

EXISTENCE OF GLOBAL WEAK SOLUTIONS TO COMPRESSIBLE ISENTROPIC FINITELY EXTENSIBLE NONLINEAR BEAD-SPRING CHAIN MODELS FOR DILUTE POLYMERS

JOHN W. BARRETT AND ENDRE SÜLI

ABSTRACT. We prove the existence of global-in-time weak solutions to a general class of models that arise from the kinetic theory of dilute solutions of nonhomogeneous polymeric liquids, where the polymer molecules are idealized as bead-spring chains with finitely extensible nonlinear elastic (FENE) type spring potentials. The class of models under consideration involves the unsteady, compressible, isentropic, isothermal Navier–Stokes system in a bounded domain Ω in \mathbb{R}^d , $d = 2$ or 3 , for the density ρ , the velocity \underline{u} and the pressure p of the fluid, with an equation of state of the form $p(\rho) = c_p \rho^\gamma$, where c_p is a positive constant and $\gamma > \frac{3}{2}$. The right-hand side of the Navier–Stokes momentum equation includes an elastic extra-stress tensor, which is the sum of the classical Kramers expression and a quadratic interaction term. The elastic extra-stress tensor stems from the random movement of the polymer chains and is defined through the associated probability density function that satisfies a Fokker–Planck-type parabolic equation, a crucial feature of which is the presence of a centre-of-mass diffusion term. We require no structural assumptions on the drag term in the Fokker–Planck equation; in particular, the drag term need not be corotational. With a nonnegative initial density $\rho_0 \in L^\infty(\Omega)$ for the continuity equation; a square-integrable initial velocity datum \underline{u}_0 for the Navier–Stokes momentum equation; and a nonnegative initial probability density function ψ_0 for the Fokker–Planck equation, which has finite relative entropy with respect to the Maxwellian M associated with the spring potential in the model, we prove, *via* a limiting procedure on certain discretization and regularization parameters, the existence of a global-in-time bounded-energy weak solution $t \mapsto (\rho(t), \underline{u}(t), \psi(t))$ to the coupled Navier–Stokes–Fokker–Planck system, satisfying the initial condition $(\rho(0), \underline{u}(0), \psi(0)) = (\rho_0, \underline{u}_0, \psi_0)$.

Keywords: Kinetic polymer models, FENE chain, compressible Navier–Stokes–Fokker–Planck system, variable density, nonhomogeneous dilute polymer

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1. INTRODUCTION

This paper establishes the existence of global-in-time weak solutions to a large class of bead-spring chain models with finitely extensible nonlinear elastic (FENE) type spring potentials, — a system of nonlinear partial differential equations that arises from the kinetic theory of dilute polymer solutions. The solvent is a compressible, isentropic, viscous, isothermal Newtonian fluid confined to a bounded Lipschitz domain $\Omega \subset \mathbb{R}^d$, $d = 2$ or 3 , with boundary $\partial\Omega$. For the sake of simplicity of presentation, we shall suppose that Ω has a ‘solid boundary’ $\partial\Omega$; the velocity field \underline{u} will then satisfy the no-slip boundary condition $\underline{u} = \underline{0}$ on $\partial\Omega$. The equations of continuity and balance of linear momentum have the form of the compressible Navier–Stokes equations (cf. Lions [45], Feireisl [31], Novotný & Straškraba [58], or Feireisl & Novotný [32]) in which the *elastic extra-stress* tensor $\underline{\tau}$ (i.e., the polymeric part of the Cauchy stress tensor) appears as a source term in the conservation of momentum equation:

Given $T \in \mathbb{R}_{>0}$, find $\rho : (\underline{x}, t) \in \Omega \times [0, T] \mapsto \rho(\underline{x}, t) \in \mathbb{R}$ and $\underline{u} : (\underline{x}, t) \in \bar{\Omega} \times [0, T] \mapsto \underline{u}(\underline{x}, t) \in \mathbb{R}^d$ such that

$$\begin{aligned}
 (1.1a) \quad & \frac{\partial \rho}{\partial t} + \nabla_x \cdot (\underline{u} \rho) = 0 && \text{in } \Omega \times (0, T], \\
 (1.1b) \quad & \rho(\underline{x}, 0) = \rho_0(\underline{x}) && \forall \underline{x} \in \Omega, \\
 (1.1c) \quad & \frac{\partial(\rho \underline{u})}{\partial t} + \nabla_x \cdot (\rho \underline{u} \otimes \underline{u}) - \nabla_x \cdot \underline{\underline{S}}(\underline{u}, \rho) + \nabla_x p(\rho) = \rho \underline{f} + \nabla_x \cdot \underline{\underline{\tau}} && \text{in } \Omega \times (0, T], \\
 (1.1d) \quad & \underline{u} = 0 && \text{on } \partial\Omega \times (0, T], \\
 (1.1e) \quad & (\rho \underline{u})(\underline{x}, 0) = (\rho_0 \underline{u}_0)(\underline{x}) && \forall \underline{x} \in \Omega.
 \end{aligned}$$

It is assumed that each of the equations above has been written in its nondimensional form; ρ denotes a nondimensional solvent density, \underline{u} is a nondimensional solvent velocity, defined as the velocity field scaled by the characteristic flow speed U_0 . Here $\underline{\underline{S}}(\underline{u}, \rho)$ is the Newtonian part of the viscous stress tensor defined by

$$(1.2) \quad \underline{\underline{S}}(\underline{u}, \rho) := \mu^S(\rho) \left[\underline{\underline{D}}(\underline{u}) - \frac{1}{d} (\nabla_x \cdot \underline{u}) \underline{I} \right] + \mu^B(\rho) (\nabla_x \cdot \underline{u}) \underline{I},$$

where \underline{I} is the $d \times d$ identity tensor, $\underline{\underline{D}}(\underline{u}) := \frac{1}{2} (\nabla_x \underline{u} + (\nabla_x \underline{u})^T)$ is the rate of strain tensor, with $(\nabla_x \underline{u})(\underline{x}, t) \in \mathbb{R}^{d \times d}$ and $(\nabla_x \underline{u})_{ij} = \frac{\partial u_i}{\partial x_j}$. The shear viscosity, $\mu^S(\cdot) \in \mathbb{R}_{>0}$, and the bulk viscosity, $\mu^B(\cdot) \in \mathbb{R}_{\geq 0}$, of the solvent are both scaled and, generally, density-dependent. In addition, p is the nondimensional pressure satisfying the isentropic equation of state

$$(1.3) \quad p(\rho) = c_p \rho^\gamma,$$

where $c_p \in \mathbb{R}_{>0}$ and the constant γ is such that $\gamma > \frac{3}{2}$.

Remark 1.1. *For the sake of simplicity of the exposition we focus here on the classical isentropic equation of state (1.3). Our analysis applies, without alterations, to some other familiar equations of state, such as the (Kirkwood-modified) Tait equation of state*

$$p(\rho) = A_0 \left(\frac{\rho}{\rho_*} \right)^\gamma - A_1,$$

where $\gamma > \frac{3}{2}$, A_0 and A_1 are constants, $A_0 - A_1 = p_*$ is the equilibrium reference pressure, and ρ_* is the equilibrium reference density. For distilled water, $\gamma \in [5.16, 7.11]$ (the actual value being dependent on the ambient temperature); for carbon tetrachloride (at 30°C) $\gamma = 12.54$, for glycerine (at 20°C) $\gamma = 9.80$; cf. Table 1.1 on p.4 of [56] for values of γ , A_0 and A_1 for other substances.

On the right-hand side of (1.1c), \underline{f} is the nondimensional density of body forces and $\underline{\underline{\tau}}$ denotes the elastic extra-stress tensor. In a *bead-spring chain model*, consisting of $K + 1$ beads coupled with K elastic springs to represent a polymer chain, $\underline{\underline{\tau}}$ is defined by a version of the *Kramers expression* depending on the probability density function ψ of the (random) conformation vector $\underline{q} := (\underline{q}_1^T, \dots, \underline{q}_K^T)^T \in \mathbb{R}^{Kd}$ of the chain (see equation (1.11) below), with \underline{q}_i representing the d -component conformation/orientation vector of the i th spring. The Kolmogorov equation satisfied by ψ is a second-order parabolic equation, the Fokker–Planck equation, whose transport coefficients depend on the velocity field \underline{u} , and the hydrodynamic drag coefficient appearing in the Fokker–Planck equation is, generally, a nonlinear function of the density ρ . The domain D of admissible conformation vectors $D \subset \mathbb{R}^{Kd}$ is a K -fold Cartesian product $D_1 \times \dots \times D_K$ of balanced convex open sets $D_i \subset \mathbb{R}^d$, $i = 1, \dots, K$; the term *balanced* means that $\underline{q}_i \in D_i$ if, and only if, $-\underline{q}_i \in D_i$. Hence, in particular, $\underline{0} \in D_i$, $i = 1, \dots, K$. Typically D_i is the whole of \mathbb{R}^d or a bounded open d -dimensional ball centred at the origin $\underline{0} \in \mathbb{R}^d$ for each $i = 1, \dots, K$. When $K = 1$, the model is referred to as the *dumbbell model*.

Let $\mathcal{O}_i \subset [0, \infty)$ denote the image of D_i under the mapping $\underline{q}_i \in D_i \mapsto \frac{1}{2} |\underline{q}_i|^2$, and consider the *spring potential* $U_i \in C^1(\mathcal{O}_i; \mathbb{R}_{\geq 0})$, $i = 1, \dots, K$. Clearly, $0 \in \mathcal{O}_i$. We shall suppose that $U_i(0) = 0$

and that U_i is unbounded on \mathcal{O}_i for each $i = 1, \dots, K$. The elastic spring-force $\underline{F}_i : D_i \subseteq \mathbb{R}^d \rightarrow \mathbb{R}^d$ of the i th spring in the chain is defined by

$$(1.4) \quad \underline{F}_i(\underline{q}_i) := U'_i(\tfrac{1}{2}|\underline{q}_i|^2) \underline{q}_i, \quad i = 1, \dots, K.$$

The partial Maxwellian M_i , associated with the spring potential U_i , is defined by

$$M_i(\underline{q}_i) := \frac{1}{\mathcal{Z}_i} e^{-U_i(\frac{1}{2}|\underline{q}_i|^2)}, \quad \mathcal{Z}_i := \int_{D_i} e^{-U_i(\frac{1}{2}|\underline{q}_i|^2)} d\underline{q}_i, \quad i = 1, \dots, K.$$

The (total) Maxwellian in the model is then

$$(1.5) \quad M(\underline{q}) := \prod_{i=1}^K M_i(\underline{q}_i) \quad \forall \underline{q} := (\underline{q}_1^T, \dots, \underline{q}_K^T)^T \in D := \bigtimes_{i=1}^K D_i.$$

Observe that, for $i = 1, \dots, K$,

$$(1.6a) \quad M(\underline{q}) \nabla_{\underline{q}_i} [M(\underline{q})]^{-1} = -[M(\underline{q})]^{-1} \nabla_{\underline{q}_i} M(\underline{q}) = \nabla_{\underline{q}_i} (U_i(\tfrac{1}{2}|\underline{q}_i|^2)) = U'_i(\tfrac{1}{2}|\underline{q}_i|^2) \underline{q}_i,$$

and, by definition,

$$(1.6b) \quad \int_D M(\underline{q}) d\underline{q} = 1.$$

Example 1.1. In the Hookean dumbbell model $K = 1$, and the spring force is defined by $\underline{F}(\underline{q}) = \underline{q}$, with $\underline{q} \in D = \mathbb{R}^d$, corresponding to $U(s) = s$, $s \in \mathcal{O} = [0, \infty)$. More generally, in a Hookean bead-spring chain model, $K \geq 1$, $\underline{F}_i(\underline{q}_i) = \underline{q}_i$, corresponding to $U_i(s) = s$, $i = 1, \dots, K$, and D is the Cartesian product of K copies of \mathbb{R}^d . The associated Maxwellian is

$$M(\underline{q}) = M_1(\underline{q}_1) \cdots M_K(\underline{q}_K) = \frac{1}{\mathcal{Z}} e^{-\frac{1}{2}|\underline{q}|^2},$$

with $|\underline{q}|^2 := |\underline{q}_1|^2 + \cdots + |\underline{q}_K|^2$ and $\mathcal{Z} := \mathcal{Z}_1 \cdots \mathcal{Z}_K = (2\pi)^{Kd/2}$. Hookean dumbbell and Hookean bead-spring chain models are physically unrealistic as they admit arbitrarily large extensions. \diamond

A more realistic class of models assumes that the springs in the bead-spring chain have finite extension: the domain D is then taken to be a Cartesian product of K *bounded* open balls $D_i \subset \mathbb{R}^d$, centred at the origin $\underline{0} \in \mathbb{R}^d$, $i = 1, \dots, K$, with $K \geq 1$. The spring potentials $U_i : s \in [0, \frac{b_i}{2}) \mapsto U_i(s) \in [0, \infty)$, with $b_i > 0$, $i = 1, \dots, K$, are in that case nonlinear and unbounded functions, and the associated bead-spring chain model is referred to as a FENE (finitely extensible nonlinear elastic) model; in the case of $K = 1$, the corresponding model is called a FENE dumbbell model.

Here we shall be concerned with finitely extensible nonlinear bead-spring chain models, with $D := B(\underline{0}, b_1^{\frac{1}{2}}) \times \cdots \times B(\underline{0}, b_K^{\frac{1}{2}})$, where $b_i > 0$, $i = 1, \dots, K$, $K \geq 1$, and $B(\underline{0}, b_i^{\frac{1}{2}})$ is a bounded open ball in \mathbb{R}^d of radius $b_i^{\frac{1}{2}}$, centred at $\underline{0} \in \mathbb{R}^d$. We shall adopt the following structural hypotheses on the spring potentials U_i and the associated partial Maxwellians M_i , $i = 1, \dots, K$.

We shall assume that $D_i = B(\underline{0}, b_i^{\frac{1}{2}})$ with $b_i > 0$ for $i = 1, \dots, K$, and that for $i = 1, \dots, K$ there exist constants $c_{ij} > 0$, $j = 1, 2, 3, 4$, and $\theta_i > 1$ such that the spring potential $U_i \in C^1[0, \frac{b_i}{2})$ and the associated partial Maxwellian M_i satisfy

$$(1.7a) \quad c_{i1} [\text{dist}(\underline{q}_i, \partial D_i)]^{\theta_i} \leq M_i(\underline{q}_i) \leq c_{i2} [\text{dist}(\underline{q}_i, \partial D_i)]^{\theta_i} \quad \forall \underline{q}_i \in D_i,$$

$$(1.7b) \quad c_{i3} \leq \text{dist}(\underline{q}_i, \partial D_i) U'_i(\tfrac{1}{2}|\underline{q}_i|^2) \leq c_{i4} \quad \forall \underline{q}_i \in D_i.$$

It follows from (1.7a,b) that (if $\theta_i > 1$, as has been assumed here,)

$$(1.8) \quad \int_{D_i} \left[1 + [U_i(\tfrac{1}{2}|\underline{q}_i|^2)]^2 + [U'_i(\tfrac{1}{2}|\underline{q}_i|^2)]^2 \right] M_i(\underline{q}_i) d\underline{q}_i < \infty, \quad i = 1, \dots, K.$$

Example 1.2. In the FENE (finitely extensible nonlinear elastic) dumbbell model, introduced by Warner [66], $K = 1$ and the spring force is given by $\underline{F}(\underline{q}) = (1 - |\underline{q}|^2/b)^{-1} \underline{q}$, $\underline{q} \in D = B(\underline{0}, b^{\frac{1}{2}})$, corresponding to $U(s) = -\frac{b}{2} \log(1 - \frac{2s}{b})$, $s \in \mathcal{O} = [0, \frac{b}{2})$, $b > 2$. More generally, in a FENE bead spring chain, one considers $K+1$ beads linearly coupled with K springs, each with a FENE spring

potential. Direct calculations show that the partial Maxwellians M_i and the elastic potentials U_i , $i = 1, \dots, K$, of the FENE bead spring chain satisfy the conditions (1.7a,b) with $\theta_i := \frac{b_i}{2}$, provided that $b_i > 2$, $i = 1, \dots, K$. Thus, (1.8) also holds and $b_i > 2$, $i = 1, \dots, K$. Note, however, that (1.8) fails for $b_i \in (0, 2]$, i.e., for $\theta_i \in (0, 1]$, which is why we have assumed in the statement of (1.7a,b) that $\theta_i > 1$ for $i = 1, \dots, K$.

It is interesting to note that in the (equivalent) stochastic version of the FENE dumbbell model ($K = 1$) a solution to the system of stochastic differential equations associated with the Fokker–Planck equation exists and has trajectorial uniqueness if, and only if, $\frac{b}{2} \geq 1$; (cf. Jourdain, Lelièvre & Le Bris [39] for details). Thus, in the general class of FENE-type bead-spring chain models considered here, the assumption $\theta_i > 1$, $i = 1, \dots, K$, is the weakest reasonable requirement on the decay-rate of M_i in (1.7a) as $\text{dist}(\underline{q}_i, \partial D_i) \rightarrow 0$.

The governing equations of the general nonhomogeneous bead-spring chain models with centre-of-mass diffusion considered in this paper are (1.1a–e), where the extra-stress tensor $\underline{\tau}$ is defined by

$$(1.9) \quad \underline{\tau}(\psi)(\underline{x}, t) := \underline{\tau}_1(\psi)(\underline{x}, t) - \left(\int_{D \times D} \gamma(\underline{q}, \underline{q}') \psi(\underline{x}, \underline{q}, t) \psi(\underline{x}, \underline{q}', t) \, d\underline{q} \, d\underline{q}' \right) \underline{I},$$

$\gamma : D \times D \rightarrow \mathbb{R}_{\geq 0}$ is a smooth, time-independent, \underline{x} -independent and ψ -independent interaction kernel, which we shall henceforth consider to be

$$\gamma(\underline{q}, \underline{q}') \equiv \mathfrak{z},$$

where $\mathfrak{z} \in \mathbb{R}_{>0}$; thus,

$$(1.10) \quad \underline{\tau}(\psi) := \underline{\tau}_1(\psi) - \mathfrak{z} \left(\int_D \psi \, d\underline{q} \right)^2 \underline{I}.$$

Here, $\underline{\tau}_1(\psi)$ is the *Kramers expression*; that is,

$$(1.11) \quad \underline{\tau}_1(\psi) := k \left[\left(\sum_{i=1}^K \underline{C}_i(\psi) \right) - (K+1) \int_D \psi \, d\underline{q} \, \underline{I} \right],$$

where $k \in \mathbb{R}_{>0}$, with the first term in the square brackets being due to the K springs and the second to the $K+1$ beads in the bead-spring chain representing the polymer molecule. Further,

$$(1.12) \quad \underline{C}_i(\psi)(\underline{x}, t) := \int_D \psi(\underline{x}, \underline{q}, t) U'_i(\tfrac{1}{2}|\underline{q}_i|^2) \underline{q}_i \underline{q}_i^T \, d\underline{q}, \quad i = 1, \dots, K.$$

The probability density function ψ is a solution of the Fokker–Planck (forward Kolmogorov) equation

$$(1.13) \quad \begin{aligned} & \frac{\partial \psi}{\partial t} + \nabla_{\underline{x}} \cdot (\underline{u} \psi) + \sum_{i=1}^K \nabla_{\underline{q}_i} \cdot \left(\underline{\sigma}(\underline{u}) \underline{q}_i \psi \right) \\ &= \varepsilon \Delta_{\underline{x}} \left(\frac{\psi}{\zeta(\rho)} \right) + \frac{1}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \nabla_{\underline{q}_i} \cdot \left(M \nabla_{\underline{q}_j} \left(\frac{\psi}{\zeta(\rho) M} \right) \right) \quad \text{in } \Omega \times D \times (0, T], \end{aligned}$$

with $\underline{\sigma}(\underline{v}) \equiv \nabla_{\underline{x}} \underline{v}$ and a, generally, density-dependent scaled drag coefficient $\zeta(\cdot) \in \mathbb{R}_{>0}$. A concise derivation of the Fokker–Planck equation (1.13) can be found in Section 1 of Barrett and Süli [12]. Let $\partial \overline{D}_i := D_1 \times \dots \times D_{i-1} \times \partial D_i \times D_{i+1} \times \dots \times D_K$. We impose the following boundary and

initial conditions on solutions of (1.13):

$$(1.14a) \quad \left[\frac{1}{4\lambda} \sum_{j=1}^K A_{ij} M \nabla_{q_j} \left(\frac{\psi}{\zeta(\rho) M} \right) - \sigma(u) \frac{q_i}{|q_i|} \psi \right] \cdot \frac{q_i}{|q_i|} = 0 \quad \text{on } \Omega \times \partial \bar{D}_i \times (0, T], \text{ for } i = 1, \dots, K,$$

$$(1.14b) \quad \varepsilon \nabla_x \left(\frac{\psi}{\zeta(\rho)} \right) \cdot \underline{n} = 0 \quad \text{on } \partial \Omega \times D \times (0, T],$$

$$(1.14c) \quad \psi(\cdot, \cdot, 0) = \psi_0(\cdot, \cdot) \geq 0 \quad \text{on } \Omega \times D,$$

where q_i is normal to ∂D_i , as D_i is a bounded ball centred at the origin, and \underline{n} is normal to $\partial \Omega$.

The first term in (1.11) is due to the K springs, whilst the second is due to the $K+1$ beads; see Chapter 15 in [17]. The nondimensional constant $k > 0$ featuring in (1.11) is a constant multiple of the product of the Boltzmann constant k_B and the absolute temperature T . In (1.13), $\varepsilon > 0$ is the centre-of-mass diffusion coefficient defined as $\varepsilon := (\ell_0/L_0)^2/(4(K+1)\lambda)$ with L_0 a characteristic length-scale of the solvent flow, $\ell_0 := \sqrt{k_B T/H}$ signifying the characteristic microscopic length-scale and $\lambda := \frac{\zeta_0 U_0}{4H L_0}$, where $\zeta_0 > 0$ is a characteristic drag coefficient and $H > 0$ is a spring-constant. The nondimensional parameter $\lambda \in \mathbb{R}_{>0}$, called the Deborah number (and usually denoted by De), characterizes the elastic relaxation property of the fluid, and $\underline{A} = (A_{ij})_{i,j=1}^K$ is the symmetric positive definite *Rouse matrix*, or connectivity matrix; for example, $\underline{A} = \text{tridiag}[-1, 2, -1]$ in the case of a (topologically) linear chain, depicted in Fig. 1; see, Nitta [57]. Concerning these scalings and notational conventions, we remark that the factor $\frac{1}{4\lambda}$ in equation (1.13) above appears as a factor $\frac{1}{2\lambda}$ in the Fokker–Planck equation in our earlier papers [11, 13, 15, 10].

Two remarks are in order at this point concerning the extra-stress tensor (1.10).

Remark 1.2. *By comparing the last term in the expression (A) in Table 15.2-1 on p.156 in [17] with the second term in (1.11), we observe that Bird et al. write K (i.e., $(N-1)$ in their notation) instead of our prefactor $(K+1)$ (i.e., N in terms of their notation), in the context of incompressible spatially homogeneous flows (i.e., the velocity field is assumed to be independent of \underline{x} there). As we shall explain in equations (1.18)–(1.23) below, the analysis in the present paper applies to a general class of extra stress tensors (1.10), with Kramers type expressions of the form*

$$(1.15) \quad \underline{\tau}_1(\psi) := k \left[\left(\sum_{i=1}^K \underline{C}_i(\psi) \right) - \mathfrak{k} \int_D \psi \, d\underline{q} \, \underline{I} \right],$$

with $\mathfrak{k} \in \mathbb{R}$ and $\mathfrak{z} > 0$. Since the actual value of \mathfrak{k} is of no particular relevance in our analysis, we took $\mathfrak{k} = K+1$ in the second term in (1.15), yielding (1.11), as this choice simplifies the expressions that arise in the course of the proof. As will be shown below, if $\mathfrak{k} \geq K+1$ and $\mathfrak{z} \geq 0$ a formal energy identity holds, and if $\mathfrak{k} < K+1$ and $\mathfrak{z} > 0$, one can still prove a formal energy inequality. In contrast with this, in the case of incompressible flows, \mathfrak{k} and \mathfrak{z} are of no relevance in the analysis and can be any two real numbers; indeed, in [11, 13, 14], $\mathfrak{k} = K$ and $\mathfrak{z} = 0$ were the values used in (1.15) and (1.10), respectively, resulting in $\underline{\tau}(\psi)$ that is the classical Kramers expression for the extra stress tensor.

