

A STRUCTURAL INVESTIGATION OF SOME COMPOUNDS SHOWING
CHARGE-TRANSFER PROPERTIES

A Thesis submitted for the
Degree of Doctor of Philosophy
in the University of Oxford



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I wish to thank all those who have helped me in the last two years, both in the laboratory and in the preparation of this thesis. I have been particularly grateful for the friendly atmosphere in the laboratory and the way that people have sacrificed much of their time to help with discussion, use of equipment, programs and operation of the computer. I shall remember the sense of community both inside and outside of the laboratory. Besides this, I particularly wish to thank my supervisor, Keith Prout, who has helped me more than I have often realised.

LAYOUT OF THIS THESIS

The thesis is divided into chapters which deal, not with the individual compounds studied in sequence, but with some particular aspect of all of them. Thus Chapters I and II are an introduction to the chemistry of the compounds. Chapters III and IV deal solely with the X-ray crystallography (data collection, structure solution, and assessment of techniques), and can be read by themselves. Chapters V and VI describe independently the results and interpretation of structure determination of two sets of compounds; either relies only on I and II. (It has been impossible to separate the solving of the structure and discussion of results in Chapter V).

All the tables and diagrams, where not included in the body of the text, have been put in at the end of each chapter. I hope that this will make it easier to refer to them continually whilst keeping the text intact. In particular structure factors are contained in a separate chapter (VII).

Since the relation of molecules to each other is important, I have tried to show this with clinographic drawings of the structures. These are to scale and are projections down the vector with direction cosines: -0.94 , -0.193 , $-.290$, in a right-handed system. The crystal axes are not arranged consistently, partly because of the differing shapes of unit cells, and partly to get the best view of the structure. The origin and axes are always labelled.

References are arranged on a chapter basis. The number refers to the table of references at the end of the chapter, unless otherwise indicated.

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SUMMARY

SUMMARY

The technique of X-ray diffraction has been used to investigate compounds which show charge-transfer properties in the solid state. It is shown how the results of structure determination can, under favourable conditions, give information about charge-transfer processes.

Two sets of compounds have been studied, both of which contain metal-chlorine bonds. The first is a series of mixed-valence cobalt hexammine chlorocuprates in which the average oxidation state of the copper can take any value between 1 and 2. The structure of the single-valence (copper (I)) compound, $[\text{Co}(\text{NH}_3)_6]_4\text{Cu}_5\text{Cl}_{17}$ was solved and data was also collected for a compound with $\text{Cu(I)}/\text{total Cu} \sim 0.25$. This was shown to have a structure to the single-valence copper (I) and copper (II) compounds and probable arrangements for the Cu(I) and Cu(II) ions are presented. It is shown that linear Cu-Cl-Cu bridges are almost certainly present and that these can account for the charge-transfer spectra. A new chlorocuprate anion, $[\text{Cu}_5\text{Cl}_{16}]^{11-}$ was discovered, which has $23(T)$ symmetry.

The second series contains a common cation, paraquat (pq^{2+}), which is N,N'-dimethyl-4,4'-bipyridylum. Four salts of this, pqCoCl_4 , pqPdCl_4 , pqCu_2Cl_6 , and pqCu_2Cl_4 have been studied, and their structures are discussed in terms of charge-transfer. It appears that short N-Cl contacts, as little as 3.25 Å, are the main feature of the structures, and these are almost perpendicular to the plane of the aromatic ring. There is also some evidence for an attraction between the nitrogen and chlorine atoms.

The compound pqFeCl_4 , which contains radicals, has been shown to be isostructural with pqCoCl_4 , and the charge-transfer process is discussed in terms of pq^{2+} and $\text{pq}^{\cdot+}$. Electron density calculations for these species support the hypothesis that in these compounds the Nitrogen atom is necessary for charge-transfer. Another chlorocuprate ion, $(\text{CuCl}_2^-)_n$, was discovered, and the coordination of metal chloride anions is discussed.

Since the arguments used to relate X-ray results to electronic processes depend on knowing the accuracy of structural parameters, an analysis of the precision of these parameters was undertaken. It was shown that the estimated standard deviations are useful criteria of accuracy.

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CHAPTER I.

CHARGE-TRANSFER and

ITS INVESTIGATION BY X-RAY STRUCTURE ANALYSIS OF CRYSTALS

INTRODUCTION

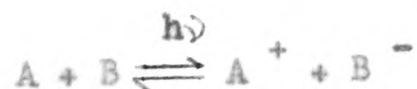
CRYSTAL STRUCTURE AND ELECTRONIC PROCESSES

The study of charge-transfer in a crystal has the advantage that the geometrical relationship of the entities involved is fixed and can be found, and that there is no variation with time. In contrast, the situation in a liquid fluctuates rapidly and at any moment is made up of many different configurations of unequal energies. Although for the complete description of charge-transfer in a solid the geometrical structure must be known, X-ray analysis is of no direct benefit in the investigation of electronic processes. It can yield an average view of the electron density in a crystal, but this is as yet too inaccurate to be described meaningfully in terms of changes in valence electron structure. An alternative approach was adopted.

We assume that charge-transfer processes in a crystal will change the electronic structure in a way that will distort the geometry of the system. Since there is no way of predicting exactly what this geometry would be in the absence of charge-transfer, the identification of distortion must be partly subjective. The risk in this process is lessened by investigating a series of similar compounds to see if they all show the same effect, and if so, to interpret any variation. The variation of geometry of complexes with atomic number is a well-known method of investigating electronic properties (cf. Reference 1). Besides the possible weakness of argument in this process, there is the difficulty that it is usually impossible to analyse accurately enough the forces in any crystal structure. Only in an isomorphous series of compounds is the variation of intermolecular forces likely to be insignificant. The combination of other physical measurements, such as conductivity or absorption spectra, with X-ray structures is often very helpful.

Charge-transfer

(1) Charge-transfer (C-T) can be loosely defined as the process:-



where A and B (not necessarily neutral) may or may not be bonded. It is an oxidation-reduction process (usually reversible) where the energy is supplied by radiation, which when in the visible region results in the compound being coloured. (In this definition can also be included partial electron transfer, or a change in polarisation of a system. However, the crystals studied may involve complete transfer since A^+ and B^- are known to be stable under moderate conditions. It is interesting that some charge-transfer compounds are thought to have a ground state where $A^+ + B^-$ is the main configuration.)

Though oxidation-reduction must involve a change in electronic configuration, it is often more clearly visualised as involving something other than simple electron-transfer. Redox mechanisms have been widely studied in solution and, for transition metal compounds, appear to take place by two or three main mechanisms. The terminology of these is adopted for solid-state charge-transfer, although it is clear that the crystalline state imposes severe restrictions on atomic and possibly electronic movements.

i) (One) Electron-transfer is well-known in both inorganic and organic reactions, and most of the systems investigated later involve only unit changes in formal oxidation state. The Franck-Condon principle implies that reaction will be efficient only if the changes in geometry are small. This process is often called "outer-sphere"(2) as it uses the outer valence orbitals and must often take place over relatively large distances. It is basically independent of the co-ordination sphere, but can be helped in certain cases by electron delocalisation onto ligands. An example of this is the $\text{IrCl}_6^{3-} \dots \text{IrCl}_6^{2-}$ system, where delocalisation which helps electron-transfer has been shown by solid state e.s.r. (3). X-ray diffraction can give little direct evidence of "outer-sphere" mechanisms, but may help to deduce the orientation of orbitals which might be involved.

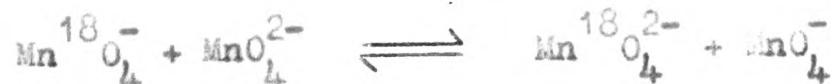
ii) Inner-sphere mechanism. If electron-transfer is so slow that ligand

replacement is a faster process, oxidation is possible by atom-transfer, as in:-



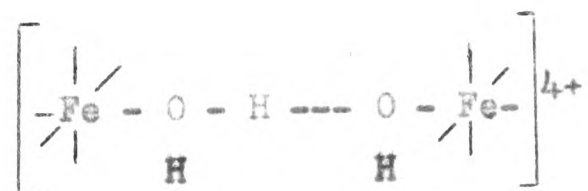
It is clear that only ligands capable of forming bridges can act, such as halogens (but not ammonia) and in charge-transfer, it is likely that intermediates involving bridges will be important.

The distinction between inner- and outer-sphere mechanisms is often not clear-cut. For example, although the equilibrium:



involves no atom transfer, it is greatly catalysed in solution by $\text{Co}(\text{NH}_3)_6^{3+}$ or Cs^+ (5). This implies that ion-pairs are important, and that polarisable species, such as Cs^+ or K^+ (see Chapter VI) (which are normally thought of as relatively "inert") can aid electron-transfer. In an ionic crystal where positive and negative charges alternate, this effect may well be significant. A reaction which occurs by both outer- and inner-sphere mechanisms simultaneously is the oxidation of $\text{Co}(\text{CN})_5^{3-}$ by $\left[\text{Co}(\text{NH}_3)_5\text{X} \right]^{2+}$ to $\left[\text{CoX}(\text{CN})_5 \right]^{3-}$. The rate, above a certain value, depends on X, but otherwise is constant (6).

iii) Hydrogen-atom transfer. Inner-sphere reactions usually involve atoms linked directly to the metal. In the equilibrium between Fe^{2+} and Fe^{3+} in aqueous solution, the reaction is catalysed much more by hydroxyl than by chloride. This has led to the proposal (7) of the bridged state:

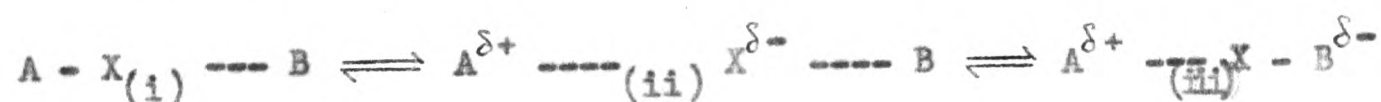


although this does not necessarily imply complete transfer. Where hydrogen bonds are formed in the crystal to ligands (e.g. NH_3 , H_2O ,

or Cl) of transition metals, this mechanism is a possibility, especially as hydrogen atoms may be labile in the crystal.

Charge-transfer in crystals

Although the preceding discussion has been of reactions in solution, the principles should apply in a limited degree to charge-transfer equilibrium in crystals. Electron transfer involving large changes in geometry is helped in solution by thermal excitation and this process is therefore rather unlikely in the crystal (unless excited states are involved). Similarly distinct atom transfer is not possible and a bridged intermediate will be more likely to involve electronic rather than atomic movement. An inner-sphere mechanism can be written:



which need not involve appreciable movement of X.

To investigate charge-transfer in crystals, we choose a system where (i) and (iii) are of roughly comparable energies. It can then be assumed that the geometry of B is not energetically too unfavourable for the formation of B^+ (or an excited state of B^+). It is possible also that state (ii) might be of lower energy than (i) or (iii) and that its existence would be suggested by bond length considerations. (In crystallographic terms, it is most unlikely that substantial amounts of (i) and (iii) exist in the same crystal, even in the presence of radiation.)

THE COMPOUNDS STUDIED

Metal-chlorine bonds were chosen as the common feature of all the compounds in which intermolecular (interionic) C-T was studied. The advantages of this are firstly that metal-chlorine systems (in the absence of C-T) have been well investigated and marked change in geometry is likely to be significant. Secondly, metal-chlorine-metal bridges are well-known and the ionisation potential of chloride is low enough to suggest that these could be important. Chloride is also roughly in the middle of the spectrochemical series suggesting that its crystal field demands on transition metal ions will not be too severe. (For example, both square planar and tetrahedral NiCl_4^{2-} ions have been observed in crystals.) It is relatively small, and its very variable co-ordination number with d^8 , d^9 and d^{10} ions suggests that this geometry might be influenced by C-T effects.

Mixed-valence compounds

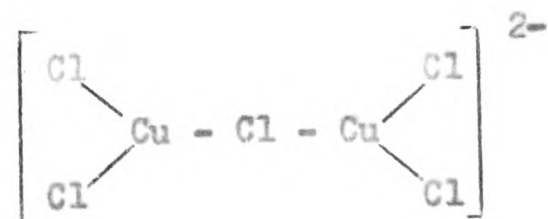
If a crystal contains two adjacent valency states of a metal close to each other and in roughly similar environments, there is often a C-T equilibrium, which may lie close to centre. The whole crystal might approximate to a reaction intermediate and unusual geometrical features should be significant.

Most mixed-valence systems are either 1:1 (or similar simple formulae) or exist over only a small range of variable stoichiometry. There is an exception in the mixed-valence chlorocuprate(I,II) salts of cobalt(III)hexammine which appear to extend completely between the formulae $[\text{Mam}_6]_4\text{Cu(I)}_5\text{Cl}_{17}$ and $(4x)[\text{Mam}_6]\text{Cu(II)Cl}_5$ (where $\text{Mam}_6 = \text{Co}(\text{NH}_3)_6^{3+}$ or $\text{Cr}(\text{NH}_3)_6^{3+}$). There appear to be no intermediate compounds of fixed stoichiometry but well formed crystals showing charge-transfer appear to be formed with any desired ratios of valences.

The formulae of the end compounds appear to be definitely established (and are confirmed by X-ray analysis (8), (Chapter V). Analytical results for the intermediate compounds are indeterminate in deciding between two possible formulations for the . This is because, in going from the copper(II) compound to the copper(I) one, a fifth

copper atom is introduced for each group of 4 original ones. Mori originally suggested that this was not introduced until the ratio $\{\text{Cu(I)}/\text{total Cu}\} > 0.5$; Day and Culpin mentioned the alternative possibility that it was introduced in proportion to the $\{\text{Cu(I)}/\text{total Cu}\}$ percentage. The intermediates can be described by the unwieldy formula; $[\text{Mam}_6]_4 (\text{Cu(II)}_{1-x} \text{Cu(I)}_x)_n \text{Cl}_{12+2n-nx}$, where x is the mole fraction of $\{\text{Cu(I)}/\text{total copper}\}$ (and is used to name the intermediates as " $x\%$ Cu(I)"). The parameter $n, 4\{\text{Cu}/\text{total Cu}\}$, is 4 for $x = 0$, 5 for $x = 1$, and is presumably a function of x , although this might depend on the method of preparation. The situation is depicted in Figure II.1., where Culpin's analyses are shown and the formulations of the intermediates are shown by lines (which need not be straight, and this is suggested in Chapter V). Mori suggested a compound $[\text{Mam}_6]_2 \text{Cu}_2 \text{Cl}_9$ as an intermediate but there is no direct evidence for this.

Charge-transfer is very marked in these compounds as the compounds are dark-black except within about 10% of the end points. The dependence of the intensity of the charge-transfer absorption on mole fraction is shown in Figure I.1a. (9). The shape of the curve, which has no discontinuity of slope, suggests that there are no discrete phases (although points near the edges are ill-defined). This is also true for the single crystal and pallet conductivity, Figure I.1b. The curves are not accurate enough to relate intensity of absorption to the parameter n (see Chapter V). The crystals were made from the very dark mixture of CuCl and CuCl_2 in HCl (which has been studied by (10)⁶³ ^{63}Cu N.M.R. and absorption spectra (14, 11)). This may contain a bridged intermediate of the type:



(although no evidence was given for the actual co-ordination of the copper).

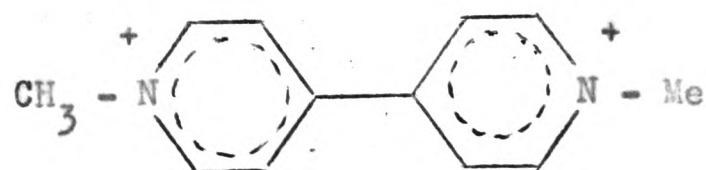
The system is very suitable for X-ray analysis. The compounds

crystallise in two space-groups ($Fd\bar{3}c$ for $0(+)\% \text{ Cu(I)}$ and $Fd\bar{3}$ for $100(-)\% \text{ Cu(I)}$), the latter being a sub-group of the first, (Part II work (12)). This means that there is no crystallographic necessity for any distinct phases and even the copper(II) and copper(I) compounds are very similar. The chlorocuprate(II) structure had been solved (8) and the positions of the corresponding heavy atoms in the chlorocuprate(I) were identical (Part II). Even though this had not been solved, its unusual stoichiometry could be explained satisfactorily in terms of space-group symmetry positions. The chlorocuprate(I) was therefore investigated more completely and data was also collected for an intermediate.

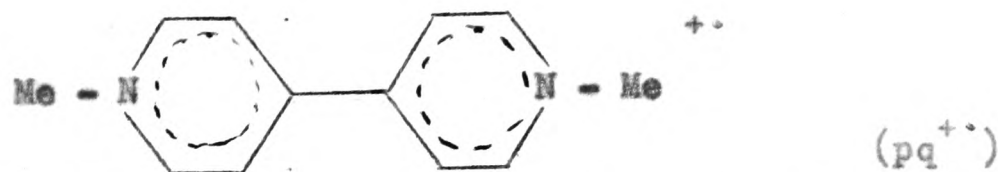
Since there is a possibility that a mixed valence compound might be sufficiently ordered to give a superlattice a similar compound for which this was likely was investigated as a parallel. Mixed lead(II) and (IV) chlorides in HCl solution give $\text{Coam}_6\text{Pb(II)(IV)Cl}_6$, which is probably $[\text{Coam}_6]_2\text{Pb(II)Cl}_6 \cdot \text{Pb(IV)Cl}_6$. The strong charge-transfer (although possibly aided by the cobalt(III)hexamine ions) is probably outer-sphere but X-ray analysis could still show if the two types of octahedra are different and regularly ordered.

Paraquat compounds

The combination of A and A^+ (as with the chlorocuprate(I,II)) is not the only method of trying to create a charge-transfer intermediate in a crystal. If A and B^+ can be chosen so that their ionisation potentials are fairly similar, and if the change in lattice energy from AB to A^+B^+ is small, then the equilibrium should be appreciable. Ionisation potentials (which are not usually available) can be crudely paralleled by electrode potentials and a particular suitable electron acceptor is paraquat, with an E.P. for reduction of 456 mV (13). Paraquat (pq^{2+}), or dimethyl viologen, is the N,N'-4,4' dipyridylum dication:-



which is easily reduced to the intense violet radical monocation:



Macfarlane (14) has made many salts of paraquat with metalchloride complex anions, of which the most interesting was $pqFeCl_4$, prepared from $pqCl_2$ and $FeCl_2$. This was blue and its absorption spectrum, Figure I.2, showed strong evidence of bands due both to pq^+ and charge-transfer. (This diagram also shows charge-transfer in other paraquat halides.) There was also an e.s.r. signal whose g-value could be interpreted as arising from a combination of pq^+ and $FeCl_4^{2-}$. Mossbauer spectra, however, showed that the percentage of $FeCl_4^{2-}$ was probably small. This strongly suggested that X-ray analysis of the structure might give some indication of the charge-transfer path. A compound had also been prepared from $Cu(I)Cl$ and $pqCl_2$ which appeared from e.s.r. evidence to contain Cu^{2+} ions and was very darkly coloured. This could have close connections with the other chlorocuprates.

There is obviously charge-transfer between paraquat and single halide ions as shown by the spectra of the crystalline halides, Figure I.3. This implies that MCl_4 anions would be worth investigating even when the central metal ion (e.g., Co^{2+} , Pd^{2+}) is not easily oxidisable. (The simple halides have been investigated elsewhere (17)).

(Paraquat is, incidentally, a well-known weed killer which probably acts, in conjunction with chlorophyll, to produce radicals by a photochemical process. Paraquat chloride solution can be changed to radical monocation by irradiation.)

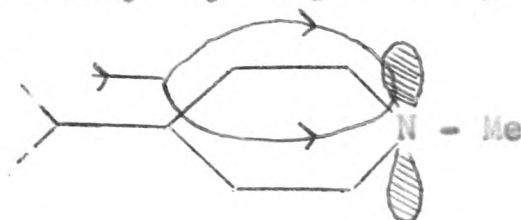
Charge-transfer involving organic systems

The movement of electrons has so far been considered as to and from definite sites, although delocalisation of (d-) electrons is certainly possible even in simple systems. When the electron accepting (or donating) power arises from no specific atom, but from a molecule as a whole, then a different mechanism must be postulated. The best known examples of charge-transfer involving delocalised (aromatic)

systems are the benzene-iodine complex and the $\pi - \pi$ complexes of flat aromatic molecules. The latter has been described by the principle of "maximum overlap" of π -orbitals (14) and is exemplified by the complex of copper oxinate and tetracyanobenzene (Chapter IV). X-ray structure analysis can give little information, except the extent of π -overlap. Since delocalised π -systems exist not only in paraquat but also in the metal anions (e.g., PdCl_4^{2-}), this sort of overlap might be possible and would be suggested by part of the anion lying close and parallel to the paraquat ion.

The benzene-iodine system and similar complexes have been shown (15) to involve a specific donor/acceptor associated with the middle of an aromatic system. The benzene-bromine and benzene- CCl_4 complexes are perhaps the most pertinent as they contain bonds with a terminal halogen directed nearly perpendicularly at the centre of the ring, and this can be called $n - \pi$ charge-transfer. Its geometric requirements are due to a combination of suitable orbitals, but it is difficult to predict which these will be. Corresponding $\text{M} - \text{Cl} \rightarrow \pi$ -system arrangements would be evidence for this type.

Organic molecules containing a donor atom (e.g., N, O) can form charge-transfer complexes which do not depend on delocalisation (15). An example is $(\text{Me})_3\text{N} \cdots \text{I}_2$, which can be written $(\text{Me})_3\text{N}^{\delta+} \cdots \text{I} \cdots \text{I}^{\delta-}$, where stability is due to the donor power of the nitrogen and the polarisability of the iodine. A similar complex occurs with pyridine where no appreciable use of the π -electrons is made. (It is possible that by back-donation the lone-pair electrons on the iodine atom nearest to the nitrogen could interact with the ring π -system.) Hassel (15) notes several very short $\text{N} \cdots \text{Cl}$ (2.31\AA) and $\text{O} \cdots \text{Cl}$ (2.50\AA) distances although some were appreciably longer ($\sim 3.2\text{\AA}$). Paraquat cannot be directly fitted into this scheme as it has no lone-pair, but the π -electron density may be partially concentrated at the (positive) nitrogen atoms:



This would allow the electron accepting/donating power to act through the nitrogen atom even through originating from the whole system. (A more rigorous basis for this is described in Chapter VI from theoretical and experimental calculations of the \bar{n} -electron density in pq^{2+} and pq^{+} .)

ASSESSMENT OF THE ACCURACY OF X-RAY STRUCTURES

This geometrical method of investigating the electronic structure of compounds relies on the differences in bonds and angles of two structures being significant. The estimation of errors in parameters assumes that errors in F_{obs} are random, which is still valid for reproducible errors provided they do not affect the parameters systematically. This is reasonably true for most positional parameters derived from data of moderate accuracy. The use of these estimated standard deviations (e.s.d.'s) depends on their practical justification, which is rarely, if ever, undertaken as it involves repeated measurement of F_{obs} data and refinement of parameters. In particular it is interesting to know if different experimenters get results agreeing within their e.s.d.'s. This would involve the duplication of the measurement of complete three-dimensional data and is not usually attempted.

Duplicate data from a molecular complex

Because of a slight misunderstanding, Wright and I determined the structure of copper oxinate.2tetracyanobenzene independently and simultaneously. I have analysed the differences between the two structures and compared them with the e.s.d.'s (Chapter IV). The structure is included solely for the data comparison, etc. (though a paper on the chemistry is in the press (IV.8) and it is hoped to include a reprint in this thesis).

Methods of data collection

Photographic methods of data collection can only compare with counter measurements if much time and care is spent on them. The data compared above suggested that one-axis Weissenberg photographic data collection is a quick method whose error is accurately represented by its e.s.d.'s. Though very widely used (IV.1), little criticism of its merits and demerits has been made and I hope to have shown (Chapter III) that its use with heavy atom compounds is worthwhile. In particular, the time saved (which is valuable if a series of compounds is being investigated) outweighs the loss in accuracy.

Errors in bonds and angles

The systems studied provided a good basis for an empirical assessment of e.s.d.'s of bonds and angles. A series of compounds was taken (from the literature and this thesis) involving oxinate complexes in environments where their bond lengths and angles should be crystallographically indistinguishable. A statistical comparison was made on the basis that bonds and angles are independent.

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(*This review also describes Huckel calculation of π -electron densities and has a plot of bond order vs bond length. Both were used in Chapter VI).

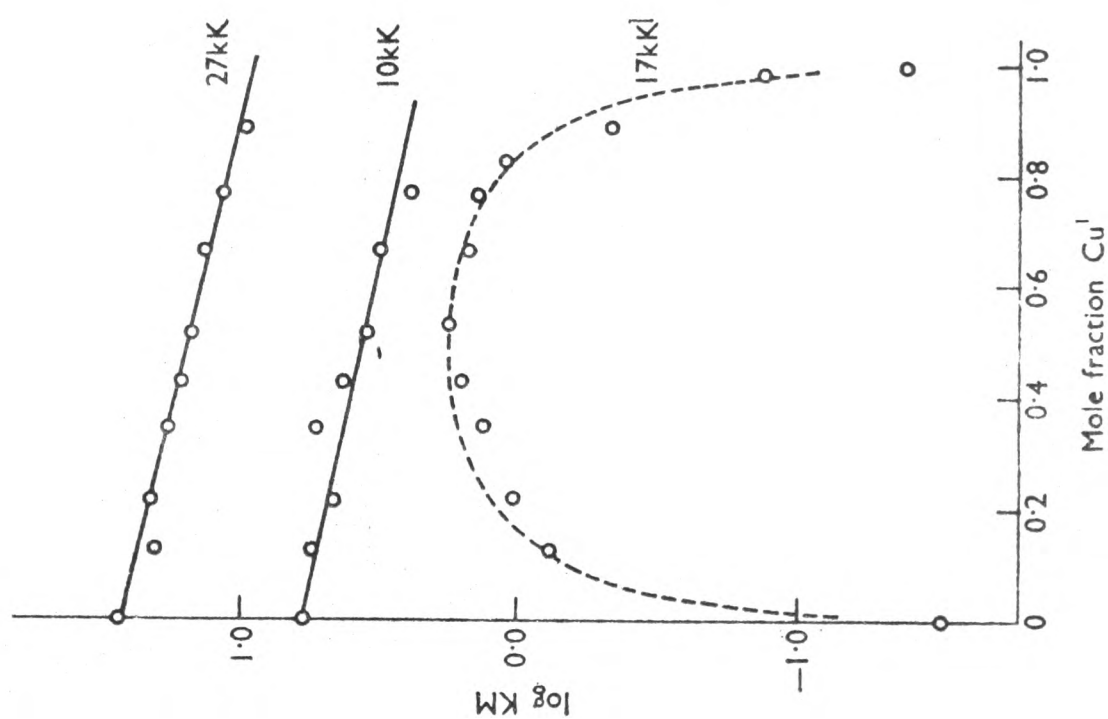


FIGURE 1. Absorption intensity and composition of mixed-valence chlorocuprates.

a

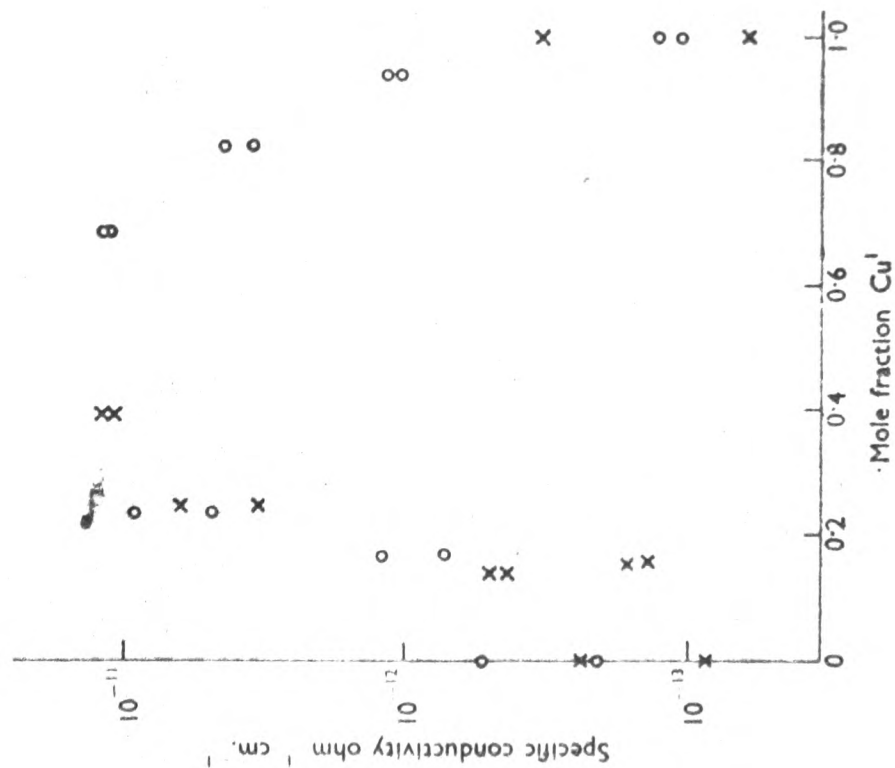


FIGURE 2. Conductivities of mixed-valence chlorocuprates. X = single crystals; O = pellets.

b

Figs. I. a. and b. Absorption and conductivities of mixed-valence Chlorocuprates(I, II)

Fig. I.3. Reflectance Spectra of solid pq halides

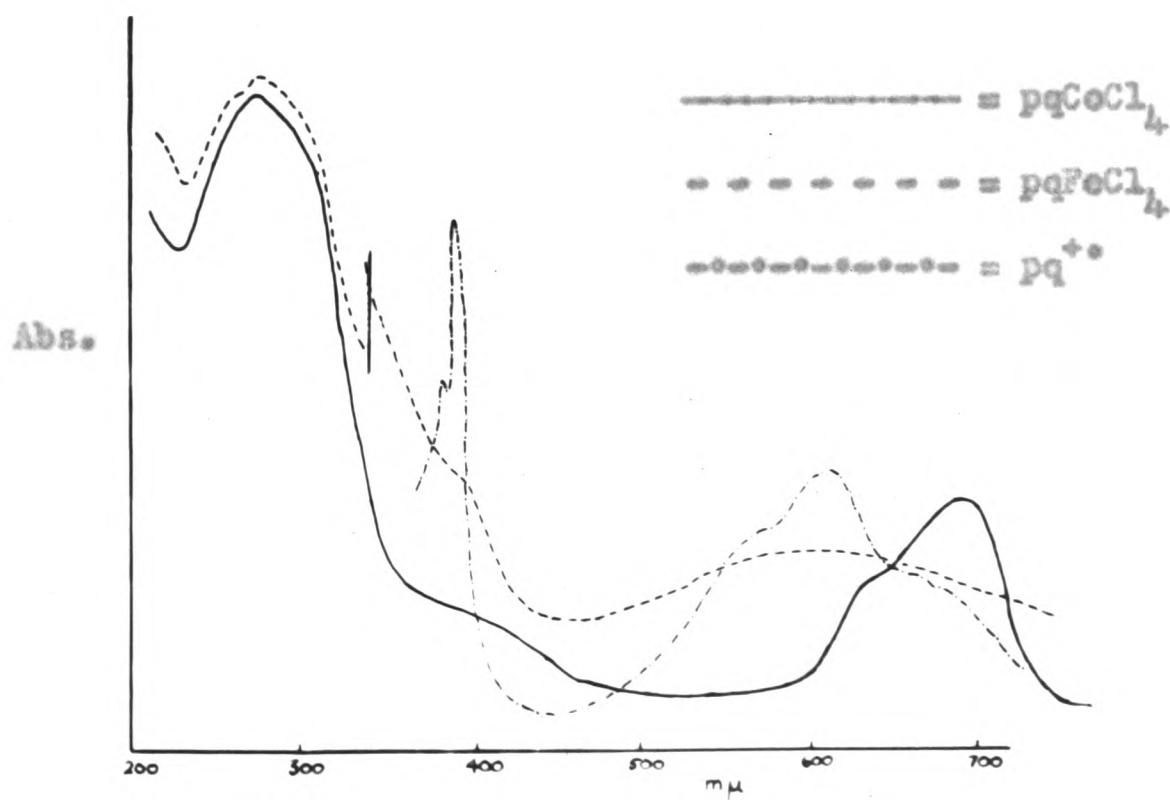
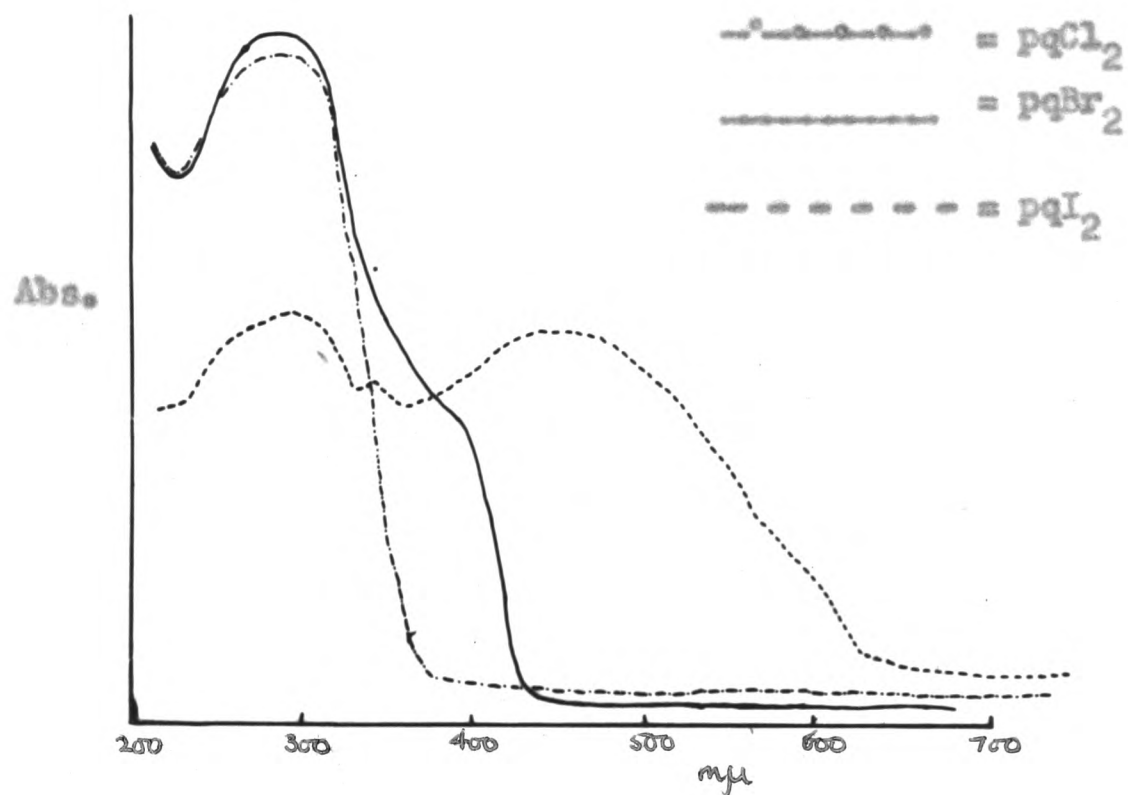


Fig. I.2. Reflectance Spectra of paraquat compounds

CHAPTER II.

PREPARATION AND PROPERTIES OF THE COMPOUNDS STUDIED

PREPARATION AND PROPERTIES OF THE COMPOUNDS

COBALT HEXAMINE CHLOROCUPRATES(I,II)

(My Part II Thesis work is indicated where necessary.)

These compounds were originally prepared by Mori (1) who deduced the formulæ of the end compounds. Day (2) and later Culpin (3) repeated the process and it was these crystals which were used in the structure analysis. They obtained the same end compounds as Mori but Culpin mentioned the possibility of the compounds lying on the lower line in Figure II.1 whereas Mori suggests the upper. The analytical results are clearly no use in differentiating these two cases and density measurements (Part II) are hampered by inhomogeneity of sample and the small variation of M.Wt. (1580 for $4\text{Coam}_6\text{CuCl}_5$ to 1537 for $[\text{Coam}_6]_4\cdot\text{Cu}_5\text{Cl}_{17}$).

Culpin also gave evidence that his method of preparation involving slow crystallisation from a hot solution, under N_2 , gave a different analysis from Day's fast cooling. There is thus a possibility that the bulk sample produced is not homogeneous. It was clear however that individual crystals of the mixed-valence preparations contained copper(I) and copper(II) as they were all darkly coloured or black. Single crystal rotation photographs (Part II) showed that crystals taken from samples of 0 - 50% Cu(I) had space group $\text{Fd}\bar{3}\text{c}$; crystals from samples over 80% Cu(I) belonged to $\text{Fd}\bar{3}(\text{m})$. A crystal from a 67% Cu(I) sample showed marked streaking of high angle reflections implying some disorder.

Day and Smith (4) have recently stated that compounds with composition $\sim 50\%$ Cu(I) are in fact mixtures. They give no new evidence for this and do not indicate whether the mixture is macroscopic or a solid solution. This is hard to reconcile with the X-ray evidence of the high percentage intermediates, whose individual crystals are definitely not in $\text{Fd}\bar{3}\text{c}$.

In the later X-ray analysis of a crystal from a 25% Cu(I) sample, no reliance was placed on this figure (Chapter V), but a substantial

proportion of Cu(I) was assumed as the crystal was virtually black. The X-ray data suggest that this figure is, in fact, correct for the crystal studied.

PARAQUAT COMPOUNDS

Macfarlane (5) had prepared and kindly provided paraquat salts of several metal halide anions, but of the chlorides only $pqCoCl_4$ had been prepared as crystals. Powder photographs (on a Guinier de Wolff camera) showed that salts with anions $FeCl_4^{2-}$, $MnCl_4^{2-}$, $ZnCl_4^{2-}$ and $CuCl_4^{2-}$ were very similar to $pqCoCl_4$. Isomorphism appears to be good whilst probably allowing some latitude in bonds and angles.

The chloroferrate(II) was a dark blue hygroscopic powder unstable in air and giving small blue crystals when a methanolic solution was slowly evaporated. These were not single but oscillation and Weissenberg photographs had very similar patterns to $pqCoCl_4$. This was disappointing, especially as the compound contains radicals, but the chlorocobaltate(II) should have a very similar structure and this was undertaken.

Pure chlorocuprates. There was evidence from the e.s.r. signal of the complex $pqCl_2 + CuCl$ that some copper(II) was present. However only a dark brown powder had been prepared and it looked as if some oxidation might have occurred and no good analysis was given. It was decided to try to prepare crystals of this compound, of the corresponding copper(II) compounds, and, in view of the e.s.r. evidence, to attempt the preparation of mixed valence compounds.

A suitable solvent was half concentrated hydrochloric acid (which had also been used in the preparation of the mixed-valence cobalt hexamine chlorocuprates). Paraquat dichloride recrystallised from ethanol was provided by I.C.I. Research Station, Jealots Hill, Berkshire. Analar copper(I) chloride was used but as it was green it was ground up with dil. HCl saturated with SO_2 , and then filtered. This reducing atmosphere was used whenever it was desired to prevent oxidation.

