an extremely low surface brightness (LSB) galaxy. However, the stripping scenario is less favoured according to galaxy merger simulations and the morphology of the H $\alpha$  blob. With the current data, we can not distinguish whether this H $\alpha$  blob is ejected gas due to a past AGN outburst, or a special category of ‘ultra-diffuse galaxy’ (UDG) interacting with MaNGA 1-24145 that further induces the gas inflow to fuel the AGN in MaNGA 1-24145.

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## 1. INTRODUCTION

It has been known that the observed number density of satellite galaxies in the Local Group and our own Milky Way is orders of magnitude lower than the predictions from the cosmological simulations, the so-called ‘missing satellite problem’ (Klypin et al. 1999; Moore et al. 1999). One of the popular solutions to this problem is that low-mass halos fail to form stars efficiently such that they are below the detection limit of most imaging surveys. Under this paradigm, ‘dark’ galaxies are gas-rich galaxies that do not emit sufficient optical light due to their low efficiency in forming stars, but are thought to be building blocks of normal star-forming galaxies. This type of galaxy is an ideal laboratory to study early stages of star formation, which could lead to understanding how the star formation is triggered.

It is in general very difficult to identify dark galaxies since they are extremely faint in the optical. Very few dark candidates have been identified and confirmed to date. The ALFALFA (Arecibo Legacy Fast ALFA) HI survey (Giovanelli et al. 2005) has discovered on the order of  $\sim 200$  HI sources not associated with apparent optical counterparts, but most of them are likely to have tidal origins (Haynes et al. 2011). A pilot study suggested that none of remaining objects they have explored is a dark galaxy after cross-checking with data in other wavelengths (Cannon et al. 2015).

Recently, van Dokkum et al. (2015) have identified a new class of low surface brightness (LSB) dwarf galaxies in the Coma cluster, often referred to as ‘ultra-diffuse’ galaxies (UDGs), using the Dragonfly Telescope Array (Abraham & van Dokkum 2014). These UDGs not only have a low surface brightness ( $24 - 26$  mag arcsec $^{-2}$ ), but also show extended structures with size comparable to  $L^*$  galaxies. Although these UDGs are found in special environments and may not necessarily represent the global dwarf population, the discovery of this type of galaxy suggests we might have missed numerous faint dwarfs due to observation limitations in the past.

In addition to the aforementioned efforts, the optical Integral Field Unit (IFU) observations open a plausible window to probe the dwarf populations. Normally, emission lines ionized by star-forming regions, AGNs, or shocks are stronger than the stellar continuum and hence can be easily detected with reasonable integration time. With the large area covered by IFU, the structures of ionized gas can be probed out to several tens of kpc. Isolated ionized gases that are separate from a nearby galaxy or a quasar have been readily studied in IFU as well as spectroscopic observations (Fu & Stockton 2007a,b, 2008; Lintott et al. 2009; Husemann et al. 2010; Keel et al. 2012; Cheung et al. 2016b). Although the majority of those ionized gas are suggested or inferred to have external origins, such as the result of gas accretion or minor mergers, their nature remains unknown.

Here we report the discovery of a giant  $H\alpha$  blob which does not have any optical counterpart in deep CFHT *gri* images. However, the morphology and emission line analyses suggest that it could either be ejected gas due to past AGN activity or a special type of UDGs (or ‘dark’ galaxies). In §2, we describe the multi-wavelength data for this system. We present the main results in §3. Section 4 discusses the plausible origins of this system and the important implications of our results. Conclusions are given in §5. Throughout this paper we adopt the following cosmology:  $H_0 = 100h$  km s $^{-1}$  Mpc $^{-1}$ ,  $\Omega_m = 0.3$  and  $\Omega_\Lambda = 0.7$ . We use a Salpeter IMF when deriving

the star formation rate from various observables. We adopt the Hubble constant  $h = 0.7$  when calculating rest-frame magnitudes. All magnitudes are given in the AB system.

## 2. DATA

## 2.1. Optical integral field data

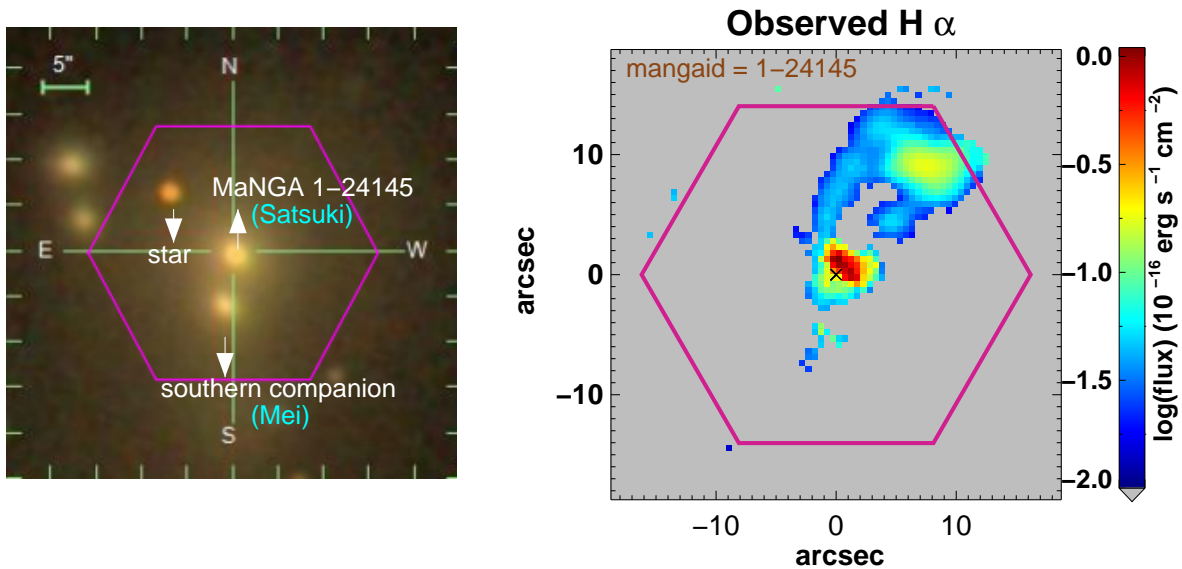
This system, MaNGA 1-24145 ( $z = 0.0322$ , RA = 258.84693, DEC = 57.43288,  $M_* \sim 10^{11} M_\odot$ <sup>27</sup>), was observed in the first 1392 galaxies, as part of the on-going SDSS-IV/MaNGA survey (Bundy et al. 2015; Drory et al. 2015; Law et al. 2016; Yan et al. 2016a,b; SDSS Collaboration et al. 2016). MaNGA is an IFU program to survey for 10k nearby galaxies with a spectral resolution varying from  $R \sim 1400$  at 4000 Å to  $R \sim 2600$  at 9000 Å. The survey uses the BOSS spectrographs (Smee et al. 2013) on the 2.5m Sloan Foundation Telescope (Gunn et al. 2006). The median full width at half maximum (FWHM) of the MaNGA point spread function (PSF) of the datacube is  $\sim 2.5''$ . The MaNGA data used were reduced using the MPL-4 version of the MaNGA data reduction pipeline (Law et al. 2016). The spectral line fitting is carried out using Pipe3D pipeline (Sánchez et al. 2016a). The stellar continuum was first modelled with a linear combination of 12 single stellar population (SSP) templates that were extracted from the MILES project (Sánchez-Blázquez et al. 2006; Vazdekis et al. 2010; Falcón-Barroso et al. 2011). The best-fit stellar continuum is then subtracted from the reduced data spectrum for the emission line measurements. Details of the fitting procedures are described in Sánchez et al. (2016b). To ensure reliable measurements, we restrict our analysis to spaxels where the error-to-flux ratio of the line fitting is less than 1 in subsequent analyses.

To correct for dust reddening, we follow the method described in the Appendix of Vogt et al. (2013) to compute the reddening correction by using the Balmer decrement at each spaxel of the IFU cube. An extinction law with  $R_V = 4.5$  (Fischera & Dopita 2005) is used. The star formation rate (SFR) is then estimated based on this extinction corrected  $H\alpha$  flux.