Remark 1.3. *Our second remark concerns the quadratic modification to $\underline{\tau}_1(\psi)$ appearing as the second term in (1.10). By defining the polymer number density*

$$(1.16) \quad \varrho(\underline{x}, t) := \int_D \psi(\underline{x}, \underline{q}, t) \, d\underline{q}, \quad (\underline{x}, t) \in \Omega \times [0, T],$$

formally integrating (1.13) over D and using the boundary condition (1.14a), we deduce that

$$(1.17a) \quad \frac{\partial \varrho}{\partial t} + \nabla_x \cdot (\underline{u} \varrho) = \varepsilon \Delta_x \left(\frac{\varrho}{\zeta(\rho)} \right) \quad \text{on } \Omega \times (0, T],$$

together with the boundary and initial conditions (which result from integrating (1.14b,c) over D)

$$(1.17b) \quad \varepsilon \nabla_x \left(\frac{\varrho}{\zeta(\rho)} \right) \cdot \underline{n} = 0 \quad \text{on } \partial \Omega \times (0, T] \quad \text{and} \quad \varrho(\underline{x}, 0) = \int_D \psi_0(\underline{x}, \underline{q}) \, d\underline{q} \quad \text{for } \underline{x} \in \Omega.$$

In order to avoid potential confusion concerning our notational conventions, we draw the reader's attention to the fact that, here and throughout the rest of the paper, the symbol ρ signifies the density of the solvent, while the symbol ϱ denotes the polymer number density, as defined in (1.16).

If $\nabla_x \cdot \mathfrak{u} \equiv 0$ and $\varrho(\cdot, 0)$ is constant, and either $\varepsilon = 0$ or $\zeta(\rho)$ is independent of ρ (and therefore identically equal to a constant), then $\varrho(\mathfrak{x}, t)$ is constant ($\equiv \varrho(\cdot, 0)$), for all $(\mathfrak{x}, t) \in \Omega \times (0, T]$, and hence

$$\int_D \psi(\mathfrak{x}, \mathfrak{q}, t) d\mathfrak{q} = \int_D \psi_0(\mathfrak{x}, \mathfrak{q}) d\mathfrak{q} \in \mathbb{R}_{>0} \quad \text{for all } (\mathfrak{x}, t) \in \Omega \times (0, T];$$

in other words, the polymer number density is constant. This conservation property then guarantees complete control of $\varrho(\mathfrak{x}, t) = \int_D \psi(\mathfrak{x}, \mathfrak{q}, t) d\mathfrak{q}$ (when $\mathfrak{z} = 0$, and a fortiori, for $\mathfrak{z} > 0$) in terms of the initial probability density function ψ_0 in the course of the weak compactness argument upon which the proof of existence of weak solutions rests (cf. [11] for the analysis in the case of $\varepsilon > 0$, constant ζ and $\mathfrak{z} = 0$). In particular, the time derivative $\frac{\partial \psi}{\partial t}$ can be bounded in a sufficiently strong norm to enable the application of an Aubin–Lions–Simon type compactness theorem that ensures strong convergence of the sequence of approximations to the probability density function ψ .

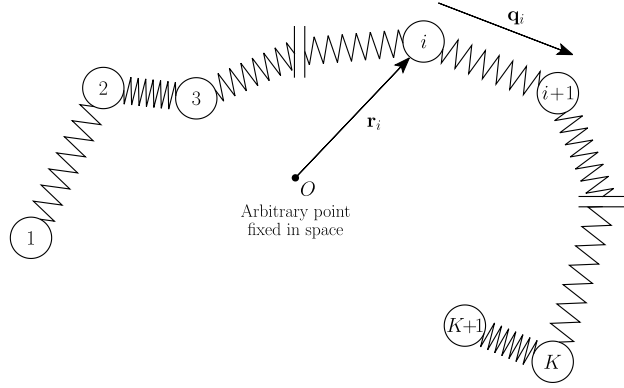
In the case of nonhomogeneous flows with density-dependent drag, the situation is more complicated because of the consequential weaker control (in the sense of norms) on the function ϱ : while, thanks to (1.17a,b) and the assumed homogeneous Dirichlet boundary condition on \mathfrak{u} , it is still true that

$$\int_{\Omega} \varrho(\mathfrak{x}, t) d\mathfrak{x} = \int_{\Omega \times D} \psi(\mathfrak{x}, \mathfrak{q}, t) d\mathfrak{x} d\mathfrak{q} = \int_{\Omega \times D} \psi_0(\mathfrak{x}, \mathfrak{q}) d\mathfrak{x} d\mathfrak{q} = \int_{\Omega} \varrho(\mathfrak{x}, 0) d\mathfrak{x}$$

for all $t \in (0, T]$, the availability of such a weak conservation property only further complicates the analysis. In [14] (in the case of $\nabla_x \cdot \mathfrak{u} \equiv 0$ and $\mathfrak{z} = 0$) a more involved argument, based on a combination of compensated compactness and Vitali's theorem, was therefore used in order to deduce strong convergence of the sequence of approximations to ψ .

In the present paper, in order to focus on the essential new difficulty — the lack of the divergence-free property of \mathfrak{u} — we shall suppose that the drag coefficient in the Fokker–Planck equation is identically equal to a constant, which we shall henceforth, without loss of generality, assume to be equal to 1, i.e., $\zeta(\rho) \equiv 1$. As in the majority of contributions to the mathematical analysis of the compressible Navier–Stokes equations to date, we shall also assume that the shear and bulk viscosity coefficients are independent of the density ρ , i.e., that $\mu^S \in \mathbb{R}_{>0}$ and $\mu^B \in \mathbb{R}_{\geq 0}$. For the analysis of the compressible Navier–Stokes equations in the case of a particular family of density-dependent shear and bulk viscosities, we refer the reader to [19, 20]. Setting $\mathfrak{z} = 0$ in this context results in the loss of a bound on the $L^1(0, T; \mathfrak{X}')$ norm of the time derivative of the probability density function ψ , for any reasonable choice of the function space \mathfrak{X} . Failure to control $\frac{\partial \psi}{\partial t}$ or a time-difference of ψ in even such a weak sense brings into question the meaningfulness of the model in the case of compressible flows for solutions of as low a degree of regularity as is guaranteed by the formal energy bound in the case of $\mathfrak{z} = 0$, see Remarks 4.1 and 4.2 below. Motivated by the papers of Constantin [22], Constantin et al. [23] and Bae & Trivisa [4], we have therefore included the quadratic term in (1.10), with $\mathfrak{z} > 0$. As we shall show later on, inclusion of the quadratic term into (1.10) does not destroy energy balance thanks to the fact that the polymer number density function ϱ satisfies the initial-boundary-value problem (1.17a,b), and has the beneficial effect of guaranteeing $L^\infty(0, T; L^2(\Omega)) \cap L^2(0, T; H^1(\Omega))$ norm control for ϱ , rendering the time derivative of the probability density function finite in the norm of $L^2(0, T; \mathfrak{X}')$, with a suitable choice of \mathfrak{X} as a Maxwellian-weighted Sobolev space of sufficiently high order. This then enables the application of Dubinskii's extension of the Aubin–Lions–Simon compactness theorem (cf. [28] and Barrett & Süli [16]).

From the physical point of view (1.2), (1.10)–(1.12) can be seen as a decomposition of the Cauchy stress π as the sum of a contribution from the solvent, π_s , and the polymeric extra stress,

FIGURE 1. A (topologically) linear bead-spring chain with K springs and $K + 1$ beads.

π_p , resulting from the presence of the polymer molecules, which are idealized here as bead-spring-chains (cf. eq. (13.3-1) in [17]):

$$\underline{\pi} = \underline{\pi}_s + \underline{\pi}_p = (\underline{S}(\underline{u}, \rho) - p_s \underline{I}) + \left(k \sum_{i=1}^K \underline{C}_i(\psi) - p_p \underline{I} \right),$$

where $p_s = p = c_p \rho^\gamma$ (with $c_p > 0$, $\gamma > \frac{3}{2}$, and ρ denoting the density) is the fluid pressure, and $p_p = k(K+1)\varrho + \frac{1}{2}\varrho^2$ (with k denoting a constant multiple of the product of the Boltzmann constant and the absolute temperature, $\frac{1}{2} > 0$ and ϱ signifying the polymer number density) is the polymeric contribution to the total pressure, defined as $p_s + p_p$. The expression $k(K+1)\varrho + \frac{1}{2}\varrho^2$ can be viewed as a quadratic truncation of a virial expansion of the polymeric pressure p_p in terms of the polymer number density ϱ .

Definition 1.1. The collection of equations and structural hypotheses (1.1a–e)–(1.14a–c) together with the assumption that the Rouse matrix A is symmetric and positive definite (as is always the case, by definition,) will be referred to throughout the paper as model (P), or as the compressible FENE-type bead-spring chain model with centre-of-mass diffusion. It will be assumed throughout the paper that the shear viscosity, $\mu^S \in \mathbb{R}_{>0}$, the bulk viscosity, $\mu^B \in \mathbb{R}_{\geq 0}$, and the drag coefficient, $\zeta \in \mathbb{R}_{>0}$, are independent of the density ρ . For the ease of exposition we shall set $\zeta \equiv 1$.

For a survey of recent developments concerning the mathematical analysis of Navier–Stokes–Fokker–Planck systems in the case of *incompressible* polymeric flows the reader is referred to the introductory sections in our papers [6, 7, 8, 11, 12, 13, 14]. The convergence of finite element approximations of Navier–Stokes–Fokker–Planck systems is discussed in [5, 9, 15], where we also review the literature on the mathematical theory of numerical algorithms for these equations. For further details concerning the mathematical analysis of kinetic models of polymers in the case of incompressible fluids we refer the reader to Renardy [62], Lions & Masmoudi [46], E, Li & Zhang [29] and Li, Zhang & Zhang [42], Jourdain, Lelièvre & Le Bris [39], Constantin [22], Du, Yu & Liu [27], Yu, Du & Liu [68], Zhang & Zhang [69], Lions & Masmoudi [47], Masmoudi [51], [52], Otto & Tzavaras [59], and the references therein.

A noteworthy feature of equation (1.13) in the model (P) compared to classical Fokker–Planck equations for bead-spring-chain models for dilute polymers appearing in the literature is the presence of the \underline{x} -dissipative centre-of-mass diffusion term $\varepsilon \Delta_x \psi$ on the right-hand side of the Fokker–Planck equation (1.13). We refer to Barrett & Süli [7] for the derivation of (1.13) in the case of $K = 1$ and constant ρ ; see also the article by Schieber [63] concerning generalized dumbbell models with centre-of-mass diffusion, and the recent paper of Degond & Liu [25] for a careful justification of the presence of the centre-of-mass diffusion term through asymptotic analysis. In the case of variable density, viscosity and drag, the derivation of the Fokker–Planck equation with centre-of-mass diffusion appears in [12] (cf. also [13, 14]).

For a survey of macroscopic models of *compressible* viscoelastic flow, the reader is referred to the paper by Bollada & Phillips [18]. The existence and uniqueness of local strong solutions and the existence of global solutions near equilibrium for macroscopic models of three-dimensional compressible viscoelastic fluids was considered in [34, 61, 60, 35, 36, 37]. The existence of measure-valued solutions to non-Newtonian compressible, isothermal, monopolar fluid flow models was studied by Nečasová in [53, 55]; for bipolar isothermal non-Newtonian compressible fluids related analysis was pursued in [54]. In a series of papers (cf. [48, 49, 50]) Mamontov developed a priori estimates for two- and three-dimensional compressible nonlinear viscoelastic flow problems and studied the existence of solutions. Zhikov & Pastukhova [70] proved the existence of global weak solutions to a class of compressible viscoelastic flow models with p -Laplacian structure. There is also a substantial literature in chemical engineering on the use of the compressible Oldroyd-B system in modelling bubble dynamics in compressible viscoelastic liquids (cf., for example, [21]). Closer to the subject of the present paper, Bae & Trivisa [4] have established the existence of global weak solutions to Doi's rod-model in three-dimensional bounded domains. The model concerns suspensions of rod-like molecules in compressible fluids and involves the coupling of a Fokker–Planck type equation with the compressible Navier–Stokes system. Also, Jiang, Jiang & Wang [38] have studied the existence of global weak solutions to the equations of compressible flow of nematic liquid crystals in two dimensions.

Our objective in the present paper is to prove the existence of global-in-time weak solutions to the general class of models of compressible viscoelastic flow, labelled above as Problem (P) (cf. Definition 1.1), that arise from the kinetic theory of dilute solutions of nonhomogeneous polymeric liquids where polymer molecules are idealized as bead-spring chains, with finitely extensible nonlinear elastic (FENE) type spring potentials, involving the coupling of the unsteady, isentropic, compressible Navier–Stokes equations with the Fokker–Planck equation. Despite their importance, in the present paper we shall, for the sake of simplicity, neglect all thermal effects and will focus instead on mechanical properties of the fluid in the isothermal setting. With a nonnegative initial density $\rho_0 \in L^\infty(\Omega)$ for the continuity equation; a square-integrable initial velocity datum \underline{u}_0 for the Navier–Stokes momentum equation; and a nonnegative initial probability density function ψ_0 for the Fokker–Planck equation, which has finite relative entropy with respect to the Maxwellian M associated with the spring potential in the model, we prove, *via* a limiting procedure on certain discretization and regularization parameters, the existence of a global-in-time bounded-energy weak solution $t \mapsto (\rho(t), \underline{u}(t), \psi(t))$ to the coupled Navier–Stokes–Fokker–Planck system, satisfying the initial condition $(\rho(0), \underline{u}(0), \psi(0)) = (\rho_0, \underline{u}_0, \psi_0)$. Since the argument is long and technical, we give a brief overview of the main steps of the proof.

At the heart of the proof is a formal energy identity, which we shall now derive under the assumption that \underline{u} , ρ , ψ and ϱ are sufficiently smooth, and, at least for our purposes in this introductory section, ρ is nonnegative, and ψ and ϱ are positive. Instead of (1.11), used in the rest of the paper, we shall make use of the more general formula (1.15), in order to explain the admissible range of \mathfrak{k} alluded to in Remark 1.2 as well as our reasons for choosing $\mathfrak{k} = K + 1$ in (1.11). By taking the $L^2(\Omega)$ inner product of equation (1.1a) first with $\frac{1}{2}|\underline{u}|^2$ and then with $P'(\rho)$, where $P(\rho) = \frac{p(\rho)}{\gamma-1}$ and so, on noting (1.3), $\rho P'(\rho) - P(\rho) = p(\rho)$, and then taking the $\underline{L}^2(\Omega)$ inner product of equation (1.1c) with \underline{u} , we deduce upon partial integration and noting the homogeneous Dirichlet boundary condition on \underline{u} and (1.10), (1.12) and (1.15) that

$$\begin{aligned}
 & \frac{d}{dt} \int_{\Omega} \left[\frac{1}{2} |\underline{u}|^2 + P(\rho) \right] d\mathbf{x} + \mu^S \int_{\Omega} \left| \underline{D}(\underline{u}) - \frac{1}{d} (\nabla_x \cdot \underline{u}) \underline{I} \right|^2 d\mathbf{x} + \mu^B \int_{\Omega} |\nabla_x \cdot \underline{u}|^2 d\mathbf{x} \\
 &= \int_{\Omega} \rho \underline{f} \cdot \underline{u} d\mathbf{x} - \int_{\Omega} \underline{\tau} : \nabla_x \underline{u} d\mathbf{x} \\
 &= \int_{\Omega} \rho \underline{f} \cdot \underline{u} d\mathbf{x} - k \sum_{i=1}^K \int_{\Omega \times D} \psi U'_i(\tfrac{1}{2} |q_i|^2) q_i q_i^T : \nabla_x \underline{u} dq d\mathbf{x} \\
 &+ k \mathfrak{k} \int_{\Omega \times D} \psi (\nabla_x \cdot \underline{u}) dq d\mathbf{x} + \mathfrak{z} \int_{\Omega} \left[\int_D \psi dq \right]^2 (\nabla_x \cdot \underline{u}) d\mathbf{x}.
 \end{aligned}
 \tag{1.18}$$

Recalling our convention that $\zeta \equiv 1$ (cf. Definition 1.1), it follows by taking the $L^2(\Omega)$ inner product of (1.17a) with ϱ , partial integration, noting the boundary condition (1.17b) and that \underline{u} satisfies a homogeneous Dirichlet boundary condition on $\partial\Omega$, that

$$(1.19) \quad \frac{d}{dt} \int_{\Omega} \varrho^2 d\tilde{x} + 2\varepsilon \int_{\Omega} |\nabla_x \varrho|^2 d\tilde{x} = - \int_{\Omega} \varrho^2 (\nabla_x \cdot \underline{u}) d\tilde{x}.$$

By noting (1.16) and substituting (1.19) into the last term on the right-hand side of (1.18), we deduce that

$$(1.20) \quad \begin{aligned} & \frac{d}{dt} \int_{\Omega} \left[\frac{1}{2} \rho |\underline{u}|^2 + P(\rho) + \mathfrak{z} \varrho^2 \right] d\tilde{x} \\ & + \mu^S \int_{\Omega} \left| \underline{D}(\underline{u}) - \frac{1}{d} (\nabla_x \cdot \underline{u}) \underline{I} \right|^2 d\tilde{x} + \mu^B \int_{\Omega} |\nabla_x \cdot \underline{u}|^2 d\tilde{x} + 2\varepsilon \mathfrak{z} \int_{\Omega} |\nabla_x \varrho|^2 d\tilde{x} \\ & = \int_{\Omega} \rho \underline{f} \cdot \underline{u} d\tilde{x} - k \sum_{i=1}^K \int_{\Omega \times D} \psi U'_i(\tfrac{1}{2} |q_i|^2) q_i q_i^T : \nabla_x \underline{u} dq d\tilde{x} \\ & + k \mathfrak{k} \int_{\Omega \times D} \psi (\nabla_x \cdot \underline{u}) dq d\tilde{x}. \end{aligned}$$

Next, the Fokker–Planck equation (1.13) (with $\zeta \equiv 1$) is multiplied by $\log \frac{\psi}{M}$, integrated over $\Omega \times D$, and integrations by parts are performed, using the boundary conditions (1.1d), (1.14a) and (1.14b), in the second and third term on the left-hand side and the two terms on the right-hand side; hence, by letting $\mathcal{F}(s) = s(\log s - 1) + 1$ for $s > 0$ and $\mathcal{F}(0) := \lim_{s \rightarrow 0+} \mathcal{F}(s) = 1$,

$$(1.21) \quad \begin{aligned} & \frac{d}{dt} \int_{\Omega \times D} M \mathcal{F} \left(\frac{\psi}{M} \right) dq d\tilde{x} + \int_{\Omega \times D} \psi (\nabla_x \cdot \underline{u}) dq d\tilde{x} \\ & - \sum_{i=1}^K \int_{\Omega \times D} \psi U'_i(\tfrac{1}{2} |q_i|^2) q_i q_i^T : \nabla_x \underline{u} dq d\tilde{x} + K \int_{\Omega \times D} \psi (\nabla_x \cdot \underline{u}) dq d\tilde{x} \\ & + 4\varepsilon \int_{\Omega \times D} M \left| \nabla_x \sqrt{\frac{\psi}{M}} \right|^2 dq d\tilde{x} + \frac{1}{\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \int_{\Omega \times D} M \nabla_{q_j} \sqrt{\frac{\psi}{M}} \cdot \nabla_{q_i} \sqrt{\frac{\psi}{M}} dq d\tilde{x} = 0. \end{aligned}$$

We now multiply (1.21) by k and add the resulting identity to (1.20) to deduce that

$$(1.22) \quad \begin{aligned} & \frac{d}{dt} \int_{\Omega} \left[\frac{1}{2} \rho |\underline{u}|^2 + P(\rho) + \mathfrak{z} \varrho^2 + k \int_D M \mathcal{F} \left(\frac{\psi}{M} \right) dq \right] d\tilde{x} + k(K+1-\mathfrak{k}) \int_{\Omega} \varrho (\nabla_x \cdot \underline{u}) d\tilde{x} \\ & + \mu^S \int_{\Omega} \left| \underline{D}(\underline{u}) - \frac{1}{d} (\nabla_x \cdot \underline{u}) \underline{I} \right|^2 d\tilde{x} + \mu^B \int_{\Omega} |\nabla_x \cdot \underline{u}|^2 d\tilde{x} + 2\varepsilon \mathfrak{z} \int_{\Omega} |\nabla_x \varrho|^2 d\tilde{x} \\ & + 4\varepsilon k \int_{\Omega \times D} M \left| \nabla_x \sqrt{\frac{\psi}{M}} \right|^2 dq d\tilde{x} + \frac{k}{\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \int_{\Omega \times D} M \nabla_{q_j} \sqrt{\frac{\psi}{M}} \cdot \nabla_{q_i} \sqrt{\frac{\psi}{M}} dq d\tilde{x} \\ & = \int_{\Omega} \rho \underline{f} \cdot \underline{u} d\tilde{x}. \end{aligned}$$

Unless $\mathfrak{k} = K+1$, one further step is necessary, in order to deal with the second term on the left-hand side. To this end, we return to (1.17a), with $\zeta(\rho) \equiv 1$, take the $L^2(\Omega)$ inner product with $\mathcal{F}'(\varrho)$, and integrate by parts in the second and third term, analogously to (1.19), to deduce that

$$\frac{d}{dt} \int_{\Omega} \mathcal{F}(\varrho) d\tilde{x} + 4\varepsilon \int_{\Omega} |\nabla_x \sqrt{\varrho}|^2 d\tilde{x} = - \int_{\Omega} \varrho (\nabla_x \cdot \underline{u}) d\tilde{x}.$$

Substitution of this into the second term in (1.22) then yields

$$\begin{aligned}
& \frac{d}{dt} \int_{\Omega} \left[\frac{1}{2} \rho |\underline{u}|^2 + P(\rho) + \mathfrak{z} \varrho^2 + k(\mathfrak{k} - (K+1)) \mathcal{F}(\varrho) + k \int_D M \mathcal{F}\left(\frac{\psi}{M}\right) d\tilde{q} \right] d\tilde{x} \\
& + \mu^S \int_{\Omega} \left| \underline{D}(\underline{u}) - \frac{1}{d} (\nabla_x \cdot \underline{u}) I \right|^2 d\tilde{x} + \mu^B \int_{\Omega} |\nabla_x \cdot \underline{u}|^2 d\tilde{x} \\
& + 2\varepsilon \mathfrak{z} \int_{\Omega} |\nabla_x \varrho|^2 d\tilde{x} + 4\varepsilon k(\mathfrak{k} - (K+1)) \int_{\Omega} |\nabla_x \sqrt{\varrho}|^2 d\tilde{x} \\
& + 4\varepsilon k \int_{\Omega \times D} M \left| \nabla_x \sqrt{\frac{\psi}{M}} \right|^2 d\tilde{q} d\tilde{x} + \frac{k}{\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \int_{\Omega \times D} M \nabla_{q_j} \sqrt{\frac{\psi}{M}} \cdot \nabla_{q_i} \sqrt{\frac{\psi}{M}} d\tilde{q} d\tilde{x} \\
(1.23) \quad & = \int_{\Omega} \rho \underline{f} \cdot \underline{u} d\tilde{x},
\end{aligned}$$

which is the desired (formal) energy identity that represents the starting point for our proof of existence of global weak solutions to problem (P). The integral over Ω of the expression in the square brackets in the first line of (1.23) is the total energy, and the sum of the terms on the second, third and fourth line of (1.23) is the dissipation of the total energy; recall that the Rouse matrix $A = (A_{ij})_{i,j=1}^K$ is, by definition, symmetric and positive definite. While, as an energy identity, (1.23) is meaningful for all $\mathfrak{z} \geq 0$ and all $\mathfrak{k} \geq K+1$, it will transpire in the course of the proof that $\mathfrak{z} > 0$ is necessary in order to ensure control of $\frac{\partial \psi}{\partial t}$ or of a time-difference of ψ (cf. Remarks 1.3 and 4.2). Once \mathfrak{z} has been chosen to be positive, the two terms in (1.23) that include the factor $(\mathfrak{k} - (K+1))$ are of lower order and contribute no additional information; we have therefore, for the sake of simplicity, set $\mathfrak{k} = K+1$ in (1.15), yielding (1.11). If $\mathfrak{k} < K+1$ and $\mathfrak{z} > 0$, then upon moving the second term on the left-hand side of (1.22) to the right, Gronwall's inequality yields a formal energy *inequality*; since the ultimate outcome is no different from the one with $\mathfrak{k} = K+1$ and $\mathfrak{z} > 0$, we shall not discuss this case further.