Paraquat chlorocuprate(I) was prepared by dissolving 1 mmole of paraquat dichloride in a small amount of acid, heating and adding a hot saturated solution of copper(I)chloride (2 mmoles) in acid, all under nitrogen. The colour darkened immediately and the solution was left in a large amount of hot water to cool slowly. Beautiful thick black needles were formed which were relatively stable in air, but decomposed to copper(I) oxide when washed with water and appeared liable to surface oxidation. The crystals are prisms using the $\{110\}$ faces and many have empty channels running down inside. The formula $(C_{12}N_2H_{14})Cu_2Cl_4$ was suggested by C,H,N analysis and confirmed by density measurements and X-ray analysis.

Paraquat chlorocuprate(II) was prepared by the same process, without the nitrogen, atmosphere using copper(II) chloride dihydrate. Large deep-red needles were produced which were green in reflected light (cf. $Me_2NH_2CuCl_3$, see Chapter VI). The formula $(C_{12}N_2H_{14})Cu_2Cl_6$ was proposed and confirmed as above. (N.B. The compound $pqCuCl_4$ (cf. Wallwork (b)) is formed only from non-aqueous solvents.)

Possible mixed-valence compounds

In an attempt to make a mixed-valence crystal, the above method was employed with a 1:1 mixture of copper(I) and copper(II) halides. Crystals were produced but appeared identical to the copper(II) compound and gave similar zero-layer Weissenberg photographs. It was observed that crystals of the copper(I) compound, if exposed to air in the presence of their mother liquor eventually oxidised to the chlorocuprate(II), and the supernatant liquid became green. An attempt to find an intermediate in this process was unsuccessful, and it is probable that it involves only dissolution, oxidation and crystallisation. The difference between the copper(I) and copper(II) compounds, as finally established by X-ray analysis (Chapter VI), is small but substantially larger than in the cobalt hexamine series and is the likely reason for the inability to crystallise an intermediate. Moreover, the e.s.r. signal from crystals of the copper(I) compound was

repeated and found to be very weak indeed which suggests no appreciable amount of Cu^{2+} ions. (The original powder used may well have been substantially oxidised.) The single crystal conductivity along the needle axis was measured by M. Price and showed no photoconductivity, which might have been present, although the crystals were semiconducting.

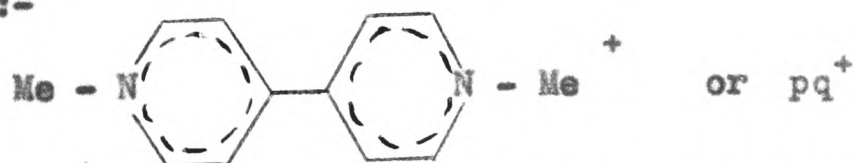
Other metal chloride salts

Crystals of other chlorides were made, including pqPdCl_4 , pqPtCl_4 and pqPbCl_4 , the last by the method above (using PbCl_2). Similar compounds with AgCl and AgBr could not be prepared but a hot solution of AgI in KI solution when added to pqI_2 solution gave dark redd crystals on cooling which were not further investigated. The lead compound crystallised as long thin pale-yellow monoclinic needles whose cell dimensions were close to those of pqCoCl_4 , and the density suggested pqPbCl_4 . Further study was abandoned because the high absorption coefficient and heavy lead atom preclude accurate determination of the light atom positions.

Crystals of pqPdCl_4 and pqPtCl_4 were made by mixing equimolar amounts of hot solutions of pqCl_2 and K_2MCl_4 ($\text{M} = \text{Pd}, \text{Pt}$). Orange rhombs crystallised quickly from the solution and the colour suggests that not much charge transfer is occurring. The compounds had virtually identical cell dimensions and were assumed to be isomorphous. The formula was confirmed by density measurement and structure analysis (Chapter VI).

Paraquat radical cation

An attempt was made to prepare crystals containing the paraquat radical ion:-

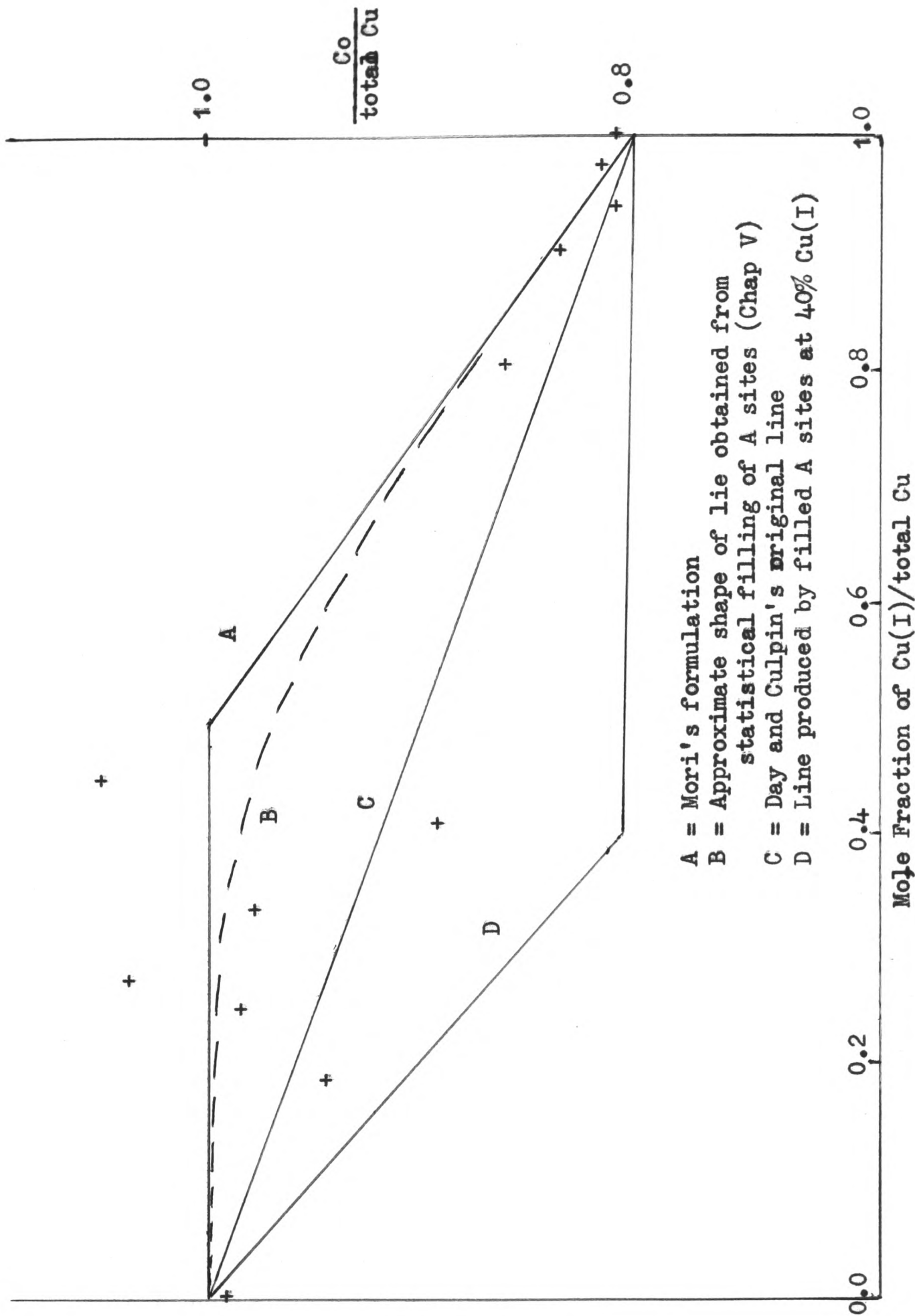


In the hope of producing a compound isostructural with the two $\text{pq}^{2+}\text{MX}_4^{2-}$ types, singly charged tetrahedral and planar AX_4^- species were used. These were BF_4^- and AuCl_4^- , and charge-transfer should be strong in the second case as the electrode potentials are similar. The

radical monocation was formed by reduction with magnesium metal of a degassed aqueous solution of paraquat dichloride under nitrogen. A very dark-blue viscous solution was obtained, and this was passed into an aqueous solution of potassium fluoroborate. Many white crystals of paraquat fluoroborate were obtained but there were some very thin ($\sim .005$ mm) dark-blue whiskers which might well contain pq^+ , but they were obviously too small for any X-ray work. A similar operation using potassium chloroaurate(III) solution produced a dark-brown intractable sludge which was not further investigated.

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(Culpin's analytical results are not ^{all} shown, for the sake of clarity. However some lie above line A and some lie on line D.)

Fig. II.1. Possible Stoichiometries for the mixed valence chlorocuprates

(A plot of the Cobalt/Copper ratio with variation of the Cu(I)/Cu ratio)

CHAPTER III.

DATA COLLECTION AND STRUCTURE SOLUTION

COLLECTION OF X-RAY DATA

(Details of compounds are given in Table III.1.)

PHOTOGRAPHY(a) Weissenberg

The accuracy of a structure determination is primarily dependent on the minimisation or corrections of errors in data collection and this is often limited by method and apparatus. The time spent on elimination of errors also increases rapidly with the precision obtained and is of doubtful value if there are appreciable approximations in the calculation of structure factors. For much of the following work only $\text{CuK}\alpha$ radiation (from rectifying but unstabilised generators) was available and was used with non integrating cameras with visual estimation of intensity. This usually leads to an 'R-factor' in the range 0.09 - 0.13 which imposes a severe restriction on the accuracy of light atom parameters ($\sigma \leq 0.07\text{\AA}$). In addition no program for calculation of anisotropic temperature factors was easily available which also limits the final accuracy attained.

One-axis data

Because of this and the nature of the compounds studied, data are not required with great accuracy and were collected from one-axis only (usually the a-axis). This technique speeds data collection by at least twice and is widely used though little justification or criticism of it has been published. We therefore note the limitations of this method, particularly when compared to the use of data from two or more axes.

- (i) Only $\sim 80\%$ of the copper sphere can be observed (if $\mu_{\text{max}} \sim 40^\circ$).
- (ii) Reflections on or near the rotation axis are unobservable or inaccurate. If, however, reflections with $\xi < 0.15$ are excluded, less than 1% of the copper sphere is lost.
- (iii) The shape of reciprocal space observed is not spherical, but only circularly symmetrical in the a-direction. Peaks which were originally spherical will be elongated along the a-axis.

- (iv) Only one independent measurement of each reflection is made.
- (v) The interlayer scalefactors are not usually available.

The first three effects are relatively unimportant as the main effect is to increase estimated standard deviations, particularly in the x-direction. An approximate formula (Cochran 12) gives $\sigma(r) \propto (R\text{-factor}) / (\text{average } \sin \theta_{\max} / \lambda \text{ for the three axes})$. For the compounds studied, this lack of data increases e.s.d.'s by $\sim 10\%$. It is noticeable that the e.s.d.'s along the x-axes are roughly this much greater than for the other two (e.g., Figure IV.2). The distortion in Fourier syntheses is small, and is unimportant unless subsequently used to predict anisotropic vibration. Two-axis data have about 65% of the copper sphere in common, and two independent measurements of the same reflection will reduce its random error and some systematic error (e.g., spot shape), as well as checking the absence of large errors. But the time needed is at least doubled and lessens the fourth objection. It is more serious when data are very limited and need as much accuracy as possible.

The absence of layer scalefactors is more serious and although some scaling can be had from one or two second axis layers, the number of common reflections is small and a different method was used. Approximate layer scalefactors were obtained from exposure times, which are inaccurate because of generator fluctuations, developing conditions and varying spot shape. They were sufficiently good, however, to locate all the atoms in the structures, and after each least-squares cycle the function $\sum w(F_{\text{obs}} k_h - F_{\text{calc}})^2$ was minimised, where k_h are the scalefactors (1).

Measurement of intensities

The contrasted reflections were estimated visually, and no spot-shape corrections (e.g., 2) were made as the integrated density was measured as far as possible. Large spots close to the axis were omitted (3) and no correction for $\alpha_1 - \alpha_2$ splitting was made. Originally interfilm factors (in a 5-film pack) were determined experimentally

giving for the four compound (Table III.1) zero-layer factors of 2.7, 2.6, 3.0, 2.9. As the value should be 2.93 (4) for $\text{CuK}\alpha$ there is a serious error in the first two which will give low values for large intensities. (This effect may well have been interpreted in many compounds as "extinction".) This is discussed in Chapter IV for the first compound. The error means that the intensity scale is effectively non-linear and is perhaps best minimised by making another or by using only a small portion of it and taking the published transmission factor (f). A term $f^{\sec\mu}$ must be used for higher layers. (This was done only for the last two compounds.)

Lorentz and polarisation corrections were applied (5) but not those for absorption. The crystals chosen were roughly cubical with $\mu a < 2$. There will be some errors because of this, but relative error between intensities is unlikely to average more than 10% and absolute error will be largely taken up by temperature factors.

(b) Precession method

This was only used, with molybdenum radiation, for the data from the 25% Cu(I) compound. This is face-centred cubic, $a = 21.8\text{\AA}$, and 5 layers of data about the 011 axis, 3 about the 001 were collected, with up to three exposures of different times for each layer. Peak-heights were measured with a microdensitometer and assumed to be proportional to the integrated intensities. L_p corrections were applied and the many common reflections on different layers gave good layer scales. The final 'merging R-factor' for all common reflections (vide infra) was 0.14 for F^2 .

COUNTER METHODS

(a) Linear Diffractometer

An early model of the Arndt-Phillips linear diffractometer (6) was available for short periods (~ 2 weeks). It had mainly been used for collecting data on protein crystals although some other data of rather low accuracy had been collected two years previously (13). Equi-inclination geometry was used which permits a Bragg angle of 30°

in any direction. However, there appeared to be inaccuracies in the ω -settings at $\mu = 30^\circ$, probably due to wear on the slides.

(1) Cobalt hexamine chlorocuprate(I)

For this molybdenum radiation with balanced Zr/Sr filters (7) was used, although there appeared to be little white radiation underneath the peak that was not proportional to the peak height. For systematic absences the count with the strontium filter was consistently larger ($\sim 20\%$) than that with Zr, and an empirical correction was made. Most reflections had an ω -spread of $\sim \frac{1}{2}^\circ$, but $1\frac{1}{2}^\circ$ was finally used to allow for slight drifts in ω , and the background was relatively low. The machine was easily set up to omit all systematic absences and, between layers, reference reflections on the zero layer were remeasured. These showed no statistically significant variation with time, implying stability of both crystal and equipment. Data processing included L_p corrections and rejection or remeasurement of reflections with significantly different backgrounds, or smaller than 3 s.d.'s of the total count.

As the system was cubic and 9 layers were collected, there was a useful check on the accuracy of measurement and also the layer scalefactors. The face-centering leads to two non-intersecting sets of data, one with all indices odd and the other with all indices even. Interlayer scalefactors can be compared within these sets but not between them, and the only solution would be data collection about a diagonal axis (as in the precession photography). Assuming that both sets are on the same scale, which is reasonable for counter data, scalefactors (k_h) from common reflections are as follows:-

Table III.2

Layer scalefactors for Cobalt hexamine chlorocuprate(I)

Layer (h)	0	1	2	3	4	5	6	7	8
Scale (k_h)	1.14	1.13	1.03	0.87	1.01	1.00	0.91	0.99	0.97

Table III.3 gives all "odd" reflections in the positive octant with more than one measurement of $|F_{hkl}^2 \mathcal{C}|$ (after scalefactor applications. Comparison with $|F_{khl}^2 \mathcal{C}|$ shows conclusively that in

general $F_{hkl}^2 \neq F_{khl}^2$ and the Lane group is $m\bar{3}$ (as opposed to $m\bar{3}m$). (It had been impossible to establish this from earlier photography, and is a demonstration of counter data for the determination of space group.) The "merging R-factor", defined as $R(a,b) = \frac{\sum |F_a^2 - F_b^2|}{\frac{1}{2} \sum (F_a^2 + F_b^2)}$ where a and b are equivalent reflections, was 0.098; this is significantly less than for the precession data and for the data compared in Chapter IV. The same function for $a = hkl$, $b = khl$ was much higher, 0.36, although low enough to suggest that F_{hkl} and F_{khl} are not completely independent.

The crystal was a near-perfect octahedron, and axial reflections ($\xi = 0$) showed a four-fold intensity variation with 360° rotation of ω . This was not more than 15% which implies that absorption is small enough to be neglected or to allow calculations from crystal geometry. Because of the inadequacy of the final structure solution (Chapter V) absorption corrections were not worthwhile. Finally, in the application of scalefactors and merging of common reflections (1), reflections where $|F_a^2 - F_b^2| > 0.2(F_a^2 + F_b^2)$ were rejected.

(ii) Paracuat chlorocuprate(II)

This data was collected with $\text{CuK}\alpha$ radiation (filtered) which limits the resolution to $d_{\min} \sim 1.5\text{\AA}$. This was only done because molybdenum radiation was not available. Only a week could be taken and a crystal had previously been aligned photographically about a non-unique axis (a). There was no time to remount it and, as there are no reflections with $\xi = 0$, there may have been significant setting errors. Serious ω -errors were encountered on higher layers ($\mu > 25^\circ$) though this may well have been due to mechanical errors. As the backgroup was small, this was partially overcome by increasing the ω -scan range to $4\frac{1}{2}^\circ$ and using a larger counter aperture. Even so, most reflections on the highest layer ($h = 4$) had to be collected manually and some were obviously mis-set (compared with F_{calc}). The data is admittedly very poor but it was able to produce a rough structure, chemically plausible, whose gross features were interesting and where time and energy were insufficient for a full photographic

determination.

(b) 3-circle diffractometer

Data for $\text{CuO} \cdot x_2 \cdot 2\text{TCNB}$ were obtained by Wright on a manually operated XRD-5 3-circle spectrogoniometer (14). It was his first experience of the machine and a few F_{obs} were very much less than F_{calc} or F_{obs} from photographs, and these reflections were excluded. Reflections at the edge of the data ($\theta \sim 50^\circ$) were also systematically weaker. $\text{CuK}\alpha$ radiation, with stationary-crystal technique (8), empirical conversion to integrated intensities (9) and a scintillation counter with pulse-height discrimination was used.

DETERMINATION OF STRUCTURES

The usual criteria for a structure solution being satisfactory are those of chemical reasonability and the agreement of observed and calculated intensities. Disagreement (a poor 'R-factor') may be due to errors in measurement of data and also approximation (e.g., the use of U_{iso}) in the calculation of structure factors. When the data may be inaccurate, chemical factors, such as interatomic distances or known geometry, are important.

The structures were all centrosymmetric and were solved by Patterson and heavy-atom techniques and also by postulating chemically reasonable structures. The heavy-atom method is particularly favourable if chlorine atoms can be located from the Patterson function as this usually removes any false symmetry. An indication of its power is shown in Figure III.1 which represents, for 4 structures, electron density in the plane of the paraquat ion phased on all the metal and chlorine atoms in each structure.

With at least the metal + chlorine atoms located, full-matrix least squares refinement was carried out. The program (10) allowed only for U_{iso} and a choice of 3 weighting schemes. Unit weights were used initially until refinement was nearly complete and then weights w , where $\sqrt{w} = F^*/F_{obs}$ unless $F^* > F_{obs}$ when $w = 1$. F^* was chosen so that $\sum w\Delta^2$ was approximately independent of F_{obs} , which worked well for photographic data so that systematic errors in Δ are minimised. Convergence was rapid and refinement was stopped when parameter shifts were less than their e.s.d.'s.

Determination of Individual Structures

There was a nice transition in the structure solutions from the mainly automatic Patterson and heavy atom methods to a chemical approach combined with the use of space-group positions. The outlines of the steps taken for each structure are described and summarised in Tables III, 4 - 7, as are final parameters. Diagrams of the structures,

which may clarify their solution, are given in Chapter VI.

CuOx₂.2TCNB was assumed to be in $P\bar{1}$ with the copper atom at the origin and structure factors dominated by its contribution. A map with all structure factors positive showed the 25 light atoms well resolved with only two spurious peaks. Only 90% of the final structure factors were negative which agrees well with theoretical predictions (e.g., Parthasarathy).

Paraquat PdCl₄, with only four molecules for the 16-fold general positions of the Patterson function showed all the necessary vectors for a square planar PdCl₄²⁻ ion at (a) and evidence, from Pd to light atom vectors, for a twisted pq²⁺ ion at (b). This latter was verified by a Fourier synthesis, on Pd and Cl only, which gave the new electron density as in Figure III.1d. No evidence for the non-centrosymmetric space-group Iba2 was seen.

Paraquat CoCl₄, in Pnab (an alternative setting of Pbcn), must have both ions on either a centre of symmetry or two-fold axis. The latter is the only real chemical possibility for CoCl₄²⁻ and the Patterson function was interpreted to give the orientation and position of these tetrahedral species. A subsequent Fourier synthesis is shown in part in Figure III.1a where the paraquat ion is on a two-fold axis and there are no significant spurious peaks.

Paraquat Cu₂Cl₆, in P2₁/n, was analysed with low resolution data ($d_{\min} = 1.5\text{\AA}$) which means that all peaks are poorly defined. Because of this there were two possible solutions for the position of the copper atom from the Patterson and vectors from chlorine atoms were not reliably identified. The first solution was chemically unlikely and Fourier syntheses based on it did not show reasonable chlorine atoms. A map on the second position, however, showed two well-resolved chlorine atoms and when these were introduced into the calculation, the third was clearly shown.

Calculated structure factors appeared encouraging at this stage and another Fourier synthesis gave Figure III.1c, clearly a paraquat ion at low resolution. Final structure factor agreement was not good

($R = 0.128$) for counter data but there is little reason to doubt the structure and it is supported by chemical evidence (Chapter VI). Some of the outer reflections on higher layers appeared to be mis-set as F_{obs} was much lower than F_{calc} (these were removed). The temperature factors were ill-determined due to the lack of data and some were just negative, possibly due to absorption. No reliance is put on light atom bond lengths (and these are not listed) but the approximate planarity of the paraquat and the five-co-ordination of the copper atom are clear. It is most unlikely that a significantly different structure would be discovered by a more accurate determination.

Paraquat Cu_2Cl_4 crystallised in $P4_2/n$ (unique axis c) which has 8 general positions. As there are 4 formula units in the cell, the paraquat ion is likely to be at a centre of symmetry or possibly on one of the two diads. Figure III.2 shows the symmetry elements in $P4_2/n$ (and also some of the false symmetry, marked *, which is particular to this problem). The chlorine atoms probably lie in two sets of general positions.

The Patterson function had large peaks only at $\frac{1}{2}\frac{1}{2}0$, $\frac{1}{2}\frac{1}{2}0$ (and symmetry-related positions) which strongly suggest that two or more special positions are occupied by the copper atoms. The cell dimensions (particularly the c -axis) and chemical considerations suggest a similarity with the last compound which has chains parallel to the short axis (Chapter VI). Two independent chains of $(\text{CuCl}_2)_n$ can be made by placing 2 atoms in each of $\frac{1}{4}\frac{1}{4}\frac{1}{4}(a)$ etc. and $\frac{1}{4}\frac{1}{4}\frac{3}{4}(b)$ etc. (both of symmetry $\bar{4}$) and the other 4 on the 4_2 screw at $\frac{1}{4}\frac{3}{4}z(c)$ etc. Assuming that the largest Patterson peaks are copper-copper vectors, we have $z \sim \frac{1}{4}$. (N.B. The origin is taken at $\bar{1}$ and all calculations are done on this basis.)

Two infinite chains of edge-linked CuCl_4 tetrahedra can be made by putting the 16 chlorine atoms in two sets of general positions (incidentally leaving no space for the paraquat ion on the diad axis). The placing of the chlorine atoms was very difficult, for, although the

orientation of both sets of tetrahedra could be clearly found from the Cu-Cl and intramolecular Cl-Cl vectors, there was no indication of their relative positions in the unit cell. In particular an isolated chain of edge linked (regular) tetrahedra involves both the symmetry elements 4_2 and $\bar{4}$ which creates further ambiguity. The only definite key at this stage, the identification and unravelling of the very large number of intermolecular Cl-Cl vectors, was quite impossible.

A Fourier synthesis on the copper atoms is useless as is shown by the false symmetry elements (*) in Figure III.2. Firstly there is a substructure with $c' = c/2$ so that copper atoms do not contribute to reflections with $l = 2n + 1$. This would be true even if the z co-ordinate of (c) were not $\frac{1}{4}$ but the fortuity of this introduces a second problem in false mirror planes (normally n -glides) at $z = 0, \frac{1}{4}$ etc. and also a false face centering for the sub cell $a' = b' = a/2$. These three independent symmetry operators create an ambiguity of origin, at P or Q, but having made an arbitrary choice we automatically define the 4_2 and $\bar{4}$ (but not which tetrahedra lie on which).

The Patterson function showed two sets of chlorine atoms at translations of 0.093, 0.067, $\pm \frac{1}{4}$ etc. (denoted by (A)) and of 0.017, 0.111, $\pm \frac{1}{2}$ (B) from copper atoms, making approximately regular tetrahedra, at undefined positions. There are 4 possible ways of arranging (A) and (B); two resulting from the ambiguity of which is on the 4_2 axis and this is doubled as there are two relative z -heights ($\pm \frac{1}{4}$) for the pair. A trial-and-error method of structure factors and least squares would be very time-consuming, very possibly unreliable and a different approach was adopted.

A Fourier synthesis was carried out with the B set attached to the copper atoms on the 4_2 screw. This axis was chosen because the chance value of $z = \frac{1}{4}$ leaves the false mirror planes unbroken (Position X in Figure III.2). This map gave a clear indication of the absolute z -co-ordinate of the A set and x -, y - parameters consistent with the Patterson. Both sets were included in a least squares

refinement of all the atoms and gave good agreement, $R < 0.3$. This solution, though seemingly correct, may not be the only one and the process was repeated with (A) used initially. This map showed not only the atoms put in but 'ghost' peaks closely related to the first solution.

Another Fourier synthesis, based on the first set of positions gave the electron density in Figure III.1b, a clearly resolved paraquat ion, which is additional indication of the correctness of the heavy atom positions. Refinement continued, but may be slightly meaningless for the z-parameters of the copper(c) and chlorine atoms because of the strong local supersymmetry of the heavy atoms. The correlation matrix (see Chapter IV) was calculated and showed no large covariances involving any positional parameters. The chlorine bridges appear to be fairly symmetric and this is quite acceptable. It had been hoped that the data, which appeared reasonably good, would have given better agreement ($R = 0.119$) but this might be due to errors in refinement.

This is the only compound solved from one-axis data that has a possible degeneracy of layer scales with heavy atom position, as the others have heavy atoms fixed in this direction (a-axis). The copper atoms make a contribution of $4f_{Cu} (1 + \cos 2\pi lz)$ to even layers implying that small errors in the scale factors may be reflected in a slightly wrong value of z. This function can be rewritten as $8f_{Cu} \cos^2 \pi lz$ (where $z \sim 0$) and errors are likely to be important only if they are roughly similar to $\cos^2 1$.

The cobalt hexamine chlorocuprates (I) and (I,II) cannot be discussed without detailed comparison of their structures and the chemical evidence. A separate Chapter (V) is devoted to this.

Results

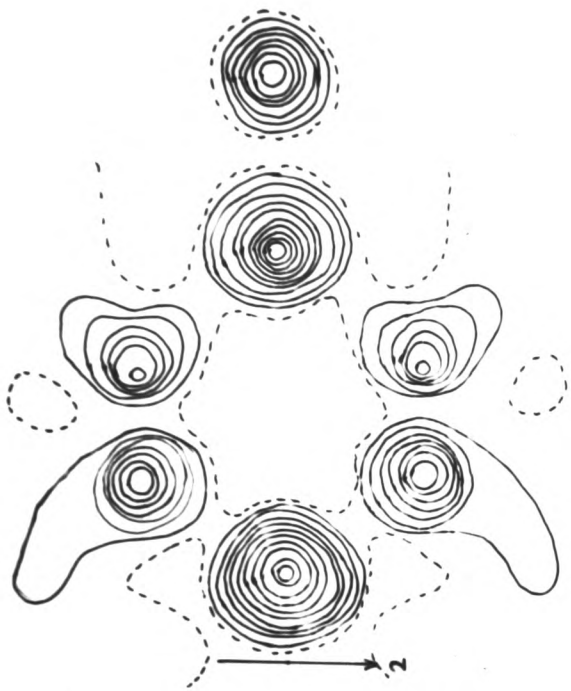
The parameters of the paraquat compounds are listed in Tables III, 4, 5, 6, 7 which also give crystal data. The bond lengths and angles in the paraquat rings are given in Figure III.3a,b,c (without $pqCu_2Cl_6$)

and those of the anions in Chapter VI. Parameters, bonds and angles of $\text{CuOx}_2 \cdot 2\text{TCNB}$ are discussed in Chapter IV.

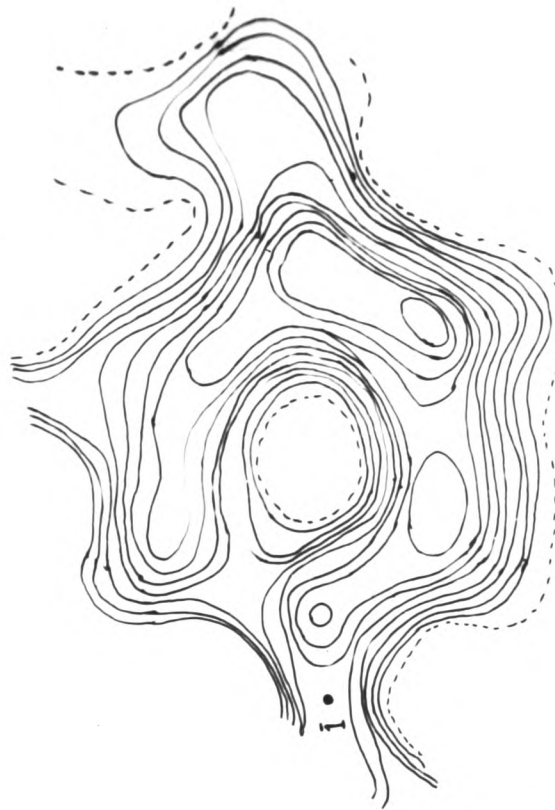
Scattering factors are those in INTERNATIONAL TABLES Vol. II (II) and anomalous dispersion corrections have been applied where necessary. Where necessary, ionic scattering factors have been used. Fourier syntheses in tetragonal space groups were calculated in monoclinic sub-groups (e.g., $P2/n$) with structure factors for $\bar{k}\bar{h}\bar{l}$ being included.

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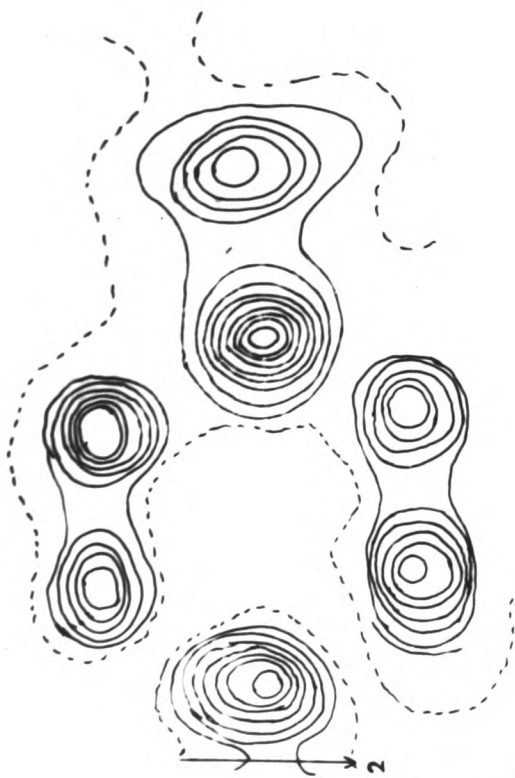
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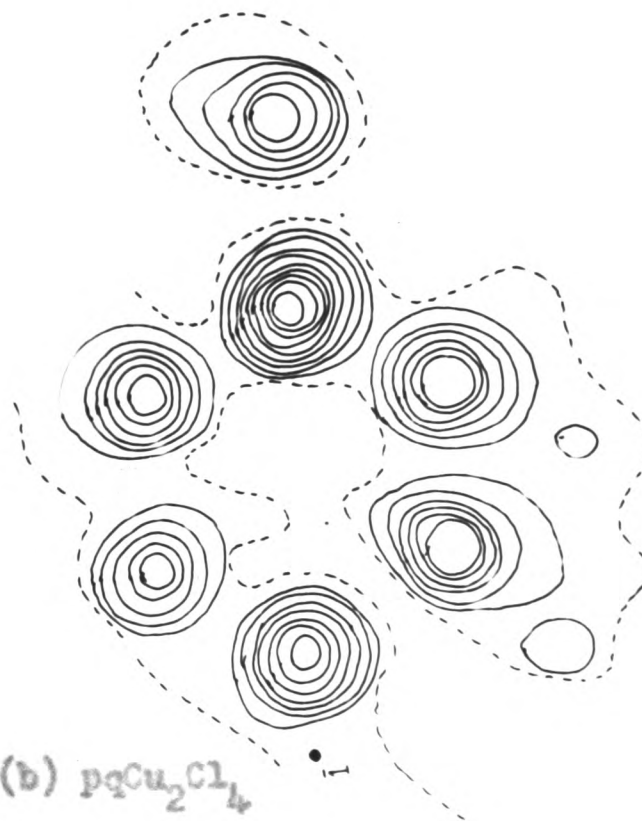
(a) $pqpAcCl_4$



(c) $pqpCu_2Cl_6$



(a) $pqpCoCl_4$



(b) $pqpCu_2Cl_4$

Fig. III.1. Electron Density (heavy atom-phased) in pq compounds

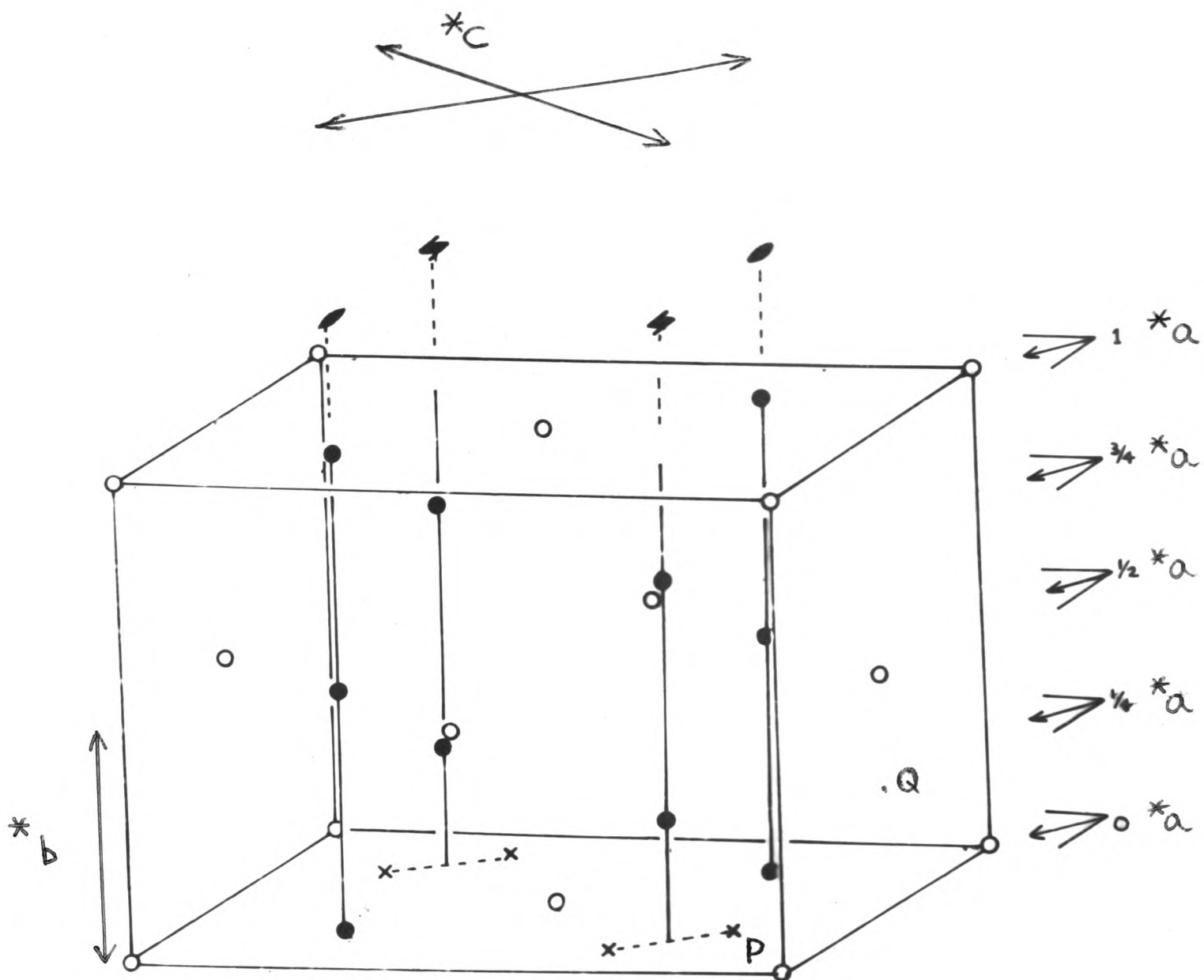


Fig. III.2. Symmetry Elements (and false symmetry) in $P4_2/n$.

N.B. False symmetry is as follows:

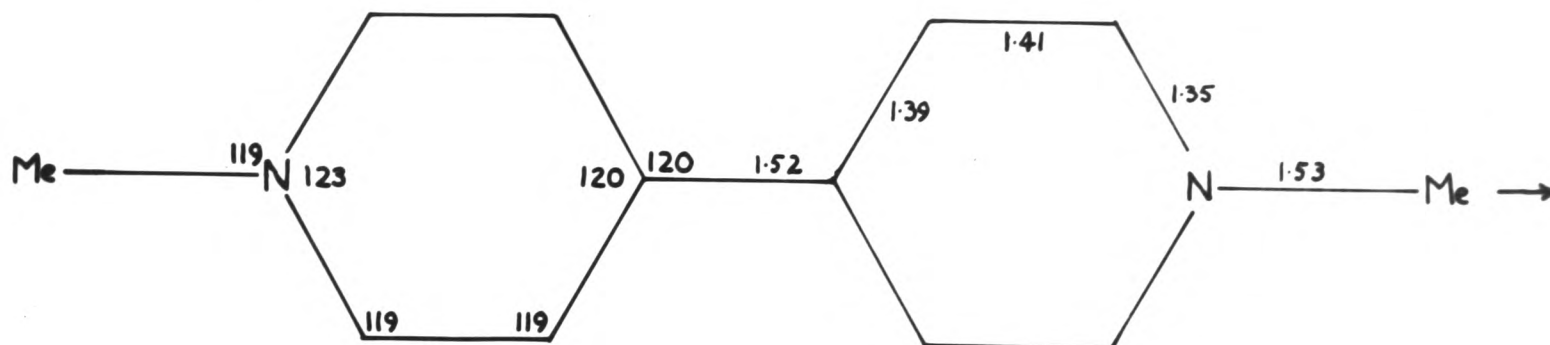
- *a, Mirror planes at $nc/4$
- *b, Subcell along c-axis
- *c, C-face-centering
- *d, Ambiguity of origin, P or Q? (P was finally chosen.)