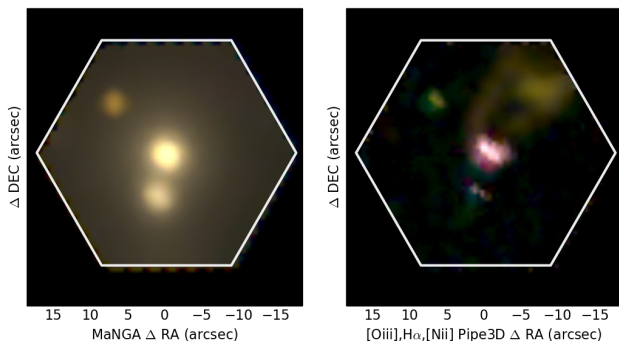
Figure 1 shows the SDSS *gri* composite image (left panel), and the  $H\alpha$  flux map (right panel) of this system (MaNGA 1-24145). The SDSS image indicates that MaNGA 1-24145 (nicknamed ‘Satsuki’) has a companion galaxy (nicknamed ‘Mei’) located in the lower left to MaNGA 1-24145. The stellar absorption lines suggest that these two galaxies are at similar redshifts with the line-of-sight velocities differed by  $\sim 200$  km s $^{-1}$ . In addition, both galaxies are round and red according to the SDSS images. Therefore, these two galaxies likely form a dry (gas-poor) merger system. The compact source in the upper-left corner of the image is a foreground star according to the MaNGA spectrum, and hence it is irrelevant to the dry merger system we are probing for this work.

What is striking about this system is that in the upper-right corner of the  $H\alpha$  map, there exists a giant  $H\alpha$  blob (RA = 258.84314; DEC = 57.43529) with which however does not have any optical counterpart in the SDSS images. The size of the  $H\alpha$  blob is  $\sim 3.2$  kpc in radius. This  $H\alpha$  blob (nicknamed ‘Totoro’) is 7.7 kpc away from Satsuki, and is connected to the  $H\alpha$  emission of Satsuki through tail-like structures. To ensure that the  $H\alpha$  emission at the position of Totoro is not due to artifacts in the data, we check the MaNGA spectra at var-

<sup>27</sup> based on the NASA-Sloan Atlas catalog: <http://www.nsatlas.org>



**Figure 1.** Left: The SDSS *gri* composite image of MaNGA 1-24145 with the MaNGA hexagonal bundle field of view (FoV) overlaid. This system was observed with the 127 fibre bundle of MaNGA, so this hexagon is  $\sim 32.5''$  in diameter. The data extend to regions a bit outside the hexagon because of the dithering. Three distinct objects are visible within the bundle, including two elliptical galaxies (Satsuki and Mei) and one foreground star. Right: the observed H $\alpha$  flux map from the MaNGA observations.



**Figure 2.** Left: The *ugi* composite image reconstructed using the MaNGA continuum with the MaNGA hexagonal bundle field of view (FoV) overlaid. Right: The [OIII]+H $\alpha$ + [NII] composite image. The flux scales of the three lines are adjusted in order to highlight the H $\alpha$  blob.

ious spaxels in the region of Totoro. We found that multiple emission lines, including H $\alpha$ , [NII] 6584, [SII] 6717, 6731, and [OIII] 5007 lines are clearly detected in those regions, suggesting that the H $\alpha$  blob is a real feature. Figure 2 displays the reconstructed *ugr* image from the MaNGA continuum (left panel) and the [OIII] + H $\alpha$  + [NII] composite image (right panel). It can be seen that the reconstructed continuum map from the MaNGA data is as smooth as SDSS. The three upper panels of Figure 3 show the MaNGA spectra of the central pixels for Satsuki, Mei, and Totoro, respectively. The two elliptical galaxies are mainly composed of old stellar populations, consistent with the morphology classification of being early-type galaxies, whereas the blob is dominated by emission lines.

## 2.2. Optical imaging data

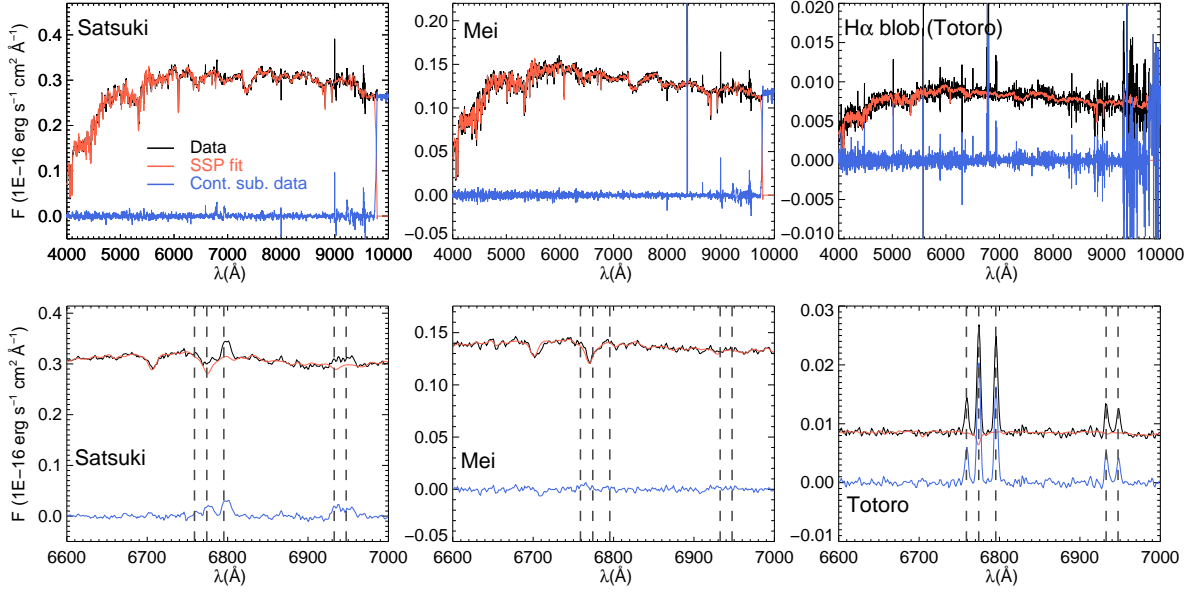
The imaging data come from two sources: the DR12 release of the SDSS photometric survey (York et al. 2000), which reaches  $r \sim 22$ , and deeper observations with

CFHT/MegaCam in *g*, *r* and *i*. The latter combines archival data downloaded from the CADC server and the data taken in 2015 summer (PI: Lihwai Lin) with the Director Discretionary time (DDT) program. All the MegaCam data were processed and stacked via MegaPipe (Gwyn 2008). The final images have  $5\sigma$  limiting mag of 25.7, 26.2, and 25.2 mag (1" aperture in radius) and surface brightness  $5\sigma$  limit of 26.4, 26.9, and 25.9 mag arcsec $^{-2}$  in *g*, *r*, and *i*, respectively.

## 2.3. Radio Continuum

We observed this system with the Karl G. Jansky Very Large Array (VLA) in the A configuration on 2015 August 20 using the C-band receiver tuned to 4–6 GHz ( $\lambda = 7.5$ –5.0 cm). The on-source time is 42 minutes; we observed 3C147 for flux and bandpass calibrations, and J0920+4441 for phase calibration. Data reduction was carried out with CASA (McMullin et al. 2007) using the following steps: (1) standard calibration using the VLA Data Reduction Pipeline (Chandler et al., in prep); (2) removal of any portions of the data corrupted by strong radio frequency interference; and (3) imaging with the task CLEAN. The imaging parameters are the following: MT-MFS deconvolver with nterms of 2, 0.''06 pixel size, and Briggs weighting with robust parameter of 0.5. The final image has a 0.''39  $\times$  0.''35 synthesized beam and rms noise at the pointing center of 7  $\mu$ Jy beam $^{-1}$ .

The radio flux of this source is about  $37 \pm 13 \mu$ Jy at 5 GHz. Assuming that this emission is synchrotron-dominated and so following a power law  $S \propto \nu^\alpha$  with a spectral index  $\alpha = -0.7$ , we derive the radio luminosity of Satsuki to be  $2.2 \times 10^{20}$  WHz $^1$  at 1.4 GHz. This luminosity implies a star-formation rate of 0.12  $M_\odot$  yr $^{-1}$  using the Bell (2003) radio SFR indicator, converted to the Salpeter IMF, which is a factor of 2.5 greater than the limit implied by H $\alpha$  emission at the location of the radio source (see Table 1). Therefore, it is likely that this radio point source is a faint AGN.