The main idea of the proof is to construct a sequence of approximating solutions, whose existence one can prove. The sequence of approximating solutions to the solution of problem (P) will be defined through suitable truncations (cut-off), regularizations and temporal semidiscretization. After deriving (truncation-, regularization-, discretization-) parameter independent bounds on the sequence of approximating solutions, we shall use a variety of compactness arguments to pass to limits in the parameters. The derivations of the various bounds on sequences of approximating solutions mimic the derivation of the formal energy identity (1.23) outlined above. The present paper combines a number of techniques developed in our earlier work on the existence of global weak solutions to Navier–Stokes–Fokker–Planck systems, in the case of incompressible flows of dilute polymers, with techniques that were introduced in the groundbreaking contributions of Lions [45] and Feireisl [31] to the mathematical theory of compressible Navier–Stokes equations. For a detailed overview of the latter, we point the reader to the, more recent, monograph of Novotný & Straškraba [58], which will be our main source of reference for technical results from the mathematical theory of compressible Navier–Stokes equations; for the convenience of the reader we have included the most relevant ones of these in Appendices A–G, at the end of the paper.

Available results in the literature concerning the existence of weak solutions to the compressible Navier–Stokes equations are based on considering sequences of approximating problems that are defined by spatial Galerkin discretization, using bases of smooth functions. Since, as in our previous work on incompressible Navier–Stokes–Fokker–Planck systems, we are motivated by constructive considerations, ultimately aimed at developing convergent numerical algorithms for the problem, we shall adopt a different approach. Our construction of a sequence of approximating problems for the compressible Navier–Stokes–Fokker–Planck system will be based on discretizing the problem with respect to the temporal variable. For a similar approach, in the case of a coupled compressible Navier–Stokes–Cahn–Hilliard system, we refer the reader to the work of Abels and Feireisl [1] and to Remark 3.1 below, which explains the aspects in which our temporal approximation differs from the one in [1]. The inclusion of the Fokker–Planck equation into the analysis is nontrivial, the main hurdle being to ensure that the presence of the extra stress term

τ_{\approx} (cf. (1.9)) on the right-hand side of the Navier–Stokes momentum equation does not destroy the Lions–Feireisl compactness argument for the compressible Navier–Stokes equations. We have only been able to achieve this for $\mathfrak{z} > 0$. As the same requirement on \mathfrak{z} has been found to be necessary in the, related, Doi model for suspensions of rod-like molecules in a compressible fluid, considered by Bae and Trivisa [4], we are confident that the condition $\mathfrak{z} > 0$ is not a byproduct of our time-discrete approach to the proof of existence of weak solutions. It is more likely that the classical Kramers expression $\tau_{\approx 1}$, appearing as the first term in (1.9), which was originally derived in the case of a homogeneous incompressible solvent, needs to be supplemented in the case of a nonhomogeneous compressible solvent by additional correction terms (e.g. to account for the fact that compressibility may impact on the size of the excluded volume surrounding a polymer molecule immersed in the solvent), such as the ‘interaction term’ appearing as the second term in (1.9).

The proof of the central theorem in the paper, Theorem 6.1, stating the existence of global bounded-energy weak solutions to problem (P), consists of six steps, which are outlined below.

Step 1. Following the approach in Barrett & Süli [8, 11, 12, 13, 14] and motivated by recent papers of Jourdain, Lelièvre, Le Bris & Otto [40] and Lin, Liu & Zhang [43] (see also Arnold, Markowich, Toscani & Unterreiter [2], and Desvillettes & Villani [26]) concerning the convergence of the probability density function ψ to its equilibrium value $\psi_{\infty}(\underline{x}, \underline{q}) := M(\underline{q})$ (corresponding to the equilibrium value $\underline{u}_{\infty}(\underline{x}) := \underline{0}$ of the velocity field in the case of constant density) in the absence of body forces \underline{f} , we observe that if $\frac{\psi}{M}$ is bounded above then, for $L \in \mathbb{R}_{>0}$ sufficiently large, the third term in (1.13), referred to as the *drag term* is equal to

$$(1.24) \quad \sum_{i=1}^K \nabla_{q_i} \cdot \left(\underline{\sigma}(\underline{u}) \underline{q}_i M \beta^L \left(\frac{\psi}{M} \right) \right)$$

(recall that, by hypothesis, $\zeta \equiv 1$), where $\beta^L \in C(\mathbb{R})$ is a cut-off function defined as

$$(1.25) \quad \beta^L(s) := \min\{s, L\}.$$

It then follows that, for $L \gg 1$, any solution ψ of (1.13), such that $\frac{\psi}{M}$ is bounded above by L , also satisfies

$$(1.26) \quad \begin{aligned} & \frac{\partial \psi}{\partial t} + \nabla_{\underline{x}} \cdot (\underline{u} \psi) + \sum_{i=1}^K \nabla_{q_i} \cdot \left(\underline{\sigma}(\underline{u}) \underline{q}_i M \beta^L \left(\frac{\psi}{M} \right) \right) \\ &= \varepsilon \Delta_{\underline{x}} \psi + \frac{1}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \nabla_{q_i} \cdot \left(M \nabla_{q_j} \left(\frac{\psi}{M} \right) \right) \quad \text{in } \Omega \times D \times (0, T]. \end{aligned}$$

We impose the following boundary and initial conditions:

$$(1.27a) \quad \left[\frac{1}{4\lambda} \sum_{j=1}^K A_{ij} M \nabla_{q_j} \left(\frac{\psi}{M} \right) - \underline{\sigma}(\underline{u}) \underline{q}_i M \beta^L \left(\frac{\psi}{M} \right) \right] \cdot \frac{\underline{q}_i}{|\underline{q}_i|} = 0$$

$$(1.27b) \quad \varepsilon \nabla_{\underline{x}} \psi \cdot \underline{n} = 0 \quad \text{on } \partial\Omega \times D \times (0, T],$$

$$(1.27c) \quad \psi(\cdot, \cdot, 0) = M(\cdot) \beta^L(\psi_0(\cdot, \cdot)/M(\cdot)) \geq 0 \quad \text{on } \Omega \times D.$$

The initial datum ψ_0 for the Fokker–Planck equation is nonnegative, defined on $\Omega \times D$, with

$$\int_D \psi_0(\underline{x}, \underline{q}) \, d\underline{q} \in L^\infty(\Omega),$$

and is assumed to have finite relative entropy with respect to the Maxwellian M ; i.e.

$$\int_{\Omega \times D} \psi_0(\underline{x}, \underline{q}) \log \frac{\psi_0(\underline{x}, \underline{q})}{M(\underline{q})} \, d\underline{q} \, d\underline{x} < \infty.$$

The model with cut-off parameter $L > 1$ is further regularized, by introducing a dissipation term of the form $-\alpha \Delta \rho$, with $\alpha > 0$, into the continuity equation (1.1a) and supplementing the resulting parabolic equation with a homogeneous Neumann boundary condition on $\partial\Omega \times (0, T]$. In addition, the equation of state (1.3) is replaced by a regularized equation of state, $p_\kappa(\rho) = p(\rho) + \kappa(\rho^4 + \rho^\Gamma)$, where $\kappa \in \mathbb{R}_{>0}$ and $\Gamma = \max\{\gamma, 8\}$. The resulting problem is denoted by $(P_{\kappa, \alpha, L})$.

Step 2. Ideally, one would like to pass to the limits $\kappa \rightarrow 0_+$, $\alpha \rightarrow 0_+$, $L \rightarrow +\infty$ to deduce the existence of solutions to (P). Unfortunately, such a direct attack at the problem is fraught with technical difficulties. Instead, we shall first (semi)discretize the problem $(P_{\kappa, \alpha, L})$ by an implicit Euler type scheme with respect to t , with step size Δt . This then results in a time-discrete version $(P_{\kappa, \alpha, L}^{\Delta t})$ of $(P_{\kappa, \alpha, L})$.

Step 3. By using Schauder's fixed point theorem, we will show in Section 3 the existence of solutions to $(P_{\kappa, \alpha, L}^{\Delta t})$. In the course of the proof, for technical reasons, a further cut-off, now from below, with a cut-off parameter $\delta \in (0, 1)$, is required. In addition, a fourth-order hyperviscosity term is added to the Navier–Stokes momentum equation (1.1c). We shall let δ pass to 0 to complete the proof of existence of solutions to $(P_{\kappa, \alpha, L}^{\Delta t})$ in the limit of $\delta \rightarrow 0_+$ in Section 3; cf. Lemma 3.4.

Step 4. In Section 4 we then go on to derive bounds on the sequence of solutions to problem $(P_{\kappa, \alpha, L}^{\Delta t})$; in particular, we develop various bounds on the sequence of weak solutions to $(P_{\kappa, \alpha, L}^{\Delta t})$ that are uniform in the time step Δt and the cut-off parameter L , and thus permit the extraction of weakly convergent subsequences, as $L \rightarrow +\infty$ and $\Delta t \rightarrow 0_+$, with $\Delta t = o(L^{-1})$, when $L \rightarrow +\infty$. The weakly convergent subsequences will then be shown to converge strongly in suitable norms. This then allows us to pass to the limit as $L \rightarrow +\infty$, with $\Delta t = o(L^{-1})$. The main result of Section 4 is Theorem 4.1, which summarizes the outcome of this limiting process.

Step 5. Section 5 is concerned with passage to the limit $\alpha \rightarrow 0_+$ with the parabolic regularization parameter that was introduced into the continuity equation in Step 1. The main result of Section 5 is Theorem 5.1, which summarizes the outcome of this limiting process.

Step 6. Finally, in Section 6 we pass to the limit $\kappa \rightarrow 0_+$ with the regularization parameter that was introduced into the equation of state in Step 1, which then leads to our main result, Theorem 6.1, stating the existence of global bounded-energy weak solutions to problem (P).

2. THE POLYMER MODEL $(P_{\kappa, \alpha, L})$

Let $\Omega \subset \mathbb{R}^d$ be a bounded open set with a Lipschitz-continuous boundary $\partial\Omega$, and suppose that the set $D := D_1 \times \cdots \times D_K$ of admissible conformation vectors $\underline{q} := (q_1^T, \dots, q_K^T)^T$ in (1.13) is such that D_i , $i = 1, \dots, K$, is an open ball in \mathbb{R}^d , $d = 2$ or 3 , centred at the origin, with boundary ∂D_i and radius $\sqrt{b_i}$, $b_i > 2$; let

$$(2.1) \quad \partial D := \bigcup_{i=1}^K \partial \bar{D}_i, \quad \text{where} \quad \partial \bar{D}_i := D_1 \times \cdots \times D_{i-1} \times \partial D_i \times D_{i+1} \times \cdots \times D_K.$$

Collecting (1.1a–e), (1.10), (1.11), (1.12), (1.26) and (1.27a–c), we then consider the following regularized initial-boundary-value problem, dependent on the following given regularization parameters $\kappa > 0$, $\alpha > 0$ and $L > 1$. As has been already emphasized in the Introduction, the centre-of-mass diffusion coefficient $\varepsilon > 0$ is a physical parameter and is regarded as being fixed throughout.

$(P_{\kappa, \alpha, L})$ Find $\rho_{\kappa, \alpha, L} : (\underline{x}, t) \in \Omega \times [0, T] \mapsto \rho_{\kappa, \alpha, L}(\underline{x}, t) \in \mathbb{R}$ and $\underline{u}_{\kappa, \alpha, L} : (\underline{x}, t) \in \bar{\Omega} \times [0, T] \mapsto \underline{u}_{\kappa, \alpha, L}(\underline{x}, t) \in \mathbb{R}^d$ such that

$$(2.2a) \quad \frac{\partial \rho_{\kappa, \alpha, L}}{\partial t} + \nabla_{\underline{x}} \cdot (\underline{u}_{\kappa, \alpha, L} \rho_{\kappa, \alpha, L}) - \alpha \Delta_{\underline{x}} \rho_{\kappa, \alpha, L} = 0 \quad \text{in } \Omega \times (0, T],$$

$$(2.2b) \quad \alpha \nabla_{\underline{x}} \rho_{\kappa, \alpha, L} \cdot \underline{n} = 0 \quad \text{on } \partial\Omega \times (0, T],$$

$$(2.2c) \quad \rho_{\kappa, \alpha, L}(\underline{x}, 0) = \rho^0(\underline{x}) \quad \forall \underline{x} \in \Omega,$$

$$(2.2d) \quad \frac{\partial(\rho_{\kappa,\alpha,L} u_{\kappa,\alpha,L})}{\partial t} - \frac{\alpha}{2} (\Delta_x \rho_{\kappa,\alpha,L}) u_{\kappa,\alpha,L} + \nabla_x \cdot (\rho_{\kappa,\alpha,L} u_{\kappa,\alpha,L} \otimes u_{\kappa,\alpha,L}) \\ - \nabla_x \cdot \underset{\sim}{S}(u_{\kappa,\alpha,L}, \rho_{\kappa,\alpha,L}) + \nabla_x p_{\kappa}(\rho_{\kappa,\alpha,L}) = \rho_{\kappa,\alpha,L} \underset{\sim}{f} + \nabla_x \cdot \underset{\sim}{\tau}(\psi_{\kappa,\alpha,L}) \quad \text{in } \Omega \times (0, T],$$

$$(2.2e) \quad u_{\kappa,\alpha,L} = 0 \quad \text{on } \partial\Omega \times (0, T],$$

$$(2.2f) \quad (\rho_{\kappa,\alpha,L} u_{\kappa,\alpha,L})(x, 0) = (\rho^0 u_0)(x) \quad \forall x \in \Omega,$$

where $\psi_{\kappa,\alpha,L} : (\underline{x}, \underline{q}, t) \in \overline{\Omega} \times \overline{D} \times [0, T] \mapsto \psi_{\kappa,\alpha,L}(\underline{x}, \underline{q}, t) \in \mathbb{R}$. Here $\underset{\sim}{S}(\cdot, \cdot)$ and $\underset{\sim}{\tau}(\cdot, \cdot)$ are given by (1.2) and (1.10), and $\rho^0(\alpha)$ is a regularization of ρ_0 , see (3.17) below. In addition, $p_{\kappa}(\cdot)$ is a regularization of $p(\cdot)$, (1.3), defined by

$$(2.3) \quad p_{\kappa}(s) := p(s) + \kappa(s^4 + s^{\Gamma}), \quad \text{where } \kappa \in \mathbb{R}_{>0} \text{ and } \Gamma = \max\{\gamma, 8\}.$$

The Fokker–Planck equation with microscopic cut-off satisfied by $\psi_{\kappa,\alpha,L}$ is:

$$(2.4) \quad \frac{\partial \psi_{\kappa,\alpha,L}}{\partial t} + \nabla_x \cdot \left(u_{\kappa,\alpha,L} M \beta^L \left(\frac{\psi_{\kappa,\alpha,L}}{M} \right) \right) + \sum_{i=1}^K \nabla_{q_i} \cdot \left[\sigma(u_{\kappa,\alpha,L}) q_i M \beta^L \left(\frac{\psi_{\kappa,\alpha,L}}{M} \right) \right] \\ = \varepsilon \Delta_x \psi_{\kappa,\alpha,L} + \frac{1}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \nabla_{q_i} \cdot \left(M \nabla_{q_j} \left(\frac{\psi_{\kappa,\alpha,L}}{M} \right) \right) \quad \text{in } \Omega \times D \times (0, T].$$

Here, for a given $L > 1$, $\beta^L \in C(\mathbb{R})$ is defined by (1.25), $\underset{\sim}{\sigma}(y) \equiv \nabla_x y$, and

$$(2.5) \quad \underset{\sim}{A} \in \mathbb{R}^{K \times K} \text{ is symmetric positive definite with smallest eigenvalue } a_0 \in \mathbb{R}_{>0}.$$

We impose the following boundary and initial conditions:

$$(2.6a) \quad \left[\frac{1}{4\lambda} \sum_{j=1}^K A_{ij} M \nabla_{q_j} \left(\frac{\psi_{\kappa,\alpha,L}}{M} \right) - \sigma(u_{\kappa,\alpha,L}) q_i M \beta^L \left(\frac{\psi_{\kappa,\alpha,L}}{M} \right) \right] \cdot \frac{q_i}{|q_i|} = 0 \\ \text{on } \Omega \times \partial \overline{D}_i \times (0, T], \quad i = 1, \dots, K,$$

$$(2.6b) \quad \varepsilon \nabla_x \psi_{\kappa,\alpha,L} \cdot \underline{n} = 0 \quad \text{on } \partial\Omega \times D \times (0, T],$$

$$(2.6c) \quad \psi_{\kappa,\alpha,L}(\cdot, \cdot, 0) = M(\cdot) \beta^L(\psi_0(\cdot, \cdot)/M(\cdot)) \geq 0 \quad \text{on } \Omega \times D,$$

where \underline{n} is the unit outward normal to $\partial\Omega$. The boundary conditions for $\psi_{\kappa,\alpha,L}$ on $\partial\Omega \times D \times (0, T]$ and $\Omega \times \partial D \times (0, T]$ have been chosen so as to ensure that

$$(2.7) \quad \int_{\Omega \times D} \psi_{\kappa,\alpha,L}(\underline{x}, \underline{q}, t) d\underline{q} d\underline{x} = \int_{\Omega \times D} \psi_{\kappa,\alpha,L}(\underline{x}, \underline{q}, 0) d\underline{q} d\underline{x} \quad \forall t \in (0, T].$$

Henceforth, we shall write

$$\widehat{\psi}_{\kappa,\alpha,L} = \frac{\psi_{\kappa,\alpha,L}}{M}, \quad \widehat{\psi}_0 = \frac{\psi_0}{M}.$$

Thus, for example, (2.6c) in terms of this compact notation becomes: $\widehat{\psi}_{\kappa,\alpha,L}(\cdot, \cdot, 0) = \beta^L(\widehat{\psi}_0(\cdot, \cdot))$ on $\Omega \times D$.

3. EXISTENCE OF A SOLUTION TO $(P_{\kappa,\alpha,L}^{\Delta t})$, A DISCRETE-IN-TIME APPROXIMATION OF $(P_{\kappa,\alpha,L})$

For later purposes, we recall the following Lebesgue interpolation result and the Gagliardo–Nirenberg inequality. Let $1 \leq r \leq v \leq s < \infty$, then, for any bounded Lipschitz domain \mathcal{O} ,

$$(3.1) \quad \|\eta\|_{L^v(\mathcal{O})} \leq \|\eta\|_{L^r(\mathcal{O})}^{1-\vartheta} \|\eta\|_{L^s(\mathcal{O})}^{\vartheta} \quad \forall \eta \in L^s(\mathcal{O}),$$

where $\vartheta = \frac{s(v-r)}{v(s-r)}$. Let $r \in [2, \infty)$ if $d = 2$, and $r \in [2, 6]$ if $d = 3$ and $\vartheta = d \left(\frac{1}{2} - \frac{1}{r} \right)$. Then, there is a constant $C = C(\Omega, r, d)$, such that

$$(3.2) \quad \|\eta\|_{L^r(\Omega)} \leq C \|\eta\|_{L^2(\Omega)}^{1-\vartheta} \|\eta\|_{H^1(\Omega)}^{\vartheta} \quad \forall \eta \in H^1(\Omega).$$

We note also the generalised Korn's inequality

$$(3.3) \quad \int_{\Omega} \left[|D(w)|^2 - \frac{1}{d} |\nabla_x \cdot w|^2 \right] dx = \int_{\Omega} |D(w) - \frac{1}{d} (\nabla_x \cdot w) I|^2 dx \geq c_0 \|w\|_{H^1(\Omega)}^2 \quad \forall w \in H_0^1(\Omega),$$

where $c_0 > 0$, see Dain [24]. We remark that the notation $|\cdot|$ will be used to signify one of the following. When applied to a real number x , $|x|$ will denote the absolute value of the number x ; when applied to a vector v , $|v|$ will stand for the Euclidean norm of the vector v ; and, when applied to a square matrix A , $|A|$ will signify the Frobenius norm, $[\text{tr}(A^T A)]^{\frac{1}{2}}$, of the matrix A , where, for a square matrix B , $\text{tr}(B)$ denotes the trace of B .

Let $\mathcal{F} \in C(\mathbb{R}_{>0})$ be defined by $\mathcal{F}(s) := s(\log s - 1) + 1$, $s > 0$. As $\lim_{s \rightarrow 0^+} \mathcal{F}(s) = 1$, the function \mathcal{F} can be considered to be defined and continuous on $[0, \infty)$, where it is a nonnegative, strictly convex function with $\mathcal{F}(1) = 0$.

We assume the following:

$$(3.4) \quad \begin{aligned} & \partial\Omega \in C^{2,\theta}, \theta \in (0, 1); \quad \rho_0 \in L_{\geq 0}^\infty(\Omega); \quad u_0 \in L^2(\Omega); \\ & \psi_0 \geq 0 \text{ a.e. on } \Omega \times D \text{ with } \mathcal{F}(\hat{\psi}_0) \in L_M^1(\Omega \times D) \text{ and } \int_D \psi_0(\cdot, q) dq \in L_{\geq 0}^\infty(\Omega); \\ & \mu^S \in \mathbb{R}_{>0}, \mu^B \in \mathbb{R}_{\geq 0}; \quad \text{the Rouse matrix } A \in \mathbb{R}^{K \times K} \text{ satisfies (2.5);} \\ & p, p_\kappa \in C^1(\mathbb{R}_{\geq 0}, \mathbb{R}_{\geq 0}) \text{ are defined by (1.3) and (2.3);} \\ & f \in L^2(0, T; L^\infty(\Omega)) \quad \text{and} \quad D_i = B(0, b_i^{\frac{1}{2}}), \quad \theta_i > 1, \quad i = 1, \dots, K, \quad \text{in (1.7a,b).} \end{aligned}$$

We introduce $P, P_\kappa \in C^1(\mathbb{R}_{\geq 0}, \mathbb{R}_{\geq 0})$, for $\kappa > 0$, such that

$$(3.5) \quad \begin{aligned} & s P'_{(\kappa)}(s) - P_{(\kappa)}(s) = p_{(\kappa)}(s) \quad \text{and} \quad P_{(\kappa)}(0) = P'_{(\kappa)}(0) = 0 \\ \Rightarrow \quad & P(s) = \frac{p(s)}{\gamma - 1} = \frac{c_p}{\gamma - 1} s^\gamma \quad \text{and} \quad P_\kappa(s) = P(s) + \kappa \left(\frac{s^4}{3} + \frac{s^\Gamma}{\Gamma - 1} \right). \end{aligned}$$

Here, and throughout, the subscript “ (\cdot) ” means with and without the subscript “ \cdot ”. We adopt a similar notation for superscripts.