(B) tetrahedra lie round the ● on the 4_2 screw.

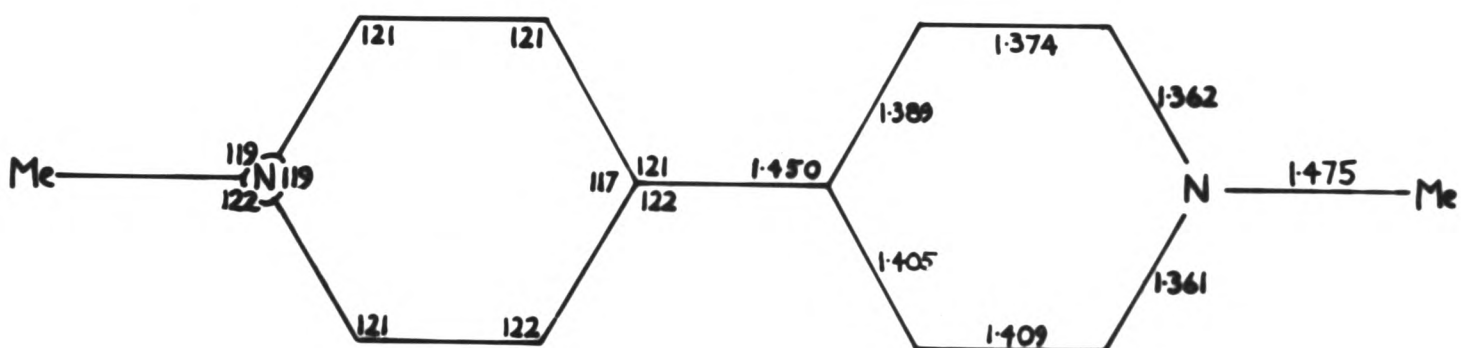
● Represents copper atoms

○ " centres of symmetry, $\bar{1}$. There are a similar set, not shown, at a translation $(\frac{1}{2}\frac{1}{2}\frac{1}{2})$.

Paraquat $PdCl_4$



Paraquat $CoCl_4$



Paraquat Cu_2Cl_4

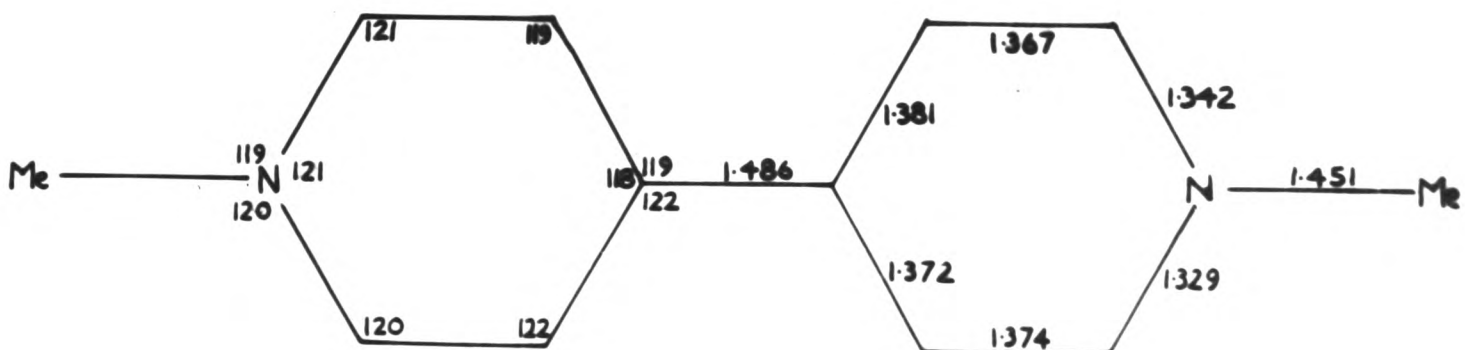


Fig. III.3e

Bond lengths and angles in the paraquat rings.

Table III.3

Intensities(F^2) for $[\text{Co}(\text{NH}_3)_6]_4\text{Cu}_5\text{Cl}_{17}$, comparing F_{hkl}^2 with F_{khl}^2

Layer No.	h	k	h	k	h	k	h	k	h	k	h	k
	1	3	1	5	1	7	3	5	3	7	5	7
hk 1							102				64	
kh 1							124				28	
hk 3	28	25	125		192	215					83	
kh 3			97		203							134
hk 5	97	102	69	75	32	32	230	221	141	134	2	
kh 5	125	124			64				84			
hk 7	203	190	64	65			84	83			6	4
kh 7	192		32	29			141					
hk 9	25	24	16	14	77	105	7	6	58	62		1
kh 9	15		6	6	147	172		10		27		
hk 11	20	16	55	52	44	68	93	87	34	35	24	24
kh 11	13	12	43	42	101	122	104	118	22	17	12	9
hk 13	52	50	14	13	33	31		18	12	10		
kh 13	57	56	9	6	28	25	15	12	16	18	3	
hk 15		53	7	10	3	7	52	52	20	23	2	2
kh 15	67	61	16	20	9	11	42	35	15	11	2	2
hk 17	9	7	10	13	26	30			3	3	11	13
kh 17	10	7	13	14	33	38			2	3	21	25
hk 19	6	9	5	5		5	39	39	17	20	10	9
kh 19	14	11	8	13	7	6	45	65	23	31	4	6
hk 21	7	8			18		18	18		9	6	6
kh 21	7				18	13	12	12		5	4	4

The values are (count $\times Lp^{-1} \times 0.1 \times k_h$) where k_h are layer scales, given in the text.

The numbers at the head of each column indicate the layer on which data was collected.

Horizontally related numbers in each "box" should agree, but vertically related ones (corresponding to reflection about $\{110\}$) need not.

(Estimated Standard Deviations are mainly between 1 and 2)

Table III.4.

Details of Paraquat CoCl₄

Full Name: N,N'-dimethyl-4,4'-bipyridylium tetrachlorocobaltate(II)

Formula: C₁₂N₂H₁₄CoCl₄ Spacegroup: Pnab, No. 60. Z= 4

Dimensions of crystal (mm): 0.18 x 0.15 x 0.14 $\mu = 187 \text{ cm}^{-1}$ (CuK α)

Cell parameters: a=7.70 b= 16.04 c=12.40 (\AA)

Number of reflections: 926 (from a-axis) Final R= 0.117

Parameters refined, including scales: 38 + 7 scales.

Final parameters ($\times 10^4$)

	x/a	σ_x	y/b	σ_y	z/c	σ_z	Uiso	σ_U
Co	2500	-	2179	2	0	-	259	8
Cl(1)	747	6	1364	2	-990	2	398	10
Cl(2)	1091	5	3038	2	1146	2	296	9
N	78	16	895	7	2636	8	303	26
C(2a)	589	21	165	9	3091	11	390	33
C(2b)	503	20	1630	8	3108	11	373	32
C(3a)	1529	20	161	8	4016	10	367	31
C(3b)	1466	19	1650	8	4046	10	331	29
C(4)	2000	17	901	7	4517	9	288	27
C(Me)	-993	23	916	10	1675	12	476	38

Bond Distances and Angles in the Anion.

Co - Cl(1) = 2.26 \AA

Co - Cl(2) = 2.28 \AA

Cl(1) - Co - Cl(1) = 109.4 $^\circ$

Cl(2) - Co - Cl(2) = 105.6 $^\circ$

Cl(1) - Co - Cl(2) = 106.5, 115.8 $^\circ$

Cl(1) - N = 3.82 \AA

Shortest contact in the " molecular vice " is Cl(2) - C(2a) = 3.35 \AA

Table III.5.

Details of Paraquat PdCl₄

Full Name: N,N'-dimethyl-4,4'-bipyridylium tetrachloropalladate(II)

Formula: C₁₂N₂H₄PdCl₄ Spacegroup; Ibam, No.72. Z= 4

Dimensions of crystal (mm): 0.21 x 0.17 x 0.12 $\mu = 149 \text{ cm}^{-1}$

Cell parameters: a= 8.43 b= 13.50 c= 13.41 (Å)

Number of reflections: 595 (from a-axis) Final R= 0.097

Parameters refined including scales: 20 + 7 scales

Final Parameters (x10⁴)

	x/a	σ_x	y/b	σ_y	z/c	σ_z	U _{iso}	σ_U
Pd	0	-	0	-	0	-	193	7
Cl(1)	-2376	6	850	3	0	-	340	12
Cl(2)	1351	5	1500	3	0	-	256	10
N	0	-	2398	12	2500	-	277	32
C(4)	0	-	4436	12	2500	-	378	36
C(3)	4378	17	1076	8	1693	9	274	25
C(2)	4404	17	2124	8	1701	9	310	26
C(Me)	0	-	1266	16	2500	-	378	46

Bonds and Angles in the Anion.

Cl(1) - Pd = 2.31 Å

Cl(2) - Pd = 2.32 Å

Cl(2) - N = 3.78 Å

Cl(1) - Pd - Cl(2) = 89.5°

Table III.6.

Details of Paraquat Cu_2Cl_4

Full Name: N,N'-dimethyl-4,4'-bipyridylum tetrachlorodicuprate(I)

Formula: $\text{C}_{12}\text{N}_2\text{H}_4\text{Cu}_2\text{Cl}_4$ Spacegroup: $P4_2/n$, No.86. Z= 4

Dimensions of crystal (mm.): 0.32 x 0.25 x 0.17 $\mu = 96 \text{ cm}^{-1}$ (CuK α)

Cell Parameters: a= b= 16.28 c= 5.90 (Å)

Number of reflections: 1199 (from c-axis) Final R = 0.119

Parameters refined, including scales: 40 + 6 scales.

Final Parameters ($\times 10^4$)

	x/a	σ_x	y/b	σ_y	z/c	σ_z	U_{iso}	σU
Cu(1)	2500	-	2500	-	2500	-	618	8
Cu(2)	2500	-	2500	-	7500	-	586	7
Cu(3)	2500	-	7500	-	-452	4	557	6
Cl(1)	3426	1	3164	1	5031	3	355	6
Cl(2)	8617	1	2335	1	-2102	3	363	6
N	9901	3	2007	3	2023	10	322	12
C(4)	-15	4	424	4	467	12	293	13
C(3a)	445	4	1035	4	-554	13	346	14
C(3b)	9467	5	637	5	2206	16	477	19
C(2a)	407	4	1815	4	303	13	365	15
C(2b)	6435	5	573	5	-2037	15	450	18
C(Me)	7849	5	138	5	-2195	15	423	16

Table III.7.

Details of Paraquat Cu₂Cl₆

Full Name:- N,N'-dimethyl-4,4'-bipyridylium hexachlorodicuprate(II)

Formula: C₁₂N₂H₄Cu₂Cl₆ Spacegroup: P2₁/n, No.14. Z=2

Dimensions of crystal(mm.): 0.27x0.20x0.15. $\mu = 101 \text{ cm}^{-1}$ (CuK α)

Cell parameters: a= 6.29 b=10.74 c= 13.02 (Å) $\beta = 98.65^\circ$

Number of reflections: 309 (from a-axis) Final R = 0.126

Parameters refined, including scales: 44 + 1 scale.

Final Parameters (x10⁴)

	x/a	σ_x	y/b	σ_y	z/c	σ_z	Uiso	σ_u
Cu	2549	9	564	6	443	5	80	33
Cl(1)	4435	15	4478	11	3927	8	-140	42
Cl(2)	9473	16	3433	12	4621	9	-42	44
Cl(3)	3378	15	1455	11	1987	9	-110	41
N	263	50	3743	35	1306	28	-110	110
C(4)	4042	67	4710	45	280	38	88	154
C(3a)	2605	67	1144	45	4774	36	37	145
C(3b)	8584	77	40	51	3654	42	64	166
C(2a)	643	70	3456	46	237	38	-64	148
C(2b)	1713	73	4595	49	1759	36	-35	149
C(Me)	8328	57	3182	41	1846	30	-174	119

CHAPTER IV.

ANALYSIS OF THE ACCURACY OF STRUCTURES

(A comparison of various parameters and their estimated errors
from identical or closely-related structures)

COMPARISON OF DATA AND PARAMETERS FROM SIMILAR STRUCTURES

Wright and I collected data, and solved the structure of, $\text{CuOx}_2 \cdot 2\text{TCNB}$ independently, and, as both of our methods of data collection are widely used, the data and parameters are compared in the light of their estimated standard deviations (e.s.d). This showed the use of one-axis data, still very widely used (1), to be a reasonable method in this case, even though layer scale factors were refined against the calculated structure. Very little has been published about the duplication of three-dimensional data collection for a particular compound and the comparison is usually not in terms of atomic parameters. The comparison appeared successful and was extended to an empirical investigation of the published e.s.d.'s for bonds and angles in similar structures. A short description of the statistical method used is given so that the notation is defined. It is clear that only the estimated precision is analysed and this cannot in any way appraise the accuracy of the answer.

Statistics of X-ray Analysis (cf 2, 3, 4)

Counting statistics (5) give a prediction of the uncertainty of an intensity measurement, but (although reasonable if F_{obs} is small) the actual disagreement with F_{calc} is usually several times greater (6). This is due to reproducible errors in F_{obs} and F_{calc} (although these often cause random errors in the parameters (p_i) if the problem is well-overdetermined; observations (m) \gg parameters (n)). This consequence of the Central Limit Theorem allows us to calculate an s.d., $\sigma(p_i)$, for each parameter by assuming that the difference ($F_{\text{obs}} - F_{\text{calc}}$) is random. The $\sigma(p_i)$ will contain both truly random and reproducible errors (but not systematic errors which are absorbed by parameter shifts).

An estimate of $\sigma(p_i)$ is c_{ii}^{-1} , a diagonal term from the inverse matrix, but if absolute weights are not known, this is modified to $c_{ii}^{-1} \sum w \Delta^2 / (m - n)$, which holds only for a full-matrix treatment. If

the off-diagonal terms of this matrix are non-zero, the parameters are not independent, and statistical analysis is difficult; this was not the case in the present work as is shown later. The second formulation for $\sigma(p_i)$ was used and was determined from a full-matrix (4).

The probability of two measurements p_{i1} and p_{i2} , with s.d.'s $\sigma(p_{i1})$ and $\sigma(p_{i2})$, being different is related to $t_i = \Delta_i / \sigma(\Delta_i)$, where $\Delta_i = p_{i1} - p_{i2}$, and $\sigma(\Delta_i) = \{\sigma^2(p_{i1}) + \sigma^2(p_{i2})\}^{1/2}$. The probabilities that a randomly generated t_i exceeds 1.645, 2.33 or 3.00 are 5, 1, 0.1% respectively. This use of t_i is useful where only single parameters are to be compared, but for several independent t_i their overall distribution must be considered. In particular, for n random t_i , $\sum_{i=1}^n t_i^2 \sim n$; the probabilities that it is substantially higher or lower are those for χ^2 (with n degrees of freedom) exceeding a particular value (3).

If these probabilities are unallowably high or low, then $\sigma(\Delta_i)$ must be an under- or overestimate of the actual uncertainty or reproducibility. In particular, a low value of $\sum t_i^2$ suggests appreciable reproducible error common to both determinations (which might be due to absorption, extinction, and particularly errors in F_{calc} , as the same method was used for both sets of data). A high value is probably due to an error being absorbed in parameter shifts and not being reflected in $\sigma(\Delta_i)$. A reasonable value of $\sum t_i^2$ may, of course, be made up of equal amounts of these effects.

The duplicated structure determination of the complex copper oxinate 2TCNB

The structure was originally determined because it was hoped to study a series of similar compounds, and particularly to relate single-crystal spectra to the simple molecular geometry of the triclinic system. The chemistry of the compound is not uninteresting and is described in a paper to be published (8). If a preprint is available it is hoped to include this in the thesis, but the compound is included solely for the analysis of its X-ray data and parameters.

Data collection and treatment

The counter data collection (8) has already been described but some extra details of the photographic method are given. Seven layers, each consisting of two settings 180° different in ω were taken about the a -axis (nkl). The two settings on any one layer were exposed for equal times and assumed to be on the same scale as spots common to both appeared equal and this was confirmed by comparison with calculated structure factors. A zero layer photograph (h0l) was taken, but not used for layer scale factors, which were refined only after all atoms had been included and after every second round of least-squares.

Polarisation (and Lorentz), but not absorption, corrections were made. This is justifiable as an axial reflection ($\chi = 90^\circ$) on the diffractometer showed about 6% variation with β . Wright's data (subscripted W) was refined to $R \sim .102$, subsequently improved by calculation in Oxford where comparison with my data (M) suggested removal of some inaccurate reflections of high θ . Both sets of data were then treated identically, with a full normal matrix and isotropic temperature factors (except for assumed values for hydrogen atoms). Final difference maps showed nothing significant except marked anisotropic thermal motion for the copper atoms. The residuals were $R_W = .082$, $R_M = .102$, neither including unmeasured reflections. It was found that all the phases were identical for the two structures which is not surprising as the weakest $F_{obs M}$ were not measured and the refinement was similar in each case.

Comparison of the determinations

Comparisons are made of the parameters, with individual t_i , (Table IV.1) and some bonds and angles with individual t_i for the oxinate rings only (Table IV.2). Common structure factors are not listed, but $|F_{obs W}|$ and $|F_{calc W}|$ are published (8) and $|F_{obs M}|$ with $|F_{calc M}|$ are in Table VII.1.

Intensity data

The $|F_{\text{obs}}|$ data cannot be compared unless they are on correct relative scales, but neither set involved any absolute scale determination. Two possible methods to determine the relative scales (K) are to minimise a function of the sort $\sum w \{ |F_{\text{obs}}|_M^k - K |F_{\text{obs}}|_W \}^2$, or to take $K = \frac{GM}{GW}$, where these are the scale factors obtained from refinement. In the first case, layer scales can be refined for F_M , and can be compared with those obtained from exposure times and the structure refinement.

The first method gives a merging R factor (Chapter III) for $|F_M|$ to $|F_W|$ of 0.079 (0.148 for F^2) and the second 0.09 (0.168 for F^2). The parameter K refined to 1.03 GM/GW (where $1/G$ is used to put F_{obs} on an absolute scale). This 3% difference between the structures is just compatible with the e.s.d.'s for GM (0.8%) and GW (0.6%), but it was found that the scales have high covariance with the temperature factors of the copper atom. This high value may also result from an effect shown in Figure IV.1, which shows plots of $\langle F_{\text{obs}M}/F_{\text{obs}W} \rangle$, $\langle F_{\text{obs}M}/F_{\text{calc}M} \rangle$ and $\langle F_{\text{obs}W}/F_{\text{calc}W} \rangle$ for various ranges of F . The first suggests strongly that $|F_{\text{obs}}|_M$ is systematically lower than $|F_{\text{obs}}|_W$ for high reflections, almost certainly due to the low film factors. Without the counter data, this might well have been ascribed to 'extinction', and it is surprising to what an extent the parameters are able to compensate for this effect (as the second curve shows). This compensation appears to lead to no systematic differences between the parameters obtained from both sets of data, with the possible exception of the temperature factor of the copper atom.

The layer scales determined for $(F_{\text{obs}})_M$ by 3 different methods are listed in Table IV.3.

Table IV.3

Values of k_h to be applied to $(F_{obs})_M$

Layer	From Exposure Times _M	From $ F_{calc} _M$	From $ F_{obs} _W$
0	1.0	0.97	0.94
1	1.0	1.12	1.11
2	1.0	0.93	0.91
3	1.0	0.96	0.96
4	1.0	1.22	1.25
5	1.0	0.81	0.88

The layer scales obtained from exposure times are poor, but as they are not wildly wrong, they can be refined against F_{calc} to give surprisingly good agreement with the third column. If we assumed that no layer scaling is necessary for the 3-circle data, then the method employed for one-axis photographic data is within the limits of acceptable error. Certainly a large amount of second-axis data would be needed to better this agreement. This is not a general justification of the process used, but suggests that it is valid if there is a heavy atom whose position cannot affect layer scales. It may well result in an incorrect temperature factor for this atom, and precludes meaningful anisotropic refinement (). It is also likely to be invalid for more accurate determinations of parameters where vibration must be accurately analysed.

Parameters

$\sum_{i=1}^{25} t_i^2$ were calculated for the four sets of parameters: x-, y-, z- co-ordinates and isotropic temperature factors, U, and are given in Table IV.4.

Table IV.4

$$\sum_{i=1}^{25} t_i^2 \text{ for various parameters} = \sum_{i=1}^{25} \left\{ (p_i)_M - (p_i)_W \right\}^2 / \left\{ \sigma^2(p_i)_M + \sigma^2(p_i)_W \right\}$$

Parameters	$\sum_{i=1}^{25} t_i^2$	Probability that $\chi_{25}^2 > \sum_{i=1}^{25} t_i^2$
25 x-co-ordinates	10.0	0.99
25 y-co-ordinate	18.2	0.80
25 z-co-ordinates	30.9	0.20
25 U-parameters	32.3	0.15

(t_i for the copper atom was 8.8, which is very improbable, and may well be due to the high covariance between G and U_{Cu} or the lack of anisotropic refinement.)

It is clear that there is good correlation between the observed and estimated difference with the exception of x-co-ordinates. This is the axis for which 6 layer scales were refined and it is possible that the numbers of degrees of freedom should be reduced from 25 to 19. The probability that $\chi_{19}^2 > 10.0$ is only 0.80 which is insignificant. Because of the slightly lower resolution in the x-direction, the e.s.d.'s of these parameters (M) are $\sim 10\%$ greater than for other axes. Overestimation of these might account for lower t_i values but even so $\sum_{i=1}^{25} t_i^2 > 12$ which is still significantly low. The good agreement of the temperature factors, depicted in Figure IV.3, implies that X-ray absorption in the two crystals is roughly equivalent.

It is interesting to compare the standard deviations of $(p_i)_M$ and $(p_i)_W$. There is very good correlation (Figure IV.2) and this is perhaps explained by Figure IV.4 where $\frac{1}{2}(\bar{\sigma}_M + \bar{\sigma}_W)$ for any atom (in Å) is plotted against $\frac{1}{2}(U_M + U_W)$. The positional s.d. of an atom, estimated from a Fourier synthesis (2), is proportional to the curvature of that atom, made up of experimental errors, form factor and temperature factor effects. Figure IV.4 suggests that $\sigma(r) \sim 0$ when $U = -0.03$ and, if the scattering factors for light atoms can be written as (2):

$$f(s) = Z \exp(-ps^2) \quad \text{where } s = \sin \theta / \lambda$$

then $p \sim 2.5 \text{Å}^{-2}$.

The validity of this analysis of e.s.d.'s has relied on neglect of any covariances, $\text{cov}(p_i, p_j)$, which are estimated by the off-diagonal terms c_{ij}^{-1} of the inverted normal matrix. The correlation matrix (10), where $N_{ij} = c_{ij}^{-1} / \sqrt{c_{ii}^{-1} c_{jj}^{-1}}$ was calculated for (M) data. The leading diagonal must be unity but of the other 5151 terms, few had significant moduli. $N(G, U_{\text{Cu}}) = 0.83$, $N(G, U_i) \sim 0.2$, $N(U_{\text{Cu}}, U_i) \sim 0.15$ and for the x - y - z - parameters of the same atom, $N(x_i, y_i) \sim N(x_i, z_i) = 0.13$ and $N(y_i, z_i) \sim 0.19$, possibly correlated with the deviations of the axes from orthogonality. The remaining ~ 5000 terms contained only 16 terms > 0.1 and their average was ~ 0.02 .

The analysis has supported the idea that e.s.d.'s for a well overdetermined structure (three-dimensional) are meaningful and that isotropic temperature factors have some significance. It would be risky to extrapolate this to analyses of much higher accuracy.

Previous comparisons of sets of data for a single compound

The above results are interesting mainly because of the lack of similar published comparisons. None of the few I shall mention are exactly similar.

The reproductivity of identical counter measurements from the same crystal (CaF_2) by different operators showed these could have "merging R's" on F^2 of 2 - 13%. Systematic errors in $\sin \theta$ could lead to a ΔB of 0.3 (18). Previous experiments on NaCl showed that $|F^2|$ data came from the same population (19). No analysis of positional parameters could, of course, be done.

Photographs are also capable of high accuracy (20) but comparisons of good photographic data are rare. Two examples of relatively low accuracy are given.

Independent refinement of di-paraanthracene from two one-axis sets of photographic data about different axes showed that within the (high) e.s.d.'s for bond lengths (0.02\AA) and B's (0.4\AA^2) the two sets of data were consistent, although both were of limited resolution (21).

A comparison of two-axis $\text{CuK}\alpha$ photographic data and limited MoK

data from one axis (collected on this linear diffractometer) can be made from an analysis of a bromoterpenoid (22). Two sets of parameters were independently refined but not compared (and both R-factors were, very high). Taking these values, we get for 47 parameters:- $\sum_{i=1}^{47} t_i^2$ for x- y- z- and B-parameters = 62, 282, 75, 138. The temperature factors might agree better if absorption corrections had been made as B_{Cu} was usually greater than B_{Mo} . Values of 62 and 75 are just allowable, but the e.s.d's along the rotation axis (y) are underestimated by a factor of at least 2. The data were much more limited in that direction than for normal one-axis photographic data. It should be noted that the analysis was for a non-centrosymmetric space-group and did not use a full matrix approximation.

COMPARISON OF RELATED STRUCTURES

If the e.s.d.'s of a structure are reliable, it is interesting to compare its parameters with those of a related one, but this cannot be done directly except for isomorphous structures. However, a molecular entity of n atoms can be defined by $3n-5$ co-ordinates, with a suitable choice of axes. These parameters are usually the reported bonds and angles, though these do not always have individual e.s.d.'s. An ideal set of structures for this is a series whose members each include a compound common to all, which is in an identical chemical state throughout. Two examples of this are molecular complexes with a common donor (metal oxinates) for which 8 reasonable structures are available and ionic compounds with a common cation (paraquat) - 6 structures. If we assume that the bond lengths and angles are unaffected by small changes in electronic configuration or molecular packing, then they should be identical within the estimated error. The absence of sizeable distortions in compounds with non-crystallographic symmetry suggests that packing forces can often be neglected.

It is often assumed that if in two structures no bond lengths differ by more than 3 e.s.d.'s, then these structures are identical. Although this argument can, with caution, be used for a single bond, for assessing two structures as a whole, the deviations must be taken together. It would be best to refer these to a common origin by minimising $\sum_{n \text{ atoms}} w(\Delta \underline{r})^2$, where $\Delta \underline{r} = \underline{r}_a - \underline{r}_b$ and then to evaluate $\sum t_i^2$, where $t_i = (\Delta \underline{r})_i / (\text{e.s.d. along the line } a \rightarrow b)$. This is tedious and the simpler comparison of bonds and angles has been made. Although the atoms are assumed independent of one another, it is impossible to choose bonds and angles to satisfy this. The analysis should however show whether the e.s.d.'s of bonds and angles are of the correct order, or are useful for comparing structures.

Metal oxinates (MOx_2)

Apart from our determination(s) of $\text{CuOx}_2 \cdot 2\text{TCNB}$, there are also reasonable ($0.08 \leq R \leq 0.14$) determinations of the α - and β - forms of

CuOx_2 (11,12), CuOx_2 .picryl azide (13), CuOx_2 .TCNQ (14) and three palladium compounds, PdOx_2 (15), PdOx_2 .TCNB (16) and PdOx .Chloranil (17). It is assumed that all bonds and angles between light atoms will be unaffected by electronic effects even from Cu^{2+} to Pd^{2+} . As the compounds are planar (all published structures obey two-dimensional Euclidean geometry) they can be defined by $2n - 3$ parameters, (19 if $n = 1$).

Those chosen are shown, along with the tabulated values (p_{ij}), $i = 1 \rightarrow 19$, for each structure (j) in Table IV.5. Also tabulated are the respective weighted means, \bar{p}_i , with standard deviations, $\sigma(\bar{p}_i)$, for all copper structures, and similarly for palladium ones. These are calculated from individual $\sigma(p_{ij})$ if these are given, and $t_{ij}^2 = \Delta_{ij}^2 / (\sigma^2(p_{ij}) + \sigma^2(\bar{p}_i))$ where $\Delta_{ij} = \bar{p}_i - p_{ij}$; where only an average e.s.d., $\bar{\sigma}(p_j)$, was published, $\sum t_i^2$ was taken as $\sum \Delta_{ij}^2 / (\bar{\sigma}(p_j))^2$. Although p_i are independent, an average $\sigma(p_i)$ was calculated as all the e.s.d.'s are of the same order. None of the 19 p_i are linearly related to any combination of others, but are not independent in relation to movement of atoms.

We consider only the three questions:-

- Is any particular compound significantly different from the mean of the others with the same metal ($\sum t_{ij}^2$ for any j)?
- Are the mean parameters of the copper and palladium oxinates different?
- Are the mean bond lengths of all oxinate structures compatible with those calculated by Huckel theory (23)?

(a) Of the 8 compounds only CuOx_2 .PA has large $\sum t_i^2$. This does not necessarily mean that the e.s.d.'s are too small and the effect may be due to the non-independence of parameters. (A large shift of C(9) alone could account for the two greater deviations (in angles 18 and 19)). In general, however, the results imply that the published s.d.'s are likely to be wrong by less than a factor of two in either direction, and can be usefully employed to compare structural features

within this accuracy.

(b) Palladium and copper oxinates have identical light-atom bond lengths angles, which agrees with their similar electronic structures. (Table IV.6.)

(c) The average deviation of r_{Cu} from r_{Huc} is $\sim 2\sigma(r_{Cu})$ which is good considering the approximations in the M.O. calculation. (Table IV.7.)

Paraquat structures (pg^{2+})

All the paraquat structures, including the first three (24) were planar (except $pgPdCl_4$) and have $\bar{1}$ or 2 symmetry. Two questions were:-

- (a) Are all the ions compatible with the mean and does this have mmm symmetry? (It must be at least mm2 or 2/m)
- (b) Can the bond lengths be interpreted in terms of valence theory?

The results are shown in Table IV.8, where mmm symmetry is assumed, and imply that all structures are compatible with this (as there are no extreme $\sum t_i^2$). The 4-4', 4-3 and 3-2 bonds differ significantly and the M-O bond orders are computed from them using an empirical curve (23). These can be compared with those calculated in Chapter VI and show reasonable agreement.

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Fig. IV.2. Comparison of $(p_1)_M$ and $(p_1)_W$ for $\text{CuO}_{x_2} \cdot 2\text{TCNB}$

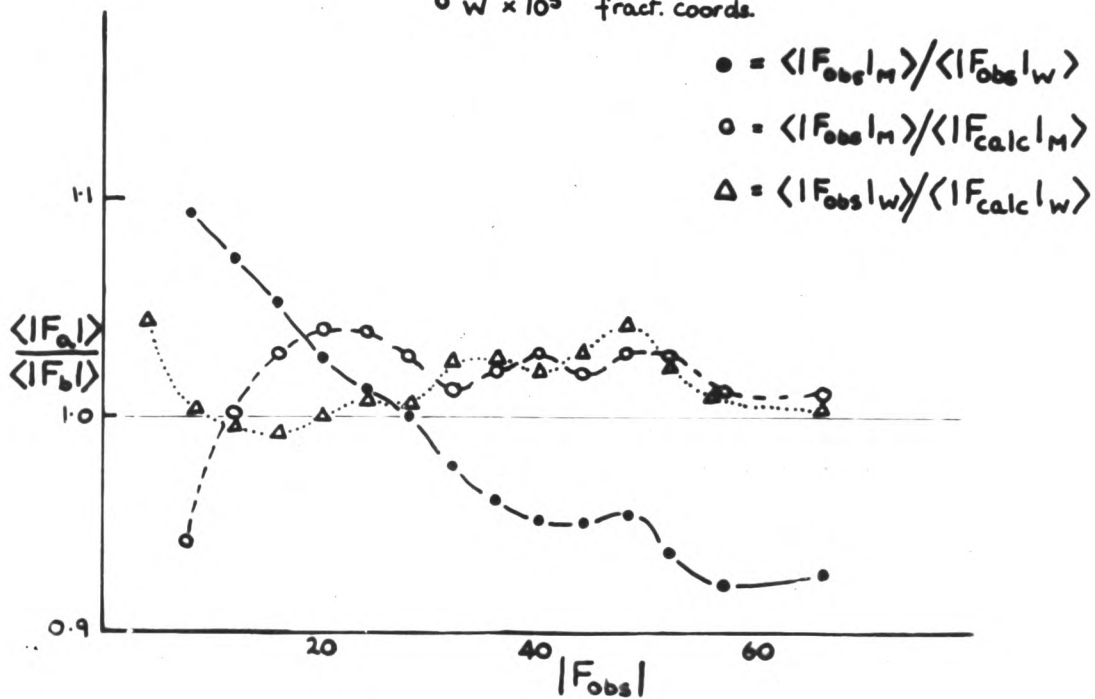
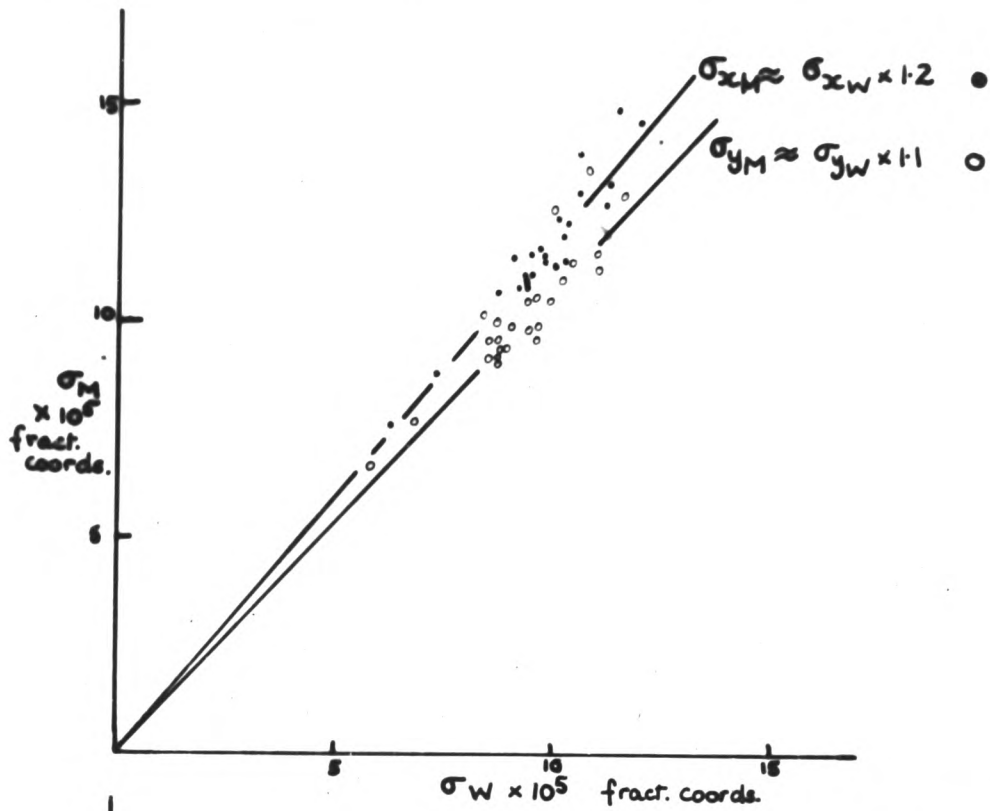


Fig. IV.1. Comparison of various F values from $\text{CuO}_{x_2} \cdot 2\text{TCNB}$

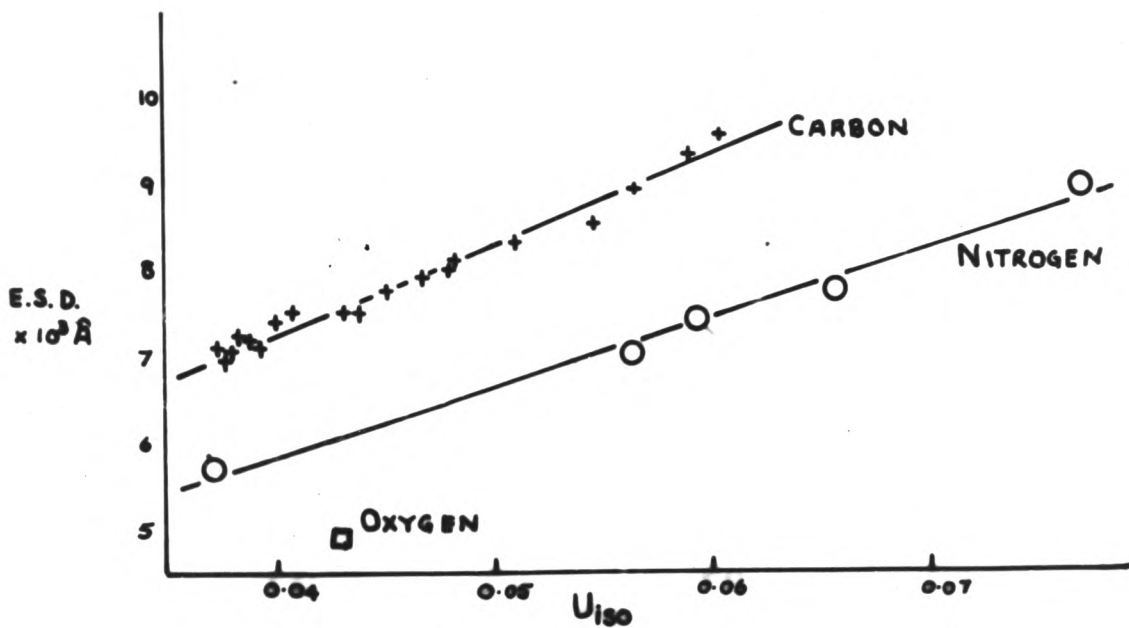


Fig. IV.4: Comparison of U_{150} with the mean e.s.d. for that atom in \AA . (for $\text{CuOx}_2 \cdot 2\text{TCNB}$)

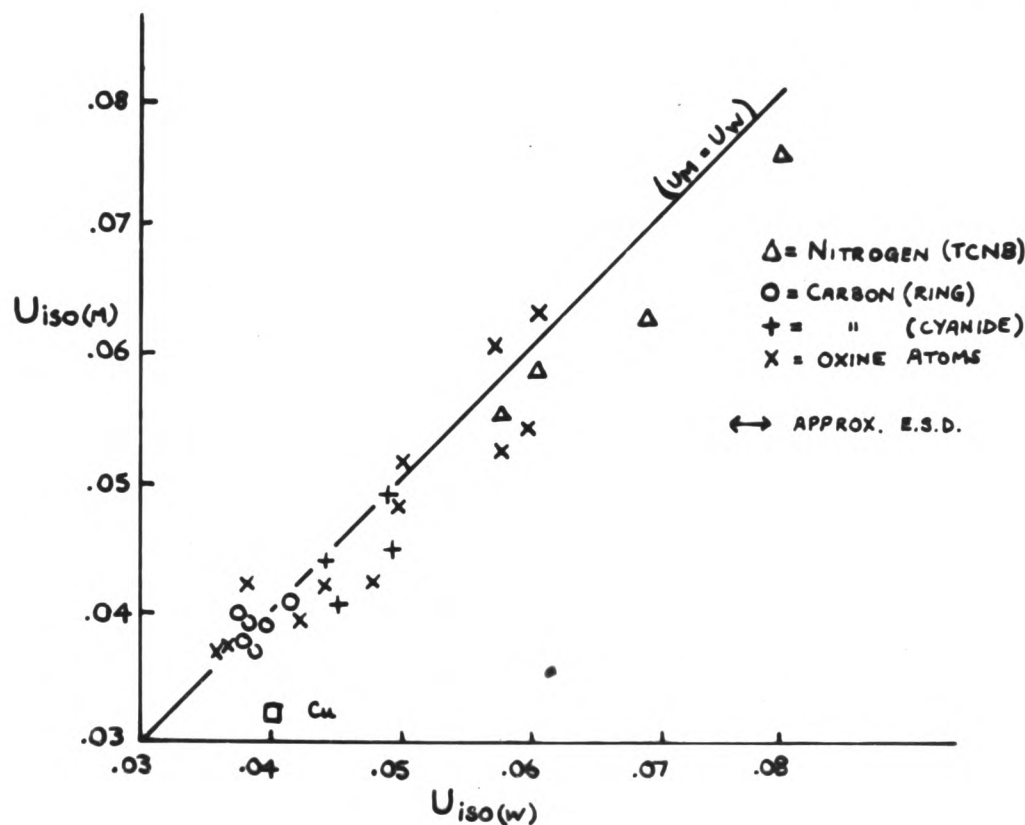


Fig. IV.3: Comparison of $U_{150}(M)$ with $U_{150}(W)$ for $\text{CuOx}_2 \cdot 2\text{TCNB}$

Table IV.1

Atomic co-ordinates and thermal parameters ($\times 10^4$) for $\text{CuOx}_2 \cdot 2\text{TCNB}$
with estimated standard deviations

For each atom the upper value is $P_{i(W)}$, the lower $P_{i(M)}$ and hydrogen atoms are omitted.