**Figure 3.** Upper panels: The black curves represent the MaNGA spectra of the central spaxel of the main galaxy ‘Satsuki’ (left), the southern companion ‘Mei’ (middle), and the  $H\alpha$  blob ‘Totoro’ (right). The red curves are the best-fitted SSP model spectra for the stellar continuum. The blue curves show the residual spectra that are used for the emission line fitting. Lower panels: The zoom-in spectra around the  $H\alpha$  line. The five dashed line denote the observed wavelengths of the [NII] 6584,  $H\alpha$ , [NII] 6584, [SII] 6717, and [SII] 6731 lines, respectively.

#### 2.4. HI

This source was observed as part of the HI-MaNGA programme at the Robert C. Byrd Green Bank Telescope (GBT), which is obtaining HI 21cm observations of a large sample of MaNGA galaxies (AGBT16A\_095, PI: K. Masters). This target was observed on 2016 Feb. 5 for 3 sets of 5 min ON/OFF pairs using the VEGAS spectrometer with a bandwidth of 23.44 MHz, centred on the frequency of 21cm emission redshifted to  $cz=9653 \text{ km s}^{-1}$ . At this frequency the FWHM of the GBT beam is  $9'$ . No HI emission was detected in this volume to a rms of 1.58 mJy (after smoothing to  $5.15 \text{ km s}^{-1}$  velocity resolution). Assuming a velocity spread of 100-400  $\text{km s}^{-1}$  this non-detection sets an  $1-\sigma$  upper limit of  $8.9\text{--}9.2 \times 10^8 M_{\odot}$  for the HI mass of this system<sup>28</sup>.

#### 2.5. X-ray

Satsuki is located within the field of view  $\sim 47$  ks *Chandra* ACIS observation OBSID 4194 (PI: Trevor Ponman), which occurred on 2003 September 17. This pointing of this observation was centered on nearby galaxy NGC 6338 (258.84256, +57.407) [J2000],  $\sim 2'$  South of the MaNGA source of interest. Before analysis, this dataset was reprocessed using the *Chandra\_repro* task in CIAO v4.8 (Fruscione et al. 2006) using CALDB v4.6.3.

The exposure corrected image of this field (Figure 4) was generated using the ‘fluximage’ command, and indicates diffuse X-ray emission coincident with the dry merger system, as also shown in Pandge et al. (2012), which performed a study on a nearby BCG based on the same X-ray dataset. The spectrum of the X-ray emission coincident with the  $H\alpha$  blob of interest was then extracted using the CIAO script *specex-*

*tract* using a source region of a  $22'' \times 18''$  ellipse centered at 17:15:23.7, +57.26:05 which encompasses all of the emission. We fit this X-ray spectrum with a single absorbed APEC model – the emission spectrum from a collisionally-ionized diffuse gas – assuming the redshift  $z = 0.032202$  measured from optical spectroscopy using XSPEC v12.8.2e (Arnaud 1996). This fit results in a reasonable reduced  $\chi^2$  (1.45 for 88 d.o.f), though somewhat under-predicts the flux  $> 5 \text{ keV}$ . As we will mention in Section 3.5, this dry merger system is part of a small group. The derived X-ray temperature is  $1.26 \pm 0.06 \text{ keV}$ , consistent with the temperature on group scale (Kettula et al. 2013). No point-like source is found within Satsuki or Totoro, indicating that there is no strong X-ray AGN present in this system.

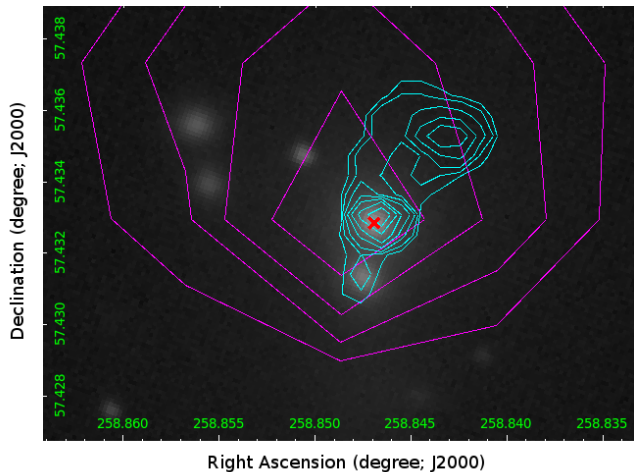
### 3. RESULTS

#### 3.1. The optical morphology

To ensure that the absence of the optical counterpart at the position of Totoro as shown in Figure 1 is not due to relatively shallow depth of SDSS imaging, we carried out a follow-up observation for this system with CFHT/MegaCam and combined it with the archival data (see section 2.2). Figure 5 displays the *gri* composite image of this galaxy. It is clear that there are extended stellar halos surrounding the two galaxies. Again at the position of the  $H\alpha$  blob, no apparent optical continuum is revealed (see the right panel of Figure 5).

To more clearly see the underlying surface brightness structure at the location of Totoro, we build detailed photometric models for this merging system using GALFIT v3.0.2 (Peng et al. 2002, 2010). The MecaCam *g*-band image is used here for this purpose because it may better reveal the continuum of Totoro if has had recent star formation (see §4). Since we are not interested in the overall structures of this complex merging system, we will only focus on a  $400 \times 400$  pixels region centered on the MaNGA bundle. We construct a mask

<sup>28</sup> It is intended that HI-MaNGA data will be released as an SDSS Value Added Catalogue in a future data release from SDSS. In addition the raw data will be publicly available via the NRAO Data Archive at <https://archive.nrao.edu> a year following observations.



**Figure 4.** The Chandra X-ray contours (magenta) and extinction-corrected MaNGA H $\alpha$  contours (cyan) overlaid on the SDSS  $r$ -band image (background image) of MaNGA 1-24145. The red cross marks the position of the VLA point-source detection. The X-ray contours correspond to 44%, 30%, 21%, and 15% of the peak value, respectively, whereas the H $\alpha$  contours correspond to 100%, 60%, 37%, 23%, 14%, 8%, 5%, 3%, and 2% of the peak value, respectively.

image to exclude isolated objects around the merging system from the fitting. The PSF of the image was modelled using the SExtractor and PSFex routines, and the PSF image used in the modelling was extracted using the central coordinate of the MaNGA bundle. Meanwhile, on top of the extended envelope, there are 11 objects (most are galaxies) that can not be easily masked out. We model them separately, then subtract them from the image. All of these smaller objects can be well modelled by a single- or double-Sérsic model locally with the help of an additional Sérsic and sky components to account for the envelope in the background. As shown in the panel (B) of Figure 6, they have been removed smoothly from the input image without any significant residual pattern. Using this “cleaned” image as input, we model the two merging galaxies along with their “common envelope” together using different combinations of Sérsic components. We started with three Sérsic components (one for each galaxy, additional one for the envelope), and gradually build up the complexity by adding more Sérsic component. As we simply want to achieve smooth residual map to study the underlying structure, the number of the Sérsic components and the detailed parameters are not a concern as long as each component behaves normally (e.g. Gu et al. 2013). For each object (including the envelope), all Sérsic components are constrained to have the same center, and only symmetric Sérsic components are used here. After visualizing the residual of initial model, it becomes clear that, on the lower part of the image, there is an additional surface brightness enhancement (caused by the merging process) that are not well fit by the simple model. Eventually, the best model we achieved include 7 Sérsic components plus a tilted-plane sky background component. Both of the merging galaxies, and the extended envelope are described by 2 Sérsic components; while the surface brightness in the south is modelled using single Sérsic component. All components behave regularly in term of size and shape, except for the central component for the brighter galaxy, no component has Sérsic index larger than 2.0.

The panels (C) and (D) of Figure 6 show the model image and the residual maps of this model reconstruction, respectively. It clearly reveals a rich system of shells and tidal tails

around the main MaNGA galaxy, indicating an on-going interaction between these two galaxies. Although the residual map is not perfectly smooth, no optical counterpart is found at the position of Totoro.

Reducing the number of components used, or changing the initial guess will not affect the above conclusion. Adding more components could not further improve the residual map, but results in ill-behaved Sérsic components. Given the complex nature of this merging system, we also try to invoke the asymmetric features in GALFIT, especially the 1st (global lopsidedness) and 4th (boxiness of the isophote) Fourier components (see Peng et al. 2010 for details). The residual map shows improvements around the tidal features, but does not change the conclusion that there is no apparent optical counterpart for Totoro<sup>29</sup>.