In (3.4), $L_M^r(\Omega \times D)$, for $r \in [1, \infty)$, denotes the Maxwellian-weighted L^r space over $\Omega \times D$ with norm

$$\|\varphi\|_{L_M^r(\Omega \times D)} := \left\{ \int_{\Omega \times D} M |\varphi|^r dq dx \right\}^{\frac{1}{r}}.$$

Similarly, we introduce $L_M^r(D)$, the Maxwellian-weighted L^r space over D . Letting

$$(3.6) \quad \|\varphi\|_{H_M^1(\Omega \times D)} := \left\{ \int_{\Omega \times D} M \left[|\varphi|^2 + |\nabla_x \varphi|^2 + |\nabla_q \varphi|^2 \right] dq dx \right\}^{\frac{1}{2}},$$

we then set

$$(3.7) \quad X \equiv H_M^1(\Omega \times D) := \left\{ \varphi \in L_{\text{loc}}^1(\Omega \times D) : \|\varphi\|_{H_M^1(\Omega \times D)} < \infty \right\}.$$

It is shown in Appendix C of [10] (with the set X denoted by \widehat{X} there) that

$$(3.8) \quad C^\infty(\overline{\Omega \times D}) \text{ is dense in } X.$$

We have from Sobolev embedding that

$$(3.9) \quad H^1(\Omega; L_M^2(D)) \hookrightarrow L^s(\Omega; L_M^2(D)),$$

where $s \in [1, \infty)$ if $d = 2$ or $s \in [1, 6]$ if $d = 3$. In addition, we note that the embeddings

$$(3.10a) \quad H_M^1(D) \hookrightarrow L_M^2(D),$$

$$(3.10b) \quad H_M^1(\Omega \times D) \equiv L^2(\Omega; H_M^1(D)) \cap H^1(\Omega; L_M^2(D)) \hookrightarrow L_M^2(\Omega \times D) \equiv L^2(\Omega; L_M^2(D))$$

are compact if $\theta_i \geq 1$, $i = 1, \dots, K$, in (1.7a,b); see Appendix D of [10].

We recall the Aubin–Lions–Simon compactness theorem, see, e.g., Simon [64]. Let \mathfrak{X}_0 , \mathfrak{X} and \mathfrak{X}_1 be Banach spaces with a compact embedding $\mathfrak{X}_0 \hookrightarrow \mathfrak{X}$ and a continuous embedding $\mathfrak{X} \hookrightarrow \mathfrak{X}_1$. Then, for $\varsigma_i \in [1, \infty)$, $i = 0, 1$, the embedding

$$(3.11) \quad \{ \eta \in L^{\varsigma_0}(0, T; \mathfrak{X}_0) : \frac{\partial \eta}{\partial t} \in L^{\varsigma_1}(0, T; \mathfrak{X}_1) \} \hookrightarrow L^{\varsigma_0}(0, T; \mathfrak{X})$$

is compact. We recall also a generalization of the Aubin–Lions–Simon compactness theorem due to Dubinskii [28], see also Barrett & Süli [16]. Prior to stating Dubinskii’s theorem we introduce the necessary prerequisites.

Let \mathfrak{X} be a linear space over the field \mathbb{R} of real numbers, and suppose that \mathfrak{M} is a subset of \mathfrak{X} such that

$$(3.12) \quad \lambda \eta \in \mathfrak{M} \quad \forall \lambda \in \mathbb{R}_{\geq 0}, \quad \forall \eta \in \mathfrak{M}.$$

In other words, whenever η is contained in \mathfrak{M} , the ray through η from the origin of the linear space \mathfrak{X} is also contained in \mathfrak{M} . Note in particular that while any set \mathfrak{M} with property (3.12) must contain the zero element of the linear space \mathfrak{X} , the set \mathfrak{M} need not be closed under summation. The linear space \mathfrak{X} will be referred to as the *ambient space* for \mathfrak{M} . Suppose further that each element η of a set \mathfrak{M} with property (3.12) is assigned a certain real number, denoted by $[\eta]_{\mathfrak{M}}$, such that:

- (i) $[\eta]_{\mathfrak{M}} \geq 0$; and $[\eta]_{\mathfrak{M}} = 0$ if, and only if, $\eta = 0$; and
- (ii) $[\lambda \eta]_{\mathfrak{M}} = \lambda [\eta]_{\mathfrak{M}}$ for all $\lambda \in \mathbb{R}_{\geq 0}$ and all $\eta \in \mathfrak{X}$.

We shall then say that \mathfrak{M} is a *seminormed set*. A subset \mathfrak{B} of a seminormed set \mathfrak{M} is said to be *bounded* if there exists a positive constant K_0 such that $[\eta]_{\mathfrak{M}} \leq K_0$ for all $\eta \in \mathfrak{B}$. A seminormed set \mathfrak{M} contained in a normed linear space \mathfrak{X} with norm $\|\cdot\|_{\mathfrak{X}}$ is said to be *embedded in \mathfrak{X}* , and we write $\mathfrak{M} \subset \mathfrak{X}$, if there exists a $K_0 \in \mathbb{R}_{>0}$ such that

$$\|\eta\|_{\mathfrak{X}} \leq K_0 [\eta]_{\mathfrak{M}} \quad \forall \eta \in \mathfrak{M}.$$

Thus, bounded subsets of a seminormed set are also bounded subsets of the ambient normed linear space the seminormed set is embedded in. The embedding of a seminormed set \mathfrak{M} into a normed linear space \mathfrak{X} is said to be *compact* if from any bounded, infinite set of elements of \mathfrak{M} one can extract a subsequence that converges in \mathfrak{X} .

Theorem 3.1 (Dubinskii’s compactness theorem). *Suppose that \mathfrak{M} is a semi-normed set that is compactly embedded into a Banach space \mathfrak{X} , which is, in turn, continuously embedded into a Banach space \mathfrak{X}_1 . Then, for $\varsigma_i \in [1, \infty)$, $i = 0, 1$, the embedding*

$$(3.13) \quad \{ \eta \in L^{\varsigma_0}(0, T; \mathfrak{M}) : \frac{\partial \eta}{\partial t} \in L^{\varsigma_1}(0, T; \mathfrak{X}_1) \} \hookrightarrow L^{\varsigma_0}(0, T; \mathfrak{X})$$

is compact.

Let \mathfrak{X} be a Banach space. We shall denote by $C_w([0, T]; \mathfrak{X})$ the set of all functions $\eta \in L^\infty(0, T; \mathfrak{X})$ such that $t \in [0, T] \mapsto \langle \varphi, \eta(t) \rangle_{\mathfrak{X}} \in \mathbb{R}$ is continuous on $[0, T]$ for all $\varphi \in \mathfrak{X}'$, the dual space of \mathfrak{X} . Here, and throughout, $\langle \cdot, \cdot \rangle_{\mathfrak{X}}$ denotes the duality pairing between \mathfrak{X}' and \mathfrak{X} . Whenever \mathfrak{X} has a predual, \mathfrak{E} , say, (viz. $\mathfrak{E}' = \mathfrak{X}$), we shall denote by $C_{w*}([0, T]; \mathfrak{X})$ the set of all functions $\eta \in L^\infty(0, T; \mathfrak{X})$ such that $t \in [0, T] \mapsto \langle \eta(t), \zeta \rangle_{\mathfrak{E}} \in \mathbb{R}$ is continuous on $[0, T]$ for all $\zeta \in \mathfrak{E}$. We note the following result.

Lemma 3.1. *Let \mathfrak{X} and \mathfrak{Y} be Banach spaces.*

- (a) *Assume that the space \mathfrak{X} is reflexive and is continuously embedded in the space \mathfrak{Y} ; then, $L^\infty(0, T; \mathfrak{X}) \cap C_w([0, T]; \mathfrak{Y}) = C_w([0, T]; \mathfrak{X})$.*
- (b) *Assume that \mathfrak{X} has a separable predual \mathfrak{E} and \mathfrak{Y} has a predual \mathfrak{F} such that \mathfrak{F} is continuously embedded in \mathfrak{E} ; then, $L^\infty(0, T; \mathfrak{X}) \cap C_{w*}([0, T]; \mathfrak{Y}) = C_{w*}([0, T]; \mathfrak{X})$.*

Proof. Part (a) is due to Strauss [65] (cf. Lions & Magenes [44], Lemma 8.1, Ch. 3, Sec. 8.4); part (b) is proved analogously, *via* the sequential Banach–Alaoglu theorem. \square

We note from Lemma 3.1(a) above and Lemma 6.2 in Novotný & Straškraba [58] (or Lemma E.1 in Appendix E) that if $\{\eta_n\}_{n \in \mathbb{N}}$ is such that

$$(3.14a) \quad \|\eta_n\|_{L^\infty(0,T;L^r(\Omega))} + \left\| \frac{\partial \eta_n}{\partial t} \right\|_{L^\varsigma(0,T;W_0^{1,\nu}(\Omega)')} \leq C, \quad r, \varsigma, \nu \in (1, \infty),$$

then there exists a subsequence (not indicated) of $\{\eta_n\}_{n \in \mathbb{N}}$ and an $\eta \in C_w([0, T]; L^r(\Omega))$ such that

$$(3.14b) \quad \eta_n \rightarrow \eta \quad \text{in } C_w([0, T]; L^r(\Omega)).$$

Throughout we will assume that (3.4) hold, so that (1.8) and (3.10a,b) hold. We note for future reference that (1.12) and (1.8) yield that, for $\varphi \in L_M^2(\Omega \times D)$,

$$(3.15) \quad \begin{aligned} \int_{\Omega} |C_i(M\varphi)|^2 dx &\underset{\sim}{=} \int_{\Omega} \left| \int_D M \varphi U'_i q_i q_i^T dq \right|_{\sim}^2 dx \\ &\leq \left(\int_D M (U'_i)^2 |q_i|^4 dq \right)_{\sim} \left(\int_{\Omega \times D} M |\varphi|^2 dq dx \right)_{\sim} \\ &\leq C \left(\int_{\Omega \times D} M |\varphi|^2 dq dx \right)_{\sim}, \quad i = 1, \dots, K, \end{aligned}$$

where C is a positive constant.

We state a simple integration-by-parts formula.

Lemma 3.2. *Let $\varphi \in H_M^1(D)$ and suppose that $B \in \mathbb{R}^{d \times d}$; then,*

$$(3.16) \quad \int_D M \sum_{i=1}^K (B q_i) \cdot \nabla_{q_i} \varphi dq = \int_D M \varphi \left[\left(\sum_{i=1}^K U'_i \left(\frac{1}{2} |q_i|^2 \right) q_i q_i^T \right) - K I \right] : B dq.$$

Proof. By Theorem C.1 in Appendix C of Barrett & Süli [10], the set $C^\infty(\overline{D})$ is dense in $H_M^1(D)$; hence, for any $\widehat{\varphi} \in H_M^1(D)$ there exists a sequence $\{\widehat{\varphi}_n\}_{n \geq 0} \subset C^\infty(\overline{D})$ converging to $\widehat{\varphi}$ in $H_M^1(D)$. As $M \in C^1(\overline{D})$ and vanishes on ∂D , the same is true of each of the functions $M\widehat{\varphi}_n$, $n \geq 1$. By replacing $\widehat{\varphi}$ by $\widehat{\varphi}_n$ on both sides of (3.16), the resulting identity is easily verified by using the classical divergence theorem for smooth functions, noting (1.6a) and that $M\widehat{\varphi}_n$ vanishes on ∂D . Then, (3.16) itself follows by letting $n \rightarrow \infty$, recalling the definition of the norm in $H_M^1(D)$ and hypothesis (1.8). \square

We now formulate our discrete-in-time approximation of problem $(P_{\kappa, \alpha, L})$ for fixed parameters $\kappa, \alpha \in (0, 1]$ and $L > 1$. For any $T > 0$ and $N \geq 1$, let $N \Delta t = T$ and $t_n = n \Delta t$, $n = 0, \dots, N$. To prove existence of a solution under minimal smoothness requirements on the initial data, recall (3.4), we regularize the initial data in terms of the parameters $\alpha, \Delta t$ and L . For $\rho_0 \in L_{\geq 0}^\infty(\Omega)$, we assign to it the function $\rho^0(\alpha) \in H^1(\Omega)$, appearing in (2.2c,f), defined as the unique solution of

$$(3.17) \quad \int_{\Omega} \left[\rho^0 \eta + \alpha \nabla_x \rho^0 \cdot \nabla_x \eta \right] dx = \int_{\Omega} \rho_0 \eta dx \quad \forall \eta \in H^1(\Omega).$$

Hence,

$$(3.18a) \quad \rho^0(\cdot) \in [0, \|\rho_0\|_{L^\infty(\Omega)}],$$

$$(3.18b) \quad \text{and} \quad \rho^0 \rightarrow \rho_0 \quad \text{weakly-}\star \text{ in } L^\infty(\Omega), \quad \text{strongly in } L^2(\Omega), \quad \text{as } \alpha \rightarrow 0_+.$$

Therefore, by (3.1) and (3.18a,b), also

$$(3.18c) \quad \rho^0 \rightarrow \rho_0 \quad \text{strongly in } L^p(\Omega), \quad \text{as } \alpha \rightarrow 0_+, \quad p \in [1, \infty).$$

Similarly, for $u_0 \in L^2(\Omega)$ we assign to it the function $u^0 = u^0(\alpha, \Delta t) \in H_0^1(\Omega)$, defined as the unique solution of

$$(3.19) \quad \int_{\Omega} \left[\rho^0 u^0 \cdot v + \Delta t \nabla_x u^0 : \nabla_x v \right] dx = \int_{\Omega} \rho^0 u_0 \cdot v dx \quad \forall v \in H_0^1(\Omega).$$

Hence, it follows from (3.19) and (3.18a) that there exists a $C \in \mathbb{R}_{>0}$, independent of Δt , L , α and κ , such that

$$(3.20a) \quad \int_{\Omega} \left[\rho^0 |\tilde{u}^0|^2 + \Delta t |\tilde{\nabla}_x \tilde{u}^0|^2 \right] d\tilde{x} \leq \int_{\Omega} \rho^0 |\tilde{u}_0|^2 d\tilde{x} \leq C,$$

$$(3.20b) \quad \text{and} \quad \int_{\Omega} \rho^0 (\tilde{u}^0 - \tilde{u}_0) \cdot \tilde{v} d\tilde{x} \rightarrow 0 \quad \forall \tilde{v} \in L^2(\Omega), \quad \text{as } \Delta t \rightarrow 0_+.$$

Analogously, we shall assign a certain ‘smoothed’ initial datum, $\hat{\psi}^0 = \hat{\psi}^0(L, \Delta t) \in H_M^1(\Omega \times D)$, to the given initial datum $\hat{\psi}_0 = \frac{\psi_0}{M}$ such that

$$(3.21) \quad \begin{aligned} \int_{\Omega \times D} M \left[\hat{\psi}^0 \varphi + \Delta t \left(\tilde{\nabla}_x \hat{\psi}^0 \cdot \tilde{\nabla}_x \varphi + \tilde{\nabla}_q \hat{\psi}^0 \cdot \tilde{\nabla}_q \varphi \right) \right] d\tilde{q} d\tilde{x} \\ = \int_{\Omega \times D} M \beta^L(\hat{\psi}_0) \varphi d\tilde{q} d\tilde{x} \quad \forall \varphi \in H_M^1(\Omega \times D). \end{aligned}$$

For $r \in [1, \infty)$, let

$$(3.22) \quad Z_r := \{\varphi \in L_M^r(\Omega \times D) : \varphi \geq 0 \text{ a.e. on } \Omega \times D\}.$$

It is proved in the Appendix of [15] that there exists a unique $\hat{\psi}^0 \in H_M^1(\Omega \times D)$ satisfying (3.21); furthermore, $\hat{\psi}^0 \in Z_2$,

$$(3.23a) \quad \int_{\Omega \times D} M \mathcal{F}(\hat{\psi}^0) d\tilde{q} d\tilde{x} + 4 \Delta t \int_{\Omega \times D} M \left[|\tilde{\nabla}_x \sqrt{\hat{\psi}^0}|^2 + |\tilde{\nabla}_q \sqrt{\hat{\psi}^0}|^2 \right] d\tilde{q} d\tilde{x} \leq \int_{\Omega \times D} M \mathcal{F}(\hat{\psi}_0) d\tilde{q} d\tilde{x}$$

and

$$(3.23b) \quad \hat{\psi}^0 = \beta^L(\hat{\psi}^0) \rightarrow \hat{\psi}_0 \quad \text{weakly in } L_M^1(\Omega \times D), \quad \text{as } L \rightarrow \infty, \quad \Delta t \rightarrow 0_+.$$

Finally, by choosing $\varphi(\tilde{x}, \tilde{q}) = \varrho^0(\tilde{x}) \otimes 1(\tilde{q})$ in (3.21), where $\varrho^0(\tilde{x}) := \int_D M(\tilde{q}) \hat{\psi}^0(\tilde{x}, \tilde{q}) d\tilde{q}$, $\tilde{x} \in \Omega$, yields, on noting (1.25) and (3.4), that

$$(3.24) \quad \begin{aligned} \frac{1}{2} \left[\|\varrho^0\|_{L^2(\Omega)}^2 + \int_{\Omega} \left(\int_D M \left(\hat{\psi}^0 - \beta^L(\hat{\psi}_0) \right) d\tilde{q} \right)^2 d\tilde{x} \right] + \Delta t \|\tilde{\nabla}_x \varrho^0\|_{L^2(\Omega)}^2 \\ = \frac{1}{2} \int_{\Omega} \left(\int_D M \beta^L(\hat{\psi}_0) d\tilde{q} \right)^2 d\tilde{x} \leq \frac{1}{2} \int_{\Omega} \left(\int_D M \hat{\psi}_0 d\tilde{q} \right)^2 d\tilde{x} \leq C. \end{aligned}$$

Next, we define

$$(3.25) \quad V := \{w \in H_0^1(\Omega) : w \in L^\infty(\Omega), \tilde{\nabla}_x \cdot w \in L^\infty(\Omega)\}$$

and

$$(3.26) \quad Y^n := L^2(t_{n-1}, t_n; H^1(\Omega)) \cap H^1(t_{n-1}, t_n; H^1(\Omega)') \cap L^\infty(t_{n-1}, t_n; L^{\frac{\Gamma}{2}}(\Omega)) \cap L_{\geq 0}^{\frac{4\Gamma}{3}}(\Omega \times (t_{n-1}, t_n)).$$

We recall also that, for all $v, w \in \tilde{H}^1(\Omega)$,

$$(3.27) \quad (\tilde{v} \otimes \tilde{v}) : \tilde{\nabla}_x w = [(\tilde{v} \cdot \tilde{\nabla}_x) w] \cdot \tilde{v} = -[(\tilde{v} \cdot \tilde{\nabla}_x) \tilde{v}] \cdot w + (\tilde{v} \cdot \tilde{\nabla}_x)(\tilde{v} \cdot w) \quad \text{a.e. in } \Omega.$$

Noting the above, our discrete-in-time approximation of $(P_{\kappa, \alpha, L})$ is then defined as follows.

($\mathbf{P}_{\kappa, \alpha, L}^{\Delta t}$) Let $N \in \mathbb{N}_{\geq 1}$ and set $\Delta t := T/N$; let, further, $\rho_{\kappa, \alpha, L}^0 := \rho^0 \in L_{\geq 0}^\infty(\Omega)$, $u_{\kappa, \alpha, L}^0 := u^0 \in H_0^1(\Omega)$ and $\hat{\psi}_{\kappa, \alpha, L}^0 := \hat{\psi}^0 \in Z_2$. For $n = 1, \dots, N$, and given $(\rho_{\kappa, \alpha, L}^{n-1}, u_{\kappa, \alpha, L}^{n-1}, \hat{\psi}_{\kappa, \alpha, L}^{n-1}) \in L_{\geq 0}^\Gamma(\Omega) \times \tilde{H}_0^1(\Omega) \times Z_2$, find

$$(3.28) \quad \rho_{\kappa, \alpha, L}^{[\Delta t], n} \in Y^n \text{ with } \rho_{\kappa, \alpha, L}^n(\cdot) := \rho_{\kappa, \alpha, L}^{[\Delta t], n}(\cdot, t_n) \in L_{\geq 0}^\Gamma(\Omega), \quad u_{\kappa, \alpha, L}^n \in H_0^1(\Omega) \text{ and } \hat{\psi}_{\kappa, \alpha, L}^n \in X \cap Z_2,$$

such that $\rho_{\kappa,\alpha,L}^{[\Delta t],n}(\cdot, t_{n-1}) = \rho_{\kappa,\alpha,L}^{n-1}(\cdot)$,

$$(3.29a) \quad \int_{t_{n-1}}^{t_n} \left[\left\langle \frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t],n}}{\partial t}, \eta \right\rangle_{H^1(\Omega)} + \int_{\Omega} \left(\alpha \nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t],n} - \rho_{\kappa,\alpha,L}^{[\Delta t],n} u_{\kappa,\alpha,L}^n \right) \cdot \nabla_x \eta \, dx \right] dt = 0$$

$$\forall \eta \in L^2(t_{n-1}, t_n; H^1(\Omega)),$$

$$(3.29b) \quad \begin{aligned} & \int_{\Omega} \left[\frac{\rho_{\kappa,\alpha,L}^n u_{\kappa,\alpha,L}^n - \rho_{\kappa,\alpha,L}^{n-1} u_{\kappa,\alpha,L}^{n-1}}{\Delta t} - \frac{1}{2} \frac{\rho_{\kappa,\alpha,L}^n - \rho_{\kappa,\alpha,L}^{n-1}}{\Delta t} u_{\kappa,\alpha,L}^n \right] \cdot w \, dx \\ & + \frac{1}{2} \int_{\Omega} \rho_{\kappa,\alpha,L}^{n-1} \left[[(u_{\kappa,\alpha,L}^{n-1} \cdot \nabla_x) u_{\kappa,\alpha,L}^n] \cdot w - [(u_{\kappa,\alpha,L}^{n-1} \cdot \nabla_x) w] \cdot u_{\kappa,\alpha,L}^n \right] \, dx \\ & + \int_{\Omega} S(u_{\kappa,\alpha,L}^n) : \nabla_x w \, dx - \int_{\Omega} \left(\frac{1}{\Delta t} \int_{t_{n-1}}^{t_n} p_{\kappa}(\rho_{\kappa,\alpha,L}^{[\Delta t],n}) \, dt \right) \nabla_x \cdot w \, dx \\ & = \int_{\Omega} \rho_{\kappa,\alpha,L}^n f^n \cdot w \, dx - \int_{\Omega} \tau_1(M \hat{\psi}_{\kappa,\alpha,L}^n) : \nabla_x w \, dx \\ & - 2 \int_{\Omega} \left(\int_D M \beta^L(\hat{\psi}_{\kappa,\alpha,L}^n) \, dq \right) \nabla_x \left(\int_D M \hat{\psi}_{\kappa,\alpha,L}^n \, dq \right) \cdot w \, dx \quad \forall w \in V, \end{aligned}$$

$$(3.29c) \quad \begin{aligned} & \int_{\Omega \times D} M \frac{\hat{\psi}_{\kappa,\alpha,L}^n - \hat{\psi}_{\kappa,\alpha,L}^{n-1}}{\Delta t} \varphi \, dq \, dx \\ & + \sum_{i=1}^K \int_{\Omega \times D} M \left[\frac{1}{4\lambda} \sum_{j=1}^K A_{ij} \nabla_{q_j} \hat{\psi}_{\kappa,\alpha,L}^n - [\sigma(u_{\kappa,\alpha,L}^n) q_i] \beta^L(\hat{\psi}_{\kappa,\alpha,L}^n) \right] \cdot \nabla_{q_i} \varphi \, dq \, dx \\ & + \int_{\Omega \times D} M \left[\varepsilon \nabla_x \hat{\psi}_{\kappa,\alpha,L}^n - \beta_L(\hat{\psi}_{\kappa,\alpha,L}^n) u_{\kappa,\alpha,L}^n \right] \cdot \nabla_x \varphi \, dq \, dx = 0 \quad \forall \varphi \in X; \end{aligned}$$

where, for $t \in (t_{n-1}, t_n]$ and $n = 1, \dots, N$,

$$(3.30) \quad f_{\sim}^{\{\Delta t\}}(\cdot, t) = f_{\sim}^n(\cdot) := \frac{1}{\Delta t} \int_{t_{n-1}}^{t_n} f(\cdot, t) \, dt \in L^{\infty}(\Omega).$$

It follows from (3.4) and (3.30) that

$$(3.31a) \quad \int_{t_{n-1}}^{t_n} \|f_{\sim}^{\{\Delta t\}}\|_{L^{\infty}(\Omega)}^2 \, dt \leq \int_{t_{n-1}}^{t_n} \|f\|_{L^{\infty}(\Omega)}^2 \, dt, \quad n = 1, \dots, N,$$

$$(3.31b) \quad f_{\sim}^{\{\Delta t\}} \rightarrow f_{\sim} \quad \text{strongly in } L^2(0, T; L^r(\Omega)), \quad \text{as } \Delta t \rightarrow 0_+,$$

where $r \in [1, \infty]$.