The value of t_i^2 is tabulated for each parameter.

	x/a	$t_{(x)}^2$	y/b	$t_{(y)}^2$	z/c	$t_{(z)}^2$	U(iso)	$t_{(U)}^2$
Cu	0(0) 0(0)		0(0) 0(0)		0(0) 0(0)		402(7) 323(5)	77
O(1)	761(6) 758(8)	0.1	-740(6) -746(7)	0.5	1184(4) 1178(4)	1.7	439(13) 418(12)	1.4
N(1)	-502(7) -506(9)	0.2	2140(7) 2150(8)	1.0	822(4) 814(4)	1.8	366(14) 375(13)	0.2
N(2)	7054(11) 7046(14)	0.2	2417(10) 2410(13)	0.2	5536(5) 5549(7)	2.3	783(21) 757(25)	0.6
N(3)	7221(9) 7216(11)	0.1	-2065(8) -2063(10)	0.0	3788(4) 3791(5)	0.2	577(18)	1.2
N(4)	3765(9) 3756(11)	0.4	1754(8) 1743(10)	0.7	-0015(5) -0022(6)	0.9	603(18) 583(19)	0.6
N(5)	3836(10) 3823(12)	0.7	6136(9) 6127(10)	0.4	1714(5) 1713(6)	0.0	694(20) 619(20)	7.4
C(1)	-1144(10) -1144(11)	0.0	3584(10) 3579(10)	0.1	586(5) 591(5)	0.3	477(19) 425(17)	4.0
C(2)	-1240(11) -1242(13)	0.0	5036(11) 5049(11)	0.7	1307(5) 1292(6)	3.5	574(21) 520(20)	3.3
C(3)	-707(11) -709(13)	0.0	4992(11) 4976(12)	1.0	2222(6) 2233(6)	1.6	596(22) 537(21)	3.9
C(4)	475(12) 456(15)	1.1	3204(12) 3189(13)	0.7	3464(6) 3463(7)	0.0	579(23) 604(23)	0.6
C(5)	1012(11) 1031(15)	1.0	1677(11) 1662(13)	0.7	3616(6) 3627(7)	1.4	603(22) 627(24)	0.6
C(6)	1122(11) 1128(13)	0.1	277(10) 283(11)	0.1	2880(5) 2887(6)	0.8	501(19) 520(20)	0.5
C(7)	652(9) 662(12)	0.4	444(9) 454(10)	0.5	1926(5) 1930(5)	0.3	381(17) 419(17)	2.5
C(8)	15(9) 0(11)	1.1	2064(9) 2061(9)	0.0	1758(4) 1759(5)	0.0	370(17) 379(16)	0.1

Table IV.1 (cont)

	x/a	$t^2_{(x)}$	y/b	$t^2_{(y)}$	z/c	$t^2_{(z)}$	U(iso)	$t^2_{(U)}$
C(9)	-72(10) -78(12)	0.1	3452(10) 3453(11)	0.0	2514(5) 2516(6)	0.1	484(19) 475(18)	0.1
C(10)	5419(10) 5417(11)	0.0	2591(10) 2584(10)	0.3	3378(5) 3384(5)	0.7	415(18) 401(16)	0.3
C(11)	6028(9) 6028(11)	0.0	2143(9) 2147(10)	0.1	3721(4) 3718(5)	0.1	375(17) 394(16)	0.7
C(12)	6074(9) 6077(11)	0.1	581(8) 567(9)	1.2	3093(4) 3094(5)	0.0	376(17) 379(16)	0.0
C(13)	5575(9) 5511(11)	0.1	435(9) 440(9)	0.2	2120(5) 2111(5)	1.6	384(17) 392(16)	0.1
C(14)	4903(9) 4913(11)	0.6	1893(9) 1887(9)	0.2	1775(4) 1773(5)	0.1	384(17) 373(16)	0.2
C(15)	4855(9) 4841(11)	0.9	3454(9) 3473(9)	2.2	2407(4) 2400(5)	1.0	396(17) 390(16)	0.1
C(16)	6606(10) 6605(12)	0.0	2283(9) 2277(11)	0.2	4718(5) 4737(6)	6.2	481(19) 483(18)	0.0
C(17)	6709(10) 6727(11)	1.5	-890(9) -922(10)	5.7	3472(5) 3475(5)	0.3	450(17) 413(19)	2.2
C(18)	4287(10) 4302(11)	0.9	1792(9) 1782(10)	0.0	786(5) 771(5)	4.3	443(18) 432(17)	0.2
C(19)	4268(10) 4258(12)	0.4	4940(10) 4957(10)	1.5	2036(5) 2027(6)	1.4	485(20) 451(18)	1.6

Table IV.2

Comparison of some bonds and angles from the two determinations (M,W)
of CuOx₂.2TCNB

(E.s.ds and Δ_i of bonds are in units of 10^{-3}\AA ; for angles, units of 0.1°)

<u>BONDS</u>	r_M	$10^3\sigma_M$	r_W	$10^3\sigma_W$	$10^3\Delta_i$	t_i^2
1	1.311	11	1.328	9	+17	1.4
2	1.405	13	1.412	11	+ 7	0.2
3	1.364	13	1.323	12	-41	5.4
4	1.405	13	1.423	11	+18	1.1
5	1.411	14	1.415	12	+ 4	0.0
6	1.336	15	1.321	12	-15	0.6
7	1.393	15	1.402	12	+ 9	0.2
8	1.397	13	1.392	11	- 5	0.1
9	1.423	12	1.424	10	+ 1	0.0
10	1.421	12	1.411	10	-10	0.4
11	1.319	10	1.304	8	-15	1.4
12	1.368	10	1.355	9	-13	0.9
						<u>11.7</u>

ANGLES

1	123.6	7	123.7	6	+ 1	0.0
2	115.0	8	114.8	7	- 2	0.0
3	121.3	9	120.8	8	- 5	0.2
4	118.2	8	118.5	7	+ 3	0.1
5	121.6	8	121.9	6	+ 3	0.1
6	116.6	8	116.4	6	- 2	0.0
7	125.1	8	125.8	6	+ 7	0.5
						<u>0.9</u>
						====
						12.6
						====

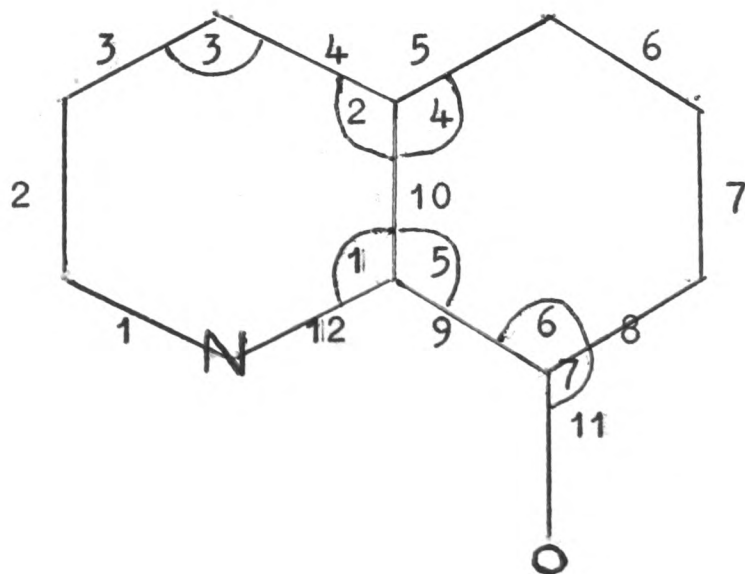


Table IV.5

Comparison of Bonds and Angles in Copper Oxinate and its complexes

BONDS	Δ_{ij} of Bonds in units of 10^{-3} \AA		Δ_{ij} of Angles in units of 0.1°		Mean(\AA)	$\sigma \sim .006$			
	α -CuOx ₂ $10^3 \times \Delta$	β -CuOx ₂ $10^3 \times \Delta$	CuOx ₂ .2TCNB(W) $10^3 \times \Delta$	CuOx ₂ .2TCNB(M) $10^3 \times \Delta$			CuOx ₂ .PA $10^3 \times \Delta$	CuOx ₂ .TCNQ $10^3 \times \Delta$	
1	+2	+12	+6	-11	1.0	-12	1.4	+6	1.322
2	-3	+12	+6	-1	0.0	+34	2.9	-11	1.406
3	+3	+3	-29	+12	0.8	+18	0.8	+16	1.352
4	-3	-18	+5	-14	1.2	+22	1.2	-3	1.418
5	-4	+2	+12	+8	0.3	-13	0.4	+10	1.403
6	+14	+6	-28	-13	0.7	+61	10.0	+23	1.349
7	+15	+29	-18	-27	3.5	0.0	0.0	-16	1.420
8	-15	+8	+2	+7	0.3	+10	0.2	-2	1.390
9	+4	-5	+1	0	0.0	-43	4.5	+1	1.423
10	-3	+5	-9	+1	0.0	+10	0.2	-6	1.420
11	+2	+7	-14	+1	0.0	+32	2.6	+3	1.318
12	-5	+12	+5	+8	0.6	+30	2.3	-8	1.360
bond									
$\sum t_i^2$		13.4	7.4	19.0	8.4		26.2		
ANGLES									
1	+2	+4	+12	+11	2.4	+3			$\sigma \sim 0.4^\circ$
2	+3	+7	-4	-2	0.0	+1			122.5
3	-2	+2	-6	-1	0.0	+1			115.2
4	-9	-11	-3	-6	0.5	+18			121.4
5	-5	+5	-8	-11	1.7	+11			118.8
6	+6	+2	-6	-4	0.2	+67			122.7
7	-7	-9	+7	+6	0.5	-89			117.0
angle									
$\sum t_i^2$		6.0	5.4	8.7	5.5	>40			
$\sum t_i^2$		19.4	12.8	27.7	13.9	>66.2			

Table IV.6

Comparison of Bonds and Angles in Palladium Oxinate and its complexes

(Δ_{ij} of bonds is in units of 10^{-2}\AA ; Δ_{ij} of angles in units 0.1°)

BONDS	PdOx ₂		PdOx ₂ .CA		PdOx ₂ .TCNB		Mean(\AA) $\sigma \sim 0.02$
	$10^2 \times \Delta$	t_i^2	$10^2 \times \Delta$	t_i^2	$10^2 \times \Delta$	t_i^2	
1	0	0.0	- 2	0.2	0	0.0	1.33
2	+ 2	0.4	+ 2	0.4	- 3	0.5	1.40
3	+ 6	4.0	- 3	1.0	- 4	1.8	1.36
4	- 3	1.0	+ 2	0.3	0	0.0	1.44
5	+ 4	1.8	- 7	5.5	+ 2	1.0	1.40
6	+ 5	2.8	- 2	0.1	- 2	0.4	1.41
7	- 1	0.1	+ 5	2.8	- 3	0.5	1.40
8	+ 4	1.8	- 2	0.4	- 2	0.4	1.40
9	+ 4	1.8	+ 1	0.1	- 5	2.7	1.44
10	+ 2	0.4	- 1	0.2	- 2	0.4	1.43
11	- 2	0.4	0	0.0	+ 2	0.4	1.31
12	- 1	0.1	+ 4	1.8	- 2	1.0	1.40
<hr/>							
bond $\sum t_i^2$		14.6		12.8		9.1	
<hr/>							
ANGLES	$10\Delta^\circ$		$10\Delta^\circ$		$10\Delta^\circ$		$\sigma \sim 1.0^\circ$
1	0	0.0	+2.0		-2.0	1.6	121.0
2	-1.0	0.3	-1.0		+2.0	1.6	116.0
3	+3.3	3.2	-2.7		-0.7	0.1	118.7
4	-0.7	0.1	+1.3		-0.7	0.1	118.7
5	-0.3	0.0	-0.3		-0.3	0.2	122.3
6	+0.3	0.0	-0.7		+0.3	0.0	116.7
7	+0.7	0.1	-1.3		+0.7	0.2	125.3
<hr/>							
angle $\sum t_i^2$		3.7		5.0		3.8	
<hr/>							
a + b $\sum t_i^2$		18.3		17.8		12.9	
<hr/>							

Table IV.7

Comparison of Bonds and Angles in Copper Oxinates,
Palladium Oxinates, and predictions from Hückel Theory

<u>BONDS</u>	CuOx ₂ (mean) $\bar{\sigma} \sim 0.007 \text{ \AA}$	t ² _{Cu/Pd}	PdOx ₂ (mean) $\bar{\sigma} \sim 0.02 \text{ \AA}$	Hückel calc ⁿ (Ref.)
1	1.322	0.1	1.33	1.32
2	1.406	0.0	1.40	1.40
3	1.352	0.1	1.36	1.38
4	1.418	0.8	1.44	1.41
5	1.403	0.0	1.40	1.41
6	1.349	6.3	1.41	1.38
7	1.420	0.7	1.40	1.41
8	1.390	0.2	1.40	1.38
9	1.423	0.5	1.44	1.41
10	1.420	0.2	1.43	1.42
11	1.318	0.1	1.31	
12	1.360	3.0	1.40	1.35
		<hr/>		
		12.0		
		<hr/>		

<u>ANGLES</u>	$\bar{\sigma} \sim 0.4^\circ$		$\bar{\sigma} \sim 1.0^\circ$
1	122.5	0.2	121.0
2	115.2	0.5	116.0
3	121.4	6.0	118.7
4	118.8	0.0	118.7
5	122.7	0.1	122.3
6	117.0	0.0	116.7
7	125.0	0.0	125.3
		<hr/>	
		6.8	
		====	
		18.8	
		<hr/>	

CHAPTER V.

STRUCTURE ANALYSIS AND DISCUSSION
of the
COBALTHEXAMMINE MIXED-VALENCE CHLOROCUPRATES(I, II)

MIXED VALENCE HALIDE COMPOUNDS

Several 1:1 mixed valence compounds containing metal halide anions have been investigated by X-rays. These have usually contained oxidation states differing by two units and among the best known are Rb_2SbCl_6 (Jensen) and $(\text{NH}_4)_2\text{SbBr}_6$ (Lawson). These both contain discrete SbX_6 octahedra and the second shows two distinct types with Sb-Br lengths of 2.56 Å and 2.80 Å, which implies Sb(V)Br_6^- and Sb(III)Br_6^{3-} ions. This difference leads to the existence of a superlattice of weak reflections, and to a very small distortion from cubic symmetry ($a = b = 10.66$ Å, $c = 2 \times 10.76 = 21.52$ Å). The strong reflections, indexed as pseudo-cubic ($\text{Pa}\bar{3}\text{m}$), have a final R of 0.078; the weak reflections, which lower the symmetry to $I4/\text{amd}$, had $R = 0.197$. This shows that relatively small changes in geometry can lead to an interpretable X-ray superlattice.

Mixed valence Lead(II) and (IV) chlorides.

A similar compound is $\text{Co}(\text{NH}_3)_6\text{PbCl}_6$ (Mori 1951), produced from Pb(II) and Pb(IV) chlorides, which shows very strong C-I, presumably outer-sphere between Pb(II)Cl_6^{4-} and Pb(IV)Cl_6^{4-} ions. As the bond distances in these are significantly different, it should be possible to see if they exist in an ordered array. The gross structure of this compound is similar to $\text{Co}(\text{NH}_3)_6\text{TlBr}_6$ with a rock salt lattice, $\text{Pa}\bar{3}$, (Wat 1950) but there are some extra powder lines (Day 1967) not belonging to this system.

Single crystals (provided by Dr. P. Day) were investigated with copper radiation, although the high absorption coefficient of ~ 600 precludes accurate work. Only zero level photographs, $h0l$, $0kl$ and hko were taken which shows the symmetry of the reciprocal lattice to be $m\bar{3}$ although there were several reflections which do not belong to $\text{Pa}\bar{3}$. These had indices

$h = 2n - \frac{1}{2}, 2n, 2n + \frac{1}{2}$ (when related to the original lattice constants) and suggests strongly a superlattice. However further analysis was very difficult as there is evidence of twinning. Reflections of large h are split, not into the usual $a_1 - a_2$ doublet, but into 3 equal spaced spots, which can only mean the existence of two axes different by $\sim 0.02 \text{ \AA}$ (in 11.5 \AA). It is unlikely that the twinned crystals are tetragonal as the reciprocal lattice has no tetrads (although these could be obscured by the twinning), and might be of 3 orthorhombic pseudo-cubic crystals related by T symmetry (23). In spite of this accident it is probably safe to assume crystallographically different PbCl_6 groups with outer-sphere charge transfer.

MIXED VALENCE CHLOROCUPRATES(I) and (II)

The structure of $\text{Cr}(\text{NH}_3)_6\text{CuCl}_5$ (Wat 1961), and the cobalt analogue (Part II), contain a rock salt lattice containing regular CuCl_5^{3-} trigonal bipyramids. The parameters, in $Fd\bar{3}m$, which have not been fully refined, are given in Table V.1. The final R was 0.19 and no structure factors have been published but comparison with the other compounds suggests this is the correct structure, and it is shown in Fig.V.1a.

TABLE V.1.

Parameters of Chromium Hexammine(III) Penta-chlorocuprate(II)

No./Cell	Atom	x/a	y/b	z/c	B	
32	Cu	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	2.0	} $a = 22.18 \text{ \AA}$
32	Cr	0	0	0	2.0	
64	Cl(axial)	0.175	0.175	0.175	2.0	assumed $\text{Cu-Cl}_{\text{ax}} = 2.32 \text{ \AA}$
96	Cl(equat.)	$\frac{1}{2}$	0.075	-0.075	2.0	} $\text{Cu-Cl}_{\text{eq}} = 2.35 \text{ \AA}$
192	H	0.07	0.05	-0.03	2.0	

An attempt to solve the structure of $\text{Co}(\text{NH}_3)_6\text{Cu}_4\text{Cl}_{17}$ from rotation photographs had been made (Part II). The space-group appeared to be $Fd\bar{3}m$ (although this has been disproved in Chapter III) and only the copper and cobalt atoms could be located. These were in identical positions to those given above, with the addition of 8 copper atoms at $\frac{1}{8}\frac{1}{8}\frac{1}{8}$ etc, but the agreement was poor, $R = 0.43$.

More orthodox data collection for $[\text{Co}(\text{NH}_3)_6]_4\text{Cu}_5\text{Cl}_{17}$ and "25%Cu(I)" has been described in Chapter III. The pure copper(I) compound is described first.

Structure Analysis of $[\text{Co}(\text{NH}_3)_6]_4\text{Cu}_5\text{Cl}_{17}$

The positioning of the heavy atoms as described above is unavoidable both from space-group symmetry and the similarity of

the intensity data to the copper(II) compound. (The cobalt atoms are actually in two equivalent 16-fold sets at 000 and $\frac{1}{2}\frac{1}{2}\frac{1}{2}$). Neither Patterson nor Fourier syntheses are useful at this point as there is too much symmetry, and cobalt and copper have nearly equal X-ray scattering powers. The positioning of the chlorine atoms must be on a trial and error basis, compatible with both the space-group and chemical criteria. These latter must preclude bunching of negative charge, and include a close similarity to the copper(II) structure. The 126 chlorine atoms in the cell can be split into 3 groups:- 32 in (d), xxx; 96 in general positions (f), xyz; and 8 in (b), $\frac{2}{8}\frac{2}{8}\frac{2}{8}$. The first group can be positioned tetrahedrally about the 8-fold copper atom (A) and bonded also to the 32-fold copper atoms (B). (See Fig. V 2). This is the only feasible coordination for copper atom (A) and it is most unlikely that this is an isolated ion. A linear Cu-Cl-Cu bridge has not otherwise been found, and the chlorine atom was placed roughly at the centre ($x/a = 0.190$). This is virtually identical with the position of one axial chlorine in the copper(II) compound. At this stage there should be enough scattering material placed to give a useful Fourier map. (The computation of these maps was impossible directly and was carried out in the non-centrosymmetric sub-group F222. Structure factors are referred to the non-centric origin in Fd3 (Table VII 7) but parameters (Table V 2) are given relative to the centric origin for ease of comparison with related structures).

The map showed acceptable peaks for all atoms introduced and two large peaks at 0.28, 0.07, -0.08, and 0.28, -0.08, 0.07, related in position by the plane 110. As the second was about half the height of the first, it is probably a "ghost" peak, produced by false symmetry. The first was used for position (f) and gave an acceptable tetrahedral coordination for copper (B). Three rounds of least-squares preceded another Fourier map which showed no ghost peak and two independent nitrogen positions giving

excellent octahedral cobalt hexamine cations. These octahedra are in virtually the same orientation as those in $\text{Cr}(\text{NH}_3)_6\text{CuCl}_5$ (e.g. Tables V.2 and V.1. The chlorine atom at $\frac{z}{8} \frac{z}{8} \frac{z}{8}$ had been included as this is the only position allowed by symmetry, but there was little evidence for it from the Fourier syntheses and its U_{iso} had risen to 0.11, very high for an ionic structure. It is in a large hole, with a minimum radius of $\sim 4 \text{ \AA}$, and might be vibrating excessively. A model was investigated with it forming the fifth coordination position of the copper(B) atoms and with statistical occupancy of 4 positions. This showed no better agreement although it is chemically acceptable as it is very close to the $\text{Cu}(\text{II})$ structure and involves penta-coordination.

Although the refinement had converged at 0.146 and the structure is chemically very acceptable, the agreement of structure factors is poor. A detailed analysis of this was carried out, firstly by comparing F_{obs} and F_{calc} and secondly F_{hkl} and F_{kh1} . For the first group, the reflections were split into those to which the heavy atoms had a large (in phase) contribution and those with none at all. The first set showed excellent agreement, $R = 0.066$, but the second set had $R = 0.22$. Although these latter are significantly weaker, this is still a serious disagreement, particularly when these weaker reflections can be measured experimentally to $\pm 5\%$ in F (Chapter III). This implies that the heavy atoms are definitely located, but the lighter atoms, particularly the chlorines, are not completely correct although they have a very close relation to the actual structure. The overall agreement is however very similar to that in the bromocantionate discussed earlier.

The discrepancy is probably explained by a comparison of F_{hkl} and F_{kh1} both for F_{obs} and F_{calc} . In effect this measures the closeness of the space group to $Fd\bar{3}m$ rather than $Fd\bar{3}$, as for the former $F_{\text{hkl}} \equiv F_{\text{h1k}}$. For a random distribution of atoms in $Fd\bar{3}$, the merging $R_{\text{hkl} \leftrightarrow \text{kh1}}$ should be ~ 0.8 , whereas it is very much

lower (Chapter III). This is represented in Fig. V.3, where the maximum value of a F_{hkl}/F_{khl} pair is plotted against the minimum. The plot for F_{obs} suggests that F_{hkl} and F_{khl} always have the same sign, whereas there is no such constraint on F_{calc} . This implies very close mirror symmetry down $\{110\}$, which is not represented in the calculated model. There are two ways by which this symmetry can be approached. Either there is a disordered structure with approximately equal amounts of two mirror-related arrangements, or the crystal consists of twinned sections related by the $\{110\}$ plane.

For the first case we can write:

$$F_{hkl} = (F(hkl,xyz)) + (1-\lambda)(F(hkl,yxz))$$

which can be rewritten as:

$$F_{hkl} = (F(hkl,xyz)) + (1-\lambda)(F(khl,xyz))$$

A model was set up for calculation with a chlorine atom at xyz with occupancy λ and an identical atom at yxz with occupancy $1-\lambda$. A value of 0.7 was taken for λ , as this was implied by the first Fourier map, and λ was refined. It converged to a value of approximately unity, reverting to the same situation as before.

The second case can be represented by:

$$F_{hkl}^2 = (F^2(hkl,xyz)) + (1-\mu)(F^2(hkl,yxz))$$

or

$$(F^2(hkl,xyz)) + (1-\mu)(F^2(hlk,xyz))$$

There is no way of evaluating μ except by writing a programme to minimise not $\sum w(F_{obs} - F_{calc})^2$ but $\sum w(F_{obs}^2 - F_{calc}^2)^2$ and introducing it as a parameter. There is no clear justification for this, especially as the present least-squares programme would need considerable modification. Data collection on several distinct crystals might also produce a crystal with better differentiation

between $|F_{hkl}|$ and $|F_{khl}|$. There is no morphological evidence as all the crystals are beautifully formed octahedra with no $\{hkl\}$ faces.

In spite of this the structure would appear to be basically correct, as it is chemically reasonable and the cobalt hexamine octahedra, which are the only definitely known chemical species, are well defined on the Fourier maps. It consists (Fig. V.3, Table V.2) of discrete $\text{Co}(\text{NH}_3)_6^{3+}$ ions with $\text{Cu}_5\text{Cl}_{16}^{11-}$ ions and, presumably, single chloride ions, vibrating excessively. This complex chlorocuprate has a central copper atom exactly tetrahedrally coordinated by 4 chlorines at 2.26 Å. The four other copper atoms have a coordination number between 3 and 4, with one long bond of 2.49 Å and 3 shorter of 2.32 Å, with the angle between these 99° instead of tetrahedral. An inviting speculation is that CuCl_4^{3-} and CuCl_3^{2-} ions occur in solution and condense together in the ratio 1:4 to give the final complex. This has $\bar{1}$ (or 23) symmetry deviating only from $\bar{4}3m$ by a rotation of the terminal CuCl_3 groups by $\sim 18^\circ$.

Structures of the Intermediate Compounds

We can see clearly how the copper(II) and copper(I) structures are related (Figs. V.1 and V.2) and how intermediate compounds could be formed over the whole range. The 32 Co and 32 Cu (B) remain in the same positions as do the amine groups. We start with the copper(II) structure which has 4 apical Cu-Cl bonds "pointing" at $\frac{1}{8} \frac{1}{8} \frac{1}{8}$. These would appear to remain in all the structures and at some stage the "hole" at $\frac{1}{8} \frac{1}{8} \frac{1}{8}$ is filled by a copper(I) ion (A), the exact point at which this starts depending on the stoichiometries of the intermediates. In the copper(I) structure the other apical chlorine has disappeared and the equatorial chlorines have moved about 0.6 Å. The formulation of the intermediates must involve fractional occupation numbers, and those between "0% Cu(I)" and "50% Cu(I)" must lie in $Fd\bar{3}c$. This means that the symmetry of the

copper atom is 32 and the overall occupancy, as seen by X-rays, of the apical positions must be equal. (The local situation may be of lower symmetry, randomly distributed). The total occupancy of these sites must decline regularly from 2 to 1 over the system as the average oxidation number of copper atoms (B) change similarly. The difference between the two stoichiometries depends only on the point of introduction of the fifth copper atom (A) and its attendant chloride ion.

The overall average electron distribution can be broken down in more than one way to local arrangements with less than space group symmetry. One such is where the CuCl_5^{3-} bipyramid changes to Cu(I)Cl_4^{3-} groups, which must be randomly oriented in two directions. These will be trigonal pyramids but it will be hard to determine if the 3 basal chlorines are coplanar with the copper atom. In the pure copper(I) phase, however, the orientation of these pyramids is ordered so we must assume a change at some stage to disorder. Assuming (as an example only) that the 50% copper(I) compound is totally disordered, then the average structure around the copper could be represented by two half occupied axial chlorine positions with either 3 fully occupied equatorial sites or 6 half-filled sites related by the diad axes.

An interesting possibility is that for compounds with small amounts of copper(I), the amount of "interstitial" copper (A) depends on the chance arrangement of the copper (B) units. For this (A) atom to be incorporated, a hole (X) (Fig. V.2) must be made for its attendant chlorine. This is impossible in the copper(II) structure and is only likely if copper(I) pyramids (B) replace copper(II) bipyramids. If enough of these are aligned correctly a large enough hole will be left at their bases. This number will be between 1 and 4 and this can be related to Fig. II.1. If it is 1, we follow a straight line to a limiting composition of $\text{Cu(II)}_3\text{Cu(I)}_2\text{Cl}_{20}^{12-}$ at 40% copper(I), whose structure is shown below, Fig. V.5.

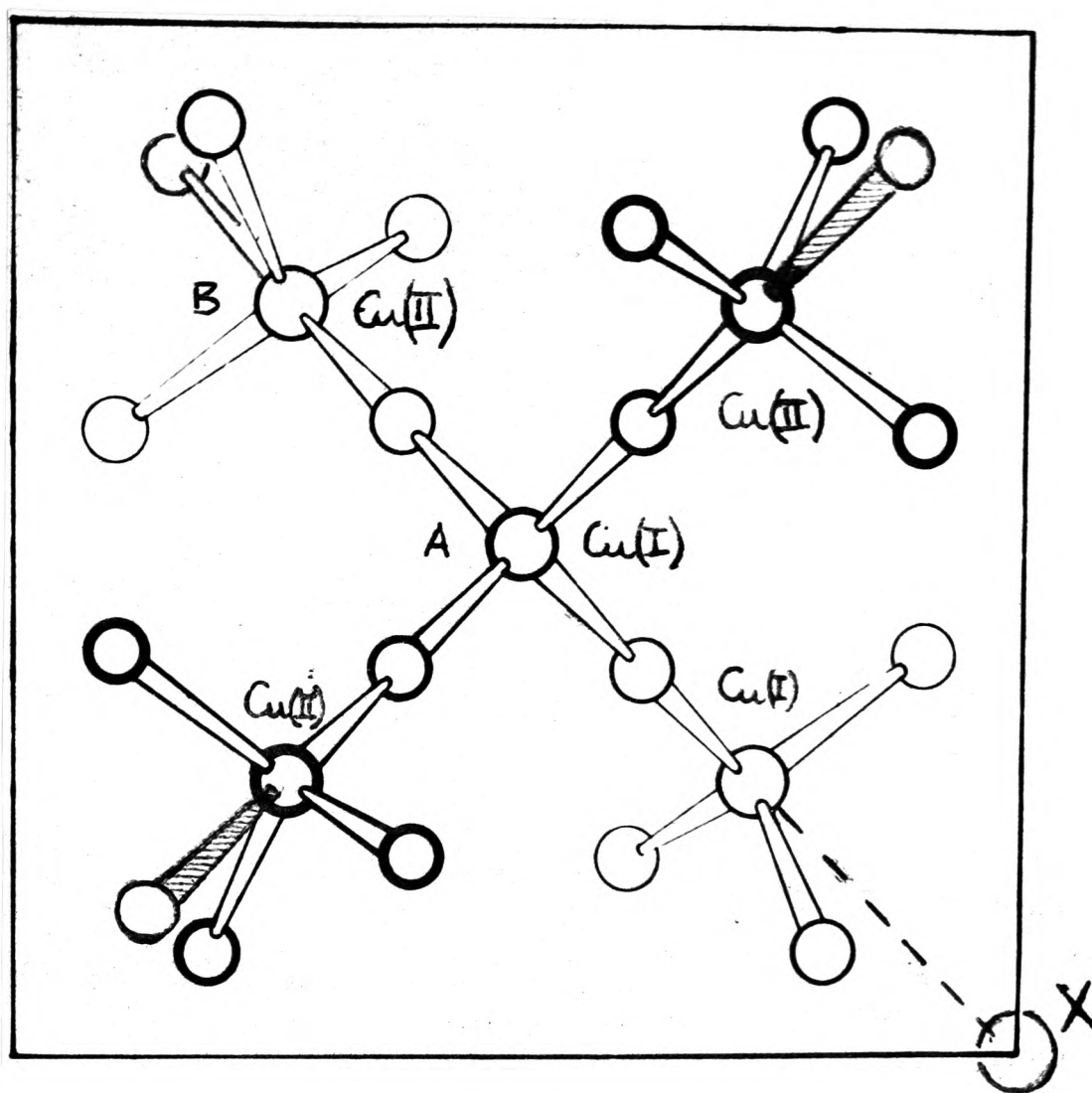


Fig. V.5.

This would be acceptable but for the chemical analyses, and the fact that the hole for the chloride ion is small and negatively charged. If however we require the absence of two apical chlorines to make a large enough hole for the introduction of the new chloride ion, and make the assumption that the arrangement of Copper(I)(B) ions is random (true if $\text{Cu(I)}/\text{Cu(II)}$ is small), we can argue as follows: For a possible Cl^- site, the chance of being surrounded by two Cu(I) ions is $4! x^2/2!2!$ where x is the fraction of copper(I). The chance that both will have their bases towards the site is $\frac{1}{2}$ of this, or $\frac{3}{2} x^2$. The corresponding formula for sites with 3 and 4 apical chlorines absent are $\frac{1}{2} x^3$ and $x^4/16$. If $x \sim 0.2$ these sites cannot be formed independently and the model becomes very involved. However we can now postulate that the sites will begin to be ordered (although no evidence for this exists). The exact arrangement of the intermediates is very comparable to the problem of hydrogen bonds in ice or ammonia.

It is much more difficult to guess at the structures of the compounds with little Cu(II), partly because the stoichiometry in that region is even less certain. The d^9 configuration of Cu^{2+} makes regular tetrahedral coordination of (B) unlikely and perhaps the best solution is a random arrangement of $[\text{Cu(I)}_4\text{Cu(II)Cl}_{17}]^{11-}$ ions with Cu(II) ions in (B) sites. This would require the shift of a chlorine atom (X) from a hole to the fifth coordination position of a (B) site.

Crystallographic Data for Intermediate Compounds

The main evidence for the absence of a superlattice in these compounds was from rotation photography (Part II) and it is possible that very weak reflections could have been obscured by the high fluorescent background. (It is significant that the reflections due to the superlattice in the chloroplumbates(II,IV) were only visible on Weissenberg films.) However the precession photographs of the 25% Cu(I) compound showed no spots not belonging to $Fd\bar{3}c$, and this is not unreasonable for compounds which have no discrete stoichiometries.

This implies that intermediate compounds must be treated as statistically disordered and that the atomic positions will be occupied by fractional amounts of atoms. The data collected on the 25% Cu(I) compound appeared good since there was a relatively large percentage of measurable reflections to which only the light atoms contribute. Since the structures of the end compounds are very similar it is unlikely that the data from an intermediate can give such information but it may answer three questions:

- (a) What is the occupancy of the (A) site at $\frac{1}{8}\frac{1}{8}\frac{1}{8}$?
- (b) What is the occupancy of the axial chlorine atoms?
- (c) Can the equatorial chlorine atoms be distinguished as filling either a single set of positions or two unrelated sites (implying distinct Cu(I) and Cu(II) atoms statistically disordered)?

The data was analysed both by Fourier and least squares methods, and three approaches were followed.

(1) A least squares calculation was done using Watanabe's parameters for the pure copper(II) compound. Although it is impossible that these are correct, the structures must be very similar and refinement converged with $R = 0.110$. These parameters are given in Table V.3 and imply that Watanabe's structure is almost certainly correct. A difference map was calculated, on which two main features appeared. These are shown in Fig. V.5., and are peaks of $3e/\text{\AA}^3$ about 0.3 \AA from the copper (B) atom and large spherical peaks of $6e/\text{\AA}^3$ just off the equatorial diads. The first implies very marked vibration of the copper atom or statistical disorder of these displaced from $4d$. The second peaks are in almost exactly the same relationship to the copper (B) atom as the 96 'equatorial' chlorine atoms in $[\text{Cu}_5\text{Cl}_{16}]^{11-}$, both $\sim 0.6 \text{ \AA}$ from the diad. As the latter atoms have final peak heights of $\sim 30e/\text{\AA}^3$, it is likely that these peaks are a pair of chlorine sites with total statistical occupation $\sim 25\%$ connected with the copper(I) atoms. There was no significant difference density either at $\frac{1}{8}\frac{1}{8}\frac{1}{8}$ or in the axial or equatorial chlorine atom positions, but this might be due to errors in temperature factors, or overall scale. The final nitrogen positions are very close to those in the copper(I) compound (and Watanabe's rough ones).

(2) In an attempt to formulate this more precisely, the following model was chosen:

- (a) Co, Cu(A) and N in standard positions (cf. copper(I) compound)
- (b) Axial and equatorial chlorines from copper(II) compound with 75% occupancy.
- (c) Two CuCl_4 tetrahedra from copper(I) compound with total occupation 25%, related by equatorial diads.

- (d) The copper atom at (A), $\frac{1}{8} \frac{1}{8} \frac{1}{8}$, with
 $0 < \text{occupancy} < 25\%$.