### 3.2. Kinematics from the MaNGA Observations

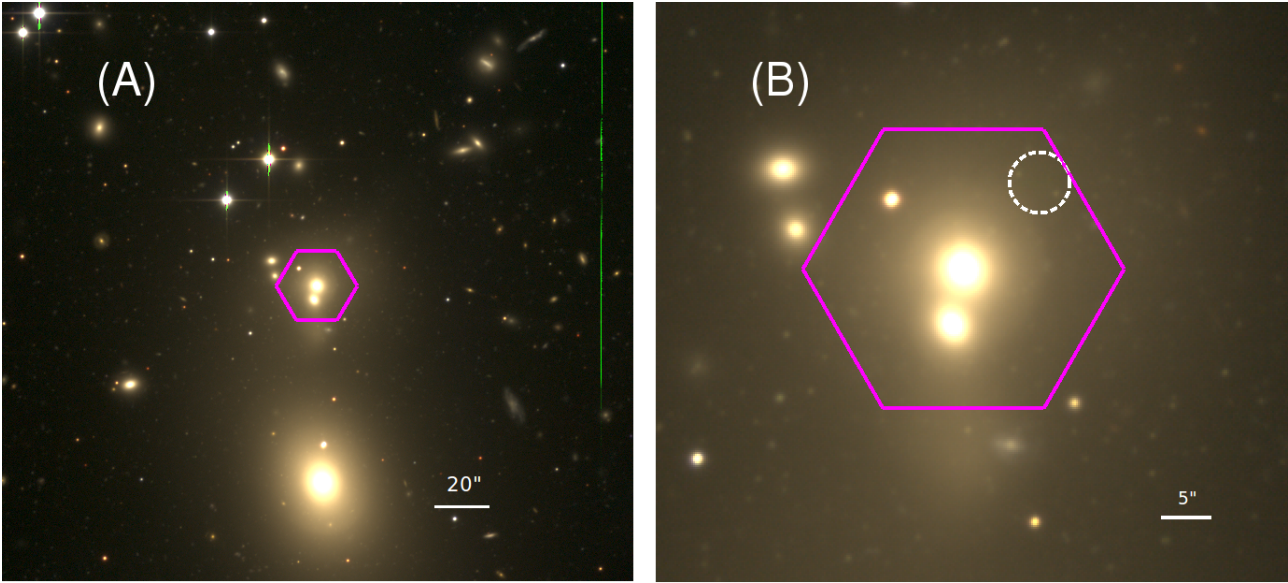
Figure 7 displays the velocity and dispersion maps for this system. While the stellar velocity field indicates that the main galaxy Satsuki is primarily pressure supported, the gas component reveals more complex structures. The central galaxy shows a weak rotation structure, while there is strong variation in the line-of-sight velocity across Totoro region, redshifted in the left tail and blueshifted in the right tail. The inconsistent velocity fields between stars and gas suggests that some part of the gas of Satsuki might have been either accreted or ejected recently. In the former case, it is similar to those early-type galaxies that exhibit misaligned gas and stellar kinematics, the so-called early-type ‘counter rotators’ (Sarzi et al. 2006; Davis et al. 2011; Chen et al. 2016; Jin et al. 2016), although counter rotators in general are defined for systems with a rotating stellar component. The gas inflow scenario is also consistent with the H $\alpha$  morphology which shows bridge (tail)-like structures that connect Totoro and Satsuki. On the other hand, the complicated velocity field can also be explained if the main galaxy Satsuki underwent a strong outburst phase, during which Totoro was expelled from Satsuki.

### 3.3. Excitation State

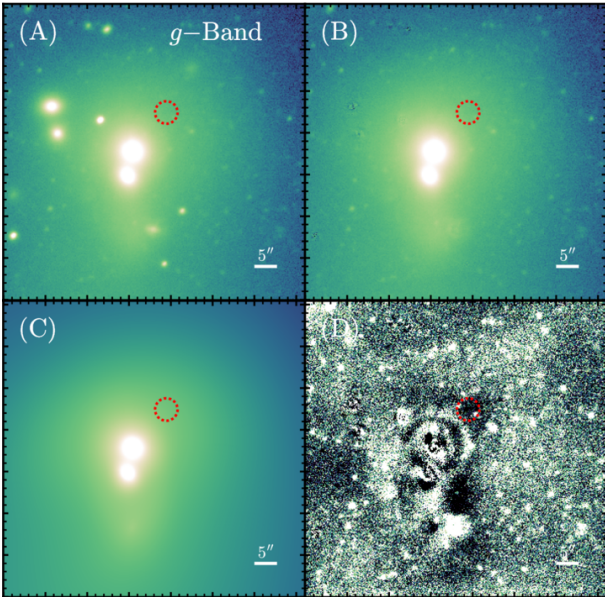
In addition to H $\alpha$  and two [NII] 6584 lines, some other weak lines such as [SII] 6717,6731 and [OIII] 5007 are also detected in both Satsuki and Totoro, allowing us to probe the ionization state for this system. Figure 8 shows the [OIII] 5007/H $\beta$  ratio, one of the frequently used ionization parameters, as a function of the distance from the main galaxy. There exists a strong [OIII] 5007/H $\beta$  gradient, decreasing from the main galaxy Satsuki (location A) to the left bridge (locations B, C, and D) that connects to Totoro (location E, F, G, H, and I). On the other hand, the [OIII] 5007/H $\beta$  ratio is nearly constant across Totoro.

The multiple line detections also allow us to classify the emission line regions into HII or AGN regimes using the standard Baldwin-Phillips-Terlevich (BPT Baldwin, Phillips & Terlevich 1981; Veilleux & Osterbrock 1987; Kauffmann et al. 2003; Kewley et al. 2006) excitation diagnostic diagrams. Here we apply three types of line diagnostics based on four line ratios, [OIII] 5007/H $\beta$ , [NII] 6584/H $\alpha$ , [SII] 6717,6731/H $\alpha$ , and [OI] 6300/H $\alpha$ . Figures 9 and 10 display the line ratio diagrams and classification maps for this system, respectively. We adopt the dividing curves suggested in

<sup>29</sup> The results still hold if we repeat the analysis using the  $i$ -band image, which is more sensitive to the stellar mass distributions.



**Figure 5.** (A) CFHT *gri* composite color image for MaNGA 1-24145 with the MaNGA hexagonal FoV overlaid. A bright BCG is  $\sim 43$  kpc away to the South. (B) A zoom-in picture of (A). The white circle marks the location of Totoro. In both panels, North is up and East is left.



**Figure 6.** (A) The original CFHT Megacam *g*-band image, (B) the *g*-band image after subtracting nearby satellite galaxies, (C) model image for (B), and (D) the residual images produced by GALFIT. The dashed white circles mark the position of Totoro. The white bars in the lower-right corner of each panel corresponds to a scale of  $5''$ .

the literature (e.g., Kewley et al. 2001; Kauffmann et al. 2003; Cid Fernandes et al. 2010) to separate various regions. All the three classifications indicate ‘LINER’-type excitations for Satsuki. As the LINER emission is extended across the entire main galaxy, this object falls in to the extended LIER (eLIER) category according to the classification scheme by Belfiore et al. (2016a,b)

On the other hand, in the regions of Totoro, the line ratios are consistent with the ‘composite’, ‘HI’, and ‘LINER’ regions when using the [NII], [SII], and [OI] diagnostics, respectively. In the remaining part of this paper, we treat Totoro as ‘composite’ regions based on the [NII], as it allows for the intermediate excitation state as opposed to the other two methods.