Remark 3.1. A possible alternative to our temporal approximation scheme (3.29a) for (2.2a), which is the weak formulation of a parabolic initial boundary-value problem posed over the time slab $\Omega \times [t_{n-1}, t_n]$, $n = 1, \dots, N$, would have been to proceed as in the work of Abels and Feireisl [1] and approximate (2.2a) by an implicit finite difference scheme with respect to t . That would have avoided the use of $\rho_{\kappa,\alpha,L}^{[\Delta t],n}$ here, but would have had the disadvantage, from the point of view of constructive considerations in numerical analysis at least, that nonnegativity of $\rho_{\kappa,\alpha,L}^n$ will have been guaranteed for $\Delta t \leq \Delta t_0$ only, where $\Delta t_0 \in (0, T]$ is sufficiently small, with the value of Δt_0 not being easily quantifiable in terms of the data and its independence of κ , α and L being less than obvious. In contrast with that, our $\rho_{\kappa,\alpha,L}^{[\Delta t],n}$, and thereby also $\rho_{\kappa,\alpha,L}^n$, will be shown to be nonnegative for all $\Delta t = \frac{T}{N}$ and all $n = 1, \dots, N$, regardless of the choice of κ , α , L and N .

We rewrite (3.29a) as

$$(3.32) \quad \int_{t_{n-1}}^{t_n} \left[\left\langle \frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t],n}}{\partial t}, \eta \right\rangle_{H^1(\Omega)} + c(u_{\kappa,\alpha,L}^n)(\rho_{\kappa,\alpha,L}^{[\Delta t],n}, \eta) \right] dt = 0 \quad \forall \eta \in L^2(t_{n-1}, t_n; H^1(\Omega)),$$

where, for all $v \in \tilde{H}_0^1(\Omega)$ and $\eta_i \in H^1(\Omega)$, $i = 1, 2$,

$$(3.33) \quad c(v)(\eta_1, \eta_2) := \int_{\Omega} \left(\alpha \nabla_x \eta_1 - \eta_1 v \right) \cdot \nabla_x \eta_2 \, dx.$$

Similarly, on noting (1.2) and (1.11), we rewrite (3.29b) as

$$(3.34) \quad b(\rho_{\kappa,\alpha,L}^n)(u_{\kappa,\alpha,L}^n, w) = \ell_b(\rho_{\kappa,\alpha,L}^{[\Delta t],n}, \hat{\psi}_{\kappa,\alpha,L}^n)(w) \quad \forall w \in V;$$

where, for all $\eta \in L_{\geq 0}^2(\Omega)$ and $w_i \in \tilde{H}_0^1(\Omega)$, $i = 1, 2$,

$$(3.35a) \quad \begin{aligned} b(\eta)(w_1, w_2) &:= \frac{1}{2} \int_{\Omega} (\eta + \rho_{\kappa,\alpha,L}^{n-1}) w_1 \cdot w_2 \, dx + \Delta t \mu^S \int_{\Omega} D(w_1) : D(w_2) \, dx \\ &\quad + \Delta t \left(\mu^B - \frac{\mu^S}{d} \right) \int_{\Omega} (\nabla_x \cdot w_1) (\nabla_x \cdot w_2) \, dx \\ &\quad + \frac{1}{2} \int_{\Omega} \rho_{\kappa,\alpha,L}^{n-1} \left[[(u_{\kappa,\alpha,L}^{n-1} \cdot \nabla_x) w_1] \cdot w_2 - [(u_{\kappa,\alpha,L}^{n-1} \cdot \nabla_x) w_2] \cdot w_1 \right] \, dx \end{aligned}$$

and, for all $\eta \in Y^n$ with $\eta(\cdot, t_n) \in L_{\geq 0}^2(\Omega)$, $\varphi \in X$ and $w \in V$,

$$(3.35b) \quad \begin{aligned} \ell_b(\eta, \varphi)(w) &:= \int_{\Omega} \left[\rho_{\kappa,\alpha,L}^{n-1} u_{\kappa,\alpha,L}^{n-1} \cdot w + \Delta t \eta(\cdot, t_n) f^n \cdot w - \Delta t k \sum_{i=1}^K C_i(M \varphi) : \nabla_x w \right] \, dx \\ &\quad + \int_{\Omega} \left(\int_{t_{n-1}}^{t_n} p_{\kappa}(\eta) \, dt + \Delta t k (K+1) \int_D M \varphi \, dq \right) \nabla_x \cdot w \, dx \\ &\quad - 2 \Delta t \mathfrak{z} \int_{\Omega} \left(\int_D M \beta^L(\varphi) \, dq \right) \nabla_x \left(\int_D M \varphi \, dq \right) \cdot w \, dx. \end{aligned}$$

It follows for fixed $u_{\kappa,\alpha,L}^{n-1} \in \tilde{H}_0^1(\Omega)$, $\rho_{\kappa,\alpha,L}^{n-1} \in L_{\geq 0}^{\Gamma}(\Omega)$ and $\eta \in L_{\geq 0}^2(\Omega)$, and the generalised Korn's inequality, (3.3), that the nonsymmetric bilinear functional $b(\eta)(\cdot, \cdot)$ is a nonsymmetric continuous coercive bilinear functional on $\tilde{H}_0^1(\Omega) \times \tilde{H}_0^1(\Omega)$. In addition, for fixed $u_{\kappa,\alpha,L}^{n-1} \in \tilde{H}_0^1(\Omega)$, $\rho_{\kappa,\alpha,L}^{n-1} \in L_{\geq 0}^{\Gamma}(\Omega)$, $\eta \in Y^n$ with $\eta(\cdot, t_n) \in L_{\geq 0}^2(\Omega)$ and $\varphi \in X$, it follows, on recalling (3.15), (2.3) and (3.26), that $\ell_b(\eta, \varphi)(\cdot)$ is a continuous linear functional on V .

It is also convenient to rewrite (3.29c) as

$$(3.36) \quad a(\hat{\psi}_{\kappa,\alpha,L}^n, \varphi) = \ell_a(u_{\kappa,\alpha,L}^n, \beta^L(\hat{\psi}_{\kappa,\alpha,L}^n))(\varphi) \quad \forall \varphi \in X,$$

where, for all $\varphi_i \in X$, $i = 1, 2$,

$$(3.37a) \quad a(\varphi_1, \varphi_2) := \int_{\Omega \times D} M \left[\varphi_1 \varphi_2 + \Delta t \varepsilon \nabla_x \varphi_1 \cdot \nabla_x \varphi_2 + \frac{\Delta t}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \nabla_{q_j} \varphi_1 \cdot \nabla_{q_i} \varphi_2 \right] \, dq \, dx$$

and, for all $v \in \tilde{H}^1(\Omega)$, $\xi \in L^{\infty}(\Omega \times D)$ and $\varphi \in X$,

$$(3.37b) \quad \ell_a(v, \xi)(\varphi) := \int_{\Omega \times D} M \left[\hat{\psi}_{\kappa,\alpha,L}^{n-1} \varphi + \Delta t \xi \left(\sum_{i=1}^K [\sigma(v) q_i] \cdot \nabla_{q_i} \varphi + v \cdot \nabla_x \varphi \right) \right] \, dq \, dx.$$

Clearly, $a(\cdot, \cdot)$ is a symmetric continuous coercive bilinear functional on $X \times X$. In addition, it is easily deduced for fixed $v \in \tilde{H}^1(\Omega)$ and $\xi \in L^{\infty}(\Omega \times D)$, on noting (3.2), that $\ell_a(v, \xi)(\cdot)$ is a continuous linear functional on X .

In order to prove existence of a solution, (3.28), to $(P_{\kappa,\alpha,L}^{\Delta t})$, $\rho_{\kappa,\alpha,L}^{[\Delta t],n}(\cdot, t_{n-1}) = \rho_{\kappa,\alpha,L}^{n-1}(\cdot)$ and (3.29a–c), which is equivalent to (3.32), (3.34) and (3.36), we require two convex regularizations of the entropy function $\mathcal{F} : s \in \mathbb{R}_{\geq 0} \mapsto \mathcal{F}(s) = s(\log s - 1) + 1 \in \mathbb{R}_{\geq 0}$, denoted by \mathcal{F}^L and \mathcal{F}_{δ}^L .

For any $L > 1$, we define $\mathcal{F}^L \in C(\mathbb{R}_{\geq 0}) \cap C_{\text{loc}}^{2,1}(\mathbb{R}_{>0})$ by

$$(3.38) \quad \mathcal{F}^L(s) := \begin{cases} s(\log s - 1) + 1, & 0 \leq s \leq L, \\ \frac{s^2 - L^2}{2L} + s(\log L - 1) + 1, & L \leq s. \end{cases}$$

Note that

$$(3.39a) \quad [\mathcal{F}^L]'(s) = \begin{cases} \log s, & 0 < s \leq L, \\ \frac{s}{L} + \log L - 1, & L \leq s, \end{cases}$$

and

$$(3.39b) \quad [\mathcal{F}^L]''(s) = \begin{cases} \frac{1}{s}, & 0 < s \leq L, \\ \frac{1}{L}, & L \leq s. \end{cases}$$

Hence, on noting the definition (1.25) of β^L , we have that

$$(3.40a) \quad \beta^L(s) = \min\{s, L\} = ([\mathcal{F}^L]''(s))^{-1}, \quad s \in \mathbb{R}_{\geq 0},$$

with the convention $\frac{1}{\infty} := 0$ when $s = 0$, and

$$(3.40b) \quad [\mathcal{F}^L]''(s) \geq \mathcal{F}''(s) = \frac{1}{s}, \quad s \in \mathbb{R}_{>0}.$$

We shall also require the following inequality, relating \mathcal{F}^L to \mathcal{F} :

$$(3.41) \quad \mathcal{F}^L(s) \geq \mathcal{F}(s), \quad s \in \mathbb{R}_{\geq 0}.$$

For $0 \leq s \leq 1$, (3.41) trivially holds, with equality. For $s \geq 1$, it follows from (3.40b), with s replaced by a dummy variable σ , after integrating twice over $\sigma \in [1, s]$, and noting that $[\mathcal{F}^L]'(1) = \mathcal{F}'(1)$ and $\mathcal{F}^L(1) = \mathcal{F}(1)$.

For $L > 1$ and $\delta \in (0, 1)$, the function $\mathcal{F}_\delta^L \in C^{2,1}(\mathbb{R})$ is defined by

$$(3.42) \quad \mathcal{F}_\delta^L(s) := \begin{cases} \frac{s^2 - \delta^2}{2\delta} + s(\log \delta - 1) + 1, & s \leq \delta, \\ \mathcal{F}^L(s), & \delta \leq s. \end{cases}$$

Hence,

$$(3.43a) \quad [\mathcal{F}_\delta^L]'(s) = \begin{cases} \frac{s}{\delta} + \log \delta - 1, & s \leq \delta, \\ [\mathcal{F}^L]'(s), & \delta \leq s, \end{cases}$$

$$(3.43b) \quad [\mathcal{F}_\delta^L]''(s) = \begin{cases} \frac{1}{\delta}, & s \leq \delta, \\ [\mathcal{F}^L]''(s), & \delta \leq s. \end{cases}$$

We note that

$$(3.44a) \quad \mathcal{F}_\delta^L(s) \leq \mathcal{F}^L(s) \quad \forall s \geq 0,$$

$$(3.44b) \quad \mathcal{F}_\delta^L(s) \geq \begin{cases} \frac{s^2}{2\delta}, & s \leq 0, \\ \frac{s}{4L} - C(L), & s \geq 0; \end{cases}$$

and that $[\mathcal{F}_\delta^L]''(s)$ is bounded below by $\frac{1}{L}$ for all $s \in \mathbb{R}$. Finally, we set

$$(3.45) \quad \beta_\delta^L(s) := ([\mathcal{F}_\delta^L]''(s))^{-1}(s) = \max\{\beta^L(s), \delta\},$$

and observe that $\beta_\delta^L(s)$ is bounded above by L and bounded below by δ for all $s \in \mathbb{R}$. Note also that both β^L and β_δ^L are Lipschitz continuous on \mathbb{R} , with Lipschitz constants equal to 1.

In addition, we regularize the bilinear functional $b(\eta)(\cdot, \cdot)$, (3.35a), on $X \times X$, by introducing the Banach space

$$(3.46) \quad \mathcal{V} := \{w \in H^2(\Omega) \cap H_0^1(\Omega) : \nabla_x \cdot w \in H^2(\Omega)\},$$

which is compactly embedded in \mathcal{V} , (3.25). Then we define for $\delta \in \mathbb{R}_{>0}$, $\eta \in L^2_{\geq 0}(\Omega)$ and $w_i \in \mathcal{V}$, $i = 1, 2$,

$$(3.47) \quad b_\delta(\eta)(w_1, w_2) := b(\eta)(w_1, w_2) + \Delta t \delta \sum_{|\lambda|=2} \int_{\Omega} \left[\frac{\partial^{|\lambda|} w_1}{\partial_{x_1}^{\lambda_1} \dots \partial_{x_d}^{\lambda_d}} \cdot \frac{\partial^{|\lambda|} w_2}{\partial_{x_1}^{\lambda_1} \dots \partial_{x_d}^{\lambda_d}} + \frac{\partial^{|\lambda|} (\nabla_x \cdot w_1)}{\partial_{x_1}^{\lambda_1} \dots \partial_{x_d}^{\lambda_d}} \frac{\partial^{|\lambda|} (\nabla_x \cdot w_2)}{\partial_{x_1}^{\lambda_1} \dots \partial_{x_d}^{\lambda_d}} \right] dx.$$

It follows that $b_\delta(\eta)(\cdot, \cdot)$ is a nonsymmetric continuous coercive bilinear functional on $\mathcal{V} \times \mathcal{V}$ for fixed $\eta \in L^2_{\geq 0}(\Omega)$ and $\delta \in (0, 1)$. We also replace $\ell_b(\eta, \varphi)(\cdot)$ by $\ell_{b,\delta}(\eta, \varphi)(\cdot)$, where $\beta^L(\varphi)$ in $\ell_b(\eta, \varphi)(\cdot)$ is replaced by $\beta_\delta^L(\varphi)$; that is, for $\delta \in (0, 1)$,

$$(3.48) \quad \ell_{b,\delta}(\eta, \varphi)(w) := \ell_b(\eta, \varphi)(w) + 2 \Delta t \mathfrak{z} \int_{\Omega} \left(\int_D M [\beta^L(\varphi) - \beta_\delta^L(\varphi)] dq \right) \nabla_x \left(\int_D M \varphi dq \right) \cdot w dx.$$

As for fixed $\eta \in Y^n$ with $\eta(\cdot, t_n) \in L^2_{\geq 0}(\Omega)$ and $\varphi \in X$, $\ell_b(\eta, \varphi)(\cdot)$ is a continuous linear functional on \mathcal{V} , it follows that $\ell_{b,\delta}(\eta, \varphi)(\cdot)$, for $\delta \in (0, 1)$, is a continuous linear functional on \mathcal{V} .

Next, we introduce

$$(3.49) \quad \Upsilon^n := L^2(t_{n-1}, t_n; H^1(\Omega)) \cap H^1(t_{n-1}, t_n; H^1(\Omega)') \cap L^\infty_{\geq 0}(\Omega \times (t_{n-1}, t_n)).$$

We note that $\Upsilon^n \hookrightarrow C([t_{n-1}, t_n]; L^2_{\geq 0}(\Omega))$ and $\Upsilon^n \subset Y^n$, (cf. (3.26)). Finally, we regularize the initial data and set

$$(3.50) \quad \rho_{\kappa,\alpha,L,\delta}^{n-1} = \beta^{\delta^{-1}}(\rho_{\kappa,\alpha,L}^{n-1}),$$

where $\beta^{\delta^{-1}}$ is given by (1.25) with $L = \delta^{-1}$.

We now consider the following regularized version of the coupled system (3.32), (3.34) and (3.36) for a given $\delta \in (0, 1)$:

Given $(\rho_{\kappa,\alpha,L}^{n-1}, u_{\kappa,\alpha,L}^{n-1}, \widehat{\psi}_{\kappa,\alpha,L}^{n-1}) \in L^\Gamma_{\geq 0}(\Omega) \times \widetilde{H}_0^1(\Omega) \times Z_2$, find $(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}, u_{\kappa,\alpha,L,\delta}^n, \widehat{\psi}_{\kappa,\alpha,L,\delta}^n) \in \Upsilon^n \times \mathcal{V} \times X$ such that $\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t_{n-1}) = \rho_{\kappa,\alpha,L,\delta}^{n-1}(\cdot)$,

$$(3.51a) \quad \int_{t_{n-1}}^{t_n} \left[\left\langle \frac{\partial \rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}}{\partial t}, \eta \right\rangle_{H^1(\Omega)} + c(u_{\kappa,\alpha,L,\delta}^n)(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}, \eta) \right] dt = 0 \quad \forall \eta \in L^2(t_{n-1}, t_n; H^1(\Omega)),$$

$$(3.51b) \quad b_\delta(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t_n))(u_{\kappa,\alpha,L,\delta}^n, w) = \ell_{b,\delta}(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}, \widehat{\psi}_{\kappa,\alpha,L,\delta}^n)(w) \quad \forall w \in \mathcal{V},$$

$$(3.51c) \quad a(\widehat{\psi}_{\kappa,\alpha,L,\delta}^n, \varphi) = \ell_a(u_{\kappa,\alpha,L,\delta}^n, \beta_\delta^L(\widehat{\psi}_{\kappa,\alpha,L,\delta}^n))(\varphi) \quad \forall \varphi \in X.$$

The existence of a solution to (3.51a–c) will be proved by using a fixed-point argument. Given $(\widetilde{u}, \widetilde{\psi}) \in \mathcal{V} \times L_M^2(\Omega \times D)$, let $(\rho^*, u^*, \psi^*) \in \Upsilon^n \times \mathcal{V} \times X$ be such that $\rho^*(\cdot, t_{n-1}) = \rho_{\kappa,\alpha,L,\delta}^{n-1}(\cdot)$,

$$(3.52a) \quad \int_{t_{n-1}}^{t_n} \left[\left\langle \frac{\partial \rho^*}{\partial t}, \eta \right\rangle_{H^1(\Omega)} + c(\widetilde{u})(\rho^*, \eta) \right] dt = 0 \quad \forall \eta \in L^2(t_{n-1}, t_n; H^1(\Omega)),$$

$$(3.52b) \quad a(\psi^*, \varphi) = \ell_a(\widetilde{u}, \beta_\delta^L(\widetilde{\psi}))(\varphi) \quad \forall \varphi \in X,$$

$$(3.52c) \quad b_\delta(\rho^*(\cdot, t_n))(u^*, w) = \ell_{b,\delta}(\rho^*, \psi^*)(w) \quad \forall w \in \mathcal{V}.$$

For fixed $v \in \mathcal{V}$, it follows that $c(v)(\cdot, \cdot)$, (3.33), is a nonsymmetric continuous bilinear functional on $H^1(\Omega) \times H^1(\Omega)$, and, moreover, for all $\eta \in H^1(\Omega)$,

$$(3.53) \quad c(v)(\eta, \eta) = \alpha \|\nabla_x \eta\|_{L^2(\Omega)}^2 + \frac{1}{2} \int_{\Omega} (\nabla_x \cdot v) \eta^2 dx \geq \alpha \|\nabla_x \eta\|_{L^2(\Omega)}^2 - \frac{1}{2} \|\nabla_x \cdot v\|_{L^\infty(\Omega)} \|\eta\|_{L^2(\Omega)}^2.$$

Hence for any fixed $\tilde{u} \in \mathcal{V}$, the existence of a unique weak solution

$$(3.54) \quad \rho^* \in L^2(t_{n-1}, t_n; H^1(\Omega)) \cap H^1(t_{n-1}, t_n; H^1(\Omega)') \hookrightarrow C([t_{n-1}, t_n]; L^2(\Omega))$$

satisfying $\rho^*(\cdot, t_{n-1}) = \rho_{\kappa, \alpha, L, \delta}^{n-1}(\cdot)$ and (3.52a) is immediate; see, for example, Wloka [67], Thm. 26.1. Further, on choosing, for $s \in (t_{n-1}, t_n]$, $\eta(\cdot, t) = \chi_{[t_{n-1}, s]} e^{-\|\tilde{\nabla}_x \cdot \tilde{u}\|_{L^\infty(\Omega)}(t-t_{n-1})} [\rho^*(\cdot, t)]_-$ in (3.52a), where, for a set $S \subset \mathbb{R}$, χ_S denotes the characteristic function of S , and recalling (3.53), we obtain that

$$(3.55) \quad e^{-\|\tilde{\nabla}_x \cdot \tilde{u}\|_{L^\infty(\Omega)}(s-t_{n-1})} \|\rho^*(\cdot, s)\|_{L^2(\Omega)}^2 + 2\alpha \int_{t_{n-1}}^s \int_{\Omega} e^{-\|\tilde{\nabla}_x \cdot \tilde{u}\|_{L^\infty(\Omega)}(t-t_{n-1})} |\tilde{\nabla}_x [\rho^*(x, t)]_-|^2 dx dt \leq 0, \quad s \in (t_{n-1}, t_n].$$

Next, we set

$$(3.56) \quad R(t) := e^{\|\tilde{\nabla}_x \cdot \tilde{u}\|_{L^\infty(\Omega)}(t-t_{n-1})} \|\rho_{\kappa, \alpha, L, \delta}^{n-1}\|_{L^\infty(\Omega)}, \quad t \in [t_{n-1}, t_n],$$

so that

$$(3.57) \quad \int_{t_{n-1}}^{t_n} \left[\left\langle \frac{\partial(\rho^* - R)}{\partial t}, \eta \right\rangle_{H^1(\Omega)} + c(\tilde{u})(\rho^* - R, \eta) \right] dt = - \int_{t_{n-1}}^{t_n} R \int_{\Omega} (\tilde{\nabla}_x \cdot \tilde{u} + \|\tilde{\nabla}_x \cdot \tilde{u}\|_{L^\infty(\Omega)}) \eta dx dt \quad \forall \eta \in L^2(t_{n-1}, t_n; H^1(\Omega)).$$

Then, similarly to (3.55), on choosing $\eta(\cdot, t) = \chi_{[t_{n-1}, s]} e^{-\|\tilde{\nabla}_x \cdot \tilde{u}\|_{L^\infty(\Omega)}(t-t_{n-1})} [\rho^*(\cdot, t) - R(t)]_+$ in (3.57) for $s \in (t_{n-1}, t_n]$, we obtain that

$$(3.58) \quad e^{-\|\tilde{\nabla}_x \cdot \tilde{u}\|_{L^\infty(\Omega)}(s-t_{n-1})} \|\rho^*(\cdot, s) - R(s)\|_{L^2(\Omega)}^2 + 2\alpha \int_{t_{n-1}}^s \int_{\Omega} e^{-\|\tilde{\nabla}_x \cdot \tilde{u}\|_{L^\infty(\Omega)}(t-t_{n-1})} |\tilde{\nabla}_x [\rho^*(x, t) - R(t)]_+|^2 dx dt \leq 0, \quad s \in (t_{n-1}, t_n].$$

On noting (3.54), extending (3.55) and (3.58) from the interval $(t_{n-1}, t_n]$ to $[t_{n-1}, t_n]$ by letting $s \rightarrow t_{n-1}$ in (3.55) and (3.58), and combining the resulting inequalities, we deduce that

$$(3.59) \quad \rho^*(\cdot, t) \in [0, R(t)] \quad \text{for } t \in [t_{n-1}, t_n], \quad \text{and so } \rho^* \in \Upsilon^n.$$

As $a(\cdot, \cdot)$ is a symmetric continuous coercive bilinear functional on $X \times X$ and $\ell_a(v, \xi)(\cdot)$ is a continuous linear functional on X for fixed $v \in H^1(\Omega)$ and $\xi \in L^\infty(\Omega \times D)$, the Lax–Milgram theorem yields the existence of a unique solution $\psi^* \in X$ to (3.52b). Similarly, for $\delta \in (0, 1)$, as $b_\delta(\eta)(\cdot, \cdot)$ is a nonsymmetric continuous coercive bilinear functional on $\mathcal{V} \times \mathcal{V}$ for fixed $\eta \in L_{\geq 0}^2(\Omega)$, and $\ell_{b, \delta}(\eta, \varphi)(\cdot)$ is a continuous linear functional on \mathcal{V} for fixed $\eta \in \Upsilon^n$ with $\eta(\cdot, t_n) \in L_{\geq 0}^2(\Omega)$ and $\varphi \in X$, the Lax–Milgram theorem yields the existence of a unique solution $\tilde{u}^* \in \mathcal{V}$ to (3.52c). Therefore the overall procedure (3.52a–c) that, for $\rho_{\kappa, \alpha, L}^{n-1} \in L_{\geq 0}^\Gamma(\Omega)$ fixed, maps $(\tilde{u}, \tilde{\psi}) \in \mathcal{V} \times L_M^2(\Omega \times D)$ to $(\rho^*, \tilde{u}^*, \psi^*) \in \Upsilon^n \times \mathcal{V} \times X$, with $\rho^*(\cdot, t_{n-1}) = \rho_{\kappa, \alpha, L}^{n-1}$, is well defined.