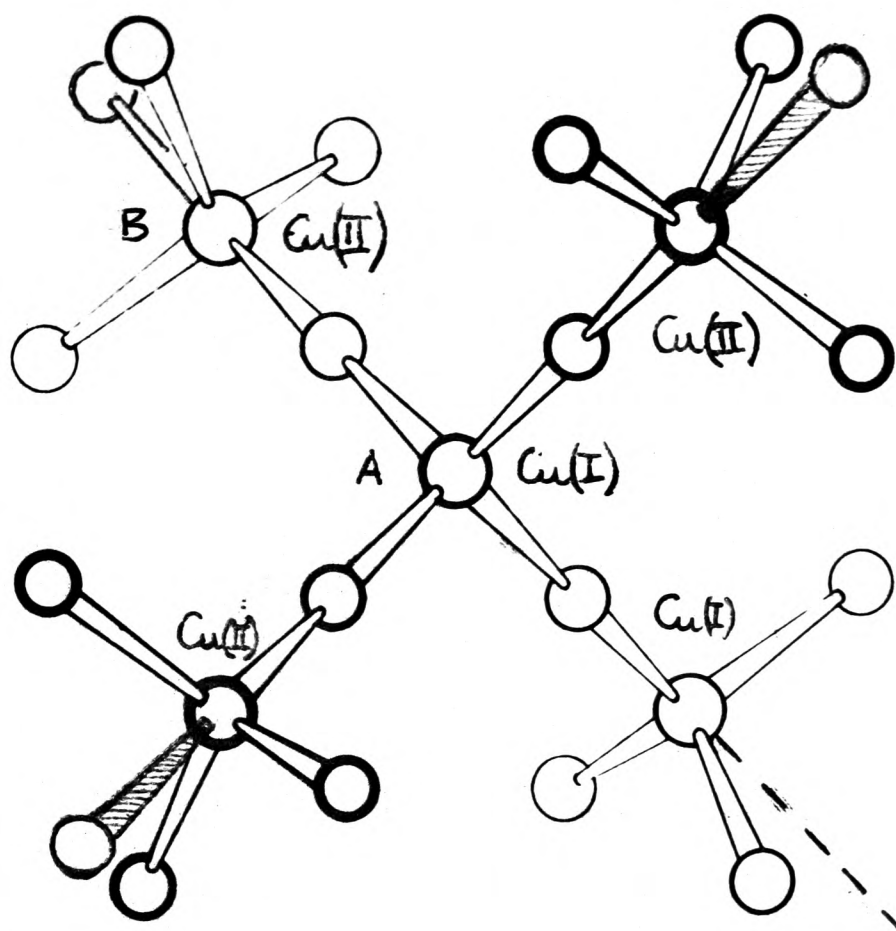
These positions were used to calculate structure factors and carry out least squares refinement. This was not successful and diverged catastrophically. It is clear that the problem has a number of near-singularities in the normal matrix and although the model is plausible, attempts to fix one or more parameters are unlikely to be useful.

(3) This model had complete chlorine atoms situated between the equatorial sites and the "copper(I) sites", 0.6 Å distant. The parameters of these atoms were refined without difficulty to $R = 0.062$. This approximation of two close disordered sites to a single atom is clearly good (although it is not impossible that, if the compounds are true C-T intermediates, the chlorine atoms occupy a position midway between the sites). The occupancy of site (A) dropped swiftly to zero.

These results suggest that the formula of the intermediate can be written $\text{Co}(\text{NH}_3)_6 \cdot \text{Cu}(\text{A})_0 \text{Cu}(\text{B})\text{Cl}(\text{eq})_3 \text{Cl}(\text{ax})_{1.75}$ where Cu(B) has an average oxidation state of 1.75. The axial chlorine atoms occupy particular sites but there are two sets of half-occupied equatorial chlorine positions, corresponding to random orientation of the CuCl_4 tetrahedra. An adequate description is the combination of 75% $\text{Co}(\text{NH}_3)_6 \text{CuCl}_5$ with 12½% of each of $\text{Co}(\text{NH}_3)_6 (\text{CuCl}_4)^\uparrow$ and $\text{Co}(\text{NH}_3)_6 (\text{CuCl}_4)^\downarrow$ where these are the two orientations.

Evidence from Physical Data

The structures of the intermediate compounds must satisfy the requirements of the absorption spectra and conductivity measurements. The spectra can be explained by intermediates of the type:-



The C-T can then take place by electron-transfer through the chlorine bridges, in a very similar fashion to the solution intermediate. Assuming that the mechanism is not affected by the Cu(I)/Cu(II) ratio, then this presumably determines the intensity of the absorption. At small ratios this should be roughly proportional to the concentration of (A) sites. The lead compound suggests that an outer sphere mechanism may be possible and electron transfer may occur between copper(I) and copper(II) in (B) sites. This might be favoured energetically as it avoids the use of the tetrahedral site (A); absorption would be proportional to x . Data from Fig. II.1. is inconclusive, but for small concentrations of copper(II) either argument leads to a linear dependence of absorption on x .

The conductivity data is more complicated as the end compounds use an ionic mechanism but the intermediates a semi-conductive one. It is possible to postulate outer sphere conductive methods which are presumably dependent solely on the random distribution of copper(I) and copper(II) units. There is a possibility however of chains being set up in the structure for which electron transfer can occur over reasonable distances. These

would be the consequence of (A) sites occurring adjacently by chance and could give inner sphere conduction. Outer sphere mechanisms might operate over the short gaps between chains. In this system there is also a chance of atom transfer (Cl) over short distances from a 5 coordinate Cu(II) to an incompletely coordinated Cu(II). The final answer must rest on more physical measurements on the compounds.

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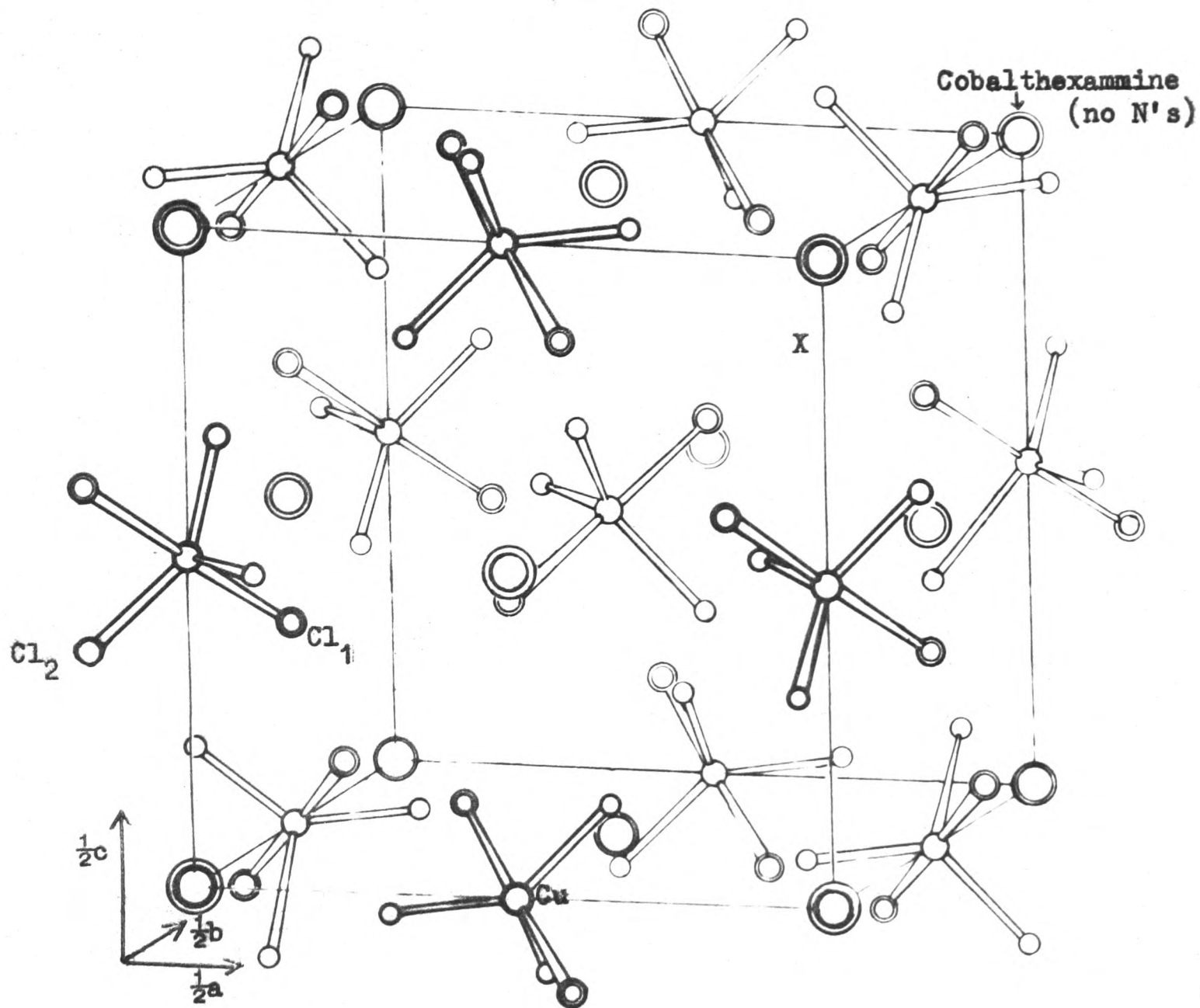


Fig. V.1. Clinographic view of Watanabe's structure

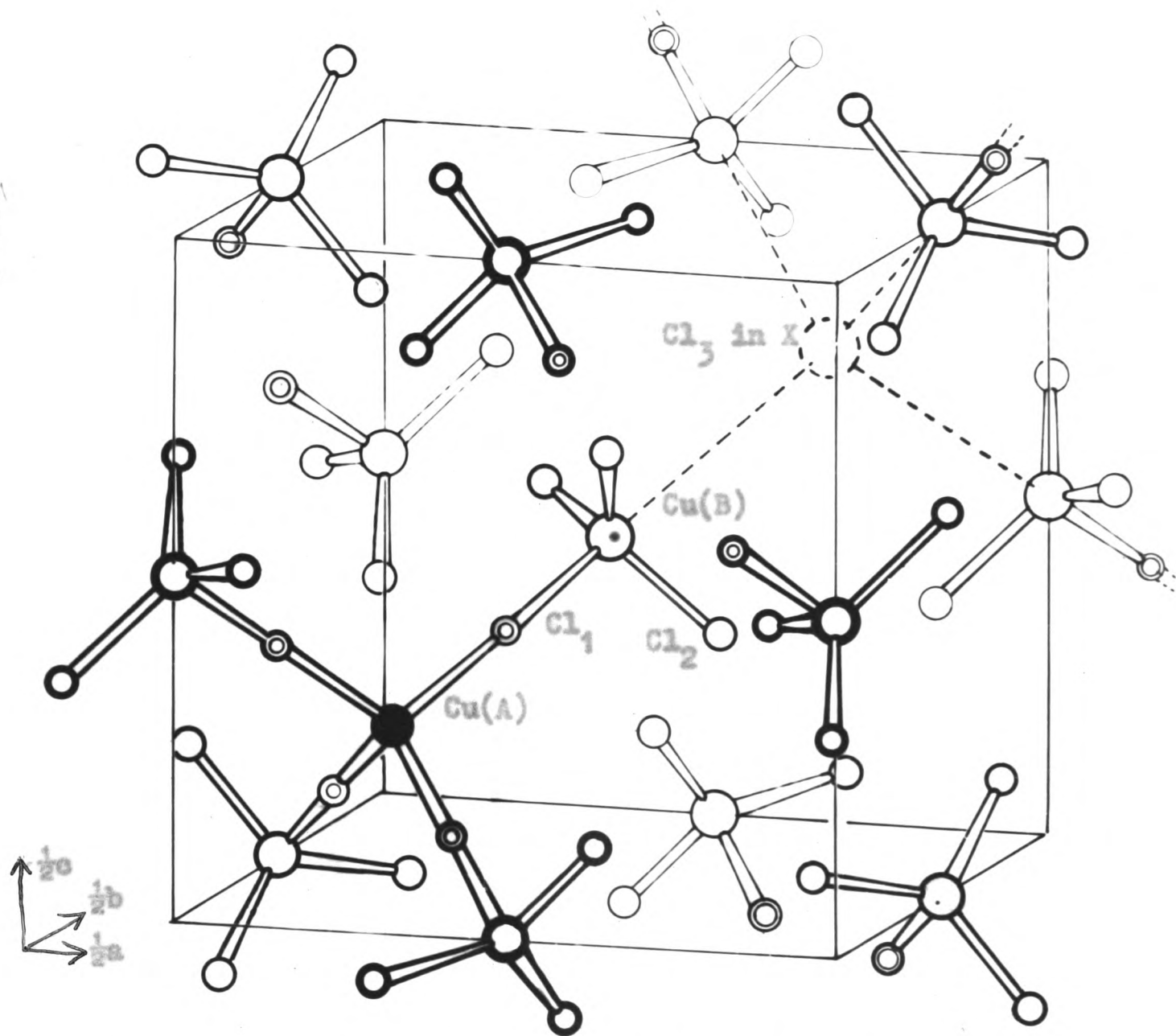
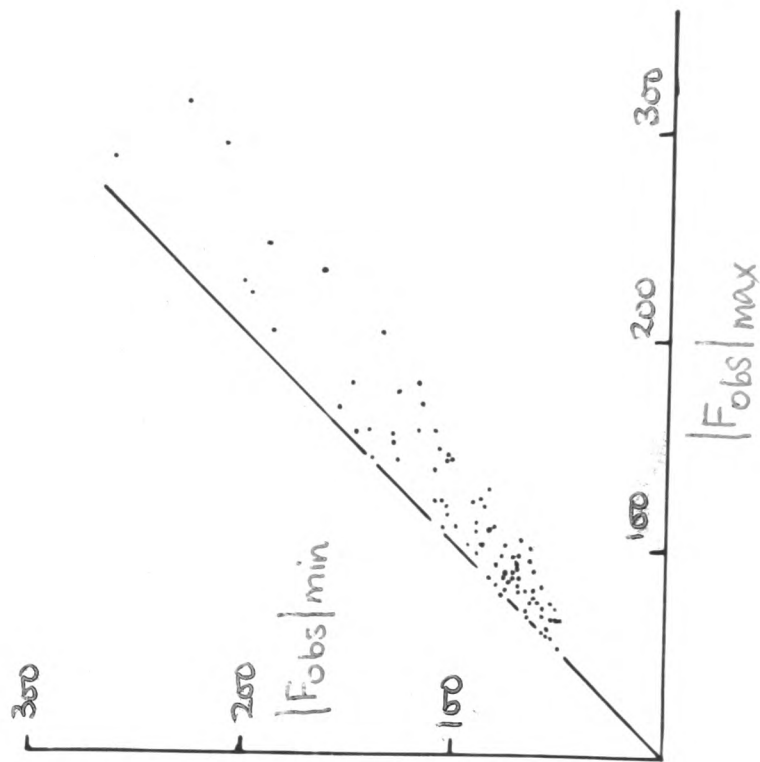
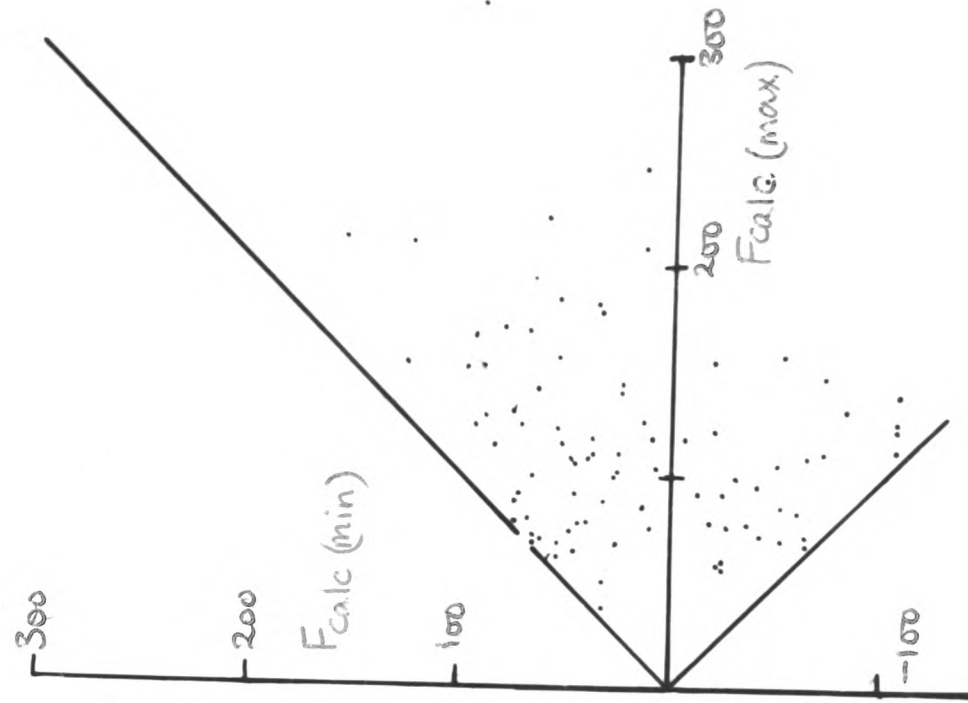


Fig. V.2. Clinographic view of Cobalthexamine Chlorocuprate(I)

(on the same scale and from the same viewpoint as the preceding diagram. The cobalthexamine ions have been omitted but are in identical positions.)



Figs V.3.a and b. Plots of the maximum of a pair F_{hkl}, F_{khl} against the minimum, for the chlorocuprate(I). Values in electrons.

Fig. V.4. Difference Electron Density in $^{63}\text{Cu(I)}\text{Cl}_2$ Chlorocuprate

Contours are at intervals of $2e/\text{\AA}^3$, and the positions of the atoms, except nitrogen and cobalt, which contribute to F_{calc} are shown, (model (1)). The equatorial atoms in model (3) were placed at P and P'. The orientation and origin of the diagram are exactly as in Fig. V.1. but the scale has been enlarged x2.

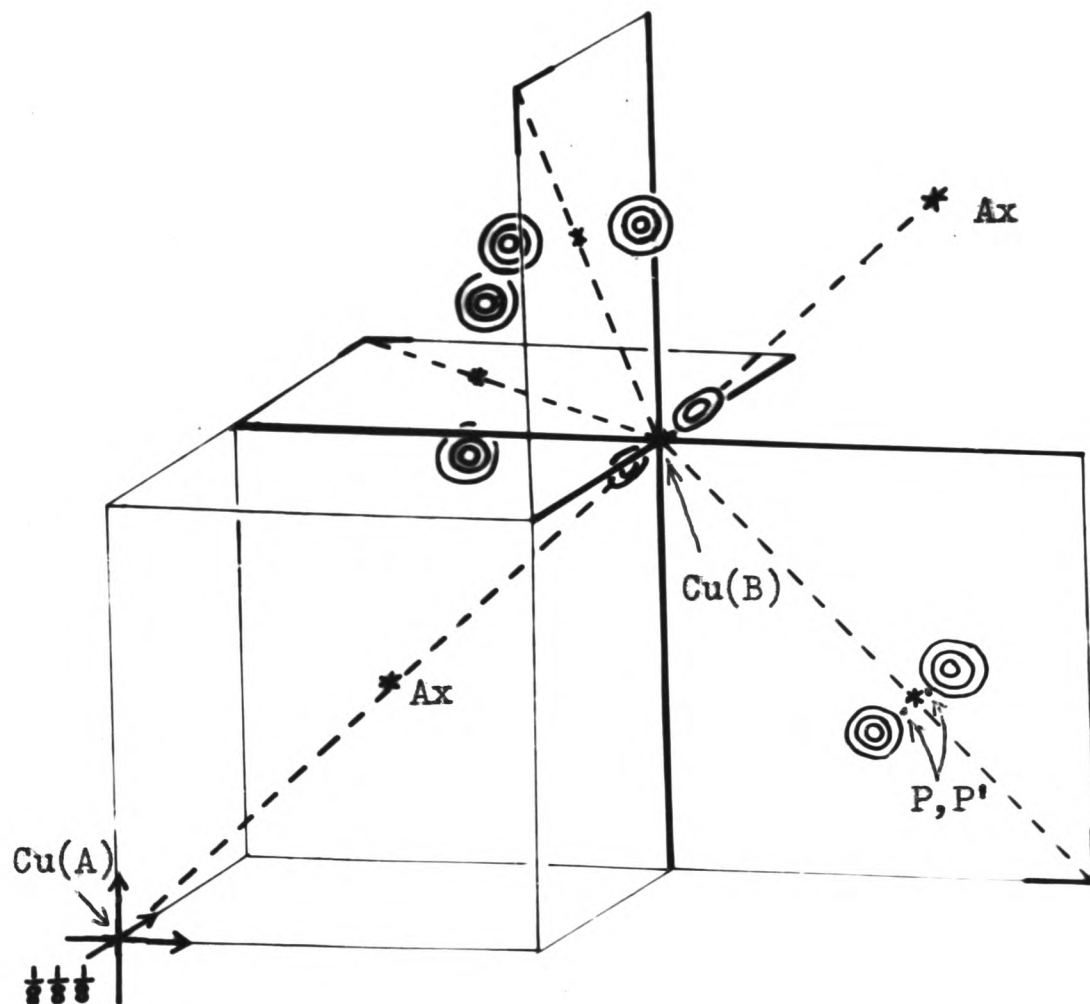


Table V.2.

Final Parameters ($\times 10^4$) for $[\text{Co}(\text{NH}_3)_6]_4\text{Cu}_5\text{Cl}_{17}$

(Referred to centrosymmetric origin)

	x/a	σ_x	y/b	σ_y	z/c	σ_z	U_{iso}	σ_U
Co(1)	0	-	0	-	0	-	105	27
Co(2)	5000	-	5000	-	5000	-	137	29
Cu(B)	2509	3	2509	3	2509	3	343	26
Cu(A)	1250	-	1250	-	1250	-	169	33
Cl(bridge)	1849	4	1849	4	1849	4	122	29
Cl(terminal)	2777	5	713	5	-748	5	343	25
Cl(interstitial)	3750	-	3750	-	3750	-	700 (constant)	
N(1)	621	13	593	13	-306	13	151	67
N(2)	5350	16	811	16	-269	16	331	86

The temperature factor of the interstitial chlorine atom was kept constant and its occupation number was refined. This converged to about 0.45 atom per site.

Other Details

Cell parameters; $a = 21.80 \text{ \AA}$ Spacegroup; $Fd\bar{3}$, No.203., origin at $\bar{3}$.

$Z = 8$ formula units Number of asymmetric units = 96.

Dimensions of crystal; Regular octahedron, side = 0.40 mm $\mu = 45 \text{ cm}^{-1} (\text{Mo})$

Number of reflections; 457 Final R = 0.146

Bonds and angles not mentioned in text;

Co(1)-N(1) = 1.98 \AA Co(2)-N(2) = 2.00 \AA

N(1)-Co(1)-N(1) = 89.6, 90.4 $^\circ$ N(2)-Co(2)-N(2) = 89.1, 90.6 \AA

Cu(A)-Cu(B) = 4.72 \AA Cu(B)-Cu(B) = 7.70 \AA

N.B. The structure factors calculated and tabulated in Chapter VII were based on the non-centrosymmetric origin at $\frac{1}{8} \frac{1}{8} \frac{1}{8}$. Their magnitudes are unaltered but their phases are different.

Table V.3.

Parameters for $\text{Co}(\text{NH}_3)_6\text{Cu(I)}_x\text{Cu(II)}_{(1-x)}\text{Cl}_n$ ($x=25\%$)

No unique solution can be given and both sets of parameters, obtained from methods (1) and (3) in the text, are given.

Final Parameters from Method (1) ($\times 10^4$)

	x/a	y/b	z/c	U_{iso}
Co	0	0	0	173
Cu(B)	2500	2500	2500	487
Cl(axial)	1880	1880	1880	540
Cl(equatorial)	2500	760	-760	590
N	704	510	-299	216

From Method (3)

Co	0	0	0	170
Cu(B)	2550	2550	2550	346
Cl(axial)	1880	1880	1880	523
Cl(equatorial)	2621	768	-756	346
N	707	511	-298	292
Cu(A)	1250	1250	1250	300(assumed)

The occupancy of the Cu(A) sites was refined and converged to a value of 0.12 per site. This is probably not significant.

Other Details

Cell parameter; $a = 21.81 \text{ \AA}$ Spacegroup; $Fd\bar{3}c$, No.228. $Z = 32$ formula units.

Origin at $\bar{3}$ for parameters, at 23 for S.F.s.

Number of reflections; 171 Final R = 0.122 (method(1)), 0.087 (method (3))

Bonds and angles from method (1):

$$\text{Cu-Cl(ax)} = 2.34 \text{ \AA} \quad \text{Cu-Cl(eq)} = 2.35 \text{ \AA} \quad \text{Co-N} = 2.01 \text{ \AA}$$

$$\text{N-Co-N} = 89.8, 90.2^\circ$$

CHAPTER VI.

STRUCTURES OF THE PARAQUAT METAL CHLORIDES
AND
DISCUSSION OF THEIR CHARGE-TRANSFER

THE STRUCTURE OF THE PARAQMAT SALTS

Although the structures of the paraquat compounds are unequivocal, mere presentation of parameters and diagrams is not very useful as this gives no very clear idea of the arrangement of the ions in space. Accordingly I have tried to select any features which are interpretable in terms of the charge-transfer and in particular to try to systematize features common to one or more structures. The presentation is thus not necessarily structure by structure, but rather by parallel discussion of common points. In the diagrams I have emphasized the features I feel important, particularly the packing of the ions. Although the compounds have been mainly studied from the viewpoint of charge-transfer, the complex anions are also interesting and are discussed separately as coordination compounds.

Gross Aspects of the Structures

All the structures have isolated paraquat ions with the mm symmetry demonstrated, (except for $pqPdCl_4$, with 222 only). The cobalt and palladium compounds have the expected tetrahedral $CoCl_4^{2-}$ and square planar $PdCl_4^{2-}$ anions but both copper anions are chain structures (along the needle axes). These can be formulated as $(Cu(I)Cl_2)_n^{2n-}$ and $(Cu_2(II)Cl_6)_n^{2n-}$. It will be remembered that the latter compound is determined only at low resolution although the heavy atoms are moderately accurate.

Structural Interpretation of the Charge-Transfer

As described in the introduction, there is no particular feature of the paraquat ion which confers on it the property of electron acceptor and the ion must be considered as a whole. There are probably about four possible geometric effects we could look for in these complexes:-

- (a) A reasonably large planar system in the anion which is parallel to and overlaps the aromatic cation.
- (b) A short contact between a specific atom, probably the nitrogen, in paraquat and a chlorine of the anion.
- (c) A direct "outer-sphere" interaction between the metal atom and the aromatic system, which would be difficult to observe unless there were some very short distances.
- (d) A chlorine atom approaching the centre of the aromatic system, with the latter preferably in line with the metal-chlorine vector.
- (e) There have been recent speculations that hydrogen atoms of phenyl groups, etc. can fill the coordination positions of a metal atom. On this unproven hypothesis it is just conceivable that a $M \cdots H_3 \equiv C-N$ system can be set up with hyperconjugation helping the "conduction" of electrons.

Features of the Structures

(All arguments involving distances or angles involve effects several times larger than $3\sigma.s.d.s$).

The main feature common to all structures is the proximity of one of the chlorine atoms to the nitrogen atom. The separation ranges from 3.23 to 3.84 Å, and the nitrogen-chlorine vector is approximately perpendicular to the paraquat ion in all four structures, although this is not too surprising. It is impossible to deny that this might be chance or that it is an accidental consequence of a more important structural feature that has not been noticed. Whether or not this close contact involves charge-transfer, however, there are two relatively convincing pieces of evidence for an attractive force between the nitrogen and chlorine.

Intermolecular Contact Distances

It is difficult and dangerous to attempt to show that any of the Cl-N distances are less than those predicted by Van der Waals radii, for these will vary significantly according to orientation and nature of hybridisation. Comparable contact distances are hard to find, but in $\pi-\pi$ charge-transfer complexes of chloranil with palladium oxinate (1) and with the oxine dimer (1), the chloranil is tilted so that chlorine-ring distances are increased. This suggests that these contacts are not much less than 3.45 Å, whereas two of the chlorine-nitrogen (paraquat) distances are \sim 3.25 Å. (These occur in $\text{p}6\text{Cu}_2\text{Cl}_4$, which shows the most intense charge-transfer of the compounds studied).

Distortion of the Anions

Much more convincing evidence of attraction is the apparent distortion of the chlorocobaltate and chlorocuprate(II) ions. Evidence from crystallographic distortion is often marginal since it is usually impossible to calculate intermolecular forces reliably enough to know how packing will affect the individual molecules.

The CoCl_4^{2-} ion in Ca_3CoCl_5 was shown to be tetrahedral (2) and later (3) to be significantly distorted from $\bar{4}3m$ symmetry to $\bar{4}2m$. This corresponds to an elongation along the $\bar{4}$ axis, with two angles of 106° and four of 111° . It was explained in terms of an electronic effect, and would be expected to be present in $\text{p}q\text{CoCl}_4$. This has only (approximate) 222 symmetry (Pnab requires only 2) and the angles are shown in Table III.4. If it is assumed that the shape of the anion is not wholly dependent on electronic effects, then the distortion from $\bar{4}2m$ can be reasonably explained as a twisting of about 10° round the diad axis. Besides distorting the bond angles (a significantly endoergic process) this moves two of the chlorine atoms nearer the nitrogens ($\text{N-Cl} = 3.82 \text{ \AA}$) which implies these may

create an attractive couple on the anion. The necessary reaction to this is provided by a molecular "vice" of two paraquat molecules "clamping" the other two chlorine atoms. Figs. VI.1 and VI.2 try to give some idea of this. The structure of the isomorphous pqCuCl_4 (4) has a distorted CuCl_4^{2-} tetrahedron with contact distances of $\text{M-Cl} = 3.50, 3.54 \text{ \AA}$. The distortion is even larger than in the cobalt compound but is readily explained by the well-known Jahn-Teller effect. This is incidentally an interesting comparison of electronic and steric energies, suggesting that at 3.5 \AA the repulsive and attractive forces of the M-Cl contact may be roughly equal, as the ion is virtually undistorted from $\bar{4}2m$ symmetry.

The case of M-Cl forces distorting the $(\text{Cu}_2\text{Cl}_6^{2-})_n$ anion in pqCu_2Cl_6 is strongly supported by the very similar structure of $(\text{Me}_2\text{NH}_2)_2\text{Cu}_2\text{Cl}_6$ (5). Both structures (Figs. VI.3. and VI.4. for pqCu_2Cl_6) have planar Cu_2Cl_6 dimers, with a long fifth Cu-Cl bond linking dimers. The 5-coordination of d^9 ions is unpredictable and, as chlorocuprate(II) anions form planar complexes (6) and 4 planar bonds with two long bonds (7), there might be a strong tendency for the dimer to remain substantially flat. In both structures there are short M-Cl contacts and the (very similar) dimers are significantly non-planar. In $(\text{Me}_2\text{NH}_2)_2\text{Cu}_2\text{Cl}_6$, although not stated, the bond must be a hydrogen bond, $\text{N}^+-\text{H} \cdots \text{Cl}^- (3.25 \text{ \AA})$. This is impossible for paraquat and implies that, regardless of the stable conformation of the anion, the energy of the $\text{M} \cdots \text{Cl}$ system (3.47 \AA) must be similar to the hydrogen bond and implies a definite attractive force. The distortion of the anion from planarity is exactly as if the relevant Cu-Cl bond was bent by an M-Cl force and is about 1° or 34° (24° only in the methylenammonium salt).

The most convincing short contacts, however, occur in pqCu_2Cl_6 , Figs. VI.5 and VI.6, and this has the most intense charge transfer, being jet-black. The nitrogen atom is surrounded on each side of the paraquat plane by chlorine atoms at 3.25 \AA and 3.26 \AA .

The Cl-N-Cl angle is 163° , and the N-Cl vectors are again nearly perpendicular to the paraquat plane. There cannot easily be distortion of the Cu-Cl bonds as the anion is a chain and each chlorine atom a bridge in it.

The planar PdCl_4^{2-} should be an ideal anion for producing a short Cl-N bond as steric hindrance is small. The crystals showed little charge-transfer and hopes were further confounded when a long N-Cl distance (3.78 \AA) and a non-planar paraquat were unexpectedly found (Figs. VI.7 and VI.8). This shows how dangerous it is to make generalizations from even half a dozen similar crystal structures. I have no explanation of this structure, particularly not what forces cause the internal rotation of the paraquat ion. This could perhaps be connected with the weak charge-transfer, but whether cause or effect is not clear; the hypothesis of charge-transfer I now present requires a planar paraquat.

This N-Cl interaction appears to be the main feature common to the structures. Besides this evidence (and the CuCl_4^{2-} structure) N-I and N-Br distances of 3.84 and 3.68 are present in pqI_2 and pqBr_2 (4). Without knowing the detailed structures of these it is difficult to know whether these represent the closest approach, but if N-Cl is taken as 3.45 \AA , these distances correlate well with ionic radii for the halide ions (III.11).

Electronic Interpretation of these Features

The existence of a N-Cl force does not per se suggest that anything but electrostatic interaction is involved. It is difficult to evaluate the partial charge on the Cl atoms but this should be high in $\text{Cu}_2\text{Cl}_4^{2-}$, as it should for N^+ . Making the assumption that this close proximity is necessary to the charge-transfer, we can investigate the possibility of the set-up approximating to a reaction intermediate. Since there are probably radicals in pqFeCl_4 , (and assuming that its structure is basically that of pqCoCl_4)

a model proposed for C-2 in this structure must be compatible with the electron distributions in pq^{2+} and pq^+ . The assumption is also made that all the structures (except possibly the palladium one) involve the same mechanism for this process.

Electronic Structure of Paraquat and its Radical

Only the ground states of the paraquat and its radical are discussed, partly because there is no evidence otherwise, and partly because the molecular orbital calculations have only been made for these. Moreover the electron spin resonance measurements quoted were taken on a solution of pq^+ (assumed to be in the ground state). The two common methods of investigating electron distribution in radicals are spin densities from e.s.r. splitting constants and direct M.O. calculation of π -electron coefficients. Splitting constants for pq^+ and related compounds have been used to calculate π -electron densities and corresponding calculations from M.O. approach have been made (8). These showed that $a_N = a_C + 0.6\beta_C$, where a is the coulomb integral and β is the resonance integral gave the best fit of observed and calculated densities. Calculations for pq^{2+} and bond orders had not been made.

A programme (11) was available for Mc calculations of π -electron densities and bond orders (I.1). Calculations were set up for both paraquat and its radicals on three models: (a) an inductive one with all resonance integrals = 1 for neighbouring atoms, a modification (b) with this integral 0.5 for the central bond, and (c) a "pseudo-atom" model to attempt to account for the effect of the methyl group. (This in effect postulates hyperconjugation between the methyl hydrogen atoms and the ring). For the first two models $a_N = a_C + 0.6\beta$ and $a_{C2} = a_C + 0.09\beta$, whereas for the last, $a_N = a_C + 1.4\beta$, $a_{C2} = a_C + 0.45\beta$. The results, shown as π -densities and π -bond orders, are given in Fig. VI.9. The first two methods give similar results except for the centre bond; although

the third has slightly different values, the qualitative picture of all three is similar. There is, of course, no account of π -densities or electron correlation in this method.

This suggests that if an electron is transferred to create a radical in the ground state, it will be most strongly involved with the nitrogen and C_4 atoms. The symmetry of the lowest available orbital for this electron is shown in Fig. VI.10. As the metal chlorine bond is always roughly parallel to the paraquat the proposed mechanism for C-T is as shown below (Fig. VI.11.)

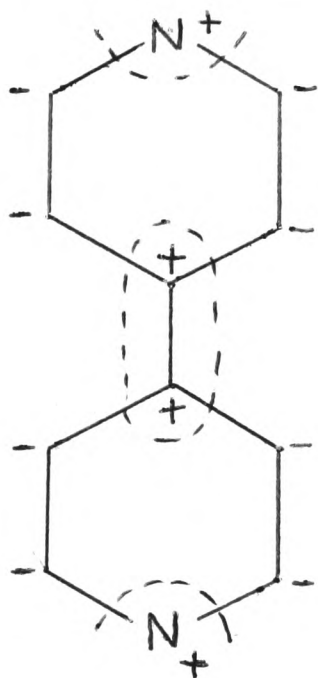


Fig. VI.10.

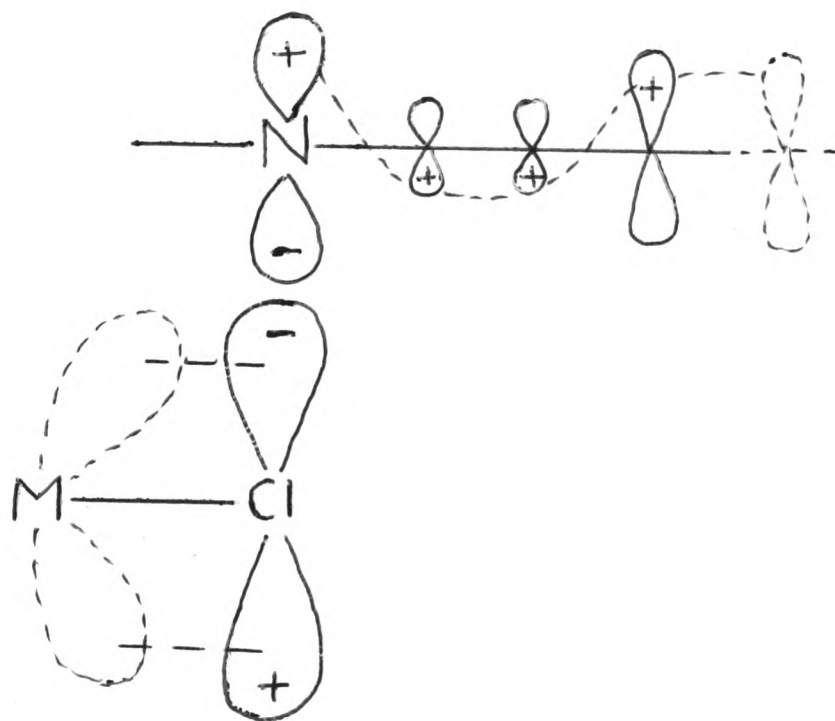
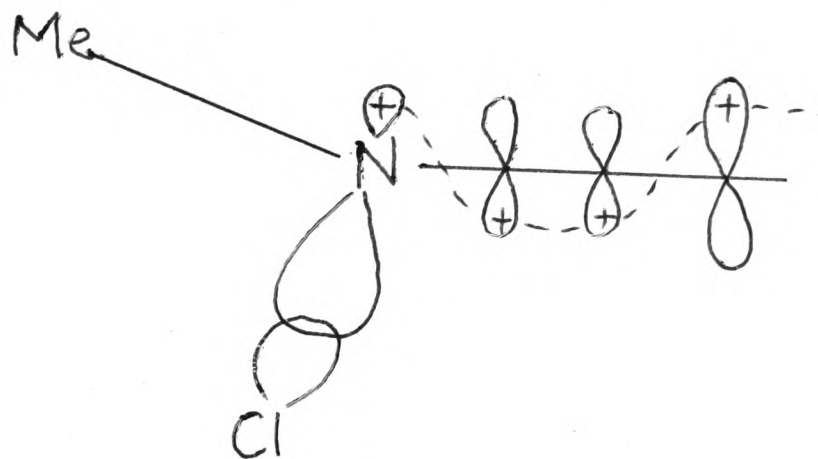


Fig. VI.11.