### 3.4. Gas metallicity

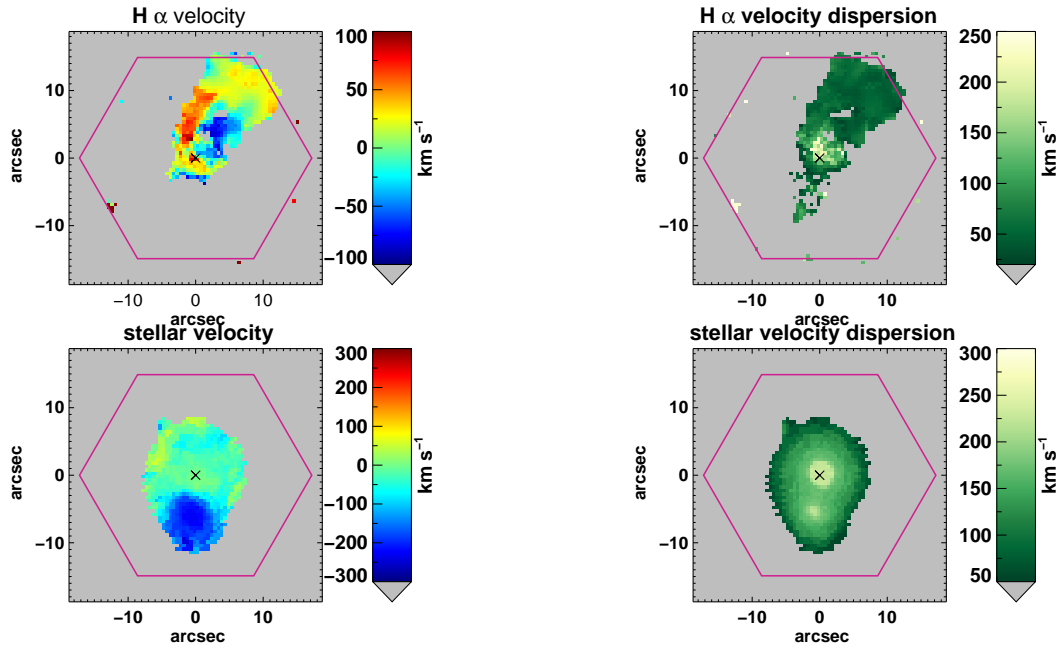
To further understand the properties of Totoro, we also measure the gas-phase metallicity ( $Z$ ) for this system. Conventionally, there are many ways to estimate the gas metallicity through emission line ratios (see Kewley & Ellison 2008). However, most of those tracers are calibrated against the HII regions, and may not be applicable to systems with ionization parameters or interstellar medium (ISM) pressure different from typical HII regions. Here, we adopt the ‘N2S2H $\alpha$ ’ calibrator that is suggested to be less sensitive to the ionization parameters (Dopita et al. 2016) to estimate the metallicity. The N2S2H $\alpha$  calibration can be expressed as the following:

$$12 + \log(\text{O}/\text{H}) = 8.77 + \text{N2S2H}\alpha \quad (1)$$

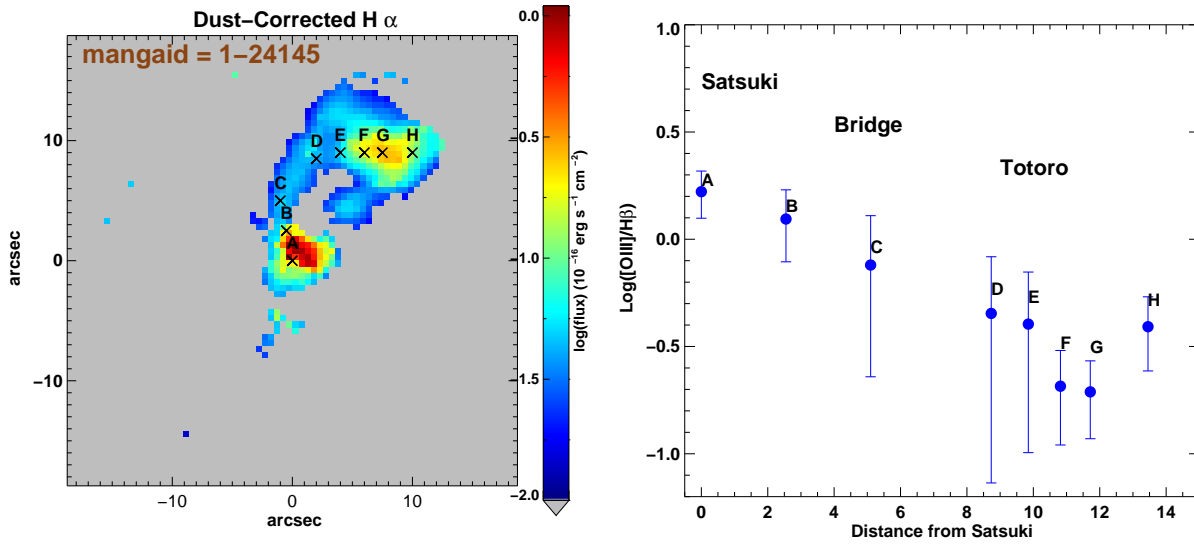
where  $\text{N2S2H}\alpha = \text{Log}([\text{NII}]/[\text{SII}]) + 0.264\text{Log}([\text{NII}]/\text{H}\alpha)$ . The [NII] to [SII] ratio is found to nearly independent of the AGN luminosity and hence provides a good metallicity indicator for even AGNs (Stern & Laor 2013). However, we note with caution that the derived metallicity may still be subject to large systematic uncertainty since the line ratios of the main galaxy and Totoro are located in the LINER and composite regions, respectively.

Figure 11 shows the  $Z(\text{N2S2H}\alpha)$  map for this system. The metallicity of the gas around Satsuki is close to solar and is  $\sim 0.3$  dex lower than the expected value ( $Z = 9.1$ ) for a massive galaxy with a stellar mass of  $10^{11} M_{\odot}$  based on the local mass–metallicity relation (Tremonti et al. 2004). The offset we see here, however, may be easily accounted for by the different metallicity calibrators adopted. On the other hand, the averaged metallicity in the H $\alpha$  region is greater for Totoro than Satsuki by a factor of 0.3 dex and is more consistent with the gas metallicity of high mass galaxies. Although the difference is significantly larger than the statistical uncertainty (0.03–0.1 dex) in the metallicity measurement, it is difficult to interpret this result given the systematic uncertainty in the metallicity measurement due to their different ionization states.