Lemma 3.3. *Let $\mathcal{T} : \mathcal{V} \times L_M^2(\Omega \times D) \rightarrow \mathcal{V} \times X$ denote the nonlinear map that takes the functions $(\tilde{u}, \tilde{\psi})$ to $(\tilde{u}^*, \psi^*) = \mathcal{T}(\tilde{u}, \tilde{\psi})$ via the procedure (3.52a–c). Then, the mapping \mathcal{T} has a fixed point. Hence, there exists a solution $(\rho_{\kappa, \alpha, L, \delta}^{[\Delta t], n}, \tilde{u}_{\kappa, \alpha, L, \delta}^n, \tilde{\psi}_{\kappa, \alpha, L, \delta}^n) \in \Upsilon^n \times \mathcal{V} \times X$ to (3.51a–c).*

Proof. Clearly, a fixed point of \mathcal{T} yields a solution of (3.51a–c). In order to show that \mathcal{T} has a fixed point, we apply Schauder’s fixed-point theorem; that is, we need to show that: (i) $\mathcal{T} : \mathcal{V} \times L_M^2(\Omega \times D) \rightarrow \mathcal{V} \times L_M^2(\Omega \times D)$ is continuous; (ii) \mathcal{T} is compact; and (iii) there exists a $C_\star \in \mathbb{R}_{>0}$ such that

$$(3.60) \quad \|\tilde{u}\|_{H^1(\Omega)} + \|\tilde{u}\|_{L^\infty(\Omega)} + \|\tilde{\nabla}_x \cdot \tilde{u}\|_{L^\infty(\Omega)} + \|\tilde{\psi}\|_{L_M^2(\Omega \times D)} \leq C_\star$$

for every $(\tilde{u}, \tilde{\psi}) \in \mathcal{V} \times L_M^2(\Omega \times D)$ and $\varkappa \in (0, 1]$ satisfying $(\tilde{u}, \tilde{\psi}) = \varkappa \mathcal{T}(\tilde{u}, \tilde{\psi})$.

(i) Let $\{\tilde{u}^{(m)}, \tilde{\psi}^{(m)}\}_{m \in \mathbb{N}}$ be such that, as $m \rightarrow \infty$,

$$(3.61) \quad \begin{aligned} \tilde{u}^{(m)} &\rightarrow \tilde{u} \quad \text{strongly in } \tilde{H}_0^1(\Omega), & \tilde{u}^{(m)} &\rightarrow \tilde{u} \quad \text{strongly in } \tilde{L}^\infty(\Omega), \\ \nabla_x \cdot \tilde{u}^{(m)} &\rightarrow \nabla_x \cdot \tilde{u} \quad \text{strongly in } L^\infty(\Omega), & \tilde{\psi}^{(m)} &\rightarrow \tilde{\psi} \quad \text{strongly in } L_M^2(\Omega \times D). \end{aligned}$$

It follows immediately from (3.61), (3.45), (1.7a) and (3.1) that

$$(3.62) \quad M^{\frac{1}{2}} \beta_\delta^L(\tilde{\psi}^{(m)}) \rightarrow M^{\frac{1}{2}} \beta_\delta^L(\tilde{\psi}) \quad \text{strongly in } L^r(\Omega \times D) \quad \text{as } m \rightarrow \infty$$

for all $r \in [1, \infty)$. In order to prove that $\mathcal{T} : \mathcal{V} \times L_M^2(\Omega \times D) \rightarrow \mathcal{V} \times L_M^2(\Omega \times D)$ is continuous, we need to show that $(\tilde{v}^{(m)}, \xi^{(m)}) := \mathcal{T}(\tilde{u}^{(m)}, \tilde{\psi}^{(m)})$ is such that, as $m \rightarrow \infty$,

$$(3.63) \quad \begin{aligned} \tilde{v}^{(m)} &\rightarrow \tilde{v} \quad \text{strongly in } \tilde{H}_0^1(\Omega), & \tilde{v}^{(m)} &\rightarrow \tilde{v} \quad \text{strongly in } \tilde{L}^\infty(\Omega), \\ \nabla_x \cdot \tilde{v}^{(m)} &\rightarrow \nabla_x \cdot \tilde{v} \quad \text{strongly in } L^\infty(\Omega), & \xi^{(m)} &\rightarrow \xi \quad \text{strongly in } L_M^2(\Omega \times D), \end{aligned}$$

where $(\tilde{v}, \xi) := \mathcal{T}(\tilde{u}, \tilde{\psi})$. We have from the definition of \mathcal{T} , recall (3.52a-c), that, for all $m \in \mathbb{N}$, $(\tilde{v}^{(m)}, \xi^{(m)}) \in \mathcal{V} \times X$ is the unique solution to

$$(3.64a) \quad a(\xi^{(m)}, \varphi) = \ell_a(\tilde{u}^{(m)}, \beta_\delta^L(\tilde{\psi}^{(m)}))(\varphi) \quad \forall \varphi \in X,$$

$$(3.64b) \quad b_\delta(\tilde{\rho}^{(m)}(\cdot, t_n))(\tilde{v}^{(m)}, w) = \ell_{b,\delta}(\tilde{\rho}^{(m)}, \xi^{(m)})(w) \quad \forall w \in \mathcal{V},$$

where $\tilde{\rho}^{(m)} \in \Upsilon^n$ is the unique solution to $\tilde{\rho}^{(m)}(\cdot, t_{n-1}) = \rho_{\kappa,\alpha,L,\delta}^{n-1}(\cdot)$ and

$$(3.65) \quad \int_{t_{n-1}}^{t_n} \left[\left\langle \frac{\partial \tilde{\rho}^{(m)}}{\partial t}, \eta \right\rangle_{H^1(\Omega)} + c(\tilde{u}^{(m)})(\tilde{\rho}^{(m)}, \eta) \right] dt = 0 \quad \forall \eta \in L^2(t_{n-1}, t_n; H^1(\Omega)).$$

It follows from (3.61) that $\|\nabla_x \cdot \tilde{u}^{(m)}\|_{L^\infty(\Omega)} \leq C_\star$ for all $m \in \mathbb{N}$. On choosing, for $s \in (t_{n-1}, t_n]$, $\eta(\cdot, t) = \chi_{[t_{n-1}, s]} e^{-C_\star(t-t_{n-1})} \tilde{\rho}^{(m)}(\cdot, t)$ in (3.65) yields, on noting (3.53), the first two bounds in

$$(3.66) \quad \begin{aligned} &\|\tilde{\rho}^{(m)}\|_{C([t_{n-1}, t_n]; L^2(\Omega))}^2 + \alpha \|\tilde{\rho}^{(m)}\|_{L^2(t_{n-1}, t_n; H^1(\Omega))}^2 \\ &+ \|\tilde{\rho}^{(m)}\|_{L^\infty(t_{n-1}, t_n; L^\infty(\Omega))}^2 + \left\| \frac{\partial \tilde{\rho}^{(m)}}{\partial t} \right\|_{L^2(t_{n-1}, t_n; H^1(\Omega)')}^2 \leq C, \end{aligned}$$

where, here and below, C is independent of m . The third bound in (3.66) follows from applying the bound (3.59) to (3.65) with \tilde{u} in (3.56) replaced by $\tilde{u}^{(m)}$ and noting (3.61). The fourth bound in (3.66) follows immediately from the first two bounds in (3.66), (3.61) and (3.65). Choosing $\varphi = \xi^{(m)}$ in (3.64a) yields, on noting (3.37a,b) and (3.61), that

$$(3.67) \quad \|\xi^{(m)}\|_{H_M^1(\Omega \times D)}^2 \leq C.$$

Choosing $w = \tilde{v}^{(m)}$ in (3.64b) yields, on noting (3.47), (3.48), (3.35a,b), $\tilde{\rho}^{(m)}(\cdot, t_n) \in L_{\geq 0}^2(\Omega)$, $\rho_{\kappa,\alpha,L}^{n-1} \in L_{\geq 0}^\Gamma(\Omega)$, (3.3), (1.6b), (3.15), (3.67), (2.3) and (3.66), that

$$(3.68) \quad \|\tilde{v}^{(m)}\|_{H^2(\Omega)}^2 + \|\nabla_x \cdot \tilde{v}^{(m)}\|_{H^2(\Omega)}^2 \leq C.$$

It follows from (3.66), (3.67), (3.68), (3.11), (3.45), (1.7a), (3.1), (3.15) and the compactness of \mathcal{V} in \mathcal{V} and $X \equiv H_M^1(\Omega \times D)$ in $L_M^2(\Omega \times D)$ that there exists a subsequence $\{(\tilde{\rho}^{(m_k)}, \tilde{v}^{(m_k)}, \xi^{(m_k)})\}_{m_k \in \mathbb{N}}$

and functions $(\tilde{\rho}, \underline{v}, \xi) \in \Upsilon^n \times \mathcal{V} \times X$ such that, as $m_k \rightarrow \infty$, for any $r \in [1, \infty)$,

$$(3.69a) \quad \tilde{\rho}^{(m_k)} \rightarrow \tilde{\rho} \quad \text{weakly in } L^2(t_{n-1}, t_n; H^1(\Omega)), \quad \text{strongly in } L^r(t_{n-1}, t_n; L^r(\Omega)),$$

$$(3.69b) \quad \frac{\partial \tilde{\rho}^{(m_k)}}{\partial t} \rightarrow \frac{\partial \tilde{\rho}}{\partial t} \quad \text{weakly in } L^2(t_{n-1}, t_n; H^1(\Omega)'),$$

$$(3.69c) \quad \tilde{\rho}^{(m_k)}(\cdot, t_n) \rightarrow \tilde{\rho}(\cdot, t_n) \quad \text{weakly in } L^2(\Omega),$$

$$(3.69d) \quad \tilde{v}^{(m_k)} \rightarrow \tilde{v} \quad \text{weakly in } \tilde{H}^2(\Omega), \quad \text{strongly in } \tilde{H}_0^1(\Omega) \cap \tilde{L}^\infty(\Omega),$$

$$(3.69e) \quad \nabla_x \cdot \tilde{v}^{(m_k)} \rightarrow \nabla_x \cdot \tilde{v} \quad \text{weakly in } H^2(\Omega), \quad \text{strongly in } L^\infty(\Omega),$$

$$(3.69f) \quad \xi^{(m_k)} \rightarrow \xi \quad \text{weakly in } H_M^1(\Omega \times D), \quad \text{strongly in } L_M^2(\Omega \times D),$$

$$(3.69g) \quad M^{\frac{1}{2}} \beta_\delta^L(\xi^{(m_k)}) \rightarrow M^{\frac{1}{2}} \beta_\delta^L(\xi) \quad \text{strongly in } L^r(\Omega \times D),$$

$$(3.69h) \quad C_i(\xi^{(m_k)}) \rightarrow C_i(\xi) \quad \text{strongly in } L^2(\Omega), \quad i = 1, \dots, K.$$

We deduce from (3.65), (3.69a,b) and (3.61) that $\tilde{\rho} \in \Upsilon^n$ is the unique solution to $\tilde{\rho}(\cdot, t_{n-1}) = \rho_{\kappa, \alpha, L, \delta}^{n-1}(\cdot)$ and

$$(3.70) \quad \int_{t_{n-1}}^{t_n} \left[\left\langle \frac{\partial \tilde{\rho}}{\partial t}, \eta \right\rangle_{H^1(\Omega)} + c(\tilde{u})(\tilde{\rho}, \eta) \right] dt = 0 \quad \forall \eta \in L^2(t_{n-1}, t_n; H^1(\Omega)).$$

Choosing $\eta = 1$ in (3.70), on noting (3.33) and (3.50), yields that

$$(3.71) \quad \int_{\Omega} \tilde{\rho}(\cdot, t_n) dx = \int_{\Omega} \rho_{\kappa, \alpha, L, \delta}^{n-1} dx \leq \int_{\Omega} \rho_{\kappa, \alpha, L}^{n-1} dx.$$

It follows from (3.64a), (3.37a,b), (3.69f), (3.61) and (3.62) that $\xi, \tilde{\psi} \in X$ and $\tilde{u} \in \mathcal{V}$ satisfy

$$(3.72) \quad a(\xi, \varphi) = \ell_a(\tilde{u}, \beta_\delta^L(\tilde{\psi}))(\varphi) \quad \forall \varphi \in C^\infty(\overline{\Omega \times D}).$$

Then, noting that $a(\cdot, \cdot)$ is a continuous bilinear functional on $X \times X$, that $\ell_a(\tilde{u}, \beta_\delta^L(\tilde{\psi}))(\cdot)$ is a continuous linear functional on X , and recalling (3.8), we deduce that $\xi \in X$ is the unique solution of (3.72) for all $\varphi \in X$. It further follows from (3.64b), (3.47), (3.48), (3.35a,b) and (3.69a,c,d,e,f,g,h) that $\underline{v} \in \mathcal{V}$ is the unique solution to

$$(3.73) \quad b_\delta(\tilde{\rho}(\cdot, t_n))(\underline{v}, w) = \ell_{b, \delta}(\tilde{\rho}, \xi)(w) \quad \forall w \in \mathcal{V}.$$

Combining (3.72) with $\varphi \in X$ with (3.73) and (3.70), we have that $(\underline{v}, \xi) = \mathcal{T}(\tilde{u}, \tilde{\psi}) \in \mathcal{V} \times X$. As (\underline{v}, ξ) is unique for fixed $(\tilde{u}, \tilde{\psi})$, the whole sequence converges in (3.69a–h), and so (3.63) holds. Therefore the mapping $\mathcal{T} : \tilde{\mathcal{V}} \times L_M^2(\Omega \times D) \rightarrow \mathcal{V} \times L_M^2(\Omega \times D)$ is continuous.

(ii) Since the embeddings $\tilde{\mathcal{V}} \hookrightarrow \mathcal{V}$ and $X \hookrightarrow L_M^2(\Omega \times D)$ are compact, we directly deduce that the mapping $\mathcal{T} : \tilde{\mathcal{V}} \times L_M^2(\Omega \times D) \rightarrow \mathcal{V} \times L_M^2(\Omega \times D)$ is compact. It therefore remains to show that (iii) holds.

(iii) Let us suppose that $(\tilde{u}, \tilde{\psi}) = \varkappa \mathcal{T}(\tilde{u}, \tilde{\psi})$; then, $(\tilde{\rho}, \tilde{u}, \tilde{\psi}) \in \Upsilon^n \times \mathcal{V} \times X$ satisfies $\tilde{\rho}(\cdot, t_{n-1}) = \rho_{\kappa, \alpha, L, \delta}^{n-1}(\cdot)$ and

$$(3.74a) \quad \int_{t_{n-1}}^{t_n} \left[\left\langle \frac{\partial \tilde{\rho}}{\partial t}, \eta \right\rangle_{H^1(\Omega)} + c(\tilde{u})(\tilde{\rho}, \eta) \right] dt = 0 \quad \forall \eta \in L^2(t_{n-1}, t_n; H^1(\Omega)),$$

$$(3.74b) \quad b_\delta(\tilde{\rho}(\cdot, t_n))(\tilde{u}, w) = \varkappa \ell_{b, \delta}(\tilde{\rho}, \tilde{\psi})(w) \quad \forall w \in \mathcal{V},$$

$$(3.74c) \quad a(\tilde{\psi}, \varphi) = \varkappa \ell_a(\tilde{u}, \beta_\delta^L(\tilde{\psi}))(\varphi) \quad \forall \varphi \in X.$$

Choosing $w = \tilde{u}$ in (3.74b) yields, as $\varkappa \in (0, 1]$, that

$$\begin{aligned}
& \frac{\varkappa}{2} \int_{\Omega} \left[\tilde{\rho}(\cdot, t_n) |\tilde{u}|^2 + \rho_{\kappa, \alpha, L}^{n-1} |\tilde{u} - u_{\kappa, \alpha, L}^{n-1}|^2 - \rho_{\kappa, \alpha, L}^{n-1} |u_{\kappa, \alpha, L}^{n-1}|^2 \right] dx \\
& + \Delta t \mu^S \int_{\Omega} |D(\tilde{u})|^2 dx + \Delta t \left(\mu^B - \frac{\mu^S}{d} \right) \int_{\Omega} |\nabla_x \cdot \tilde{u}|^2 dx \\
& + \Delta t \delta \sum_{|\lambda|=2} \int_{\Omega} \left[\left| \frac{\partial^{|\lambda|} \tilde{u}}{\partial x_1^{\lambda_1} \dots \partial x_d^{\lambda_d}} \right|^2 + \left| \frac{\partial^{|\lambda|} (\nabla_x \cdot \tilde{u})}{\partial x_1^{\lambda_1} \dots \partial x_d^{\lambda_d}} \right|^2 \right] dx \\
& \leq \varkappa \Delta t \left[\int_{\Omega} \tilde{\rho}(\cdot, t_n) f^n \cdot \tilde{u} dx - k \sum_{i=1}^K \int_{\Omega} C_i(M \tilde{\psi}) : \nabla_x \tilde{u} dx \right] \\
& + \varkappa \int_{\Omega} \left(\int_{t_{n-1}}^{t_n} p_{\kappa}(\tilde{\rho}) dt + \Delta t k (K+1) \int_D M \tilde{\psi} dq \right) \nabla_x \cdot \tilde{u} dx \\
& - 2 \varkappa \Delta t \int_{\Omega} \left(\int_D M \beta_{\delta}^L(\tilde{\psi}) dq \right) \nabla_x \left(\int_D M \tilde{\psi} dq \right) \cdot \tilde{u} dx.
\end{aligned} \tag{3.75}$$

On recalling (3.5), we choose $\eta(\cdot, t) = \chi_{[t_{n-1}, s]} P'_{\kappa}(\tilde{\rho}(\cdot, t) + \varsigma)$ in (3.74a), for any $s \in (t_{n-1}, t_n]$ and any fixed $\varsigma \in \mathbb{R}_{>0}$, to obtain, on noting (3.33), that

$$\begin{aligned}
& \int_{\Omega} P_{\kappa}(\tilde{\rho}(\cdot, s)) dx + \alpha \kappa \int_{t_{n-1}}^s \int_{\Omega} (4 \tilde{\rho}^2 + \Gamma \tilde{\rho}^{\Gamma-2}) |\nabla_x \tilde{\rho}|^2 dx dt \\
& \leq \int_{\Omega} P_{\kappa}(\tilde{\rho}(\cdot, s) + \varsigma) dx + \alpha \int_{t_{n-1}}^s \int_{\Omega} P'_{\kappa}(\tilde{\rho} + \varsigma) |\nabla_x \tilde{\rho}|^2 dx dt \\
& = \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L, \delta}^{n-1} + \varsigma) dx + \int_{t_{n-1}}^s \int_{\Omega} \tilde{\rho} \tilde{u} \cdot \nabla_x P'_{\kappa}(\tilde{\rho} + \varsigma) dx dt \\
& = \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L, \delta}^{n-1} + \varsigma) dx + \int_{\Omega} \left(\int_{t_{n-1}}^s [P_{\kappa}(\tilde{\rho} + \varsigma) - \tilde{\rho} P'_{\kappa}(\tilde{\rho} + \varsigma)] dt \right) \nabla_x \cdot \tilde{u} dx.
\end{aligned} \tag{3.76}$$

As $\rho_{\kappa, \alpha, L, \delta}^{n-1} \in L_{\geq 0}^{\infty}(\Omega)$, $\tilde{\rho} \in L^{\infty}(t_{n-1}, t_n; L_{\geq 0}^{\infty}(\Omega))$ and $\tilde{u} \in \mathcal{V}$, one can pass to the limit $\varsigma \rightarrow 0_+$ in (3.76) using Lebesgue's dominated convergence theorem to obtain, for any $s \in (t_{n-1}, t_n]$, that

$$\begin{aligned}
& \int_{\Omega} P_{\kappa}(\tilde{\rho}(\cdot, s)) dx + \alpha \kappa \int_{t_{n-1}}^s \left[\|\nabla_x(\tilde{\rho}^2)\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x(\tilde{\rho}^{\frac{\Gamma}{2}})\|_{L^2(\Omega)}^2 \right] dt \\
& \leq \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L, \delta}^{n-1}) dx + \int_{\Omega} \left(\int_{t_{n-1}}^s [P_{\kappa}(\tilde{\rho}) - \tilde{\rho} P'_{\kappa}(\tilde{\rho})] dt \right) \nabla_x \cdot \tilde{u} dx \\
& \leq \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L}^{n-1}) dx - \int_{\Omega} \left(\int_{t_{n-1}}^s p_{\kappa}(\tilde{\rho}) dt \right) \nabla_x \cdot \tilde{u} dx,
\end{aligned} \tag{3.77}$$

where we have noted (3.50) and (3.5) for the final inequality. We remark that we needed to choose $P'_{\kappa}(\tilde{\rho}(\cdot, t) + \varsigma)$, as opposed to $P'_{\kappa}(\tilde{\rho}(\cdot, t))$, in the testing procedure as $P'_{\kappa}(\tilde{\rho})$, that would appear in (3.76), may not be well-defined for $\gamma \in (\frac{3}{2}, 2)$, as we only know that $\tilde{\rho}$ is nonnegative as opposed to being strictly positive.