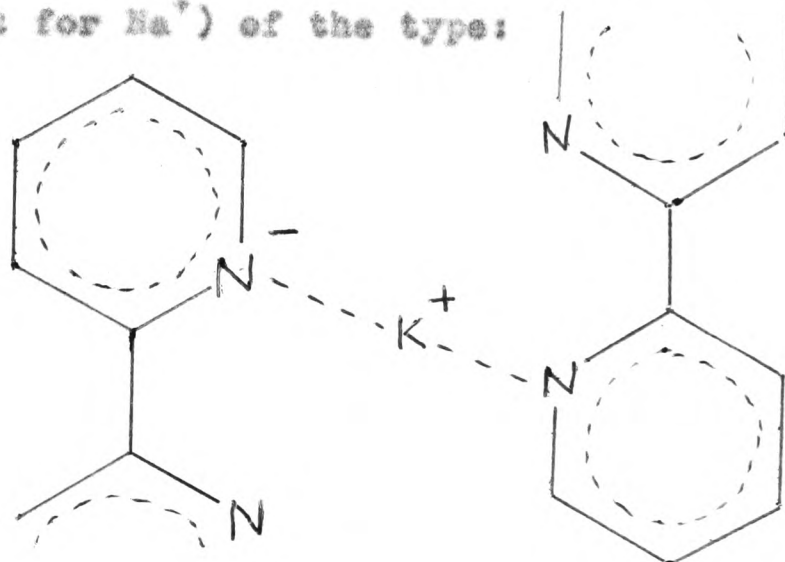
(This also accounts for C-T in the simple paraquat halides). The p orbital may be bonded to the metal which would facilitate electron transfer from a metal via the chlorine. If the unfilled orbital on the nitrogen were regarded as localized, then although C-T could take place, in a similar manner but reverse direction to that in the $py-I_2$ complex, radical formation is unlikely. It seems that the system is more similar to the benzene- I_2 complex, involving the whole

π -system but a specific atom at the point of transfer. There is a possibility that if the Cl-N distance were very short the nitrogen might have some sp^3 hybridisation and the methyl group be distorted out of the paraquat plane:-



There is also the possibility that if similar compounds occur with the diacid of 4,4'-dipyridyl, then hydrogen bonded charge-transfer could be possible.

Electron transfer between 2,2'-dipyridyl and the ion-pair $K^+ 2,2'$ -dipyridyl $^-$ has been studied by line broadening of e.s.r. spectra (12). This may well involve a bridged intermediate (for K^+ but not for Na^+) of the type:



with $\Delta E_{act} \sim 10$ kcal/mole. The electron structures of the two species will be very similar to those of paraquat (although there is the possibility of the lone pair being used). If $K^+ \cdots N$ interactions are possible, a $K^+ \cdots Cl$ analogue appears very reasonable.

The changes in bond order (and length) from pq^{2+} to pq^+ are not very great and so the Franck-Condon principle should be reasonably obeyed. The crystallographically observed bond orders, Fig. VI.9 calculated from plots of bond order against bond length (I.1), give some confirmation that the paraquat ion in the compounds is in the dipositive (ground) state.

Conformation of the Paraquat Ion

The two main forces determining the conformation of bipyridyls (and biphenyls) are steric hindrance (repulsion) between the C_3 hydrogen atoms and the resonance stabilisation, which are in opposite directions. They predict a twist from planar (θ) of 90° and 0° and gaseous electron diffraction determinations of θ as 41.6° for biphenyl (13) and 37.2 for bipyridyl (14,15) suggest that they are of the same order. Crystallographic determinations show that these type of compounds are almost all planar in the solid (16)(17), biphenyl (18), triphenyl (19), although this has marked libration of the central ring (20), and seven of the paraquat compounds. Only $pqPdCl_4$ (50°) and 4,4'-dinitrodiphenyl (33°) (21) appear to be twisted. The packing of biphenyl has been explained in terms of quadrupole moments (22) and it is likely that intermolecular forces are important in determining conformation, besides the two intramolecular effects above.

COORDINATION IN THE ANIONS

The bond lengths in CoCl_4^{2-} (Table III.4) are in good agreement with those of Mason (2.25 Å); those in PdCl_4^{2-} (Table III.5) with $(\text{NH}_4)_2\text{PdCl}_4$ (23), Pd-Cl = 2.30 Å in 4/mm. The coordination of copper(I) and (II) has been extensively studied and even with halogens there are over twenty different coordinations. Copper(II) chloride anions form square planes (6), distorted octahedra (7), distorted tetrahedra (4), trigonal bipyramids, distorted square pyramids (5) and the coordination appears to be almost wholly related to the packing. For instance, it is difficult to visualise anything but a trigonal bipyramid if the large spherical tripositive $\text{Co}(\text{NH}_3)_6^{3+}$ is used as a cation.

For copper(I) the situation is as varied. Surprisingly no crystalline CuCl_4^{3-} has been found (although $\text{Cu}(\text{CN})_4^{3-}$ is known) but chains of tetrahedra exist:- vertex linked (25), edged linked (pqCu_2Cl_4), edge and vertex-linked in Cu_2Cl_3^- (26) and vertex-linked ($\text{Cu}_5\text{Cl}_{16}^{11-}$). 3-coordination does not seem to occur with chloride ion although several with cyanide have been found. The compound Cu.CN.NH_3 (27) has 4 coordinate copper(I) with Cu-Cu distances of 2.14 Å, (cf. pqCu_2Cl_4). It seems that although electronic forces do play a part in copper(II) compounds with chlorine, to a large extent their stereochemistry, and that of copper(I), is dependent on the shape, size and charge of the positive ion.

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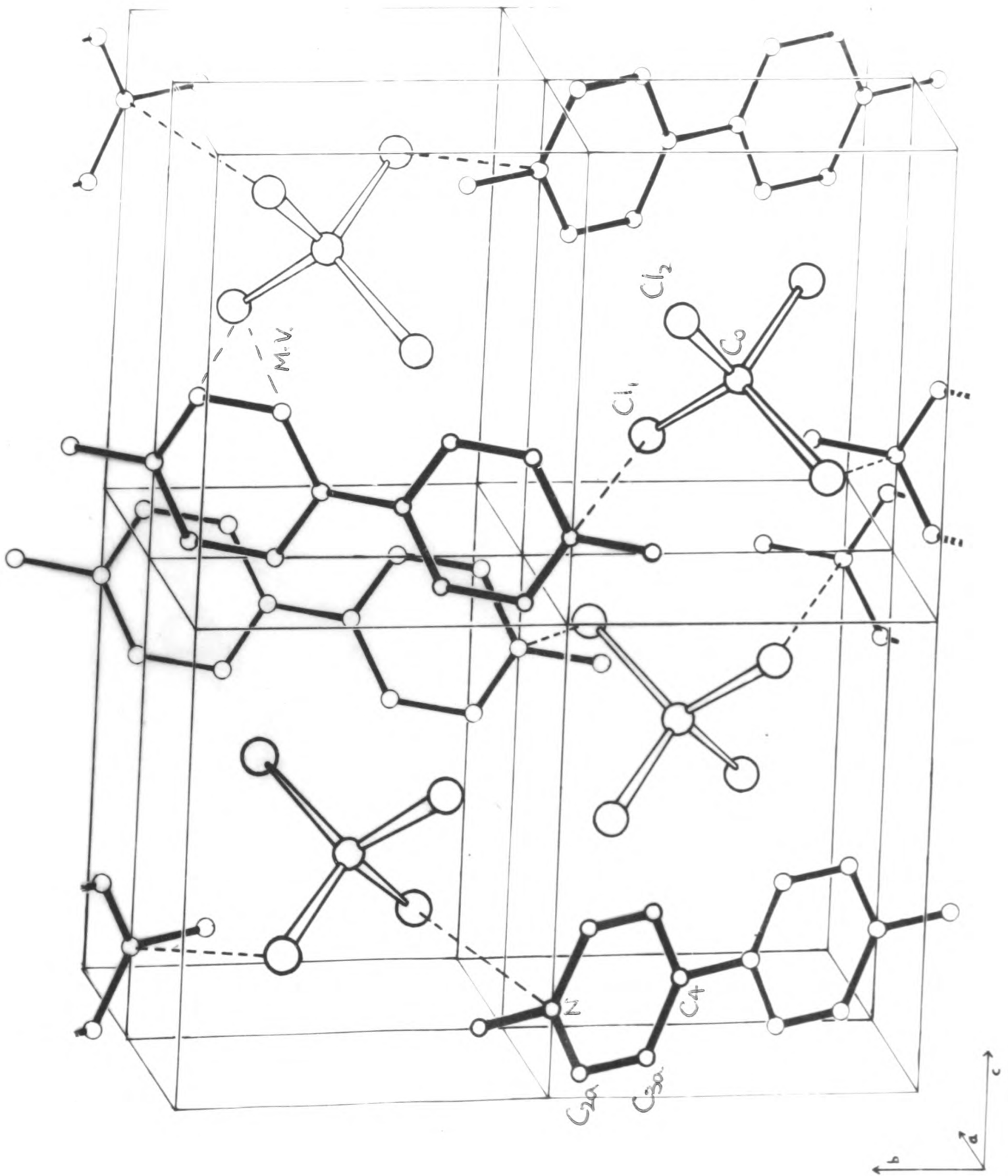


Fig. VI.1. Clinographic view of Parquat CoCl_2 (showing whole cell)

M.V. = the "molecular vice" referred to in the text

The angle shown represents the angle in projection

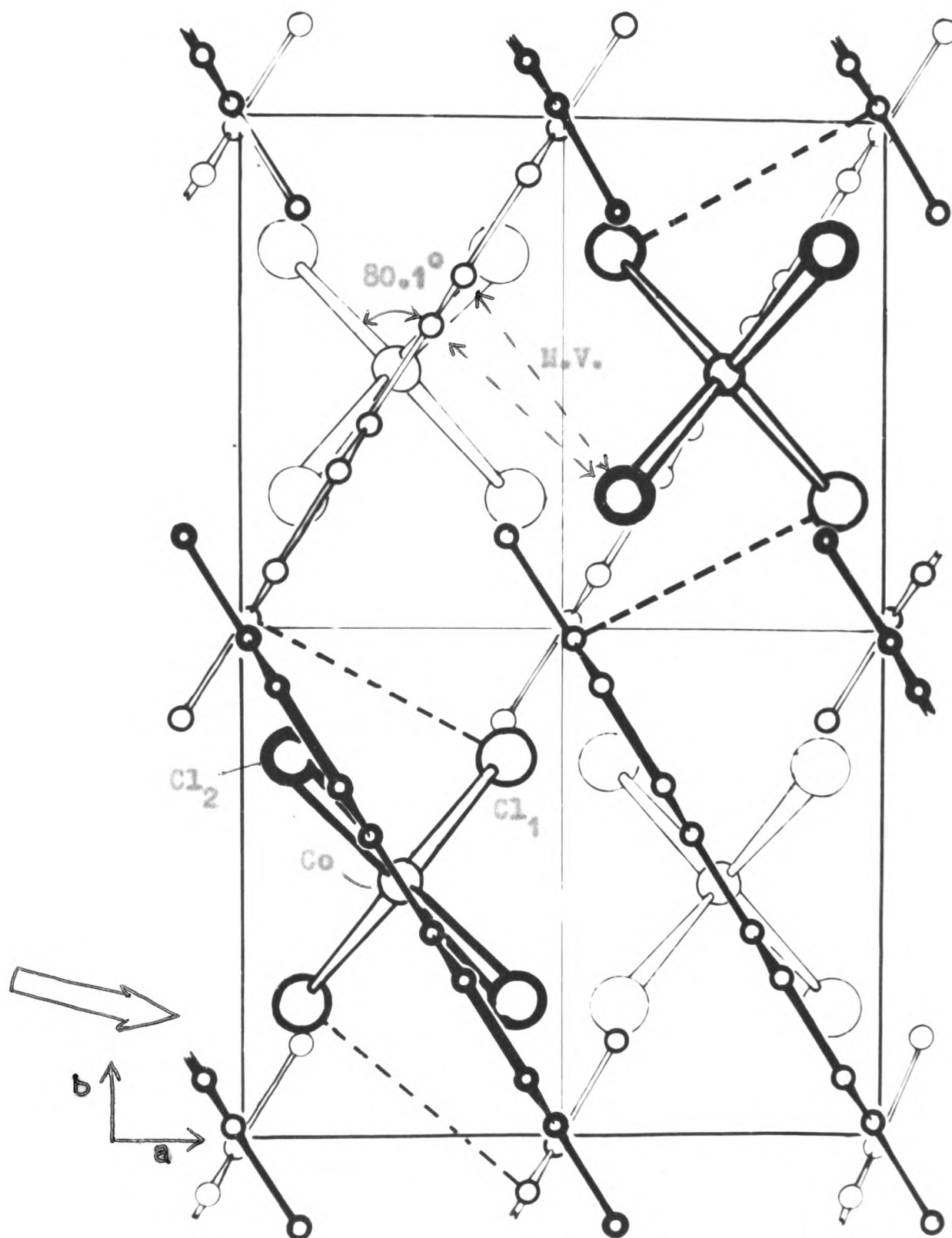


Fig. VI.2. Projection down c for Paraquat CoCl_4

(A whole cell is shown. The arrow represents the approximate angle of the clineographic view.)

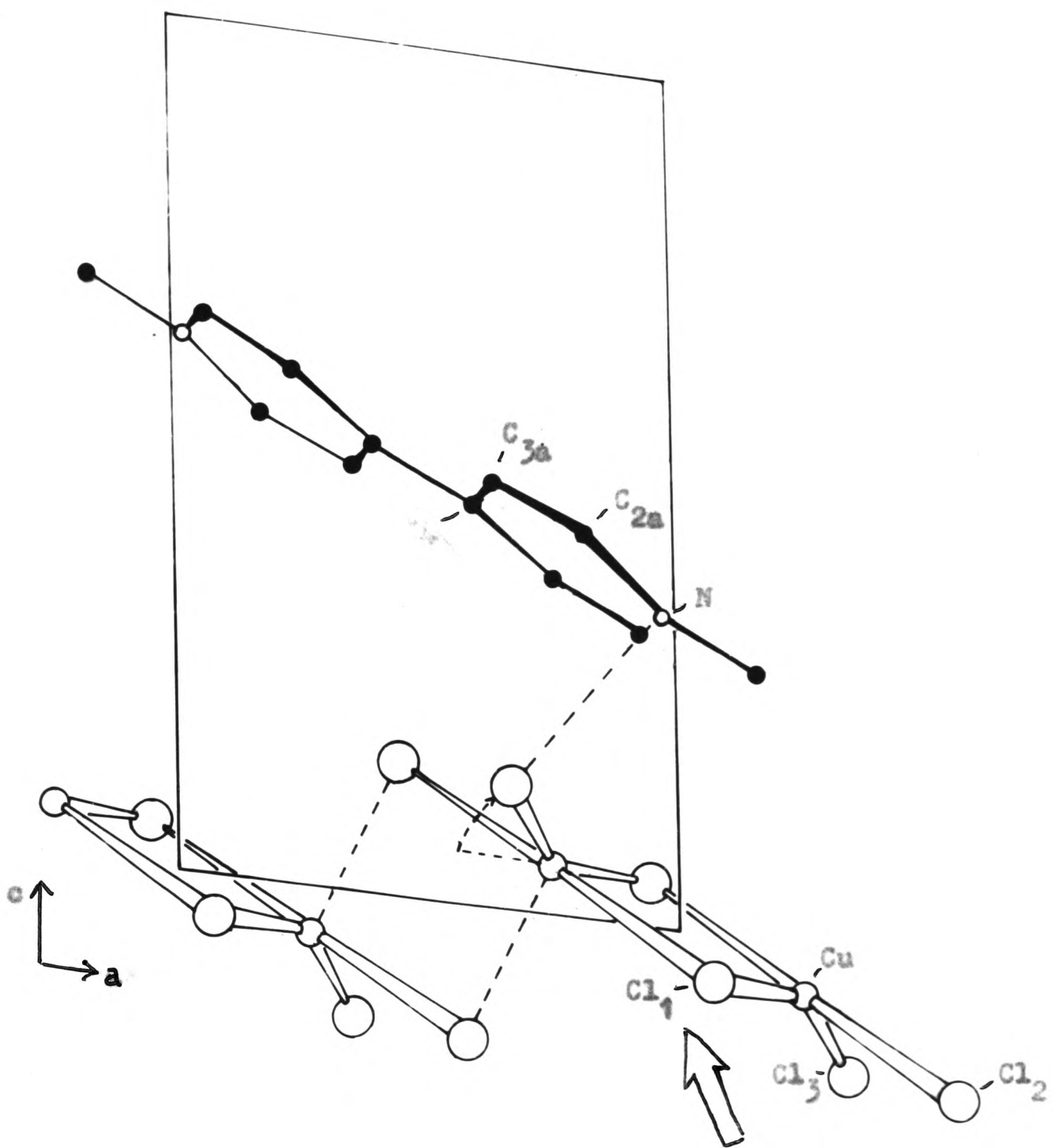


Fig. VI.4. Projection along b for Paraquat Cu_2Cl_6

(A whole cell, but not all its molecules, is shown.)

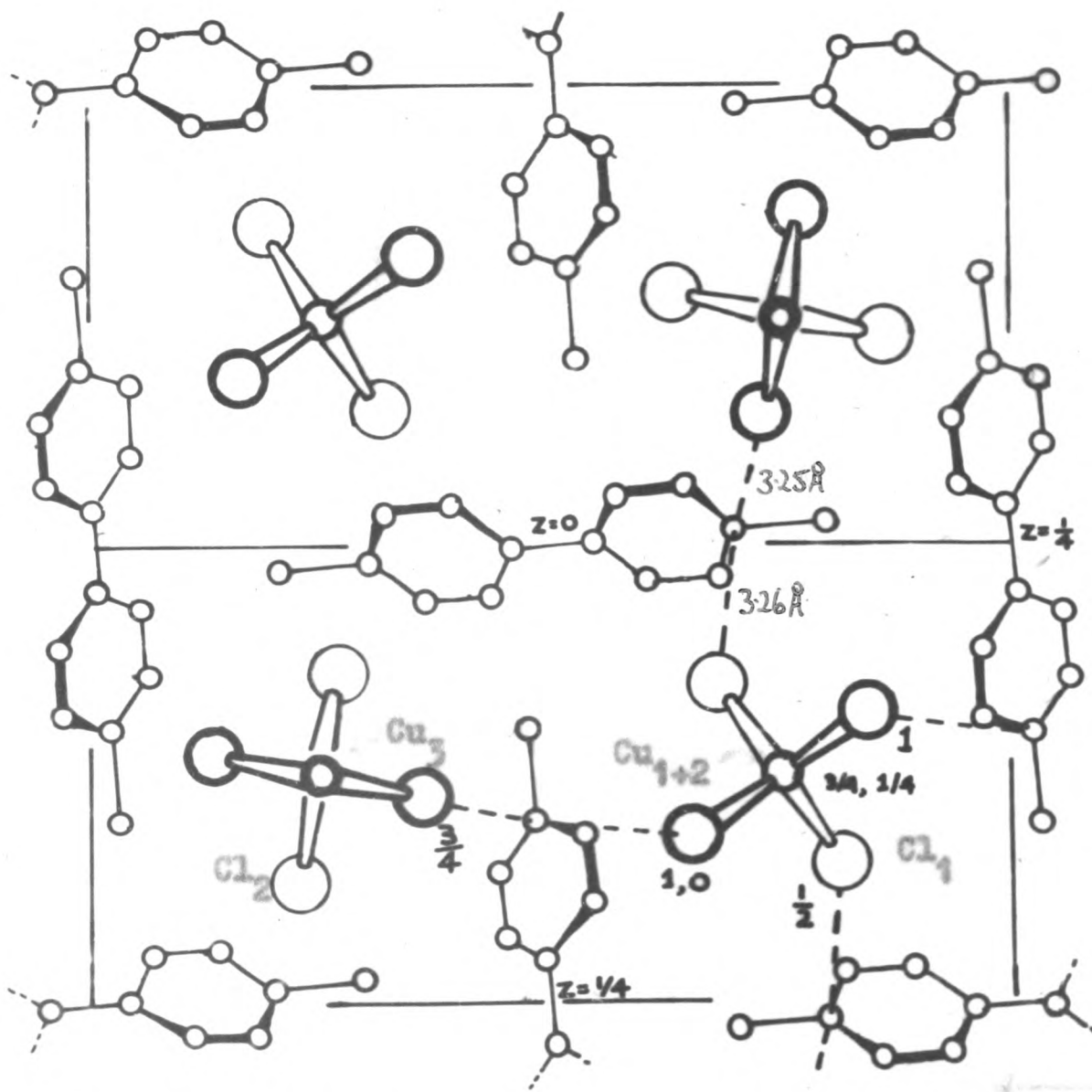


Fig. VI.5. Projection down c for Parquat Cu_2Cl_4

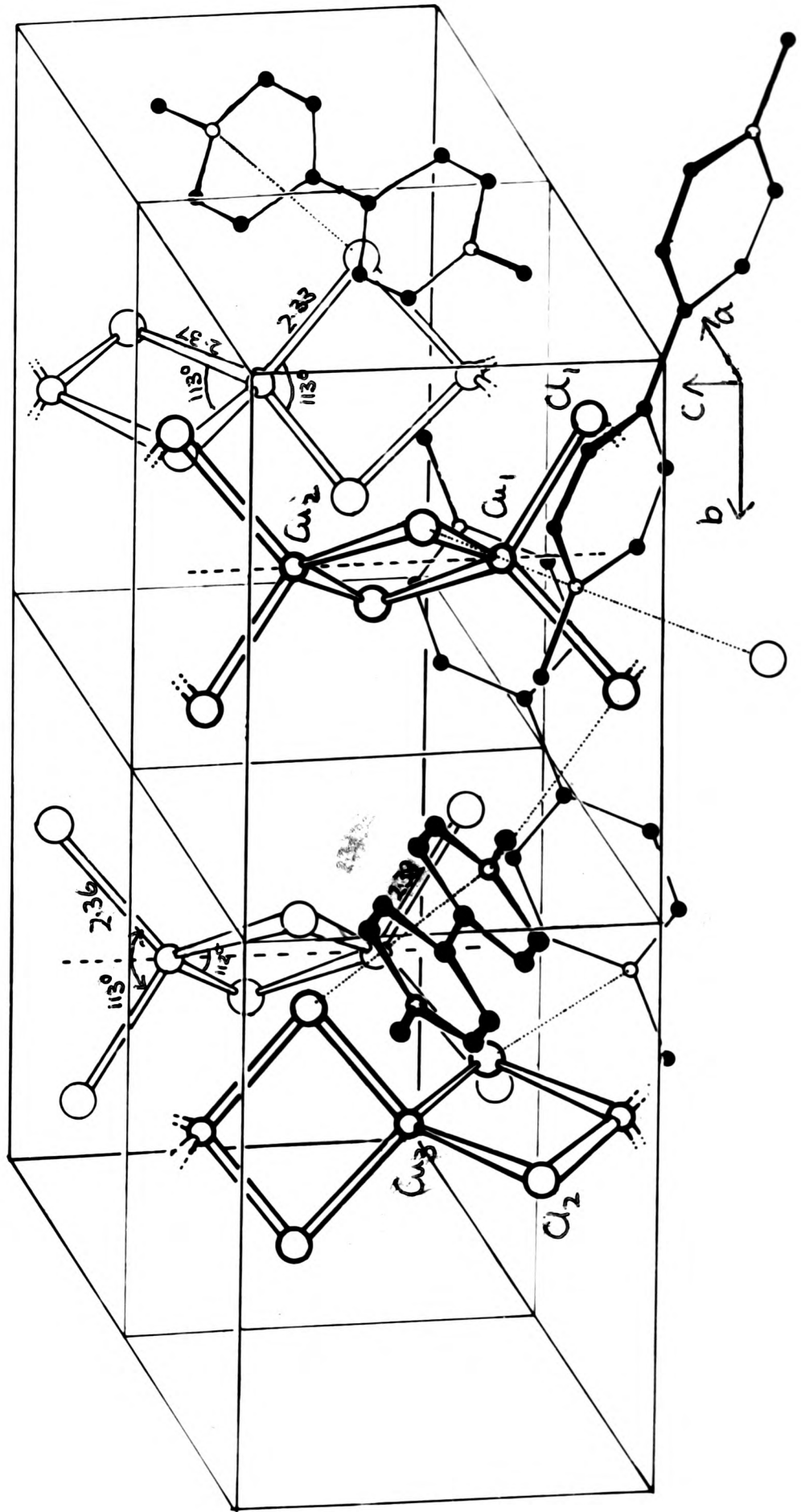


Fig. VI.6. Clinographic view of Paraguan Cu_2Cl_4
 (A whole cell, but not all its molecules, is shown)

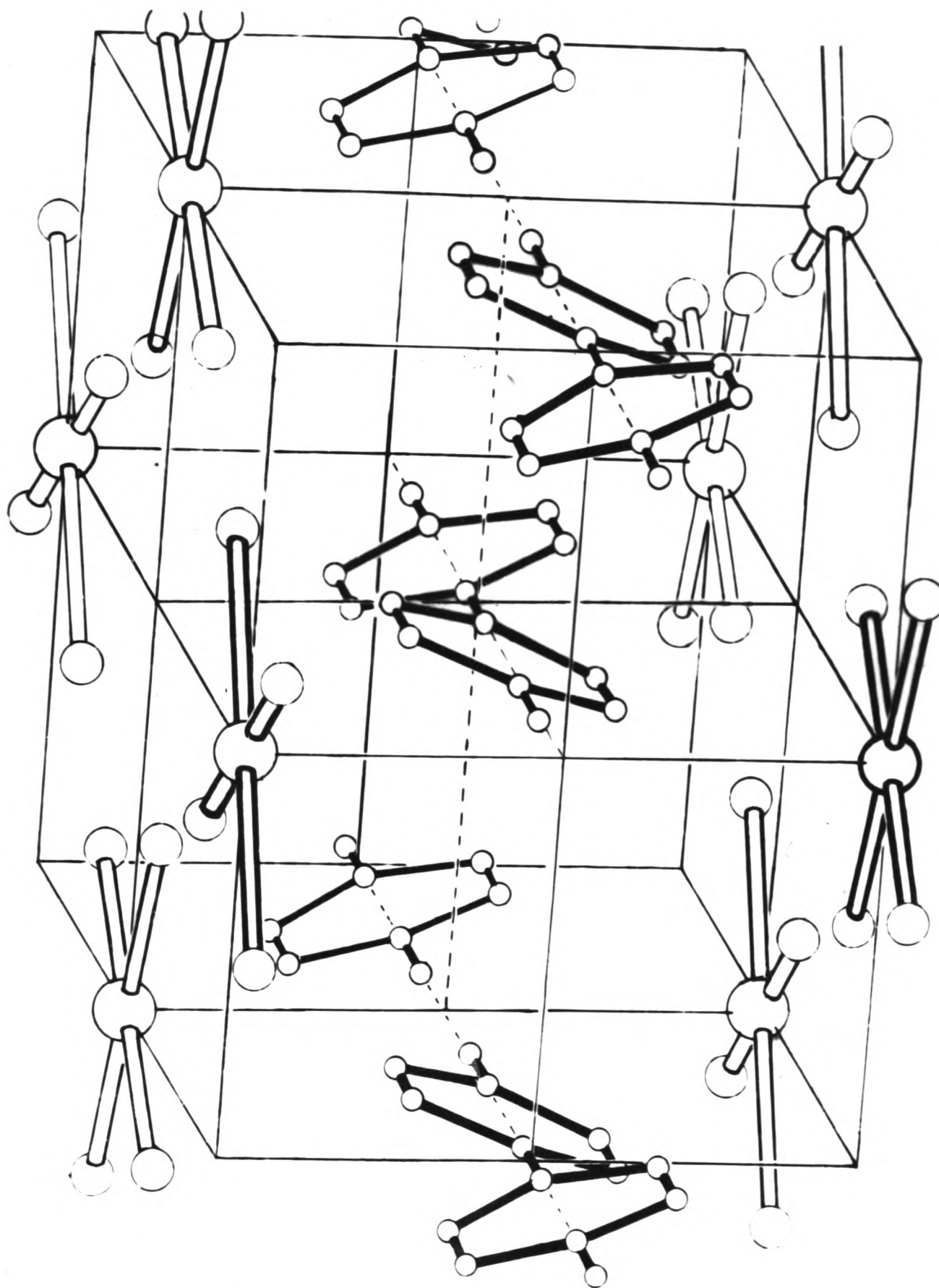


Fig. VI.7. Clinographic view of Paracquat PdCl₄
 (note position of origin is not at corner of diagram.)

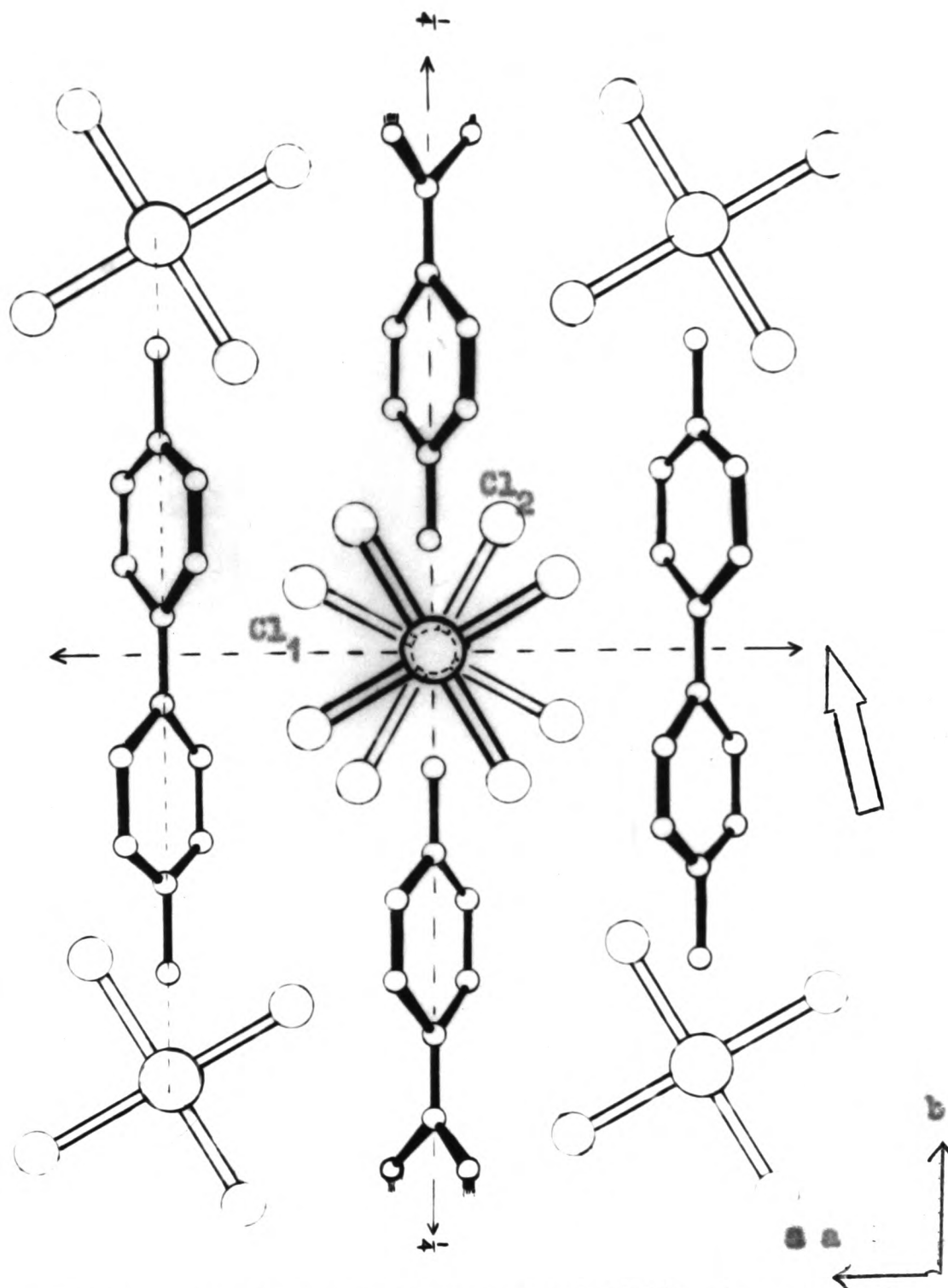


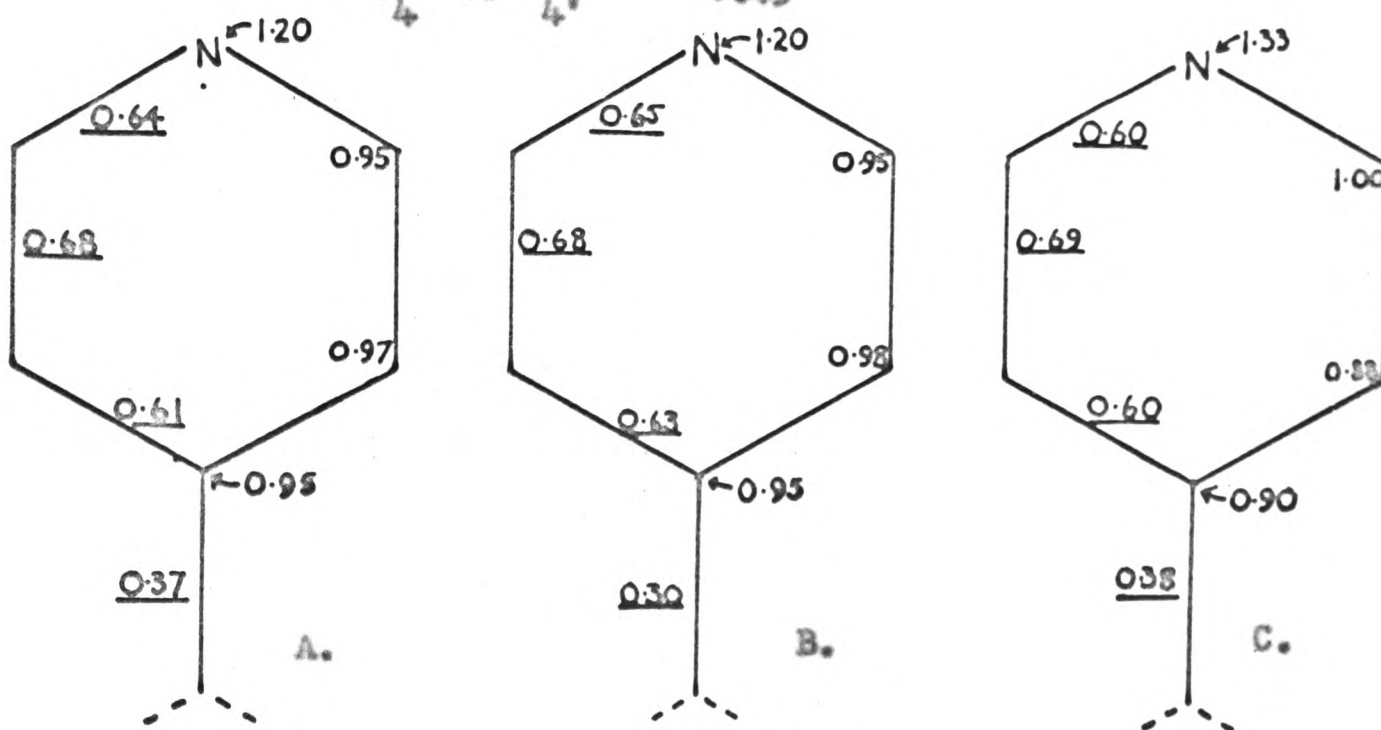
Fig. VI.8. Projection of Paraquat PdCl_4 down c axis

(Note that the origin is at the centre of the diagram.)

Crystallographically observed bond orders (I.1.):-

C ₂	to C ₃	0.72
C ₃	to C ₄	0.62
C ₄	to C _{4'}	0.43

DICATION



MONOCATION

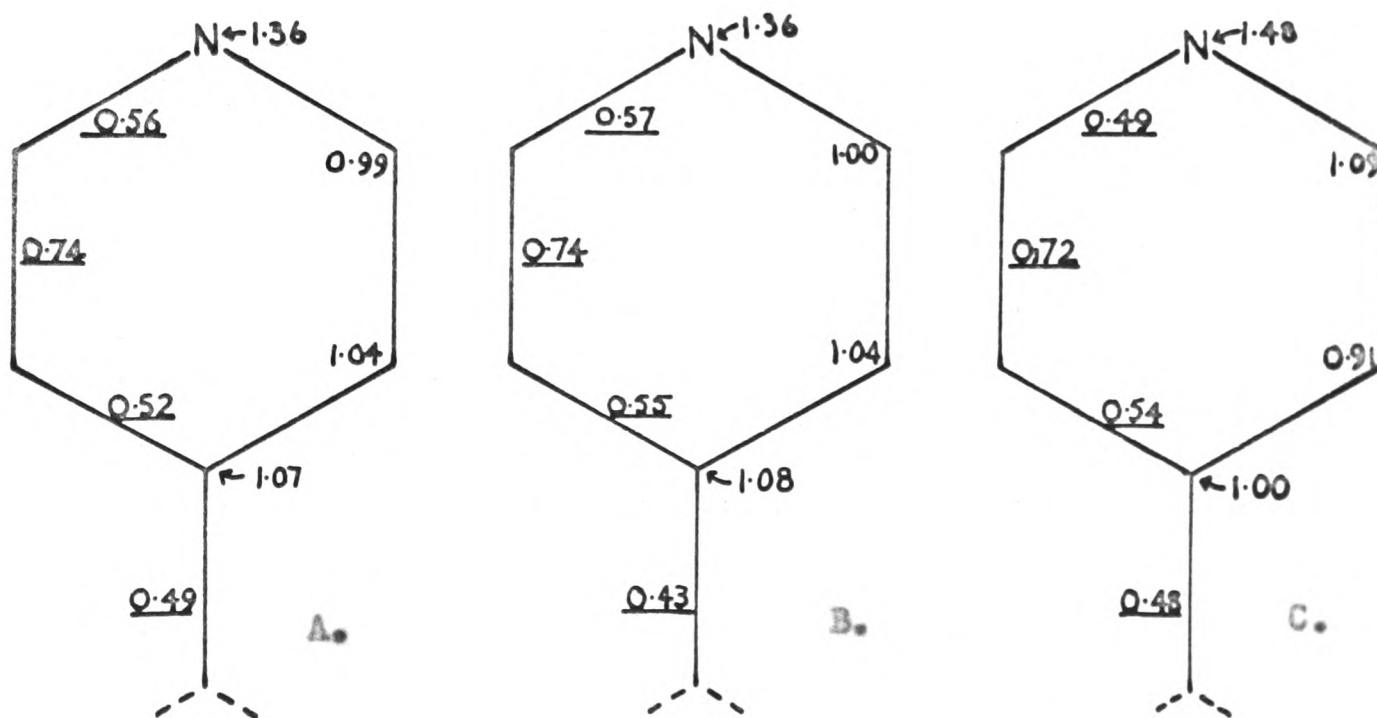


Fig. VI.9. Bond orders and charges calculated by M.O. theory (3 methods)

(for paraquat and its radical monocation)

Crystallographically observed bond orders are also given.

CHAPTER VII.

OBSERVED AND CALCULATED STRUCTURE FACTORS

TABULATION OF STRUCTURE FACTORS

The tables, except VII.3., are arranged in lines of constant h and k . These lines are vertical for the first five tables and horizontal for the last two. They are preceded by $*h$ k for that line and consist of the triplet, l F_{obs} F_{calc} . Values for F_{obs} have already been multiplied by the layer scales. The numbers given are in electrons multiplied by a factor to make the values more manageable and this factor is given for each table.

All the structures are centrosymmetric, but the cobalthexammine chlorocuprates are referred to a non-centric origin for structure factor calculation. For these the phases, or information to calculate them, are given.