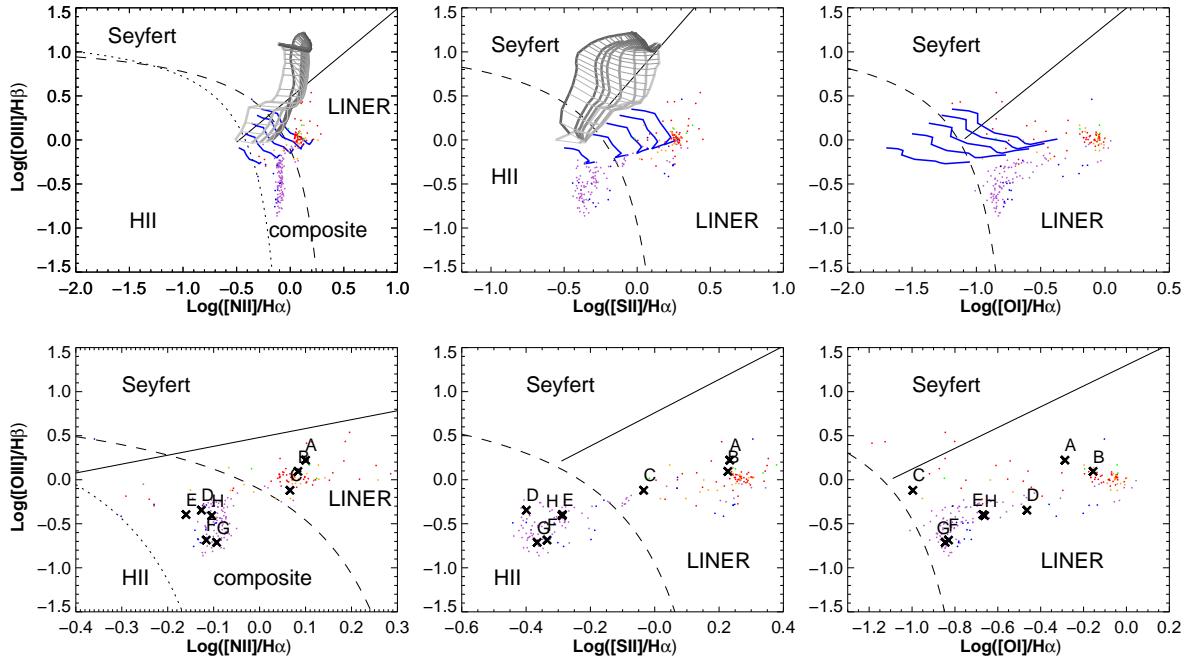
### 3.5. Environment



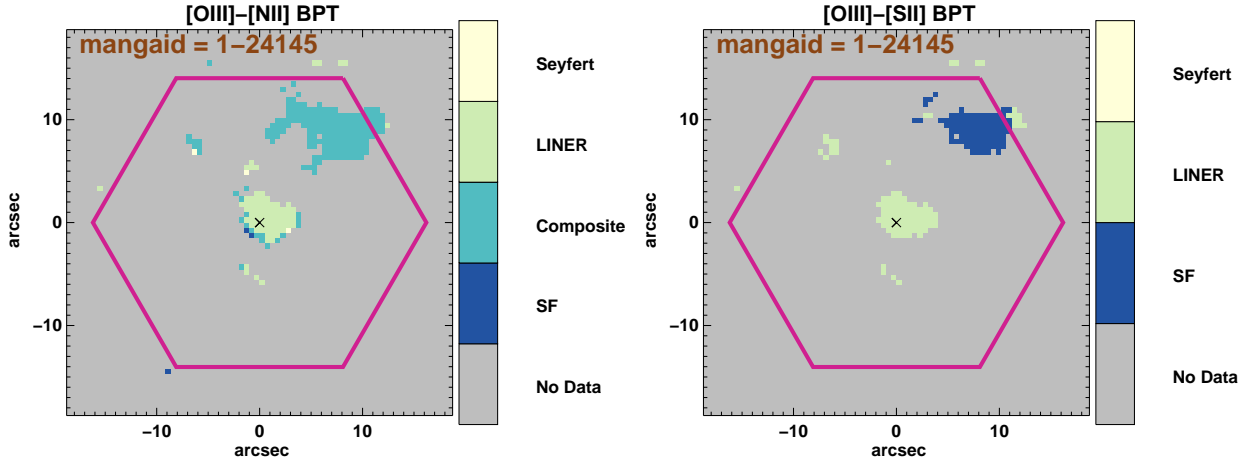
**Figure 7.** Velocity fields of MaNGA 1-24145 based on H $\alpha$  line (upper panels) and stellar components (lower panels). The left panels show the radial velocities and the right panels show the velocity dispersions.



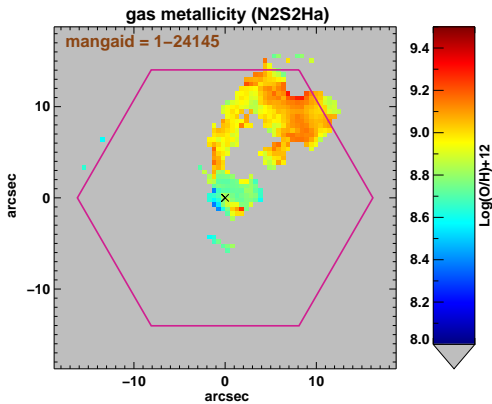
**Figure 8.** Left: The extinction-corrected H $\alpha$  map of MaNGA 1-24145. Right: The [OIII] 5007/H $\beta$  ratio as a function of distance from the center of Satsuki. The letters (A,B,C,...etc) mark different locations of this system.



**Figure 9.** Upper panels: BPT diagnostic diagrams for MaNGA 1-24145 based on the  $[\text{OIII}] 5007/\text{H}\beta$  vs.  $[\text{NII}] 6584/\text{H}\alpha$  (left panels),  $[\text{OIII}] 5007/\text{H}\beta$  vs.  $[\text{SII}] 6717,6731/\text{H}\alpha$  (middle panels), and  $[\text{OIII}] 5007/\text{H}\beta$  vs.  $[\text{OI}] 6300/\text{H}\alpha$  (right panels). Various colors on the data points indicate the physical separation from the center of Satsuki (from near to far: red, origin, green, blue, purple). The solid, dashed, and dotted lines show the classification lines suggested by Cid Fernandes et al. (2010), Kewley et al. (2001), and Kauffmann et al. (2003), respectively. The blue curves display the model predictions of the shock and photoionization mixing sequences by Ho et al. (2014). Each line corresponds to a certain shock fraction (from bottom to top: 20% to 100%). The shock velocity ranges from 100 to 300  $\text{km s}^{-1}$  (from left to right). The MAPPINGS III shock+precursor models (gray grid) with  $n = 1 \text{ cm}^{-3}$  from Allen et al. (2008) are also shown for comparison. The thick lines represent constant magnetic parameter while the thin lines display the constant shock velocity ranging from 200 to 1000  $\text{km s}^{-1}$  (bottom to top; with 25  $\text{km s}^{-1}$  intervals). Bottom panels: A zoomed-in view of the upper panels. The letters (A,B,C,...etc) mark different locations of this system, following the definitions of Figure 8.



**Figure 10.** BPT classification maps for MaNGA 1-24145 based on the [OIII] 5007/H $\beta$  vs. [NII] 6584/H $\alpha$  (left panel) and [OIII] 5007/H $\beta$  and [SII] 6717,6731/H $\alpha$  (right panel). The [OI] diagnostic map is not included here as all spaxels are classified as LINER.



**Figure 11.** The gas metallicity map of MaNGA 1-24145, computed using the method described in Dopita et al. (2016) based on the emission line ratios among [NII], [SII], and H $\alpha$ .

The environment of this system is rather complex. As already mentioned, it has an early-type companion (Mei) just 4 kpc away to the South, which makes them a possible dry merger candidate. Moreover, Satsuki is also part of the system MCG+10+24-117, a small group falling onto a galaxy cluster with NGC 6338 as a central Brightest Cluster Galaxy that is 43 kpc away to the south (see Figure 5; Pandge et al. 2012; Dupke et al. 2013).

#### 4. DISCUSSION

##### 4.1. The origin of the H $\alpha$ blob

Using the first-year MaNGA data, we discover a giant H $\alpha$  blob, Totoro, associated with a dry merger system. This object, however, does not have any optical counterparts down

to 26.9 mag arcsec<sup>-2</sup> in deep CFHT/MegaCam *g*, *r*, and *i* images. There are several possible scenarios to explain the origins of Totoro:

Scenario 1: Totoro is associated with gas tidally stripped from Satsuki during the interaction between Satsuki and Mei.

Scenario 2: Totoro is associated with the gas ram-pressure stripped during the infall of Satsuki toward the center of NGC 6338 galaxy cluster, similar to NGC4569 located in the Virgo cluster (Boselli et al. 2016).

Scenario 3: Totoro is associated with gas ejected from Satsuki by an AGN outflow during the mergers between Satsuki and Mei. If the central black hole of Satsuki is turned on during mergers, the energy can ionize the stripped gas, resulting in the H $\alpha$  emissions, similar to the known ‘Hanny’s voorwerp’ phenomenon (Lintott et al. 2009).

Scenario 4: Totoro is an UDG (or alternatively, a LSB galaxy), which falls below the detection limit (26.9 mag arcsec<sup>-2</sup> in *r*-band) of CFHT imaging data, making it a ‘dark’ galaxy. The morphology of Totoro (Figure 1), especially the (tail) bridge-like structures that connect between Satsuki and Totoro, indicates that Satsuki is likely under interaction with the hosting galaxy of Totoro, which has a comparable physical size as Satsuki.

To estimate the mass of the warm gas component, we follow the approach adopted by Cheung et al. (2016b). We first estimate the electron density  $n_e$  to be  $\sim 260$  cm<sup>-3</sup> based on the median value of the [SII]6717/[SII]6731 ratio in the region of Totoro following Equation 3 of Proxauf et al. (2014), which assumes the electron temperature  $T_e = 10,000K$ . Next, we calculate the extinction-corrected H $\beta$  luminosity,  $L_{H\beta}$ , to be  $3.1 \times 10^{39}$  erg s<sup>-1</sup>. The ionized gas mass can then be derived using the following equation (Osterbrock & Ferland

2006; Fu et al. 2012):

$$\frac{M_{\text{HI}}}{6.8 \times 10^7 M_{\odot}} = \left( \frac{L_{\text{H}\beta}}{10^{40} \text{ ergs}^{-1}} \right) \left( \frac{n_e}{1 \text{ cm}^{-3}} \right)^{-1}. \quad (2)$$

We obtain the ionized gas mass  $\sim 8.2 \times 10^4 M_{\odot}$ .