For a.e. $x \in \Omega$, let

$$\tilde{\varrho}(x) := \int_D M(q) \tilde{\psi}(x, q) dq \quad \text{and} \quad \varrho_{\kappa, \alpha, L}^{n-1}(x) := \int_D M(q) \hat{\psi}_{\kappa, \alpha, L}^{n-1}(x, q) dq. \tag{3.78}$$

Choosing $\varphi(\underline{x}, \underline{q}) = \tilde{\varrho}(\underline{x}) \otimes 1(\underline{q})$ in (3.74c) yields that

$$(3.79) \quad \begin{aligned} & \frac{\varkappa}{2} \left[\|\tilde{\varrho}\|_{L^2(\Omega)}^2 + \|\tilde{\varrho} - \varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 \right] + \Delta t \varepsilon \|\nabla_x \tilde{\varrho}\|_{L^2(\Omega)}^2 \\ &= \frac{\varkappa}{2} \|\varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 + \varkappa \Delta t \int_{\Omega} \left(\int_D M \beta_{\delta}^L(\tilde{\psi}) \, d\tilde{q} \right) \tilde{u} \cdot \nabla_x \tilde{\varrho} \, d\tilde{x}. \end{aligned}$$

Combining (3.75), (3.77) for $s = t_n$ and (3.79) yields, on noting (3.3), (3.4) and (3.78), that, for all $\varkappa \in (0, 1]$,

$$(3.80) \quad \begin{aligned} & \frac{\varkappa}{2} \int_{\Omega} \left[\tilde{\rho}(\cdot, t_n) |\tilde{u}|^2 + \rho_{\kappa, \alpha, L}^{n-1} |\tilde{u} - \tilde{u}_{\kappa, \alpha, L}^{n-1}|^2 \right] d\tilde{x} + \varkappa \int_{\Omega} P_{\kappa}(\tilde{\rho}(\cdot, t_n)) d\tilde{x} \\ &+ \varkappa \alpha \kappa \int_{t_{n-1}}^{t_n} \left[\|\nabla_x(\tilde{\rho}^2)\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x(\tilde{\rho}^{\frac{\Gamma}{2}})\|_{L^2(\Omega)}^2 \right] dt + \Delta t \mu^S c_0 \|\tilde{u}\|_{H^1(\Omega)}^2 \\ &+ \Delta t \delta \sum_{|\lambda|=2} \int_{\Omega} \left[\left| \frac{\partial^{|\lambda|} \tilde{u}}{\partial x_1^{\lambda_1} \cdots \partial x_d^{\lambda_d}} \right|^2 + \left| \frac{\partial^{|\lambda|} (\nabla_x \cdot \tilde{u})}{\partial x_1^{\lambda_1} \cdots \partial x_d^{\lambda_d}} \right|^2 \right] d\tilde{x} \\ &+ \varkappa \mathfrak{z} \left[\|\tilde{\varrho}\|_{L^2(\Omega)}^2 + \|\tilde{\varrho} - \varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 \right] + 2 \Delta t \mathfrak{z} \varepsilon \|\nabla_x \tilde{\varrho}\|_{L^2(\Omega)}^2 \\ &\leq \frac{\varkappa}{2} \int_{\Omega} \rho_{\kappa, \alpha, L}^{n-1} |\tilde{u}_{\kappa, \alpha, L}^{n-1}|^2 d\tilde{x} + \varkappa \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L}^{n-1}) d\tilde{x} + \varkappa \Delta t \int_{\Omega} \tilde{\rho}(\cdot, t_n) f^n \cdot \tilde{u} d\tilde{x} \\ &+ \varkappa \mathfrak{z} \|\varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 + \varkappa k \Delta t \left[(K+1) \int_{\Omega} \tilde{\varrho} \nabla_x \cdot \tilde{u} d\tilde{x} - \sum_{i=1}^K \int_{\Omega} C_i(M \tilde{\psi}) : \nabla_x \tilde{u} d\tilde{x} \right]. \end{aligned}$$

Choosing $\varphi = [\mathcal{F}_{\delta}^L]'(\tilde{\psi})$ in (3.74c) and noting (3.45) implies that

$$(3.81) \quad \begin{aligned} & \int_{\Omega \times D} M \left(\mathcal{F}_{\delta}^L(\tilde{\psi}) - \mathcal{F}_{\delta}^L(\varkappa \tilde{\psi}_{\kappa, \alpha, L}^{n-1}) + \frac{1}{2L} |\tilde{\psi} - \varkappa \tilde{\psi}_{\kappa, \alpha, L}^{n-1}|^2 \right) d\tilde{q} d\tilde{x} \\ &+ \frac{\Delta t}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \int_{\Omega \times D} M \nabla_{q_j} \tilde{\psi} \cdot \nabla_{q_i} ([\mathcal{F}_{\delta}^L]'(\tilde{\psi})) d\tilde{q} d\tilde{x} \\ &+ \Delta t \varepsilon \int_{\Omega \times D} M \nabla_x \tilde{\psi} \cdot \nabla_x ([\mathcal{F}_{\delta}^L]'(\tilde{\psi})) d\tilde{q} d\tilde{x} \\ &\leq \varkappa \Delta t \int_{\Omega \times D} M \left[\left(\sum_{i=1}^K \sigma(\tilde{u}) q_i \right) \cdot \nabla_{q_i} \tilde{\psi} + \tilde{u} \cdot \nabla_x \tilde{\psi} \right] d\tilde{q} d\tilde{x} \\ &= \varkappa \Delta t \int_{\Omega} \left[\sum_{i=1}^K C_i(M \tilde{\psi}) : \sigma(\tilde{u}) - (K+1) (\nabla_x \cdot \tilde{u}) \int_D M \tilde{\psi} d\tilde{q} \right] d\tilde{x}, \end{aligned}$$

where in the transition to the final line we applied (3.16) with $\underline{B} = \underline{\sigma}(\tilde{u})$ (on account of it being independent of the variable \underline{q}), and recalled (1.12). Combining (3.80) and (3.81), and noting (2.5) and (3.45), yields, for all $\varkappa \in (0, 1]$ and $\varsigma \in \mathbb{R}_{>0}$, that

$$\begin{aligned} & \frac{\varkappa}{2} \int_{\Omega} \left[\tilde{\rho}(\cdot, t_n) |\tilde{u}|^2 + \rho_{\kappa, \alpha, L}^{n-1} |\tilde{u} - \tilde{u}_{\kappa, \alpha, L}^{n-1}|^2 \right] d\tilde{x} + \varkappa \int_{\Omega} P_{\kappa}(\tilde{\rho}(\cdot, t_n)) d\tilde{x} \\ &+ \varkappa \alpha \kappa \int_{t_{n-1}}^{t_n} \left[\|\nabla_x(\tilde{\rho}^2)\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x(\tilde{\rho}^{\frac{\Gamma}{2}})\|_{L^2(\Omega)}^2 \right] dt \\ &+ k \int_{\Omega \times D} M \left(\mathcal{F}_{\delta}^L(\tilde{\psi}) + \frac{1}{2L} |\tilde{\psi} - \varkappa \tilde{\psi}_{\kappa, \alpha, L}^{n-1}|^2 \right) d\tilde{q} d\tilde{x} + \Delta t \mu^S c_0 \|\tilde{u}\|_{H^1(\Omega)}^2 \end{aligned}$$

$$\begin{aligned}
& + \Delta t \delta \sum_{|\lambda|=2} \int_{\Omega} \left[\left| \frac{\partial^{|\lambda|} \tilde{u}}{\partial_{x_1}^{\lambda_1} \cdots \partial_{x_d}^{\lambda_d}} \right|^2 + \left| \frac{\partial^{|\lambda|} (\nabla_x \cdot \tilde{u})}{\partial_{x_1}^{\lambda_1} \cdots \partial_{x_d}^{\lambda_d}} \right|^2 \right] dx \\
& + \frac{\Delta t k a_0}{4 \lambda L} \sum_{i=1}^K \int_{\Omega \times D} M |\nabla_{q_i} \tilde{\psi}|^2 dq dx + \frac{\Delta t k \varepsilon}{L} \int_{\Omega \times D} M |\nabla_x \tilde{\psi}|^2 dq dx \\
& + \varkappa \mathfrak{z} \left[\|\tilde{\varrho}\|_{L^2(\Omega)}^2 + \|\tilde{\varrho} - \varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 \right] + 2 \Delta t \mathfrak{z} \varepsilon \|\nabla_x \tilde{\varrho}\|_{L^2(\Omega)}^2 \\
& \leq \frac{\varkappa}{2} \int_{\Omega} \rho_{\kappa, \alpha, L}^{n-1} |u_{\kappa, \alpha, L}^{n-1}|^2 dx + \varkappa \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L}^{n-1}) dx + k \int_{\Omega \times D} M \mathcal{F}_{\delta}^L(\varkappa \hat{\psi}_{\kappa, \alpha, L}^{n-1}) dq dx \\
& \quad + \varkappa \mathfrak{z} \|\varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 + \varkappa \Delta t \int_{\Omega} \tilde{\rho}(\cdot, t_n) f^n \cdot \tilde{u} dx \\
& \leq \frac{\varkappa}{2} \int_{\Omega} \rho_{\kappa, \alpha, L}^{n-1} |u_{\kappa, \alpha, L}^{n-1}|^2 dx + \varkappa \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L}^{n-1}) dx + k \int_{\Omega \times D} M \mathcal{F}_{\delta}^L(\varkappa \hat{\psi}_{\kappa, \alpha, L}^{n-1}) dq dx \\
(3.82) \quad & + \varkappa \mathfrak{z} \|\varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 + \frac{\varkappa \Delta t}{2} \left[\varsigma \int_{\Omega} \tilde{\rho}(\cdot, t_n) |\tilde{u}|^2 dx + \frac{1}{\varsigma} \|f^n\|_{L^{\infty}(\Omega)}^2 \int_{\Omega} \rho_{\kappa, \alpha, L}^{n-1} dx \right],
\end{aligned}$$

where, in deriving the final inequality, we have noted (3.71). It is easy to see that $\mathcal{F}_{\delta}^L(s)$ is nonnegative for all $s \in \mathbb{R}$, with $\mathcal{F}_{\delta}^L(1) = 0$. Furthermore, for any $\varkappa \in (0, 1]$, $\mathcal{F}_{\delta}^L(\varkappa s) \leq \mathcal{F}_{\delta}^L(s)$ if $s < 0$ or $1 \leq \varkappa s$, and also $\mathcal{F}_{\delta}^L(\varkappa s) \leq \mathcal{F}_{\delta}^L(0) \leq 1$ if $0 \leq \varkappa s \leq 1$. Thus we deduce that

$$(3.83) \quad \mathcal{F}_{\delta}^L(\varkappa s) \leq \mathcal{F}_{\delta}^L(s) + 1 \quad \forall s \in \mathbb{R}, \quad \forall \varkappa \in (0, 1].$$

Hence, the bounds (3.82) and (3.83), on noting (3.44b) and, from (3.45) and (1.25), that $\beta_{\delta}^L(\cdot) \leq L$, give rise, for ς sufficiently small, to the desired bound (3.60) with C_* dependent only on $\delta, L, \Delta t, M, k, \mu^S, c_0, a_0, \underline{f}, \rho_{\kappa, \alpha, L}^{n-1}, u_{\kappa, \alpha, L}^{n-1}$ and $\hat{\psi}_{\kappa, \alpha, L}^{n-1}$. Therefore (iii) holds, and so \mathcal{T} has a fixed point, proving existence of a solution to (3.51a–c). \square

Similarly to (3.79), choosing $\varphi(\underline{x}, q) = \varrho_{\kappa, \alpha, L, \delta}^n(\underline{x}) \otimes 1(q)$ in (3.51c), where $\varrho_{\kappa, \alpha, L, \delta}^n(\underline{x}) := \int_D M(q) \hat{\psi}_{\kappa, \alpha, L, \delta}^n(\underline{x}, q) dq dx$, yields that

$$\begin{aligned}
(3.84) \quad & \frac{1}{2} \left[\|\varrho_{\kappa, \alpha, L, \delta}^n\|_{L^2(\Omega)}^2 + \|\varrho_{\kappa, \alpha, L, \delta}^n - \varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 \right] + \Delta t \varepsilon \|\nabla_x \varrho_{\kappa, \alpha, L, \delta}^n\|_{L^2(\Omega)}^2 \\
& = \frac{1}{2} \|\varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 + \Delta t \int_{\Omega} \left(\int_D M \beta_{\delta}^L(\varrho_{\kappa, \alpha, L, \delta}^n) dq \right) u_{\kappa, \alpha, L, \delta}^n \cdot \nabla_x \varrho_{\kappa, \alpha, L, \delta}^n dx.
\end{aligned}$$

Choosing $w = u_{\kappa, \alpha, L, \delta}^n$ in (3.51b) and $\eta = P'_{\kappa}(\rho_{\kappa, \alpha, L, \delta}^{[\Delta t], n})$ in (3.51a), and combining with (3.84), yields, similarly to (3.80), that

$$\begin{aligned}
& \frac{1}{2} \int_{\Omega} \left[\rho_{\kappa, \alpha, L, \delta}^{[\Delta t], n}(\cdot, t_n) |u_{\kappa, \alpha, L, \delta}^n|^2 + \rho_{\kappa, \alpha, L}^{n-1} |u_{\kappa, \alpha, L, \delta}^n - u_{\kappa, \alpha, L}^{n-1}|^2 \right] dx + \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L, \delta}^{[\Delta t], n}(\cdot, t_n)) dx \\
& + \alpha \kappa \int_{t_{n-1}}^{t_n} \left[\|\nabla_x [(\rho_{\kappa, \alpha, L, \delta}^{[\Delta t], n})^2]\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x [(\rho_{\kappa, \alpha, L, \delta}^{[\Delta t], n})^{\frac{\Gamma}{2}}]\|_{L^2(\Omega)}^2 \right] dt \\
& + \Delta t \mu^S c_0 \|u_{\kappa, \alpha, L, \delta}^n\|_{H^1(\Omega)}^2 \\
& + \Delta t \delta \sum_{|\lambda|=2} \int_{\Omega} \left[\left| \frac{\partial^{|\lambda|} u_{\kappa, \alpha, L, \delta}^n}{\partial_{x_1}^{\lambda_1} \cdots \partial_{x_d}^{\lambda_d}} \right|^2 + \left| \frac{\partial^{|\lambda|} (\nabla_x \cdot u_{\kappa, \alpha, L, \delta}^n)}{\partial_{x_1}^{\lambda_1} \cdots \partial_{x_d}^{\lambda_d}} \right|^2 \right] dx \\
& + \mathfrak{z} \left[\|\varrho_{\kappa, \alpha, L, \delta}^n\|_{L^2(\Omega)}^2 + \|\varrho_{\kappa, \alpha, L, \delta}^n - \varrho_{\kappa, \alpha, L}^{n-1}\|_{L^2(\Omega)}^2 \right] + 2 \Delta t \mathfrak{z} \varepsilon \|\nabla_x \varrho_{\kappa, \alpha, L, \delta}^n\|_{L^2(\Omega)}^2
\end{aligned}$$

$$\begin{aligned}
&\leq \frac{1}{2} \int_{\Omega} \rho_{\kappa,\alpha,L}^{n-1} |u_{\kappa,\alpha,L}^{n-1}|^2 dx + \int_{\Omega} P_{\kappa}(\rho_{\kappa,\alpha,L}^{n-1}) dx + \Delta t \int_{\Omega} \rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t_n) f^n \cdot u_{\kappa,\alpha,L,\delta}^n dx \\
&\quad + \mathfrak{z} \|\varrho_{\kappa,\alpha,L}^{n-1}\|_{L^2(\Omega)}^2 + k \Delta t (K+1) \int_{\Omega} \varrho_{\kappa,\alpha,L,\delta}^n \nabla_x \cdot u_{\kappa,\alpha,L,\delta}^n dx \\
(3.85) \quad &\quad - k \Delta t \sum_{i=1}^K \int_{\Omega} C_i(M \hat{\psi}_{\kappa,\alpha,L,\delta}^n) : \nabla_x u_{\kappa,\alpha,L,\delta}^n dx.
\end{aligned}$$

Choosing $\varphi = [\mathcal{F}_{\delta}^L]'(\hat{\psi}_{\kappa,\alpha,L,\delta}^n)$ in (3.51c), combining with (3.85) and noting (3.44a), yields, similarly to (3.82), that, for $\varsigma \in \mathbb{R}_{>0}$ sufficiently small, the solution $(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}, u_{\kappa,\alpha,L,\delta}^n, \hat{\psi}_{\kappa,\alpha,L,\delta}^n) \in \Upsilon^n \times \mathcal{V} \times X$ of (3.51a–c) satisfies

$$\begin{aligned}
&\frac{1}{2} \int_{\Omega} \left[\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t_n) |u_{\kappa,\alpha,L,\delta}^n|^2 + \rho_{\kappa,\alpha,L}^{n-1} |u_{\kappa,\alpha,L,\delta}^n - u_{\kappa,\alpha,L}^{n-1}|^2 \right] dx + \int_{\Omega} P_{\kappa}(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t_n)) dx \\
&\quad + \alpha \kappa \int_{t_{n-1}}^{t_n} \left[\|\nabla_x[(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^2]\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x[(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\frac{\Gamma}{2}}]\|_{L^2(\Omega)}^2 \right] dt \\
&\quad + k \int_{\Omega \times D} M \left(\mathcal{F}_{\delta}^L(\hat{\psi}_{\kappa,\alpha,L,\delta}^n) + \frac{1}{2L} |\hat{\psi}_{\kappa,\alpha,L,\delta}^n - \hat{\psi}_{\kappa,\alpha,L}^{n-1}|^2 \right) dq dx \\
&\quad + \Delta t \mu^S c_0 \|u_{\kappa,\alpha,L,\delta}^n\|_{H^1(\Omega)}^2 + \Delta t \delta \sum_{|\lambda|=2} \int_{\Omega} \left[\left| \frac{\partial^{|\lambda|} u_{\kappa,\alpha,L,\delta}^n}{\partial x_1^{\lambda_1} \dots \partial x_d^{\lambda_d}} \right|^2 + \left| \frac{\partial^{|\lambda|} (\nabla_x \cdot u_{\kappa,\alpha,L,\delta}^n)}{\partial x_1^{\lambda_1} \dots \partial x_d^{\lambda_d}} \right|^2 \right] dx \\
&\quad + \frac{\Delta t k a_0}{4 \lambda L} \sum_{i=1}^K \int_{\Omega \times D} M |\nabla_{q_i} \hat{\psi}_{\kappa,\alpha,L,\delta}^n|^2 dq dx + \frac{\Delta t k \varepsilon}{L} \int_{\Omega \times D} M |\nabla_x \hat{\psi}_{\kappa,\alpha,L,\delta}^n|^2 dq dx \\
&\quad + \mathfrak{z} \left[\|\varrho_{\kappa,\alpha,L,\delta}^n\|_{L^2(\Omega)}^2 + \|\varrho_{\kappa,\alpha,L,\delta}^n - \varrho_{\kappa,\alpha,L}^{n-1}\|_{L^2(\Omega)}^2 \right] + 2 \Delta t \mathfrak{z} \varepsilon \|\nabla_x \varrho_{\kappa,\alpha,L,\delta}^n\|_{L^2(\Omega)}^2 \\
&\leq \frac{1}{2} \int_{\Omega} \rho_{\kappa,\alpha,L}^{n-1} |u_{\kappa,\alpha,L}^{n-1}|^2 dx + \int_{\Omega} P_{\kappa}(\rho_{\kappa,\alpha,L}^{n-1}) dx + k \int_{\Omega \times D} M \mathcal{F}_{\delta}^L(\hat{\psi}_{\kappa,\alpha,L}^{n-1}) dq dx \\
&\quad + \mathfrak{z} \|\varrho_{\kappa,\alpha,L}^{n-1}\|_{L^2(\Omega)}^2 + \Delta t \int_{\Omega} \rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t_n) f^n \cdot u_{\kappa,\alpha,L,\delta}^n dx \\
&\leq \frac{1}{2} \int_{\Omega} \rho_{\kappa,\alpha,L}^{n-1} |u_{\kappa,\alpha,L}^{n-1}|^2 dx + \int_{\Omega} P_{\kappa}(\rho_{\kappa,\alpha,L}^{n-1}) dx + k \int_{\Omega \times D} M \mathcal{F}^L(\hat{\psi}_{\kappa,\alpha,L}^{n-1}) dq dx \\
&\quad + \mathfrak{z} \|\varrho_{\kappa,\alpha,L}^{n-1}\|_{L^2(\Omega)}^2 + \frac{\Delta t}{2} \left[\varsigma \int_{\Omega} \rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t_n) |u_{\kappa,\alpha,L,\delta}^n|^2 dx + \frac{1}{\varsigma} \|f^n\|_{L^\infty(\Omega)}^2 \int_{\Omega} \rho_{\kappa,\alpha,L}^{n-1} dx \right] \\
&\leq C,
\end{aligned}
\tag{3.86}$$

where C is independent of δ and Δt .