For table VIII.3., $pqCu_2Cl_4$, the constant indices are $*l$ h .

TABLES IN CHAPTER VII.

		<u>Value Tabulated</u>
VII.1	Copper Oxinate.2TCNB	5F
VII.2	Paraquat CoCl_4	5F
VII.3	Paraquat Cu_2Cl_4	5F
VII.4	Paraquat PdCl_4	2.5F
VII.5	Mixed-Valence (25%Cu(I)) compound	2.5F
VII.6	Paraquat Cu_2Cl_6	0.5F
VII.7	$\text{Co}(\text{NH}_3)_6 \text{Cu}_5\text{Cl}_{17}$	0.1F

Table VII.1. $\ln k, 1/\bar{v}$ obs \bar{v} calc for $\text{CuOx}_2 \cdot 2\text{TCNB}(\text{H})$

0	0		6	77	86	3	87	96	0	0	6	6	36	20	7	58	45
4	31	-18	7	41	40	4	75	77	4	57	52	7	34	24	8	100	103
7	73	-62	9	172	193	5	45	39	5	51	53	1	1	1	14	30	23
8	68	66	10	162	168	6	29	29	8	50	51	-7	22	26	-12	25	23
9	112	103	11	69	-58	7	38	39	9	36	42	-4	37	40	-11	20	28
10	57	58	15	22	26	9	43	41	10	26	32	-3	26	27	-10	38	49
11	127	118	0	0	3	10	32	40	12	28	28	-1	39	38	-7	29	33
12	45	39	-17	21	24	11	30	27	12	38	28	0	46	42	-5	44	35
13	45	44	-16	28	32	13	41	42	-14	50	41	1	49	46	-4	70	79
14	39	39	-13	32	-24	14	30	30	-6	53	55	3	43	38	-2	60	60
15	22	32	-12	28	28	16	35	34	-5	28	21	4	49	45	-1	47	33
16	40	46	-10	57	53	-16	35	34	-3	51	50	9	42	35	0	27	-23
17	34	41	-9	35	35	-14	28	26	-2	40	42	10	45	37	1	63	56
0	0		-8	53	54	-13	67	58	-1	90	80	11	43	28	2	23	22
-17	43	41	-7	112	111	-12	46	49	-1	94	94	12	23	16	3	88	66
-15	37	25	-6	87	88	-9	95	34	0	97	-96	1	47	-7	4	87	73
-14	39	38	-5	94	94	-8	41	31	1	47	-49	-5	40	45	5	76	64
-13	64	60	-4	167	169	-7	91	89	2	47	54	-4	44	49	6	6	1
-10	59	59	-3	50	51	-6	76	69	3	62	62	-1	32	22	-14	24	22
-8	41	-29	-1	155	-164	-5	90	83	4	59	53	0	67	65	-12	50	58
-7	140	148	0	17	22	-4	113	120	5	35	34	1	52	-42	-11	31	38
-5	75	82	1	73	70	-3	56	50	7	38	36	2	37	31	-10	53	58
-4	149	163	2	85	87	-2	132	135	10	22	22	3	39	38	-7	48	55
4	64	-61	3	192	225	0	91	84	11	27	20	4	44	33	-6	55	53
5	17	16	4	154	165	1	87	92	-12	27	20	5	76	62	-5	66	64
6	39	37	5	83	-72	2	100	100	-11	35	24	6	53	48	-4	85	94
7	76	74	6	33	26	3	77	76	-8	30	30	7	30	20	-3	27	28
8	116	114	7	65	66	4	73	76	-7	25	25	8	45	41	-2	29	28
9	73	70	9	95	85	5	52	55	-4	42	45	9	31	31	-1	78	76
10	179	184	10	28	-24	6	47	48	1	22	22	10	48	36	0	75	76
11	114	128	12	46	55	10	37	42	2	28	36	11	56	45	1	91	91
12	28	29	15	34	27	13	37	45	3	28	36	12	60	36	2	99	39
14	22	33	0	0	4	16	22	21	4	37	35	13	52	30	3	51	53
16	39	31	-16	30	28	-16	22	21	5	36	33	13	52	30	4	51	53
-17	29	27	-13	37	35	-15	32	27	9	24	16	11	20	22	5	24	21
-15	47	50	-12	51	-46	-14	34	31	9	24	16	11	20	22	6	24	21
-14	76	83	-11	72	61	-13	38	35	-8	29	20	-9	21	26	7	72	53
-10	48	43	-10	141	152	-12	38	35	-4	36	29	-7	38	45	8	40	38
-8	97	101	-9	56	53	-12	43	38	-3	36	29	-6	42	45	9	78	84
-7	212	222	-8	60	53	-11	27	23	-3	23	17	-5	30	22	10	20	20
-6	108	105	-7	79	77	-7	51	47	2	21	-15	-4	28	28	11	60	52
-5	38	35	-6	69	66	-5	109	106	2	45	37	-3	50	53	12	78	73
-4	35	38	-5	136	137	-4	110	100	1	45	37	-2	45	-43	13	35	30
-3	24	-8	-4	82	90	-3	52	55	-4	22	23	0	23	21	14	21	24
-2	45	-35	-3	136	-141	-2	40	55	-1	38	32	2	105	107	17	43	38
0	24	23	-1	69	68	-1	157	162	1	47	46	3	99	97	18	17	17
3	21	21	0	119	120	0	100	102	2	29	22	4	29	23	-16	17	21
5	74	83	2	143	154	3	67	68	3	37	29	5	89	85	-15	33	34
									4	32	28	6	36	23	-14	30	28

-12	27	24	12	60	49	3	130	128	0	1	2	-12	99	96	0	1	6
-11	100	114	13	34	35	4	121	123	-4	129	138	-10	40	34	-15	31	20
-10	57	71	14	47	42	5	75	-63	-3	49	-42	-9	55	54	-14	29	26
-9	36	34	15	44	47	6	49	41	-1	49	53	-8	84	79	-13	36	33
-6	25	20	16	42	35	7	74	78	0	68	62	-7	141	141	-10	63	-54
-5	51	49	17	31	28	8	75	74	1	97	105	-6	108	104	-8	46	47
-4	131	152	18	25	21	9	80	87	2	30	29	-5	104	108	-7	52	43
-3	130	148	19	25	21	10	32	38	3	101	104	-4	75	72	-6	89	75
-2	189	214	-14	23	15	11	55	52	4	40	34	-3	72	76	-5	38	32
-1	111	119	-12	21	24	12	22	29	5	26	-17	-1	121	126	-4	60	51
0	140	166	-11	113	123	13	31	29	6	138	141	0	129	138	-3	63	57
1	55	-46	-10	120	132	14	31	29	7	109	117	1	162	172	-2	63	52
2	165	174	-9	57	62	15	23	24	8	54	65	2	88	98	0	71	74
3	94	100	-8	122	106	16	23	24	9	89	91	3	93	102	1	23	26
4	119	127	-7	25	19	17	23	24	10	27	23	4	30	29	2	56	56
5	124	124	-6	43	38	18	28	26	11	27	23	5	25	21	3	33	27
6	135	141	-5	177	200	19	28	26	12	28	23	6	36	39	4	30	33
7	132	126	-4	18	19	20	48	50	13	39	39	7	25	27	5	32	36
8	60	63	-3	156	179	21	49	46	14	51	41	8	32	33	6	26	27
9	38	25	-2	114	105	22	50	51	15	27	27	9	47	60	7	26	27
11	85	79	-1	130	121	23	32	34	16	39	29	10	58	65	8	25	16
12	36	35	2	80	-67	24	99	96	17	44	41	11	40	42	9	65	67
14	35	30	3	172	177	25	69	74	18	79	81	12	35	40	10	29	23
15	50	45	4	33	-21	26	57	61	19	43	36	13	35	40	11	63	63
16	20	23	5	122	122	27	42	43	20	53	51	14	24	22	12	42	32
17	32	21	6	197	228	28	37	-41	21	119	113	15	1	5	13	28	-37
-16	22	23	7	146	146	29	37	37	22	112	105	16	27	23	14	30	-29
-13	39	30	8	122	122	30	50	46	23	76	71	17	32	32	15	27	27
-11	102	103	9	36	33	31	22	-8	24	121	124	18	32	32	16	27	27
-9	22	34	10	58	49	32	22	-8	24	121	124	19	32	32	17	27	27
-8	127	141	11	40	44	33	22	-8	24	121	124	20	32	32	18	27	27
-7	74	82	12	45	35	34	22	-8	24	121	124	21	32	32	19	27	27
-6	114	123	13	34	34	35	22	-8	24	121	124	22	32	32	20	27	27
-5	140	144	14	26	25	36	22	-8	24	121	124	23	32	32	21	27	27
-4	96	97	15	25	19	37	22	-8	24	121	124	24	32	32	22	27	27
-3	174	184	16	25	19	38	22	-8	24	121	124	25	32	32	23	27	27
-1	187	203	17	24	24	39	22	-8	24	121	124	26	32	32	24	27	27
0	292	307	18	24	27	40	22	-8	24	121	124	27	32	32	25	27	27
1	78	84	19	11	24	41	22	-8	24	121	124	28	32	32	26	27	27
2	71	65	20	38	34	42	22	-8	24	121	124	29	32	32	27	27	27
3	83	90	21	70	64	43	22	-8	24	121	124	30	32	32	28	27	27
4	81	79	22	39	33	44	22	-8	24	121	124	31	32	32	29	27	27
5	163	149	23	56	49	45	22	-8	24	121	124	32	32	32	30	27	27
6	111	112	24	144	149	46	22	-8	24	121	124	33	32	32	31	27	27
7	126	130	25	133	137	47	22	-8	24	121	124	34	32	32	32	27	27
8	94	92	26	102	-91	48	22	-8	24	121</							

Table VII.1a (cont.)

n k, 1 50 obs 5 scale

0	1	9	-5	69	59	-5	93	87	0	2	-2	-7	22	-13	-11	45	51
1	37	35	-4	75	77	-4	104	104	-9	43	40	-6	28	-13	-11	45	51
2	31	33	-3	28	30	-3	163	149	-8	42	41	-5	36	-29	-8	144	149
3	22	24	-8	77	-80	-2	25	94	-7	34	28	3	31	37	-7	185	186
4	36	37	-1	121	-105	-1	162	154	-6	102	-79	4	25	32	-6	37	29
5	15	16	0	142	141	0	84	78	-2	99	88	5	221	239	-5	133	150
6	29	21	1	113	92	2	24	24	-1	55	-38	6	84	91	-4	102	101
7	33	30	2	30	30	4	149	129	0	80	74	7	31	-22	-3	64	67
8	23	22	3	66	50	5	120	107	1	96	-79	8	59	56	-2	190	205
9	33	27	4	81	62	6	86	82	2	69	66	10	32	32	-1	105	114
10	23	23	5	84	81	7	36	26	3	139	116	12	42	42	0	158	179
11	18	17	6	123	111	8	91	-94	4	75	74	13	20	25	1	77	85
12	22	21	11	57	45	9	69	79	5	50	52	14	24	24	2	47	45
13	21	14	13	36	35	10	197	216	6	130	118	16	36	39	4	34	30
14	26	26	13	42	50	11	36	34	7	93	87	18	26	36	5	86	88
15	44	39	14	43	50	12	64	47	8	48	49	20	26	36	6	93	98
16	45	40	15	48	48	13	40	34	9	71	73	22	29	36	7	29	32
17	41	30	16	63	60	14	26	24	11	29	26	24	23	28	8	43	56
18	46	40	17	58	53	15	57	49	12	51	-56	26	22	25	9	174	201
19	42	26	18	85	92	16	32	24	13	29	26	28	46	49	10	82	103
20	46	40	19	43	29	17	23	25	14	39	44	30	83	76	11	23	20
21	42	40	20	43	29	18	23	26	15	25	33	32	47	53	12	36	26
22	22	13	21	134	132	19	49	52	16	27	30	34	108	101	13	36	26
23	30	24	22	124	128	20	53	53	17	24	-27	34	33	40	14	17	25
24	29	23	23	81	74	21	39	42	18	45	-24	33	33	-14	16	33	43
25	43	38	24	113	105	22	63	56	19	156	164	35	158	150	17	66	67
26	49	38	25	104	99	23	71	79	20	129	141	36	128	116	18	39	40
27	23	26	26	125	120	24	135	120	21	35	27	37	37	-13	10	33	39
28	28	34	27	56	49	25	81	70	22	59	55	38	62	60	11	63	70
29	31	30	28	66	59	26	172	156	23	99	111	39	52	54	12	108	121
30	72	73	29	79	85	27	139	136	24	84	80	40	149	-133	13	31	33
31	66	53	30	25	-22	28	63	-53	25	70	66	41	168	172	14	63	64
32	47	53	31	53	-59	29	176	-148	26	51	49	42	59	50	15	27	25
33	25	-12	32	49	41	30	119	118	27	49	48	43	114	117	16	128	143
34	35	32	33	66	59	31	154	143	28	121	120	44	95	95	17	30	25
35	25	24	34	48	59	32	92	96	29	22	241	54	54	-51	18	109	120
36	51	54	35	44	43	33	132	131	30	74	81	45	71	71	19	107	92
37	46	47	36	44	30	34	101	87	31	39	40	46	70	70	20	68	77
38	27	21	37	28	27	35	68	66	32	48	64	47	51	55	21	45	49
39	30	24	38	38	33	36	53	41	33	42	41	48	118	117	22	42	48
40	44	30	39	48	60	37	13	50	34	33	28	44	44	45	23	136	160
41	26	20	40	39	46	38	15	33	35	33	31	47	39	44	24	100	108
42	35	28	41	22	24	39	16	23	36	25	14	31	31	34	25	31	30
43	17	23	42	47	48	40	23	27	37	31	29	32	22	22	26	28	20
44	40	49	43	28	24	41	51	52	38	109	111	32	28	34	27	40	38
45	35	38	44	37	29	42	65	-59	39	120	127	33	50	50	28	34	30
46	42	40	45	41	53	43	40	-36	40	70	70	34	51	60	29	45	50
			46	79	74	46	117	128	41	144	132	35	32	32	30	45	50

0	4	86	4	46	53	-3	35	38	4	3	-5	-2	92	83	-7	52	49
1	86	66	5	32	38	-2	31	25	5	70	70	-1	99	104	-6	52	49
2	65	77	6	39	40	-1	33	74	6	70	95	1	16	104	-5	74	65
3	87	77	7	51	54	0	46	44	7	37	32	2	81	75	-4	152	153
4	143	133	8	57	60	1	65	53	8	101	94	3	54	47	0	95	100
5	23	20	9	50	52	2	42	36	9	86	63	4	122	113	1	8	10
6	94	-86	10	31	38	3	34	33	10	71	62	5	87	73	2	136	138
7	69	72	11	42	44	4	37	30	11	93	89	6	42	45	3	61	55
8	23	27	12	44	44	5	35	23	12	29	23	7	67	77	4	82	81
9	77	81	13	30	32	6	36	32	13	39	25	8	41	-36	5	103	99
10	122	123	14	24	24	7	38	32	14	49	33	9	56	45	6	118	105
11	82	87	15	54	53	8	39	29	15	27	25	10	67	62	7	114	114
12	43	45	16	29	30	9	33	22	16	39	30	11	26	19	8	54	57
13	47	44	17	36	38	10	53	45	17	25	23	12	28	22	9	73	73
14	23	22	18	28	34	11	16	13	18	29	-4	13	38	36	10	60	50
15	43	44	19	40	42	12	20	22	19	43	27	14	23	24	11	29	34
16	45	70	20	28	39	13	23	30	20	28	23	15	23	24	12	44	40
17	121	129	21	43	-45	14	66	72	21	28	25	16	44	55	13	40	37
18	46	43	22	43	46	15	45	44	22	52	45	17	21	18	14	31	26
19	45	48	23	30	36	16	57	55	23	98	102	18	35	37	15	31	26
20	106	96	24	26	30	17	99	91	24	94	85	19	87	84	16	14	20
21	47	48	25	18	20	18	25	23	25	43	34	20	132	126	17	30	20
22	90	-80	26	19	19	19	18	15	26	81	77	21	112	104	18	38	50
23	13	18	27	24	22	20	66	70	27	37	-31	22	109	152	19	98	102
24	66	68	28	22	22	21	68	-55	28	68	65	23	74	62	20	63	58
25	75	80	29	44	38	22	45	40	29	19	19	24	168	171	21	99	97
26	64	63	30	38	26	23	92	82	30	108	94	25	154	147	22	52	52
27	43	49	31	29	22	24	66	48	31	94	81	26	82	74	23	144	138
28	36	38	32	26	21	25	53	51	32	32	30	27	76	75	24	39	42
29	26	33	33	32	23	26	70	63	33	66	67	28	119	116	25	123	106
30	22	21	34	32	23	27	44	42	34	26	25	29	122	108	26	26	22
31	22	21	35	37	27	28	59	43	35	48	35	30	134	116	27	109	109
32	42	43	36	43	47	29	66	53	36	54	50	31	38	35	28	98	97
33	41	40	37	36	31	30	44	34	37	38	34	32	24	24	29	120	111
34	75	77	38	37	30	31	44	28	38	43	42	33	38	42	30	146	159
35	53	-43	39	46	48	32	37	22	39	34	34	34	32	25	31	93	91
36	14	21	40	38	30	33	30	22	40	17	10	35	23	20	32	77	72
37	62	57	41	38	30	34	25	25	41	20	20	36	36	26	33	50	48
38	40	38	42	19	13	35	23	31	42	28	32	37	40	40	34	53	50
39	37	37	43	25	21	36	30	28	43	30	36	38	21	21	35	62	65
40	40	37	44	31	20	37	38	42	44	35	37	39	20	11	36	52	54
41	38	39	45	44	34	38	40	43	45	34	34	40	20	-1	37	35	26
42	32	34	46	29	20	39	40	33	46	27	27	41	12	18	38	34	26
43	82	82	47	17	17	40	50	50	47	38	44	42	21	25	39	25	22
44	65	66	48	27	24	41	31	-28	48	42	42	43	70	60	40	31	22
45	39	38	49	35													

Table VII.1.(cont.)

h k, l 5th obs 5th calc

• 3	1	-10	54	58	-6	35	41	• 3	8	-3	97	53	8	105	-103
-13	40	37	-9	83	83	-5	37	40	-4	43	45	-2	109	-106	72
-12	25	26	-8	77	79	-3	33	31	-3	27	-27	-1	33	-21	54
-11	37	44	-7	63	69	-2	79	81	1	21	22	0	59	55	61
-10	29	32	-6	31	25	-1	81	83	2	26	27	1	35	31	42
-9	84	87	-5	31	38	0	76	84	3	24	30	2	60	56	18
-8	83	80	-4	16	-11	1	36	42	• 4	• 4	• 4	5	86	84	27
-7	47	-36	-3	30	-20	2	45	54	-1	34	34	6	64	61	• 4
-5	47	47	-2	49	52	3	32	32	0	50	43	7	17	12	-15
-4	42	33	-1	55	55	5	29	36	1	49	40	8	29	25	-15
-3	176	172	0	71	72	8	33	38	4	21	18	10	40	36	-12
-2	165	175	1	106	122	9	24	28	6	19	13	12	26	22	-11
-1	144	126	2	23	27	10	41	44	• 6	• 4	-8	13	34	23	-9
1	94	98	3	43	39	11	19	25	-6	11	-10	11	34	23	-9
2	52	50	4	50	53	• 3	• 6	• 6	-4	62	56	-10	• 4	• 5	-8
3	70	71	6	38	104	-15	17	23	-3	18	16	-9	27	29	-6
4	71	-66	7	82	99	-14	33	36	-2	19	18	-7	34	37	-5
5	23	26	9	47	46	-13	34	38	-1	28	28	-6	44	41	-4
6	104	117	• 3	• 4	• 4	-12	21	23	0	25	27	-5	69	69	-3
7	60	63	-17	17	28	-11	20	20	1	34	26	-3	81	-75	-2
8	64	72	-16	32	33	-8	32	38	2	36	30	-2	67	-61	-1
9	22	23	-13	43	47	-7	32	34	3	30	26	-1	129	124	0
10	58	46	-12	42	46	-6	29	30	4	30	25	0	78	62	1
15	19	22	-11	28	20	-5	47	54	7	38	29	1	29	29	3
• 3	• 2	• 2	-10	22	28	-4	26	30	8	22	11	2	59	51	4
-16	21	23	-9	52	55	-3	64	66	9	47	32	3	110	92	5
-13	25	29	-8	69	69	-1	39	39	10	44	35	4	84	76	6
-12	28	32	-7	67	68	0	35	35	• 4	• 4	-7	5	82	70	7
-10	27	32	-6	56	56	1	35	39	-9	13	17	6	31	24	8
-9	36	43	-5	22	18	2	47	57	-6	41	41	10	32	26	9
-8	95	99	-4	95	88	3	45	49	-5	28	27	• 3	• 4	• 4	10
-7	102	106	-3	22	23	8	33	33	-4	28	33	-13	18	26	11
-5	88	85	-2	66	73	9	27	28	-3	23	23	-12	38	52	12
-4	25	18	-1	45	51	• 3	• 7	• 7	-2	38	32	-11	33	40	13
-3	122	126	0	93	106	-11	38	-39	-1	38	28	-10	21	25	14
-2	102	92	1	88	96	-9	42	44	0	48	42	-7	46	52	15
0	48	-36	2	100	116	-8	20	17	1	24	14	-6	48	47	• 4
1	63	63	3	70	73	-7	40	44	2	65	49	-5	66	70	-16
2	58	-46	4	43	46	-6	12	16	3	26	24	-4	45	35	-14
3	63	62	6	26	21	-5	21	25	6	25	26	-3	45	33	-13
4	52	47	9	50	64	-4	66	71	7	30	27	-2	122	117	-11
6	62	58	10	26	30	-3	40	39	9	46	38	-1	83	73	-9
8	44	55	• 3	• 5	• 5	-1	37	39	10	69	46	0	114	101	-8
9	79	86	-14	28	33	2	39	40	• 4	• 3	-3	1	125	104	-7
10	47	60	-13	63	73	5	24	25	-8	22	22	2	86	75	-6
• 3	• 2	• 2	-10	21	27	-10	24	30	-7	22	22	4	72	60	-4
-16	22	25	-9	36	39	-7	34	37	-6	44	44	5	72	66	-3
-12	26	30	-8	83	95	-6	29	26	-5	42	42	6	99	99	-2
-11	32	33	-7	63	65	-5	30	28	-4	88	90	7	61	-61	-1
• 4	• 2	• 2	-8	57	56	9	76	84	• 4	• 5	• 5	1	19	22	4
0	51	-36	-7	53	47	10	33	44	-10	27	34	-4	5	-8	6
1	86	-60	-6	19	-22	12	31	31	-9	48	46	-4	18	-16	7
2	106	106	-5	38	41	• 4	• 3	• 3	-8	25	30	-1	43	37	8
3	113	119	-4	91	74	-16	41	49	-7	52	53	0	22	17	9
4	89	92	1	23	23	-15	31	42	-6	99	89	1	22	18	10
5	117	118	2	54	57	-12	24	26	-4	37	37	-25	22	19	11
6	77	74	4	86	95	-9	84	99	0	46	52	5	39	28	12
7	48	47	5	140	160	-8	94	105	1	39	46	6	25	16	• 5
8	83	81	6	40	49	-7	50	57	2	51	57	7	25	16	-12
9	44	35	8	25	28	-6	54	63	3	36	42	• 5	• 7	• 7	-11
10	52	50	9	22	28	-5	52	51	5	20	22	-3	23	21	-9
12	29	23	10	23	-21	-4	26	22	7	23	24	0	49	39	-8
13	22	13	12	27	36	-3	103	115	8	14	24	1	44	34	-7
14	24	21	13	30	26	-2	69	65	9	22	26	2	54	50	-6
15	20	19	15	26	22	-1	125	140	• 4	• 6	• 6	3	38	28	-5
• 4	• 1	• 1	• 4	• 1	• 1	0	29	31	-13	51	52	8	39	26	-3
-14	29	27	-13	60	68	1	39	-18	-12	64	71	9	40	31	-2
-13	62	-62	-12	54	65	2	52	54	-9	29	31	• 5	-6	31	0
-11	102	119	-11	45	52	3	57	53	-8	27	31	-8	28	31	1
-10	37	43	-10	91	87	4	49	59	-7	46	44	-7	36	40	1
-9	16	10	-9	20	-14	5	57	61	-5	30	-29	-6	33	36	2
-8	41	38	-7	77	76	7	30	47	-2	28	32	-5	47	50	3
-6	192	168	-5	63	56	9	46	55	-1	20	19	-4	30	31	5
-5	220	217	-4	56	-35	• 4	• 4	• 4	-1	28	32	-2	90	88	6
-4	65	-39	0	47	62	-16	23	32	1	28	34	-1	29	25	9
-3	79	79	1	55	59	-12	21	25	2	24	30	0	29	29	10
-2	23	18	3	30	29	-11	19	22	3	21	27	1	31	29	11
0	23	26	4	139	158	-10	28	33	4	27	34	3	34	24	12
1	25	30	5	30	33	-9	55	59	8	18	18	4	42	39	• 5
2	40	35	6	26	-14	-8	41	48	• 4	• 7	• 7	5	25	19	-14
3	46	46	7	34	40	-6	66	61	-9	17	21	6	28	20	-12
4	22	25	9	39	44	-5	90	85	-6	35	37	9	49	33	-7
5	103	112	• 4	• 2	• 2	-3	47	54	-5	23	25	10	45	26	-6
6	108	107	-13	39	39	-2	41	39	-4	18	-10	11	50	39	-5
7	40	44	-12	37	42	-1	48	57	-3	25	28	12	22	15	-4
8	55	61	-11	60	69	0	19	-3	-2	58	58	• 5	• 5	• 5	-3
9	28	35	-8	87	84	1	57	67	-1	34	36	-11	21	27	-2
10	17	12	-7	61	53	2	105	127	• 4	• 8	• 8	• 5	25	27	-1
11	43	-44	-5	58	53	3	24	29	-2	22	26	-9	25	27	0
12	31	30	-4	58	78	5	38	44	3	28	32	-8	43	55	1
13	48	50	-3	149	154	7	67	83	4	18	19	-7	37	37	1
14	14	23	-1	30	46	8	44	47	• 5	• 17	-17	-6	58	55	2
15	26	24	0	19	-3	9	23	-27	• 4	• 8	• 8	-5	79	80	4
• 4	• 0	• 0	1	107	111	11	17	22	-8	30	36	-2	47	51	5
-14	27	-23	2	45	47	-14	18	-14	-7	28	23	-1	28	-20	7
-12	102	113	3	25	23	-13	28	-26	-6	20	26	0	52	43	10
-11	64	74	4	30	32	-12	28	-26	-4	28	17	1	114	94	11
-10	52	59	6	46	48	-12	28	-26	-3	74	79	2	62	42	12
-9	57	57	8	40	45	-11	28	-26	-2	37	-28	3	77	73	13

Table VII.1.(cont.)

σ_h k_0 l 5° obs 5° calc

0	5	-2	3	40	37	0	5	5	0	6	-4	0	6	-1	0	4	94	93
-13	30	35	4	74	78	-14	5	5	0	1	50	43	-14	26	31	0	6	94
-12	50	56	5	63	73	-13	0	0	0	2	75	70	-13	25	-27	6	24	30
-11	25	32	0	5	0	-9	0	0	0	3	55	48	-12	24	34	0	26	36
-10	26	29	-16	17	19	-8	0	0	0	4	29	30	-9	38	43	0	39	43
-6	43	44	-14	32	42	-7	6	-8	0	5	55	49	-8	37	-28	0	38	-17
-5	33	34	-14	32	37	0	0	0	0	6	35	30	-6	116	131	0	90	-75
-4	69	61	-12	60	65	-2	0	0	0	7	35	-24	-5	80	93	0	150	145
-3	60	46	-10	36	36	-1	0	0	0	11	25	17	-4	61	81	0	19	3
-2	76	77	-9	80	90	0	0	0	0	0	0	0	-3	48	54	0	24	22
-1	37	-39	-8	34	32	0	0	0	0	-12	17	26	-3	41	47	0	30	23
1	23	23	-7	43	48	0	0	0	0	-8	23	29	4	64	62	0	72	75
2	60	56	-4	73	76	-2	0	0	0	-7	27	26	5	59	52	0	20	13
3	53	52	-3	125	136	-1	0	0	0	-6	29	34	6	38	30	0	0	0
4	67	68	-2	56	52	0	0	0	0	-5	42	43	7	55	49	0	13	19
5	73	68	-1	115	119	1	0	0	0	-4	36	29	8	27	17	0	19	23
6	48	53	0	28	30	5	0	0	0	-3	65	65	9	0	0	0	24	28
7	41	48	1	60	62	8	0	0	0	-2	35	32	0	0	0	0	41	46
0	5	-1	2	85	107	0	6	-4	0	-1	86	87	-14	39	-39	0	27	30
-13	61	80	3	55	57	0	6	34	0	0	74	63	-12	55	70	0	27	30
-10	26	25	4	45	48	-7	33	37	0	3	46	45	-9	29	37	0	37	32
-7	75	69	5	34	50	-4	60	63	0	2	56	56	-8	36	41	0	23	28
-6	64	47	0	5	0	-3	33	28	0	4	28	27	-7	76	78	0	70	65
-5	53	44	-16	12	18	-1	29	32	0	5	67	60	-6	69	89	0	81	82
-4	86	87	-14	21	20	1	43	49	0	6	67	60	-4	55	66	0	25	22
-3	45	43	-13	26	27	2	34	32	0	7	91	-64	-2	42	60	0	28	31
-2	108	114	-10	29	35	5	33	32	0	8	91	81	1	22	13	0	39	35
-1	38	34	-9	28	36	8	41	26	0	9	59	49	2	26	29	0	30	28
1	43	47	-7	23	25	0	6	-5	0	10	44	36	4	83	81	0	33	-16
2	50	47	-6	50	49	-10	28	37	0	11	38	29	5	74	78	0	39	38
3	70	74	-5	24	23	-7	27	37	0	0	0	0	6	32	30	0	65	54
4	69	67	-4	93	96	-6	42	50	0	11	6	29	7	37	35	0	92	92
5	66	57	-3	118	111	-5	61	72	0	-14	17	34	8	42	34	0	6	48
6	37	37	-2	98	101	-4	30	28	0	-13	16	28	10	27	-19	0	46	48
0	5	0	-1	67	83	-3	86	-103	0	-9	35	37	11	28	24	0	46	48
-13	37	43	0	5	3	-1	56	58	0	-8	32	29	12	44	36	0	32	31
-12	58	65	-10	31	38	-1	26	26	0	-6	33	39	0	6	1	0	46	52
-11	42	54	-9	63	62	2	23	15	0	-5	64	83	-13	57	61	0	37	33
-10	66	64	-8	67	67	4	80	80	0	-4	44	47	-12	27	35	0	63	72
-9	41	44	-6	51	50	5	62	52	0	-3	59	67	-11	36	56	0	29	-14
-7	96	81	-5	37	33	9	46	33	0	-2	83	88	-10	32	36	0	45	49
-6	55	55	-4	41	43	10	28	19	0	-1	93	87	-8	26	29	0	49	41
-5	64	65	-3	37	38	10	28	19	0	0	47	-39	-6	44	47	0	25	29
-4	82	83	0	5	4	11	27	15	0	3	68	65	-5	73	85	0	29	35
-3	93	102	-13	26	27	0	6	-4	0	4	49	49	-4	73	-67	0	92	91
-2	71	75	-11	24	25	-11	35	47	0	5	51	49	-3	63	-30	0	22	18
-1	74	93	-10	33	39	-6	40	46	0	7	40	35	-1	39	40	0	22	18
0	12	12	-9	38	41	-4	52	-8	0	8	64	62	0	45	44	0	26	31
1	52	67	-8	31	36	-2	68	71	0	9	22	22	1	37	31	0	28	30
2	36	30	-3	25	35	0	47	48	0	12	45	35	2	29	-2	0	41	39
										12	29	20	3	69	67	0	60	54
																4	21	23
																2	21	10

Table VII.2. (cont.)

on k₀ 1 5^F obs 5^F calc

• 3 8	• 3 15	4 176	-138	• 4 7	• 4 13	5 102	-87
• 4 77	1 152	5 49	21	9 215	3 152	7 142	261
6 146	2 87	6 195	185	10 63	4 91	8 67	-75
7 231	4 103	7 155	-182	11 82	5 117	10 133	-143
8 155	5 141	8 181	180	12 51	11 42	11 140	-116
10 249	7 88	9 177	-211	13 42	13 13	• 5 102	80
• 3 9	9 101	12 132	-120	• 4 8	0 128	1 102	3
• 4 120	• 3 17	• 4 4	• 4 4	0 114	1 78	2 120	-117
3 96	1 206	0 155	141	1 157	2 142	3 195	205
6 138	• 4 0	1 226	-211	3 241	3 108	4 88	-85
8 64	1 4	2 221	236	4 138	4 106	5 291	305
10 121	1 31	3 270	-271	7 203	6 89	6 242	214
11 97	2 89	4 110	104	8 80	7 83	7 144	140
13 104	3 394	7 127	102	9 78	8 67	8 125	141
• 3 10	5 136	8 100	62	11 88	10 54	9 83	34
• 4 112	6 298	9 92	92	12 59	• 4 15	10 66	-52
1 98	7 160	10 122	116	13 82	2 170	11 44	47
4 140	8 61	12 91	111	• 4 9	3 83	12 123	-114
6 229	9 45	14 47	64	1 168	4 151	13 84	88
8 260	10 80	• 4 5	• 4 5	3 203	6 108	• 5 251	265
7 64	11 64	1 531	569	5 184	8 96	1 154	145
8 110	12 170	3 57	-8	7 166	• 4 16	2 167	167
• 3 11	13 84	4 50	-49	9 99	2 206	3 212	238
3 193	14 189	5 107	-68	11 74	3 95	4 54	54
5 228	• 4 1	6 91	76	13 54	4 62	5 213	-204
7 90	1 115	7 178	165	• 4 10	• 4 17	6 132	119
• 3 12	3 91	8 78	74	0 74	1 4	7 160	-177
0 175	4 180	9 299	332	1 114	2 35	8 143	135
1 108	5 175	10 91	49	2 62	3 68	10 103	134
4 330	6 64	14 136	-102	3 159	4 61	11 155	141
6 220	7 81	• 4 6	• 4 6	4 270	6 59	13 52	33
7 63	8 114	0 414	-377	5 46	• 4 18	• 5 34	7
• 3 13	9 101	1 97	-74	6 80	0 139	2 136	131
1 110	11 64	2 125	-118	7 195	2 93	3 140	-125
2 110	12 104	3 117	101	11 97	3 56	4 57	79
3 319	• 4 2	4 37	14	12 71	4 47	5 241	-248
4 156	1 33	6 146	156	• 4 11	• 5 1	7 140	-130
5 253	2 243	7 230	-221	1 63	2 168	• 6 136	-106
6 62	3 156	8 194	-194	2 135	3 276	1 52	-41
7 134	4 330	10 101	-122	3 199	5 65	2 137	-130
8 57	6 294	13 55	58	4 195	6 59	3 56	-50
• 3 14	7 71	• 4 7	• 4 7	5 284	7 100	4 88	117
0 110	8 62	1 331	-331	8 97	9 158	5 256	-248
1 64	10 92	2 174	144	9 57	10 94	6 256	-248
2 127	12 169	3 146	-142	• 4 12	11 189	7 140	-130
3 219	14 123	4 69	58	2 138	13 134	8 143	-117
4 139	• 4 3	5 72	-37	4 153	• 5 2	9 67	62
7 112	1 250	6 62	57	6 149	1 269	10 57	-74
10 66	3 311	7 239	-217	9 74	2 236	11 57	-74
	3 85	8 65	-76	10 65			

• 5 7	• 5 12	6 176	135
1 69	0 135	9 92	-81
2 142	1 48	• 6 2	• 6 2
3 42	2 143	0 276	322
4 98	3 67	2 203	200
5 102	4 72	3 117	-119
6 199	5 76	4 93	58
8 116	7 91	8 199	220
10 61	10 91	10 138	138
12 100	• 5 13	• 6 3	• 6 3
• 5 8	1 155	1 138	157
0 147	2 92	2 102	103
1 113	3 79	3 101	102
2 230	4 63	5 85	60
4 165	5 129	6 110	-114
5 105	9 66	10 86	70
6 106	• 5 14	• 6 4	• 6 4
8 175	1 108	2 71	-39
10 71	2 76	4 278	-273
12 74	3 155	8 93	-70
• 5 9	4 144	• 6 5	• 6 5
1 142	5 137	2 203	-201
2 352	6 137	3 275	-268
3 103	7 126	4 181	-180
4 105	8 54	5 190	-177
5 247	• 5 15	6 114	150
6 135	1 153	10 116	-107
8 121	3 103	• 6 6	• 6 6
9 96	5 60	0 79	-62
10 109	6 76	1 98	72
12 67	7 67	2 130	96
• 5 10	• 5 16	3 209	183
0 186	0 96	6 114	98
2 170	1 102	7 115	-118
3 82	3 106	• 6 7	• 6 7
4 107	6 30	2 138	155
5 67	• 5 17	3 199	199
6 82	1 58	5 361	416
8 62	3 57	• 6 8	• 6 8
10 101	• 6 0	1 110	-68
• 5 11	2 264	3 92	-67
1 83	3 269	5 115	65
2 117	4 115	• 6 9	• 6 9
3 117	5 115	1 160	-146
5 115	6 186	3 148	-118
6 113	7 291	5 114	-124
7 62	8 163	• 6 11	• 6 11
8 41	10 232	1 186	185
10 77	• 6 1	6 118	118
11 42	1 185	• 6 12	• 6 12
	5 115	0 231	241
		2 127	117
		6 94	-57
		• 6 13	• 6 13
		1 135	-133
		2 133	109

.....

Table VII.3.