The HI observation of this system provides an upper limit of  $(8.9\text{--}9.2) \times 10^8 M_{\odot}$  for the HI mass due to the null detection. We account for He gas by applying a factor of 1.33 to this upper limit. We can also infer the  $\text{H}_2$  content of this Totoro using the  $\text{H}\alpha$  emission. The integrated  $\text{H}\alpha$  flux over spaxels within Totoro region is  $3.3 \times 10^{-15} \text{ erg s}^{-1} \text{ cm}^{-2}$ , corresponding to a luminosity of  $7.8 \times 10^{39} \text{ erg s}^{-1}$ . Assuming all the  $\text{H}\alpha$  flux results from star formation, we obtain the total star formation rate (SFR) of Totoro to be  $0.059 M_{\odot} \text{ yr}^{-1}$  using the conversion between the  $\text{H}\alpha$  luminosity and star formation rate from Kennicutt (1998). We then estimate the total  $\text{H}_2$  gas mass ( $= \text{SFR} \times t_{\text{dep}}$ ) of the  $\text{H}\alpha$  blob to be  $\sim 1.2 \times 10^8 M_{\odot}$  assuming a typical gas depletion time ( $t_{\text{dep}}$ ) of 2 Gyr. This value is also an upper limit since we have assumed that all the  $\text{H}\alpha$  fluxes are contributed by the star formation. This implies an upper limit of the entire cold gas ( $\text{HI} + \text{He} + \text{H}_2$ ) of  $\sim 1.3 \times 10^9 M_{\odot}$ .

Although Satsuki and its companion Mei are consistent with being early-type galaxies (ETGs), the amount of the cold gas in Totoro is not totally unexpected if it is part of Satsuki. For example, recent works have found that a significant fraction of massive ETGs that show signs of star formation possess cold gas with gas mass comparable to this Totoro (O’Sullivan et al. 2015; Davis et al. 2016). Therefore, it is reasonable to speculate that Totoro may originally be part (if not all) of the cold gas contained Satsuki, and then get ejected during the galaxy-galaxy interaction or because of the ram-pressure stripping (scenarios 1 to 3, respectively).

To test the galaxy interaction scenario (scenario 1), we have performed and examined a set of N-body simulations of the interaction of two Elliptical galaxies (E-E) with the code Gadget2 (Springel 2005). Different initial orbital configurations (e.g. pericentric distance) have been considered in a similar way to Peirani et al. (2010). Although we did not include any gas component in our dry merger simulations, the features in the stellar components can still be useful to understand the system as the stellar streams well trace the gas streams before 1 Gyr during galaxy interaction.

Among all the cases we have looked at, we did not see the formation of any clear stellar stream or a centrally concentrated blob-like structures during the interaction. This strongly suggested that Totoro is unlikely to be produced during an E-E interaction. In addition, based on other existing gas-rich merger simulations (see for instance Barnes & Hernquist 1996; Di Matteo et al. 2007; Peirani et al. 2010), the prominent gas (or stellar) streams are expected to be formed when at least one of merger progenitors have extended disk structures, unlikely to be the case for this dry merger system even if these two ellipticals have small amount of gas. Furthermore, such latter simulations also suggest that it is very difficult to produce such blob-like structure. Although tidal dwarf galaxies can be formed during the interaction of disk galaxies (see for instance Bournaud, Duc et Masset 2003; Duc, Bournaud et Masset 2004) in these cases, the blob-like structure is more extended with no clear star formation activity. Nevertheless, we note that the simulations we have examined have the following caveats: 1) the possible parameter

space is extensive, so it is impossible to fully explore, and 2) the stellar/AGN feedback which could drive galactic scale outflows and these physics are not included. Such advanced simulations are beyond the scope of this work. We defer more detailed comparisons with simulations to a future work.

On the other hand, the centrally concentrated blob-like structure is not expected in the ram-pressure stripped gas (scenario 2), either, which often shows the tail and/or ‘jellyfish’ like structures (e.g., Boselli et al. 2016). Moreover, galaxies that show ram-pressure stripping phenomena are mostly gas-rich late-type galaxies, unlike Satsuki.

Another explanation for Totoro is the materials ejected by an AGN located in Satsuki (scenario 3), similar to the extended emission-line regions found around quasars (Fu & Stockton 2007a,b). The outflow scenario is also consistent with the complicated velocity field seen in the gas component of this system. On the other hand, as discussed in §2, we do not detect an X-ray point source or an extended radio jet in the main galaxy, suggesting that there is no on-going strong AGN activity. However, it has been known that the AGN brightness can vary over timescales of several to  $10^5$  years (Denney et al. 2014; McElroy et al. 2016), either due to a change in the black-hole accretion rate or Tidal-Disruption Events (TDEs, e.g. Saxton et al. 2012; Merloni et al. 2015). Therefore, we can not rule out the possibility of a recent past AGN outflow, similar to the notable phenomena of Henny’s Voorwerp, in which the ionizing source has already diminished whereas a clump of ionized gas is found several kpc away. Nevertheless, it is unclear why the ejected gas would have a higher metallicity than the gas remaining in Satsuki (although it is unclear whether the metallicity difference is real given the potential systematics in our metallicity estimate).

An alternative explanation of the  $\text{H}\alpha$  blob is a separate faint galaxy, interacting with the dry merger system (scenario 4). A few morphological features, for example, the centrally concentrated  $\text{H}\alpha$  and the bridges extended from Totoro toward Satsuki, are most consistent with Totoro being a separated gas component. According to numerical simulations, these features can indeed be explained by the interaction between the main galaxy and a faint disk galaxy (scenario 1) (see for instance, Peirani et al. 2010; Cheung et al. 2016).

Recently, a class of UDGs has been identified in the Coma cluster and several low-redshift clusters (van Dokkum et al. 2015; van der Burg et al. 2016). These UDGs are surprising large in size ( $r_{\text{eff}} = 1.5 - 4.6 \text{ kpc}$ ) despite of their low stellar masses ( $< 10^8 M_{\odot}$ ). The majority of these UDGs are found to lie on the red-sequence, indicating old stellar populations and quenched star formation activities (van der Burg et al. 2016). Possible formation mechanisms include that gas is stripped by the ram pressure when falling into the cluster, which prevents subsequent star formation, or that gas is expelled due to strong stellar or supernova feedback for galaxies with halo mass between  $10^{10}\text{--}10^{11} M_{\odot}$  (Di Cintio et al. 2016). Nevertheless, it remains an open question regarding the origin of UDGs and whether these UDGs are already quenched before falling into the cluster environments. Totoro identified in this work has a comparable size ( $\sim 3.2 \text{ kpc}$ ) as the UDGs. However, the averaged surface brightness of Totoro has an upper limit of  $26.9 \text{ mag arcsec}^{-2}$ , at least 1-3 mag dimmer compared to that of the known UDGs ranging from 24 to  $26 \text{ mag arcsec}^{-2}$ . Another aspect of Totoro that is distinct from the known UDGs is that the latter are generally old in the stellar populations as

indicated from their red colors, whereas Totoro shown in this work is likely to have small amount of on-going star formation based on the BPT diagnostics but with little old stellar populations. Therefore, Totoro may represent a different category of LSB galaxy from the quiescent UDGs.

The non-detection of optical light associated with Totoro can provide an upper limit of the amount of star formation. Assuming that this H $\alpha$  cloud is a young star-forming system with age  $< 0.1$  Gyr, the flux in the optical regime is expected to be comparable to that in the UV (1500–2800 Å) in the case of no dust extinction. Following the conversion between the UV luminosity, H $\alpha$  luminosity, and the star formation rate given by (Kennicutt 1998) and adopt the surface brightness limit (26.9 mag arcsec $^{-2}$ ) from our CFHT MegaCam  $r$ -band observation, we estimate that the corresponding H $\alpha$  surface density of this Totoro is on the order of  $1.1 \times 10^{-17}$  erg s $^{-1}$  cm $^{-1}$  arcsec $^{-2}$ , which is one magnitude lower than the peak value of Totoro. By integrating it over the H $\alpha$  region, we estimate that the star formation contributes less than 37% to the total H $\alpha$  flux, otherwise we should be able to detect the optical counterpart.

Given the very low amount of star formation that can possibly occur in Totoro, we speculate that Totoro is a gas cloud that fails to form stars efficiently, and thus a ‘dark’ gas cloud. The origin of this gas cloud, whether it is associated with a satellite dark matter, or a pure gas cloud, is difficult to pin down. However, the latter scenario is unlikely since it would be difficult to maintain the kpc-scale gas cloud against gravity without invoking the existence of a dark matter halo. If Totoro is indeed hosted by a subhalo, it would be a strong evidence of the existence of ‘dark’ subhalos, which help to alleviate the missing satellite problem. Since there is no stellar component in the region of Totoro, we could roughly estimate the halo mass of Totoro based on the gas velocity dispersion using the virial theorem. Taking  $R = 3.2$  kpc and  $\sigma_{gas} = 50$  km s $^{-1}$ , we obtain the halo mass  $M_{halo} \sim 5.6 \times 10^9 M_{\odot}$ .

On the other hand, the high metallicity of Totoro (see Section 3.4) is unexpected for a typical dwarf galaxy. One possible explanation is that the gas cloud has been enriched by the surrounding environment, and may have a different evolution process from the typical known dwarf galaxies. However, we caution that the metallicity measurement presented in this work is subject to large uncertainty since the ionization source is not well constrained.

#### 4.2. Sources of ionization

Shocks can often lead to line ratios similar to those occupying the composite regions (Ho et al. 2014). In the case where there are shocks, the gas velocity dispersion ( $\sigma_{gas}$ ) is expected to be as high as several hundreds of km s $^{-1}$ . As revealed in Figure 7,  $\sigma_{gas}$  in the region of Totoro is  $\sim 50$  km s $^{-1}$ . Although it is close to the lower end of velocity dispersion distribution that have been found in typical shocked regions, we can not ruled out the shock excitation of the emission lines. To gain insights to whether shock is responsible for producing the ionized ratios seen in this system, we show the shock and photoionization mixing models from Ho et al. (2014) as blue curves in the upper panels of Figure 9. These models are produced based on the MAPPINGS IV code (Dopita et al. 2013) and span a wide range of shock fractions (from 20% to 100%) and shock velocities (from 100 to 300 km s $^{-1}$ ). As it can be seen, the models predict a greater [OIII] 5007/H $\beta$  ratio than what is observed in the data and do not cover the major-

ity of the regions occupied by the data points even at shock velocity up to  $\sim 300$  km s $^{-1}$ . In addition to the pure shock models, we also compare the data to the shock + precursor models of Allen et al. (2008) with  $n = 1$  cm $^{-3}$  and solar metallicity, shown as the gray grid in figure 9. Although the grid starts with the shock velocity of 200 km s $^{-1}$ , it is expected by extrapolation that models with lower velocity values still can not reproduce the [NII]/H $\alpha$  ratios of Totoro. These comparisons suggest that shocks are unlikely to be the dominant mechanism that is responsible for ionizing the gas blob.

There have been studies showing that AGN is able to ionize gas clouds extending to several kpc (Lintott et al. 2009; Fu & Stockton 2008), and the effect from AGN may persist even  $\sim 10^5$  yr after the central engines shut off (e.g., the ‘Hanny’s Voorewerp’; Lintott et al. 2009). If Totoro is indeed a ‘dark’ galaxy or a cloud interacting with Satsuki, the tidal field would cause a gas inflow that fuels the central black hole of Satsuki and possibly trigger the AGN activity. Although no X-ray point source is detected in the position of either Satsuki or Totoro, the detection of point-like radio source (see Sec. 2.3) in the center of Satsuki indicates the presence of a low-activity AGN. Therefore, an alternative explanation for the line ratios seen in Totoro is due to the star formation – AGN mixing effect (Davies et al. 2014a,b).

Observationally a starburst–AGN mixing sequence is often found in starbursting galaxy that hosts a central AGN, in which case the line ratio moves from Seyfert to composite to HII regions as the distance from the central AGN increases. The position of line ratios depends on the fractional contribution between AGN and star formation. According to the mixing model by Davies et al. (2014b), when the star-forming cloud is photonionized by an AGN, the line ratios can fall into the ‘composite’ region on BPT diagrams. In the case of 100% contribution from AGN, the emission line ratio falls into the ‘AGN’ region, instead of ‘composite’ region where Totoro lies. This implies that the emission line ratios seen in Totoro can not be fully attributed to pure AGN photonionization, and some level of star formation may be required. Unlike the typical star formation – AGN mixing (Davies et al. 2014a,b), our case is analogous to the star formation – LINER sequence, in which the LINER excitation is due to the low-activity AGN located in Satsuki.

#### 4.3. Comparison to similar objects in the literature

There are several similar systems reported in the literature that show offset ionized gas components, and hence it is intriguing to compare this Totoro with previous cases. For example, the famous Hanny’s Voorewerp (Lintott et al. 2009), also exhibits a separate warm gas component that is offset from the main galaxy by several kpc. However, there are several different features between Hanny’s Voorewerp and Totoro presented in this work. First, the ionized gas of Hanny’s Voorewerp has much higher excitation state and its line ratios are consistent with Seyfert, as opposed to the ‘composite’ region for Totoro. Moreover, the H $\alpha$  morphology of the Hanny’s Voorewerp is lumpy and irregular, unlike a disk-like structure seen in our case. Another well-known example is an H $\alpha$  emission line component (often referred as the ‘cap’) at a projected distance of 11 kpc northwest of M82 (Devine & Bally 1999; Lehnert et al. 1999; Stevens et al. 2003). The cap has a shell-like structure and may possibly be bow shock formed by the starburst-driven superwind. In both cases, the nearby galaxy is a late-type star

**Table 1**  
Properties of MaNGA 1-24145 (Satsuki), its southern companion (Mei), and the H $\alpha$  blob (Totoro).

Object	$z$	RA	DEC	$M_*$ ( $M_\odot$ )	$M_{\text{HII}}$ ( $M_\odot$ )	$M_{\text{H}_2}$ ( $M_\odot$ )	$M_{\text{HI}}$ ( $M_\odot$ )	$M_{\text{halo}}$ ( $M_\odot$ )	SFR ( $\text{yr}^{-1} M_\odot$ )
Satsuki	0.0322	258.84695	57.43288	$1.2 \times 10^{11}$	...	...	$< 9.2 \times 10^8$	...	$< 0.049^b$
Mei	0.0322	258.84750	57.43133	$4.0 \times 10^{10}{}^a$	...	...	$< 9.2 \times 10^8$	...	...
Totoro	0.0322	258.84314	57.43529	...	$8.2 \times 10^4$	$< 1.2 \times 10^8$	$< 9.2 \times 10^8$	$5.6 \times 10^9$	$< 0.059^b$

**Note.** — <sup>(a)</sup> This is scaled from the stellar mass of Satsuki by using the difference in their SDSS  $r$ -band magnitudes.; <sup>(b)</sup> This upper limit is derived assuming all the H $\alpha$  fluxes come from the star formation.

forming galaxy, different from Satsuki, which is an early-type galaxy. Totoro may thus represent a different category of offset ionized gas in nearby galaxies.

## 5. CONCLUSION

Here we report a discovery of a puzzling giant H $\alpha$  blob, Totoro, identified from the first-year MaNGA data. The data disfavor the scenario that Totoro is tidally-stripped gas from MaNGA 1-24145 (Satsuki) that is interacting with the southern companion (Mei), or the ram-pressure stripped gas when Satsuki falls into the center of the cluster it is located. Despite there being no X-ray point source or radio jet detected in this system, we can not rule out the possibility that Totoro is ejected from a past AGN activity in Satsuki, which likely hosts a faint AGN given its radio luminosity. On the other hand, the H $\alpha$  morphology and the lack of stellar tidal streams suggest that Totoro could also be a separate ‘dark’ galaxy (or an extremely LSB galaxy) interacting with Satsuki. The non-detection of the stellar continuum suggests Totoro is different from known dwarf populations or UDGs: it is either completely ‘dark’ or with a star formation rate that contributes  $< 37\%$  of the H $\alpha$  flux.

As for the source that powers the line excitation for Totoro, the ‘composite’ line excitation can be explained either by a star-forming cloud being excited by a low-velocity shock or by the star formation – LINER mixing effect. The shock scenario, however, is less favoured because of the low velocity dispersion observed in Totoro region. The decrease in the [OIII] 5007/H $\beta$  ratio away from Satsuki indicates that the ionizing source is possibly located inside Satsuki. Thus, the star formation – LINER mixing effect seems to be the most probable ionizing mechanism. In this scenario, the hypothesis is that the ionizing source is the low-activity AGN residing in Satsuki, being triggered by the gas inflow induced during the interaction between Satsuki and the ‘dark’ galaxy (or gas cloud). The AGN subsequently photoionizes the ‘dark’ gas cloud, which then emits the H $\alpha$  photons.

However, we have not considered the case where Totoro is tidally-stripped from Satsuki while falling into the cluster center at the same time, which results in the non-typical tidally or ram-pressure stripped gas morphology of Totoro. More sophisticated modelling considering both the effects of tidal disrupting and the orbital motions of Satsuki relative to the cluster environments, as well as future resolved atomic and molecular gas observations, such as HI and CO, are required to further understand the origin of Totoro.

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