σ_1 h, k 5σ obs 5σ calc for Paraquat Cu_2Cl_4

h	k	5σ obs	5σ calc	h	k	5σ obs	5σ calc	h	k	5σ obs	5σ calc	h	k	5σ obs	5σ calc
0	0	344	-358	1	1	246	220	4	4	326	-344	17	17	109	68
0	1	550	620	3	3	549	556	6	6	597	668	10	10	109	68
0	2	664	-738	5	5	450	-424	8	8	442	-511	11	11	60	55
0	3	388	389	7	7	129	105	10	10	258	261	12	12	65	60
0	4	193	-230	11	11	67	70	12	12	116	-149	13	13	65	-40
0	5	211	225	13	13	333	-305	14	14	152	163	14	14	66	83
0	6	110	-107	15	15	73	42	16	16	259	-278	15	15	18	18
0	7	174	132	17	17	113	108	18	18	165	165	16	16	202	238
1	0	420	-471	19	19	92	-76	19	19	143	-134	17	17	201	-253
1	1	439	398	2	2	391	452	7	7	206	178	18	18	146	204
1	2	566	-626	4	4	421	-490	11	11	90	108	19	19	19	-120
1	3	307	356	6	6	640	767	15	15	118	150	20	20	117	122
1	4	134	103	8	8	603	-571	17	17	84	-72	21	21	84	-72
1	5	209	-231	10	10	255	250	19	19	102	105	22	22	88	-117
1	6	52	43	12	12	160	-160	2	2	197	-196	23	23	93	-63
1	7	188	108	14	14	52	55	4	4	139	154	24	24	102	178
1	8	103	-91	16	16	175	-133	6	6	232	-261	25	25	102	178
2	0	546	1120	18	18	111	92	8	8	293	325	26	26	112	-125
2	1	525	-603	1	1	231	-188	10	10	114	-121	27	27	121	-145
2	2	430	463	3	3	287	-278	12	12	69	75	28	28	121	-145
2	3	724	-864	5	5	613	628	14	14	213	-227	29	29	125	-163
2	4	778	879	7	7	109	-76	16	16	84	87	30	30	125	-163
2	5	694	-729	9	9	92	-67	18	18	86	-71	31	31	127	-163
2	6	305	390	11	11	50	38	20	20	154	141	32	32	127	-163
2	7	204	-226	13	13	74	-58	22	22	104	-90	33	33	127	-163
2	8	99	114	15	15	43	58	24	24	126	-149	34	34	127	-163
2	9	54	-46	17	17	43	58	26	26	175	184	35	35	127	-163
3	0	458	521	2	2	810	-923	28	28	175	184	36	36	127	-163
3	1	204	-174	4	4	664	769	30	30	66	-62	37	37	127	-163
3	2	29	38	6	6	473	489	32	32	90	76	38	38	127	-163
3	3	307	233	8	8	140	-141	34	34	127	129	39	39	127	-163
3	4	169	-143	10	10	164	-128	36	36	222	-238	40	40	127	-163
3	5	222	-246	12	12	102	-104	38	38	173	173	41	41	127	-163
3	6	224	-231	14	14	217	215	40	40	60	-66	42	42	127	-163
3	7	160	172	16	16	182	-167	42	42	95	79	43	43	127	-163
4	0	554	-748	1	1	295	250	44	44	127	121	44	44	127	-163
4	1	492	623	3	3	209	207	46	46	124	124	45	45	127	-163
4	2	408	-499	5	5	121	103	48	48	156	-181	46	46	127	-163
4	3	664	722	7	7	121	-98	50	50	119	58	47	47	127	-163
4	4	368	-409	9	9	212	212	52	52	88	-105	48	48	127	-163
4	5	277	286	11	11	146	-160	54	54	71	-58	49	49	127	-163
4	6	238	-240	13	13	73	-42	56	56	217	-217	50	50	127	-163
4	7	262	237	15	15	75	-81	58	58	217	-217	51	51	127	-163
4	8	95	-87	17	17	32	57	60	60	111	99	52	52	127	-163
4	9	58	62	19	19	20	57	62	62	127	-148	53	53	127	-163
5	0	67	56	2	2	45	-17	64	64	55	48	54	54	127	-163
5	1	220	-228	4	4	41	60	66	66	245	-223	55	55	127	-163
5	2	87	-79	6	6	16	-66	68	68	62	71	56	56	127	-163
5	3	72	76	8	8	50	-66	70	70	144	160	57	57	127	-163
5	4	225	-234	10	10	124	150	72	72	27	32	58	58	127	-163
5	5	183	-190	12	12	46	68	74	74	214	-212	59	59	127	-163
5	6	76	-82	14	14	98	-93	76	76	203	93	60	60	127	-163
5	7	56	63	16	16	17	-93	78	78	76	69	61	61	127	-163
5	8	49	-22	18	18	44	59	80	80	31	-33	62	62	127	-163
5	9	47	-42	20	20	109	-123	82	82	177	-188	63	63	127	-163
6	0	82	-82	22	22	45	59	84	84	216	211	64	64	127	-163
6	1	207	223	24	24	18	-121	86	86	96	-99	65	65	127	-163
6	2	200	-25	26	26	74	102	88	88	233	224	66	66	127	-163
6	3	88	95	28	28	59	72	90	90	243	224	67	67	127	-163
6	4	224	127	30	30	75	-78	92	92	243	-256	68	68	127	-163
6	5	52	59	32	32	19	-112	94	94	179	164	69	69	127	-163
6	6	52	-70	34	34	104	-112	96	96	44	48	70	70	127	-163
6	7	86	-97	36	36	52	35	98	98	177	-149	71	71	127	-163
6	8	47	-52	38	38	49	48	100	100	252	-262	72	72	127	-163
6	9	43	54	40	40	33	-28	102	102	70	597	73	73	127	-163
6	10	129	140	42	42	38	44	104	104	70	-74	74	74	127	-163
7	0	70	-69	44	44	20	-28	106	106	122	-98	75	75	127	-163
7	1	70	67	46	46	29	42	108	108	69	68	76	76	127	-163
7	2	223	94	48	48	62	-57	110	110	122	115	77	77	127	-163
7	3	52	-30	50	50	0	90	112	112	172	-173	78	78	127	-163
7	4	202	-202	52	52	90	90	114	114	54	74	79	79	127	-163
7	5	89	97	54	54	200	192	116	116	20	24	80	80	127	-163
7	6	88	93	56	56	286	-279	118	118	24	28	81	81	127	-163
7	7	89	60	58	58	242	-273	120	120	256	-243	82	82	127	-163
7	8	87	-60	60	60	407	408	122	122	307	276	83	83	127	-163
7	9	45	-58	62	62	37	-20	124	124	332	-248	84	84	127	-163
7	10	92	83	64	64	122	-124	126	126	62	72	85	85	127	-163
8	0	52	-53	66	66	25	-6	128	128	105	103	86	86	127	-163
8	1	209	181	68	68	84	94	130	130	86	-90	87	87	127	-163
8	2	224	222	70	70	65	53	132	132	122	-122	88	88	127	-163
8	3	223	-233	72	72	140	121	134	134	81	87	89	89	127	-163
8	4	89	68	74	74	107	-102	136	136	81	87	90	90	127	-163
8	5	46	-39	76	76	128	-127	138	138	130	122	91	91	127	-163
8	6	44	-48	78	78	87	-43	140	140	89	54	92	92	127	-163
8	7	58	-58	80	80	89	48	142	142	69	54	93	93	127	-163
8	8	53	245	82	82	28	22	144	144	48	48	94	94	127	-163
8	9	202	-227	84	84	28	22	146	146	22	22	95	95	127	-163
8	10	52	-52	86	86	78	-4	148	148	81	-57	96	96	127	-163
9	0	86	102	88	88	89	-100	150	150	40	74	97	97	127	-163
9	1	226	-228	90	90	305	300	152	152	27	27	98	98	127	-163

Table VII.3.(cont.)

*1 h, k 5^o obs 5^o calo

3	0	16	7a	-66	a	205	-216	3	12	17	16	81	74							
1	81	69	141	184	18	28	24	4	27	-29	6	167	-159	1	99	-113	17	79	61	
2	92	71	139	-304	3	6	188	-184	7	51	-48	8	23	-37	2	23	-37	4	23	
3	144	130	4	131	-148	1	45	-12	6	143	-133	8	122	132	3	78	95	1	186	100
4	339	-360	5	180	175	2	93	58	8	77	-68	9	49	-42	5	75	80	2	171	115
5	295	-275	6	117	-137	3	48	-50	9	159	-136	10	47	-36	6	28	38	3	85	-139
6	191	169	7	237	256	4	71	-59	10	98	-87	11	51	42	7	62	-60	4	84	-138
7	22	-26	8	112	-105	5	81	-75	11	29	13	12	117	-118	8	32	34	5	99	127
8	172	165	9	55	-67	6	49	-47	12	84	70	13	21	-31	9	18	18	6	62	58
9	16	-33	10	50	58	7	132	113	13	38	-29	14	114	131	1	62	70	7	212	-257
10	153	-154	11	128	-133	8	126	-122	14	24	-24	15	23	28	2	47	-27	8	195	-224
11	40	29	12	118	124	9	29	-23	15	31	-25	16	31	-25	3	37	-44	9	220	218
12	29	12	13	41	26	10	22	51	16	18	-18	17	22	31	4	43	35	10	222	305
13	98	-101	14	67	-74	11	51	-51	17	28	31	18	28	31	5	79	-86	11	149	-154
14	120	123	15	25	-17	12	130	132	18	10	10	19	28	31	6	44	-25	12	222	-222
15	40	44	16	167	152	13	26	45	19	102	117	20	71	81	7	22	19	13	60	64
16	90	-116	17	314	-334	14	150	-143	20	232	-259	21	55	-56	8	43	58	14	172	177
17	46	-45	18	31	-30	15	34	-40	21	159	-174	22	62	-66	9	52	68	15	107	-102
18	3	1	19	187	-187	16	35	30	22	266	284	23	42	32	10	20	-36	16	68	-58
19	53	70	20	204	187	17	22	-24	23	86	89	24	54	59	11	10	-36	17	123	102
2	282	291	21	22	-30	18	22	-24	24	100	84	25	13	-15	12	10	10	18	60	64
3	309	314	22	22	-48	19	22	-24	25	41	32	26	102	108	13	164	238	19	153	194
4	200	-210	23	22	-93	20	173	-163	26	122	-102	27	100	108	14	173	250	20	147	160
5	143	-143	24	22	-69	21	86	-86	27	29	26	28	41	-55	15	151	-173	21	142	-126
6	161	164	25	22	131	22	152	137	28	102	-94	29	85	-85	16	43	-41	22	89	-122
7	51	-49	26	22	113	23	145	122	29	123	127	30	68	61	17	180	175	23	56	75
8	98	114	27	22	77	24	152	-136	30	66	61	31	66	66	18	149	135	24	145	177
9	88	93	28	22	-61	25	80	58	31	68	-77	32	37	-34	19	274	-281	25	42	40
10	106	-86	29	22	-38	26	62	69	32	120	-130	33	46	-25	20	274	-252	26	184	-187
11	87	-98	30	22	34	27	20	-35	33	70	83	34	29	-21	21	141	152	27	38	-40
12	56	-66	31	22	88	28	20	10	34	86	88	35	29	23	22	163	166	28	217	228
13	112	82	32	22	22	29	22	-19	35	123	-126	36	74	87	23	41	47	29	130	84
14	44	39	33	22	22	30	22	-28	36	271	263	37	85	85	24	86	76	30	106	113
15	19	-10	34	22	108	31	22	-69	37	229	20	38	29	20	25	61	-51	31	40	-28
16	46	-29	35	22	88	32	22	-28	38	145	-147	39	55	-68	26	14	-36	32	85	-65
17	126	140	36	22	93	33	22	-62	39	51	43	40	27	45	27	27	51	33	96	-107
18	144	-137	37	22	214	34	22	-62	40	108	104	41	27	45	28	114	161	34	96	-107
19	88	-85	38	22	-20	35	22	-62	41	56	45	42	27	45	29	114	161	35	96	-107
20	181	195	39	22	94	36	22	-62	42	101	-105	43	27	45	30	114	161	36	96	-107
21	207	-188	40	22	45	37	22	-62	43	56	45	44	27	45	31	114	161	37	96	-107
22	32	-20	41	22	11	38	22	-62	44	36	-31	45	27	45	32	114	161	38	96	-107
23	24	14	42	22	60	39	22	-62	45	33	-48	46	27	45	33	114	161	39	96	-107
24	63	72	43	22	48	40	22	-62	46	46	46	47	27	45	34	114	161	40	96	-107
25	109	-133	44	22	14	41	22	-62	47	11	23	48	27	45	35	114	161	41	96	-107
26	120	-147	45	22	15	42	22	-62	48	11	23	49	27	45	36	114	161	42	96	-107
27	39	-32	46	22	16	43	22	-62	49	11	23	50	27	45	37	114	161	43	96	-107
28	37	41	47	22	17	44	22	-62	50	11	23	51	27	45	38	114	161	44	96	-107
29	19	-7	48	22	18	45	22	-62	51	11	23	52	27	45	39	114	161	45	96	-107

14	130	-95	5	291	-303	4	12	-34	5	0	102	13	53	71	13	46	58
15	88	66	6	332	-348	5	39	-34	6	114	102	14	115	-136	14	10	58
16	111	-98	7	165	136	6	27	20	7	44	35	15	25	28	15	42	-44
17	20	-13	8	196	184	7	108	-105	8	106	-104	16	181	-157	16	42	-40
18	61	-63	9	165	-140	8	53	-69	9	45	-42	17	20	-12	17	42	-40
19	279	277	10	83	-90	9	137	131	10	59	-69	18	20	-12	18	42	-40
20	149	-129	11	14	-24	11	129	122	11	55	68	19	20	-12	19	42	-40
21	276	-232	12	34	19	12	143	-132	12	1	-32	20	20	-12	20	42	-40
22	78	-69	13	83	88	13	20	27	13	82	-70	21	20	-12	21	42	-40
23	166	141	14	65	-46	14	48	37	14	112	-114	22	20	-12	22	42	-40
24	159	119	15	192	159	15	52	-42	15	207	234	23	20	-12	23	42	-40
25	67	-48	16	97	83	16	13	-72	16	37	-86	24	20	-12	24	42	-40
26	48	-29	17	193	-163	17	66	-72	17	93	-86	25	20	-12	25	42	-40
27	99	87	18	193	-163	18	27	25	18	66	-69	26	20	-12	26	42	-40
28	152	-133	19	212	219	19	37	-30	19	42	-29	27	20	-12	27	42	-40
29	53	42	20	67	58	20	45	49	20	87	80	28	20	-12	28	42	-40
30	79	66	21	99	-85	21	83	-63	21	80	-90	29	20	-12	29	42	-40
31	34	33	22	88	-71	22	65	57	22	65	57	30	20	-12	30	42	-40
32	167	166	23	46	-41	23	105	-101	23	121	150	31	20	-12	31	42	-40
33	162	144	24	14	113	24	20	-27	24	56	70	32	20	-12	32	42	-40
34	255	-207	25	11	113	25	124	113	25	193	-189	33	20	-12	33	42	-40
35	65	-56	26	203	174	26	14	-25	26	45	-51	34	20	-12	34	42	-40
36	327	296	27	176	142	27	35	31	27	51	-33	35	20	-12	35	42	-40
37	79	-59	28	137	-124	28	60	75	28	33	8	36	20	-12	36	42	-40
38	242	-185	29	135	-121	29	54	-42	29	24	14	37	20	-12	37	42	-40
39	158	122	30	87	71	30	98	-91	30	47	39	38	20	-12	38	42	-40
40	62	53	31	236	213	31	42	-32	31	81	-74	39	20	-12	39	42	-40
41	39	29	32	174	-159	32	109	98	32	74	67	40	20	-12	40	42	-40
42	26	-18	33	171	-155	33	15	70	33	181	185	41	20	-12	41	42	-40
43	40	14	34	150	123	34	38	20	34	162	-172	42	20	-12	42	42	-40
44	76	-56	35	62	67	35	20	-70	35	21	25	43	20	-12	43	42	-40
45	23	-21	36	48	-63	36	64	-70	36	80	-85	44	20	-12	44	42	-40
46	183	179	37	31	-19	37	95	102	37	33	-43	45	20	-12	45	42	-40
47	217	-215	38	13	39	38	20	-31	38	71	-64	46	20	-12	46	42	-40
48	86	-78	39	14													

Table VII.4a

h k, l 2.5° obs 2.5° calc for Paraquat PdCl₄

0	0	0	0	12	0	12	0	1	12	7	335	298	6	134	129	0	3	5	702		
8	1186	1253	0	511	503	1	156	145	11	87	94	9	265	239	8	174	183	0	393	702	
10	815	685	2	733	703	1	156	145	11	13	94	11	195	213	10	200	177	2	620	688	
12	1172	1102	4	536	614	3	162	151	0	604	552	13	127	171	0	2	14	4	645	635	
14	303	406	6	461	471	5	106	116	2	485	417	15	100	133	0	257	226	4	569	563	
16	506	556	8	465	520	7	123	103	4	475	442	0	598	668	2	283	249	8	515	491	
0	4	1056	862	10	412	457	0	1	6	495	444	0	866	1148	4	241	220	10	442	411	
2	759	622	0	14	293	0	995	1085	8	398	364	2	742	745	6	245	221	12	312	352	
4	739	628	0	483	488	2	658	624	10	294	303	4	701	745	8	237	191	14	299	287	
6	639	544	2	467	513	4	717	754	0	281	276	6	554	584	1	162	-122	1	122	116	
8	303	326	4	424	469	6	705	715	2	278	297	8	556	632	3	155	-130	3	162	143	
10	368	371	6	421	446	8	538	562	4	238	222	10	297	346	5	134	-111	5	99	95	
12	220	258	8	389	393	10	432	424	6	278	300	12	323	436	7	110	-104	7	115	91	
14	220	258	10	412	457	12	445	508	8	208	193	14	229	-183	0	425	351	0	435	500	
16	220	258	12	145	163	14	249	278	2	629	-607	16	168	-160	2	347	282	2	240	254	
0	4	136	102	4	145	163	1	495	-438	4	532	517	4	127	-124	4	432	382	4	301	374
2	211	230	6	649	660	5	353	-365	6	261	200	6	754	764	8	356	420	6	272	287	
4	307	212	8	789	776	7	329	-334	8	406	338	8	513	538	10	314	315	8	172	181	
6	154	122	10	646	627	9	333	-365	10	503	435	10	583	645	12	330	336	10	243	253	
8	162	105	12	409	372	11	184	-223	12	237	200	12	470	503	14	314	300	12	115	125	
10	145	100	14	440	445	13	156	-197	14	149	-212	14	459	498	16	211	180	14	207	177	
12	170	111	16	294	310	0	405	337	16	694	-696	16	307	360	18	263	222	16	107	124	
0	162	244	0	304	259	2	358	363	2	168	-146	2	323	380	4	172	153	4	162	165	
2	775	888	2	115	135	4	512	506	4	403	514	4	323	-283	6	693	-782	6	522	633	
4	525	470	4	779	751	6	189	216	6	155	-169	6	220	-180	8	247	-249	8	493	481	
6	545	523	6	162	150	8	199	194	8	430	435	8	265	-253	10	579	-618	10	654	698	
8	458	403	8	307	267	10	204	195	10	446	418	10	241	-216	12	162	-171	12	322	312	
10	509	489	10	174	198	12	253	-209	12	418	388	12	127	-169	14	325	-341	14	507	573	
12	294	266	12	87	120	14	189	-128	14	282	234	14	100	-127	16	167	-234	16	304	320	
14	373	363	14	493	655	16	137	-156	16	295	198	16	344	367	18	922	1128	18	295	353	
0	691	692	0	554	534	0	359	318	0	134	173	0	418	394	2	767	739	2	167	108	
2	377	412	2	554	534	2	425	378	2	63	-60	2	134	173	4	604	564	4	295	258	
4	582	553	4	567	491	4	394	288	4	110	-102	4	134	136	6	579	636	6	162	178	
6	419	432	6	405	390	6	333	337	6	703	695	6	307	289	8	455	463	8	57	57	
8	435	426	8	278	294	8	230	240	8	751	703	8	224	219	10	414	421	10	441	422	
10	322	297	10	554	711	10	284	256	10	560	530	10	273	224	12	152	-147	12	521	475	
12	328	347	12	603	619	12	179	193	12	437	449	12	241	211	14	275	-232	14	427	391	
14	235	216	14	626	564	14	194	157	14	294	362	14	195	171	16	91	-109	16	437	421	
0	432	423	0	528	474	0	194	187	0	383	444	0	179	149	2	107	-103	2	319	329	
2	397	369	2	402	393	2	130	140	2	364	380	2	241	252	4	115	-126	4	319	329	
4	154	139	4	358	344	4	115	139	4	386	357	4	195	195	6	115	-126	6	253	259	
6	507	552	6	260	272	6	260	272	6	260	272	6	260	272	8	260	272	8	260	272	
8	238	263	8	260	272	8	260	272	8	260	272	8	260	272	10	260	272	10	260	272	
10	214	225	10	260	272	10	260	272	10	260	272	10	260	272	12	260	272	12	260	272	
12	214	225	12	260	272	12	260	272	12	260	272	12	260	272	14	260	272	14	260	272	
14	140	158	14	260	272	14	260	272	14	260	272	14	260	272	16	260	272	16	260	272	

Table VII.5.

%h k, l 2.5^o obs 2.5^o calc for the 25% (Cu(I)) m.v.compound

12	0	0	15	417	446	45	13	3	7	15	326	355	135
16	3796	3638	17	389	260	135	13	367	356	5	11		
0	3132	3194	21	309	283	135	17	225	270	13	455	381	-45
2	449	388	1	9			19	300	190	15	462	452	-135
6	318	514	11	183	164	-45	3	9		6	6		
10	455	577	13	167	123	-135	11	474	467	6	973	1131	90
14	792	840	15	237	79	-45	3	7		10	410	242	90
4	3281	3141	19	225	181	-45	21	290	127	22	309	213	0
8	2196	1996	2	2			4	4		14	474	161	-90
12	2715	2609	4	497	483	0	4	3744	3756	18	540	583	-90
16	1803	1888	6	694	930	-90	6	2832	2832	6	8		
20	1029	1013	8	730	1056	-180	8	3365	3403	8	343	321	-90
24	956	844	10	237	269	-90	10	1391	1720	10	690	644	180
10	635	779	12	259	479	0	12	2037	2243	14	410	279	0
14	513	530	14	734	680	90	14	590	688	16	382	186	90
18	318	288	22	410	325	-90	16	1514	1770	6	10		
8	3766	3724	18	270	122	90	18	545	559	10	580	479	-90
12	2234	2196	4	2			20	1318	1330	14	212	120	90
16	988	886	6	585	650	-90	4	6		18	212	76	90
20	1179	1056	8	183	63	0	6	270	193	6	12		
24	1106	1023	10	938	905	90	10	237	44	12	335	383	-90
14	857	830	10	519	380	-180	12	956	924	6	14		
18	1342	1490	14	248	273	180	8	8		14	270	612	-90
20	718	933	6	6			8	3575	3504	16	248	151	-180
18	2483	2430	8	270	384	0	12	1753	1830	18	343	619	90
16	1342	1490	10	237	210	-90	16	1465	1457	7	9		
14	857	830	14	917	982	90	20	1395	1340	11	259	295	-135
12	2483	2430	18	359	583	-90	24	745	764	15	167	294	45
10	635	779	8	2			4	10		7	11		
8	2196	1996	8	167	113	90	10	318	297	15	270	235	-45
6	318	514	10	1093	1086	0	12	640	604	8	8		
4	3281	3141	14	248	501	180	14	367	216	8	2084	2194	0
2	449	388	16	212	80	-90	18	335	132	12	1253	1384	-180
0	3132	3194	16	10			4	12		16	1732	1754	0
12	1150	1188	12	248	382	180	12	1595	1565	20	1333	1240	180
16	1498	1494	14	237	217	-90	14	343	377	8	10		
14	449	478	18	617	593	90	16	1224	1166	10	813	960	-180
10	309	44	2	14			18	374	333	8	12		
8	2196	1996	14	524	528	-90	20	892	901	12	1230	1275	0
6	318	514	7	5			24	813	688	16	1264	1077	180
4	3281	3141	9	631	544	45	4	16		20	902	762	0
2	449	388	11	130	142	135	16	944	1006	24	635	556	-180
0	3132	3194	13	485	485	45	5	7		8	14		
12	1150	1188	13	604	617	-45	9	270	161	14	318	363	-180
16	1498	1494	15	309	337	-135	11	212	158	8	16		
14	449	478	21	318	197	135	13	270	219	16	682	729	0
10	309	44	3	7			15	183	1	20	799	785	-180
8	2196	1996	9	367	301	-135	5	9		10	10		
6	318	514	11	686	634	135	11	580	684	12	335	460	0
4	3281	3141	13	502	485	45	13	502	485	10	12		
2	449	388								18	212	187	180
0	3132	3194								10	14		
										18	443	512	-90
										12	12		
										12	1842	1861	180
										18	335	178	-90
										20	535	577	-180
										12	16		
										16	622	829	-180

Table VII. 6.

Observed and Calculated Structure Factors for $pcCu_2Cl_6$ (see p. 28)

- *0,0;l=2,4,6,8;(58,72)(24,38)(23,19)(30,-27)
- *0,1;l=1 to8;(42,53)(30,-22)(32,26)(46,36)(10,-8)(25,-20)(45,-40)(4,4)
- *0,2;l=0 to8;(51,66)(4,1)(41,39)(37,28)(21,-15)(24,-21)(17,-12)(31,-26)
(4,4)
- *0,3;l=1 to8;(6,7)(13,-14)(16,-14)(8,-11)(40,-32)(49,-49)(12,10)(32,32)
- *0,4;l=0 to6;(38,-33)(10,8)(4,1)(35,-29)(25,-20)(59,-59)(13,-14)
- *0,5;l=1 to6;(12,-8)(7,8)(5,9)(59,-61)(7,-6)(51,-55)
- *0,6;l=0 to5;(22,20)(18,18)(12,-13)(20,-23)(11,11)(40,-34)
- *0,7;l=1;(24,-28)
- *1,-7;l=0 to3;(16,-18)(14,-16)(2,4)(16,14)
- *1,-6;l=1 to5;(13,15)(15,-16)(31,33)(2,2)(14,12)
- *1,-5;l=0 to5;(36,47)(3,-7)(45,49)(37,-36)(5,0)(8,-7)
- *1,-4;l=1 to7;(42,51)(23,-21)(45,47)(20,-40)(3,-5)(22,20)(29,-24)
- *1,-3;l=1 to7;(32,24)(52,58)(26,-27)(24,20)(22,-22)(12,10)(18,18)
- *1,-2;l=1 to8;(33,36)(52,-56)(36,-27)(8,5)(3,-2)(7,-9)(37,-34)(6,-7)
- *1,-1;l=0 to8;(31,35)(13,-12)(15,11)(7,2)(25,-16)(43,-40)(20,-17)
(37,-31)(12,11)
- *1,0;l=1 to8;(4,-2)(47,-48)(37,-31)(28,-24)(45,-40)(64,-66)(19,15)(12,11)
- *1,1;l=0 to8;(36,-48)(19,-13)(4,-3)(40,-35)(45,-38)(65,-68)(25,-23)
(66,-69)(36,-31)
- *1,2;l=1 to7;(10,-8)(8,6)(13,8)(49,-47)(35,-26)(52,-50)(23,21)
- *1,3;l=0 to7;(4,2)(44,44)(7,-10)(10,-7)(17,-13)(26,-20)(3,4)(12,-8)
- *1,4;l=1 to4;(36,-38)(29,25)(44,-43)(4,8)
- *1,5;l=0 to4;(35,-36)(18,16)(48,-50)(54,53)(18,-13)
- *1,6;l=1;(55,-59)
- *2,-3;l=1 to7;(12,-13)(44,-46)(48,-52)(35,-34)(30,-30)(54,-56)(15,14)
- *2,-2;l=0 to8;(31,-34)(35,-36)(11,-10)(31,-27)(39,-35)(46,-41)(19,-18)
(50,-51)(35,33)
- *2,-1;l=1 to8;(4,-6)(3,4)(5,2)(14,-14)(43,-38)(24,-21)(30,28)(21,-17)
- *2,0;l=2 to8;(3,-3)(24,16)(37,-34)(13,10)(15,-11)(14,13)(42,42)
- *2,1;l=1 to7;(31,-35)(46,47)(48,-50)(10,11)(1,0)(44,40)(5,-1)
- *2,2;l=0 to7;(16,-23)(34,34)(53,-59)(63,68)(29,-21)(44,38)(7,4)(14,12)

Table VII.6. (cont.)

*2,3;l=1to5;(36,-33)(20,-19)(38,-37)(31,31)(3,2)
*2,4;l=0to5;(19,-19)(1,2)(19,-13)(51,51)(12,10)(16,14)
*2,5;l=1to3;(13,11)(17,19)(34,36)
*2,6;l=0,1;(4,-2)(15,18)
*3,-5;l=0to6;(13,-8)(33,-35)(15,-15)(13,-14)(19,-23)(7,-8)(2,-5)
*3,-4;l=1to7;(6,-7)(1,-1)(1,1)(23,20)(31,-34)(10,7)(20,22)
*3,-3;l=0to7;(17,-23)(21,22)(2,0)(28,27)(37,-36)(38,33)(17,-16)(33,32)
*3,-2;l=1to7;(18,-21)(42,42)(34,-31)(20,-16)(35,-30)(56,61)(12,10)
*3,-1;l=2to7;(50,-41)(55,48)(29,-21)(43,37)(3,-2)(29,25)
*3,0;l=1to7;(20,13)(24,21)(28,-24)(24,23)(9,-6)(3,-1)(5,-5)
*3,1;l=0to6;(30,35)(14,10)(3,6)(48,46)(26,18)(17,15)(5,3)
*3,2;l=1to6;(41,49)(22,20)(38,36)(50,52)(4,5)(21,-19)
*3,3;l=0to6;(19,16)(9,11)(47,56)(33,35)(29,26)(21,20)(2,-2)
*3,4;l=1to4;(10,10)(27,32)(36,41)(11,12)
*3,5;l=0,1;(4,2)(19,26)
*4,-4;l=2to6;(6,-8)(14,14)(14,-16)(25,22)(9,-10)
*4,-2;l=1to7;(11,10)(27,24)(21,21)(20,16)(11,11)(6,-8)(9,5)
*4,-1;l=1to7;(64,67)(19,16)(28,21)(42,39)(13,13)(28,-28)(16,-14)
*4,0;l=0to7;(22,29)(4,-4)(52,47)(35,31)(14,12)10,10)(11,-11)(13,-16)
*4,1;l=1to6;(4,4)(21,19)(32,27)(11,7)(14,-16)(15,-13)
*4,2;l=0to6;(11,-15)(16,23)(8,10)(12,-11)(6,5)(38,-34)(3,-2)
*4,3;l=1to5;(2,4)(9,-9)(4,3)(42,-44)(19,20)
*4,4;l=0to3;(13,-17)(3,6)10,-12)(24,-26)
*4,5;l=1;(12,17)

The F values are 0.5x absolute value in electrons.

Table VII.7.

Structure Factors (x0.1) for Cobalthexammine Chlorocuprate(I).

The calculated values are based on the parameters in Table V.2. but related to a non-centrosymmetric origin at $\frac{1}{8} \frac{1}{8} \frac{1}{8}$. The phase can be calculated from the relationship:

$$\alpha = (x+2)\pi/4 \pm \pi/2, \text{ where the sign is given below,}$$

and $x = h+k+l \pmod{4}$. The table is arranged as:

$$*h,k;(1, F_{obs}, F_{calc})_1 (1, F_{obs}, F_{calc})_2 \text{ \&c.}$$

-
- *0,0;(8,18,25)(12,109,-98)(16,117,111)(24,29,24)(28,46,-40)(
 - *0,2;(6,33,32)(10,33,16)(14,34,25)(18,10,13)
 - *0,4;(4,118,166)(8,73,-73)(12,124,128)(16,79,-76)(20,59,62)(24,33,-32)
(28,35,35)(32,15,-16)
 - *0,6;(8,13,11)
 - *0,8;(4,77,-83)(8,142,185)(12,76,-74)(16,68,70)(20,66,-65)(24,54,55)
(28,11,-12)
 - *0,10;(2,28,-13)(6,53,75)(14,14,-7)(18,12,-6)(22,19,17)(26,18,22)
 - *0,12;(4,137,135)(8,79,-85)(12,77,79)(16,39,-35)(20,43,41)(24,26,-23)
 - *0,14;(2,37,38)(10,16,-14)(14,22,22)(18,18,22)
 - *0,16;(4,76,-69)(8,67,67)(12,41,-43)(16,54,59)(20,27,-25)(24,23,20)
 - *0,18;(2,11,17)(6,11,-7)(10,17,14)(14,17,19)(22,12,-6)
 - *0,20;(4,36,-34)(8,56,66)(12,25,-23)(16,21,16)
 - *0,22;(6,15,14)(10,26,28)(18,11,-7)
 - *0,24;(4,36,-34)(8,56,66)(12,25,-23)(16,21,16)
 - *0,26;(10,13,2)
 - *0,28;(4,34,31)(12,43,30)(16,19,-21)
 - *0,32;(4,14,-14)
 - *1,1;(9,34,30)(11,15,-12)(13,36,-28)(15,28,25)(17,22,-25)(19,12,12)
 - *1,3;(3,19,-27)(5,37,30)(7,52,46)(9,18,19)(11,16,-13)(13,27,-21)
(15,28,-21)(17,11,10)(19,11,-7)(21,11,10)
 - *1,5;(3,42,-42)(5,31,-29)(7,29,37)(9,14,-15)(11,26,21)(13,13,-14)
(15,10,3)(17,12,11)
 - *1,7;(3,48,-41)(5,20,5)(13,20,-19)(17,19,10)(21,15,14)
 - *1,9;(7,45,48)(9,31,-36)(11,14,15)(13,18,-18)(15,18,-21)(17,11,14)
 - *1,11;(3,13,-5)(5,23,-20)(7,38,-41)(9,17,21)(11,18,-20)(13,27,30)
(15,23,24)

Table VII.7.(cont.)

- *1,13;(3,28,24)(5,10,0)(7,18,-10)(9,14,10)(11,20,-4)(15,11,12)(17,20,-22)
- *1,15;(3,30,26)(5,15,18)(7,11,-10)(9,14,-5)(11,15,-3)(13,13,16)
(15,19,-28)(17,13,-3)
- *1,17;(3,11,11)(5,13,-11)(7,21,22)(13,15,3)(15,17,18)
- *1,19;(3,13,-16)(9,10,-13)
- *1,21;(3,10,6)(7,14,11)(9,12,-12) *1,25;(7,16,15)
- *2,2;(4,41,42)(6,26,27)(12,11,11)(14,26,-23)(16,3,29)(22,25,24)
- *2,4;(6,35,35)(8,57,-67)(12,29,28)(14,14,15)(16,15,14)(18,22,22)
(22,12,9)(28,11,-15)
- *2,6;(6,35,32)(8,31,36)(12,23,27)(14,24,-22)(24,12,12)
- *2,8;(4,44,16)(6,22,8)(10,54,63)(12,11,5)(22,12,5)
- *2,10;(4,17,19)(8,39,29)(12,22,4)(14,15,17)(20,19,19)
- *2,12;(4,22,-11)(6,16,5)(10,32,40)(14,11,7)
- *2,14;(6,26,-27)(10,14,13)(12,15,19)(18,15,14)
- *2,16;(8,15,16)(10,17,-20)(14,11,12)
- *2,18;(4,19,16)(14,11,17)(16,19,19)
- *2,20;(6,13,12)(10,12,2) *2,22;(4,14,14)
- *2,24;(6,12,11) *2,30;(6,14,-14)
- *3,3;(3,53,56)(9,15,-12)(17,16,-11)(19,15,12)(23,20,20)
- *3,5;(5,57,55)(7,35,9)(9,10,4)(11,36,-30)(13,16,-15)(15,28,-28)
(19,24,7)(21,16,21)
- *3,7;(5,45,-47)(9,29,-31)(11,23,20)(13,12,0)(15,18,22)(19,16,-7)
- *3,9;(5,12,13)(11,11,4)(13,12,8)(15,12,-10)(19,16,-11)
- *3,11;(5,40,-39)(7,17,8)(9,10,11)(11,14,-11)(15,22,23)
- *3,13;(5,14,8)(7,16,-18)(9,14,19)(13,16,15)(19,12,-3)
- *3,15;(5,24,-16)(9,9,5)(11,20,14)
- *3,17;(11,11,-6)(17,16,15)
- *3,19;(5,28,27)(7,20,18)(9,18,-14) *3,21;(1,11,10)(5,13,0)
- *4,4;(4,104,-116)(6,70,71)(8,123,151)(10,30,-25)(12,81,-81)(14,48,-49)
(16,81,83)(18,37,37)(20,51,-43)(22,10,9)(24,43,43)(28,19,-19)
- *4,6;(6,13,5)(8,26,22)(10,23,12)(12,43,-48)(14,14,16)(16,11,-11)
(20,16,19)
- *4,8;(6,29,-30)(8,103,-110)(10,11,15)(12,83,90)(14,18,11)(16,47,-48)
(18,10,12)(20,63,61)(22,15,-13)(24,29,-28)(28,22,27)
- *4,10;(6,32,36)(10,28,27)(12,11,11)(14,20,22)
- *4,12;(6,30,-30)(8,85,94)(12,63,-61)(14,25,12)(16,65,66)(18,22,-25)
(20,30,-30)(24,35,35)(26,12,11)(28,27,-20)

Table VII.7.(cont.)

- *4,14;(6,16,16)(8,20,-16)(10,16,8)(12,37,43)(20,15,-20)(24,10,-7)
- *4,16;(6,10,-6)(8,46,-44)(12,65,67)(16,42,-41)(18,12,-11)(20,34,36)
(24,20,-20)
- *4,18;(12,19,-14)(20,11,8)
- *4,20;(8,61,58)(12,30,-30)(14,13,-7)(16,36,38)(20,26,-24)(24,22,22)
- *4,22;(8,14,7)(12,11,-5) *4,28;(8,22,26)(12,22,-22)
- *4,24;(8,30,-30)(12,33,35)(16,20,-19)(20,21,22)
-
- *5,5;(5,22,24)(9,10,7)(11,35,-31)(13,27,-27)
- *5,7;(11,18,-19)(17,12,7)(19,11,-12)
- *5,9;(9,26,-18)(11,26,27)(17,13,-9)
- *5,11;(7,12,4)(9,26,25)(13,16,11)(19,11,-12)(21,11,5)
- *5,13;(11,22,27)(15,10,12) *5,15;(13,9,11)
- *5,17;(7,18,19)(9,17,-19)(11,9,8)
-
- *6,6;(6,14,-19)(8,28,28)(10,18,-18)(12,19,20)(14,25,20)(16,12,11)
- *6,8;(14,16,18)(16,15,-15)(18,23,27)(24,9,9)
- *6,10;(8,23,28)(10,11,9)(12,19,21)(16,20,13)
- *6,12;(10,16,-11)(12,21,20)(14,11,10)(16,16,16)(18,13,14)
- *6,14;(8,15,14)(14,12,-4)(18,11,-13)
- *6,16;(8,13,9)(10,28,32)(12,11,-5)(14,10,10)(18,10,-7)
- *6,18;(12,10,6)(14,12,-11) *6,20;(10,17,19)(12,10,-12)
-
- *7,7;(7,41,41)(13,28,-27)
- *7,9;(9,30,-31)(11,17,17)(13,15,14)(15,15,-16)(17,27,31)19,13,-10)
- *7,11;(15,12,15)(17,17,-19) *7,13;(9,20,24)(11,15,-14)(13,18,18)
- *7,15;(9,15,-14)(11,9,2)(15,10,10)
- *7,17;(9,19,10)(11,12,-6) *7,19;(9,16,-20)(11,18,20)
-
- *8,8;(8,83,90)(12,50,-52)(16,66,71)(20,27,-25)(24,22,22)(28,22,-24)
- *8,10;(14,19,16)(18,27,30)
- *8,12;(12,78,80)(16,51,-51)(20,40,46) *8,14;(10,26,32)
- *8,16;(12,51,-46)(14,12,10)(16,47,44)(20,27,-28)
- *8,18;(10,20,17) *8,20;(12,41,43)(16,28,-28)(20,20,22)
- *8,22;(10,12,-12)(14,12,13)(18,10,12)
- *8,24;(12,24,-25)(14,11,-7)(16,31,36)(20,15,-12)