On choosing, for any $s \in (t_{n-1}, t_n]$, $\eta(\cdot, t) = \chi_{[t_{n-1}, s]} [\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t)]^{\vartheta-1}$, for $\vartheta = 2$ and $\frac{\Gamma}{2}$, in (3.51a), we obtain, on noting (3.33), (3.50) and (3.86), that

$$\begin{aligned}
&\frac{1}{\vartheta} \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, s)\|_{L^{\vartheta}(\Omega)}^{\vartheta} + \frac{4\alpha(\vartheta-1)}{\vartheta^2} \int_{t_{n-1}}^s \|\nabla_x[(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\frac{\vartheta}{2}}]\|_{L^2(\Omega)}^2 dt \\
&= \frac{1}{\vartheta} \left[\|\rho_{\kappa,\alpha,L,\delta}^{n-1}\|_{L^{\vartheta}(\Omega)}^{\vartheta} + (\vartheta-1) \int_{t_{n-1}}^s \int_{\Omega} u_{\kappa,\alpha,L,\delta}^n \cdot \nabla_x [(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\vartheta}] dx dt \right] \\
&\leq \frac{1}{\vartheta} \left[\|\rho_{\kappa,\alpha,L}^{n-1}\|_{L^{\vartheta}(\Omega)}^{\vartheta} + \Delta t \|u_{\kappa,\alpha,L,\delta}^n\|_{L^2(\Omega)}^2 + \frac{(\vartheta-1)^2}{4} \int_{t_{n-1}}^{t_n} \|\nabla_x[(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\vartheta}]\|_{L^2(\Omega)}^2 dt \right] \\
(3.87) \quad &\leq C,
\end{aligned}$$

where C is independent of δ and Δt . On denoting by \bar{f} the mean value of the function f over Ω , it follows from a Poincaré inequality, (3.86) and (3.87) for $\vartheta = \frac{\Gamma}{2}$ that

$$\begin{aligned}
 \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^\Gamma(t_{n-1},t_n;L^\Gamma(\Omega))}^\Gamma &= \|(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\frac{\Gamma}{2}}\|_{L^2(t_{n-1},t_n;L^2(\Omega))}^2 \\
 &\leq 2\|(I - \bar{f})(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\frac{\Gamma}{2}}\|_{L^2(t_{n-1},t_n;L^2(\Omega))}^2 + 2\|\bar{f}(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\frac{\Gamma}{2}}\|_{L^2(t_{n-1},t_n;L^2(\Omega))}^2 \\
 (3.88) \quad &\leq C\|\nabla_x[(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\frac{\Gamma}{2}}]\|_{L^2(t_{n-1},t_n;L^2(\Omega))}^2 + C\Delta t\|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^\infty(t_{n-1},t_n;L^{\frac{\Gamma}{2}}(\Omega))}^\Gamma \leq C.
 \end{aligned}$$

Next, we obtain from (3.33), (3.2), (3.86) and (3.87) for $\vartheta = \frac{\Gamma}{2}$, on recalling that $\Gamma \geq 8$, that

$$\begin{aligned}
 \left| \int_{t_{n-1}}^{t_n} c(\tilde{u}_{\kappa,\alpha,L,\delta}^n)(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}, \eta) dt \right| &\leq \alpha \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^2(t_{n-1},t_n;H^1(\Omega))} \|\eta\|_{L^2(t_{n-1},t_n;H^1(\Omega))} \\
 &\quad + \left| \int_{t_{n-1}}^{t_n} \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^3(\Omega)} \|\tilde{u}_{\kappa,\alpha,L,\delta}^n\|_{L^6(\Omega)} \|\nabla_x \eta\|_{L^2(\Omega)} dt \right| \\
 &\leq C \|\eta\|_{L^2(t_{n-1},t_n;H^1(\Omega))} + C \left| \int_{t_{n-1}}^{t_n} \|\tilde{u}_{\kappa,\alpha,L,\delta}^n\|_{H^1(\Omega)} \|\eta\|_{H^1(\Omega)} dt \right| \\
 &\leq C \left[1 + \left(\Delta t \|\tilde{u}_{\kappa,\alpha,L,\delta}^n\|_{H^1(\Omega)}^2 \right)^{\frac{1}{2}} \right] \|\eta\|_{L^2(t_{n-1},t_n;H^1(\Omega))} \\
 (3.89) \quad &\leq C \|\eta\|_{L^2(t_{n-1},t_n;H^1(\Omega))} \quad \forall \eta \in L^2(t_{n-1},t_n;H^1(\Omega)).
 \end{aligned}$$

Hence, we deduce from (3.87) for $\vartheta = 2$ and $\frac{\Gamma}{2}$, (3.89), on noting (3.51a), (3.86), (3.5) and (3.88) that $\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n} \in \Upsilon^n$ is such that

$$\begin{aligned}
 \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^\infty(t_{n-1},t_n;L^{\frac{\Gamma}{2}}(\Omega))} &+ \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^2(t_{n-1},t_n;H^1(\Omega))}^2 + \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{H^1(t_{n-1},t_n;H^1(\Omega))}^2 \\
 (3.90) \quad &+ \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t_n)\|_{L^\Gamma(\Omega)}^\Gamma + \|(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\frac{\Gamma}{2}}\|_{L^2(t_{n-1},t_n;H^1(\Omega))}^2 \leq C,
 \end{aligned}$$

where C is independent of δ and Δt . Furthermore, we deduce from (3.2) and the last bound in (3.90) that

$$\begin{aligned}
 \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^\Gamma(t_{n-1},t_n;L^{3\Gamma}(\Omega))}^\Gamma &= \|(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\frac{\Gamma}{2}}\|_{L^2(t_{n-1},t_n;L^6(\Omega))}^2 \leq C \|(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^{\frac{\Gamma}{2}}\|_{L^2(t_{n-1},t_n;H^1(\Omega))}^2 \leq C.
 \end{aligned}$$

Finally, it follows from (3.1) with $v = \frac{4\Gamma}{3}$, $r = \frac{\Gamma}{2}$ and $s = 3\Gamma$ yielding $\vartheta = \frac{3}{4}$, the first bound in (3.90) and (3.91) that

$$(3.92) \quad \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^{\frac{4\Gamma}{3}}(t_{n-1},t_n;L^{\frac{4\Gamma}{3}}(\Omega))}^{\frac{4\Gamma}{3}} \leq \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^\infty(t_{n-1},t_n;L^{\frac{\Gamma}{2}}(\Omega))}^{\frac{\Gamma}{3}} \|\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}\|_{L^\Gamma(t_{n-1},t_n;L^{3\Gamma}(\Omega))}^\Gamma \leq C,$$

where C is independent of δ and Δt .

As the bounds (3.86), (3.91) and (3.92) are independent of δ , we are now ready to pass to the limit $\delta \rightarrow 0_+$ in (3.51a-c), to deduce the existence of a solution $\{(\rho_{\kappa,\alpha,L}^{[\Delta t],n}, \tilde{u}_{\kappa,\alpha,L}^n, \hat{\psi}_{\kappa,\alpha,L}^n)\}_{n=1}^N$ to $(P_{\kappa,\alpha,L}^{\Delta t})$.

Lemma 3.4. *There exists a subsequence (not indicated) of $\{(\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}, \tilde{u}_{\kappa,\alpha,L,\delta}^n, \hat{\psi}_{\kappa,\alpha,L,\delta}^n)\}_{\delta>0}$, and functions $\rho_{\kappa,\alpha,L}^{[\Delta t],n} \in Y^n$ with $\rho_{\kappa,\alpha,L}^n(\cdot) = \rho_{\kappa,\alpha,L}^{[\Delta t],n}(\cdot, t_n) \in L_{\geq 0}^\Gamma(\Omega)$, $\tilde{u}_{\kappa,\alpha,L}^n \in H_0^1(\Omega)$ and $\hat{\psi}_{\kappa,\alpha,L}^n \in X \cap Z_2$, $n = 1, \dots, N$, with*

$$(3.93) \quad \varrho_{\kappa,\alpha,L}^n(\cdot) := \int_D M(q) \hat{\psi}_{\kappa,\alpha,L}^n(\cdot, q) dq \in H^1(\Omega), \quad n = 1, \dots, N,$$

such that, as $\delta \rightarrow 0_+$,

$$(3.94a) \quad \rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n} \rightarrow \rho_{\kappa,\alpha,L}^{[\Delta t],n} \quad \text{weakly in } L^2(t_{n-1}, t_n; H^1(\Omega)), \quad \text{weakly in } H^1(t_{n-1}, t_n; H^1(\Omega)'),$$

$$(3.94b) \quad \rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n} \rightarrow \rho_{\kappa,\alpha,L}^{[\Delta t],n} \quad \text{strongly in } L^2(t_{n-1}, t_n; L^r(\Omega)), \quad \text{strongly in } L^v(\Omega \times (t_{n-1}, t_n)),$$

$$(3.94c) \quad \rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n}(\cdot, t_n) \rightarrow \rho_{\kappa,\alpha,L}^n(\cdot) \quad \text{weakly in } L^\Gamma(\Omega),$$

$$(3.94d) \quad (\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n})^\vartheta \rightarrow (\rho_{\kappa,\alpha,L}^{[\Delta t],n})^\vartheta \quad \text{weakly in } L^2(t_{n-1}, t_n; H^1(\Omega)), \quad \vartheta = 2 \text{ and } \frac{\Gamma}{2},$$

$$(3.95a) \quad \underset{\sim}{u}_{\kappa,\alpha,L,\delta}^n \rightarrow \underset{\sim}{u}_{\kappa,\alpha,L}^n \quad \text{weakly in } \underset{\sim}{H}_0^1(\Omega), \quad \text{strongly in } \underset{\sim}{L}^r(\Omega),$$

$$(3.95b) \quad \delta \frac{\partial^{|\lambda|} \underset{\sim}{u}_{\kappa,\alpha,L,\delta}^n}{\partial_{x_1}^{\lambda_1} \dots \partial_{x_d}^{\lambda_d}} \rightarrow 0 \quad \text{strongly in } \underset{\sim}{L}^2(\Omega), \quad \forall |\lambda| = 2,$$

$$(3.95c) \quad \delta \frac{\partial^{|\lambda|} (\nabla_x \cdot \underset{\sim}{u}_{\kappa,\alpha,L,\delta}^n)}{\partial_{x_1}^{\lambda_1} \dots \partial_{x_d}^{\lambda_d}} \rightarrow 0 \quad \text{strongly in } L^2(\Omega), \quad \forall |\lambda| = 2,$$

where $r \in [1, \infty)$ if $d = 2$ and $r \in [1, 6)$ if $d = 3$, and $v \in [1, \frac{4\Gamma}{3})$; and

$$(3.96a) \quad M^{\frac{1}{2}} \underset{\sim}{\nabla}_q \widehat{\psi}_{\kappa,\alpha,L,\delta}^n \rightarrow M^{\frac{1}{2}} \underset{\sim}{\nabla}_q \widehat{\psi}_{\kappa,\alpha,L}^n \quad \text{weakly in } \underset{\sim}{L}^2(\Omega \times D),$$

$$(3.96b) \quad M^{\frac{1}{2}} \underset{\sim}{\nabla}_x \widehat{\psi}_{\kappa,\alpha,L,\delta}^n \rightarrow M^{\frac{1}{2}} \underset{\sim}{\nabla}_x \widehat{\psi}_{\kappa,\alpha,L}^n \quad \text{weakly in } \underset{\sim}{L}^2(\Omega \times D),$$

$$(3.96c) \quad M^{\frac{1}{2}} \widehat{\psi}_{\kappa,\alpha,L,\delta}^n \rightarrow M^{\frac{1}{2}} \widehat{\psi}_{\kappa,\alpha,L}^n \quad \text{strongly in } L^2(\Omega \times D),$$

$$(3.96d) \quad M^{\frac{1}{2}} \beta_\delta^L(\widehat{\psi}_{\kappa,\alpha,L,\delta}^n) \rightarrow M^{\frac{1}{2}} \beta^L(\widehat{\psi}_{\kappa,\alpha,L}^n) \quad \text{strongly in } L^s(\Omega \times D),$$

$$(3.96e) \quad C_i(M \widehat{\psi}_{\kappa,\alpha,L,\delta}^n) \rightarrow C_i(M \widehat{\psi}_{\kappa,\alpha,L}^n) \quad \text{strongly in } \underset{\sim}{L}^2(\Omega), \quad i = 1, \dots, K,$$

$$(3.96f) \quad \varrho_{\kappa,\alpha,L,\delta}^n \rightarrow \varrho_{\kappa,\alpha,L}^n \quad \text{weakly in } H^1(\Omega), \quad \text{strongly in } L^r(\Omega),$$

where $s \in [1, \infty)$. Furthermore, $(\rho_{\kappa,\alpha,L}^{[\Delta t],n}, \underset{\sim}{u}_{\kappa,\alpha,L}^n, \widehat{\psi}_{\kappa,\alpha,L}^n)$ solves (3.29a–c) for $n = 1, \dots, N$. Hence, there exists a solution $\{(\rho_{\kappa,\alpha,L}^{[\Delta t],n}, \underset{\sim}{u}_{\kappa,\alpha,L}^n, \widehat{\psi}_{\kappa,\alpha,L}^n)\}_{n=1}^N$ to $(P_{\kappa,\alpha,L}^{\Delta t})$.

Proof. The weak convergence results (3.94a,c,d) follow immediately from (3.90). The strong convergence results (3.94b) follow from (3.94a), (3.11), (3.92) and the interpolation result (3.1). Hence $\rho_{\kappa,\alpha,L}^{[\Delta t],n} \in Y^n$ with $\rho_{\kappa,\alpha,L}^n(\cdot) = \rho_{\kappa,\alpha,L}^{[\Delta t],n}(\cdot, t_n) \in L_{\geq 0}^\Gamma(\Omega)$ as $\rho_{\kappa,\alpha,L,\delta}^{[\Delta t],n} \in \Upsilon^n$. The weak convergence result (3.95a) and the strong convergence results (3.95b,c) follow immediately from (3.86), and hence $\underset{\sim}{u}_{\kappa,\alpha,L}^n \in \underset{\sim}{H}_0^1(\Omega)$ as $\underset{\sim}{u}_{\kappa,\alpha,L,\delta}^n \in \mathcal{Y}$. The strong convergence result (3.95a) follows as $\underset{\sim}{H}^1(\Omega)$ is compactly embedded in $\underset{\sim}{L}^r(\Omega)$ for the stated values of r .

The weak convergence results (3.96a,b) follow from (3.86); the strong convergence result (3.96c) and the fact that $\widehat{\psi}_{\kappa,\alpha,L}^n \geq 0$ a.e. on $\Omega \times D$ follow from the fourth bound in (3.86), (3.44b) and (3.10b). Hence $\widehat{\psi}_{\kappa,\alpha,L}^n \in X \cap Z_2$. The desired results (3.96d,e) follow from (3.96c), (3.45), (1.12) and (3.15). See the proof of Lemma 3.3 in [11] for details of the results (3.96a–e). Finally, (3.96f) follows from (3.86) and (3.96c).

It follows from (3.94a–c), (3.95a–c), (3.96a–f), (3.33), (3.47), (3.48), (3.35a,b), (3.37a,b) and (3.8) that we may pass to the limit $\delta \rightarrow 0_+$ in (3.51a–c) to obtain that $(\rho_{\kappa,\alpha,L}^{[\Delta t],n}, \underset{\sim}{u}_{\kappa,\alpha,L}^n, \widehat{\psi}_{\kappa,\alpha,L}^n)$ solves (3.32), (3.34), and (3.36); that is, (3.29a–c).

Finally, as $(\rho_{\kappa,\alpha,L}^0, \underset{\sim}{u}_{\kappa,\alpha,L}^0, \widehat{\psi}_{\kappa,\alpha,L}^0) \in L_{\geq 0}^\Gamma(\Omega) \times \underset{\sim}{H}_0^1(\Omega) \times Z_2$, performing the above existence proof at each time level t_n , $n = 1, \dots, N$, yields a solution $\{(\rho_{\kappa,\alpha,L}^{[\Delta t],n}, \underset{\sim}{u}_{\kappa,\alpha,L}^n, \widehat{\psi}_{\kappa,\alpha,L}^n)\}_{n=1}^N$ to $(P_{\kappa,\alpha,L}^{\Delta t})$ with

$\rho_{\kappa,\alpha,L}^n(\cdot) = \rho_{\kappa,\alpha,L}^{[\Delta t],n}(\cdot, t_n)$, $n = 1, \dots, N$, by noting that $\rho_{\kappa,\alpha,L}^{[\Delta t],n}$ thus constructed is an element of $C([t_{n-1}, t_n]; L^2(\Omega))$, $n = 1, \dots, N$. \square

4. EXISTENCE OF A SOLUTION TO $(P_{\kappa,\alpha})$

Next, we derive bounds on the solution of $(P_{\kappa,\alpha,L}^{\Delta t})$, independent of Δt and L . Our starting point is Lemma 3.4, concerning the existence of a solution to the problem $(P_{\kappa,\alpha,L}^{\Delta t})$. The model $(P_{\kappa,\alpha,L}^{\Delta t})$ includes ‘microscopic cut-off’ in the drag and convective terms of the Fokker–Planck equation, where $L > 1$ is a (fixed, but otherwise arbitrary,) cut-off parameter. Our next objective is to pass to the limits $L \rightarrow \infty$ and $\Delta t \rightarrow 0_+$ in the model $(P_{\kappa,\alpha,L}^{\Delta t})$, with L and Δt linked by the condition $\Delta t = o(L^{-1})$, as $L \rightarrow \infty$. To that end, we need to develop various bounds on sequences of weak solutions of $(P_{\kappa,\alpha,L}^{\Delta t})$ that are uniform in the time step Δt and the cut-off parameter L , and thus permit the extraction of weakly convergent subsequences, as $L \rightarrow \infty$, through the use of a weak compactness argument. The derivation of such bounds, based on the use of the relative entropy associated with the Maxwellian M , is our main task in this section.

We define

(4.1a)

$$\rho_{\kappa,\alpha,L}^{[\Delta t]} := \rho_{\kappa,\alpha,L}^{[\Delta t],n}, \quad t \in [t_{n-1}, t_n], \quad n = 1, \dots, N, \quad \Rightarrow \quad \rho_{\kappa,\alpha,L}^{[\Delta t]}(\cdot, t_n) = \rho_{\kappa,\alpha,L}^n(\cdot), \quad n = 0, \dots, N, \quad (4.1b)$$

$$\rho_{\kappa,\alpha,L}^{\{\Delta t\}}(\cdot, t) := \frac{1}{\Delta t} \int_{t_{n-1}}^{t_n} \rho_{\kappa,\alpha,L}^{[\Delta t]}(\cdot, s) ds, \quad t \in (t_{n-1}, t_n], \quad n = 1, \dots, N.$$

Further, we define the pressure variable

$$(4.2) \quad p_{\kappa,\alpha,L}^{\{\Delta t\}}(\cdot, t) := \frac{1}{\Delta t} \int_{t_{n-1}}^{t_n} p_{\kappa}(\rho_{\kappa,\alpha,L}^{[\Delta t]}(\cdot, s)) ds, \quad t \in (t_{n-1}, t_n], \quad n = 1, \dots, N,$$

and the momentum variable

$$(4.3) \quad \tilde{m}_{\kappa,\alpha,L}^n := \rho_{\kappa,\alpha,L}^n \tilde{u}_{\kappa,\alpha,L}^n, \quad n = 0, \dots, N.$$

We then introduce the following definitions:

$$(4.4a) \quad \tilde{u}_{\kappa,\alpha,L}^{\Delta t}(\cdot, t) := \frac{t - t_{n-1}}{\Delta t} \tilde{u}_{\kappa,\alpha,L}^n(\cdot) + \frac{t_n - t}{\Delta t} \tilde{u}_{\kappa,\alpha,L}^{n-1}(\cdot), \quad t \in [t_{n-1}, t_n], \quad n = 1, \dots, N,$$

$$(4.4b) \quad \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+}(\cdot, t) := \tilde{u}_{\kappa,\alpha,L}^n(\cdot), \quad \tilde{u}_{\kappa,\alpha,L}^{\Delta t,-}(\cdot, t) := \tilde{u}_{\kappa,\alpha,L}^{n-1}(\cdot), \quad t \in (t_{n-1}, t_n], \quad n = 1, \dots, N.$$

We shall adopt $\tilde{u}_{\kappa,\alpha,L}^{\Delta t(\pm)}$ as a collective symbol for $\tilde{u}_{\kappa,\alpha,L}^{\Delta t}$, $\tilde{u}_{\kappa,\alpha,L}^{\Delta t,\pm}$. The corresponding notations $\hat{\psi}_{\kappa,\alpha,L}^{\Delta t}$, $\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,\pm}$ and $\hat{\psi}_{\kappa,\alpha,L}^{\Delta t(\pm)}$; $\rho_{\kappa,\alpha,L}^{\Delta t}$, $\rho_{\kappa,\alpha,L}^{\Delta t,\pm}$ and $\rho_{\kappa,\alpha,L}^{\Delta t(\pm)}$; $\tilde{m}_{\kappa,\alpha,L}^{\Delta t}$, $\tilde{m}_{\kappa,\alpha,L}^{\Delta t,\pm}$ and $\tilde{m}_{\kappa,\alpha,L}^{\Delta t(\pm)}$, and $\varrho_{\kappa,\alpha,L}^{\Delta t}$, $\varrho_{\kappa,\alpha,L}^{\Delta t,\pm}$ and $\varrho_{\kappa,\alpha,L}^{\Delta t(\pm)}$ are defined analogously. The notation $\rho_{\kappa,\alpha,L}^{\Delta t}$ signifying the piecewise linear interpolant of $\rho_{\kappa,\alpha,L}^{[\Delta t]}$ with respect to the variable t is not to be confused with $\rho_{\kappa,\alpha,L}^{[\Delta t]}$, itself, which denotes the function defined piecewise, over the union of time slabs $\Omega \times [t_{n-1}, t_n]$, $n = 1, \dots, N$, solving (3.29a) subject to the initial condition $\rho_{\kappa,\alpha,L}^{[\Delta t]}(\cdot, t_{n-1}) = \rho_{\kappa,\alpha,L}^{n-1}(\cdot)$, $n = 1, \dots, N$, with $\rho_{\kappa,\alpha,L}^0 := \rho^0$.

Using the above notation, (3.29a–c) summed for $n = 1, \dots, N$ can be restated in the form: find $(\rho_{\kappa,\alpha,L}^{[\Delta t]}(\cdot, t), \tilde{u}_{\kappa,\alpha,L}^{\Delta t}(\cdot, t), \hat{\psi}_{\kappa,\alpha,L}^{\Delta t}(\cdot, t)) \in H^1(\Omega) \cap L_{\geq 0}^{\frac{4\Gamma}{3}}(\Omega) \times \tilde{H}_0^1(\Omega) \times (X \cap Z_2)$, with $\frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t]}}{\partial t}(\cdot, t) \in H^1(\Omega)'$, a.e. $t \in (0, T)$, such that $\tilde{m}_{\kappa,\alpha,L}^{\Delta t}$ is defined via (4.3) and

$$(4.5a) \quad \int_0^T \left\langle \frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t]}}{\partial t}, \eta \right\rangle_{H^1(\Omega)} dt + \int_0^T \int_{\Omega} \left(\alpha \tilde{\nabla}_x \rho_{\kappa,\alpha,L}^{[\Delta t]} - \rho_{\kappa,\alpha,L}^{[\Delta t]} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \right) \cdot \tilde{\nabla}_x \eta \, dx \, dt = 0 \quad \forall \eta \in L^2(0, T; H^1(\Omega)),$$

$$\begin{aligned}
& \int_0^T \int_\Omega \left[\frac{\partial m_{\kappa,\alpha,L}^{\Delta t}}{\partial t} - \frac{1}{2} \frac{\partial \rho_{\kappa,\alpha,L}^{\Delta t}}{\partial t} u_{\kappa,\alpha,L}^{\Delta t,+} \right] \cdot w \, dx \, dt + \int_0^T \int_\Omega S(u_{\kappa,\alpha,L}^{\Delta t,+}) : \nabla_x w \, dx \, dt \\
& \quad + \frac{1}{2} \int_0^T \int_\Omega \left[[(m_{\kappa,\alpha,L}^{\Delta t,-} \cdot \nabla_x) u_{\kappa,\alpha,L}^{\Delta t,+}] \cdot w - [(m_{\kappa,\alpha,L}^{\Delta t,-} \cdot \nabla_x) w] \cdot u_{\kappa,\alpha,L}^{\Delta t,+} \right] dx \, dt \\
& \quad - \int_0^T \int_\Omega p_{\kappa,\alpha,L}^{\{\Delta t\}} \nabla_x \cdot w \, dx \, dt \\
& = \int_0^T \int_\Omega \left[\rho_{\kappa,\alpha,L}^{\Delta t,+} f^{\{\Delta t\}} \cdot w - \tau_1(M \hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) : \nabla_x w \right] dx \, dt \\
(4.5b) \quad & - 2 \mathfrak{z} \int_0^T \int_\Omega \left(\int_D M \beta^L(\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \, dq \right) \nabla_x \varrho_{\kappa,\alpha,L}^{\Delta t,+} \cdot w \, dx \, dt \quad \forall w \in L^2(0, T; V),
\end{aligned}$$

$$\begin{aligned}
& \int_0^T \int_{\Omega \times D} M \frac{\partial \hat{\psi}_{\kappa,\alpha,L}^{\Delta t}}{\partial t} \varphi \, dq \, dx \, dt + \frac{1}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \int_0^T \int_{\Omega \times D} M \nabla_{q_j} \hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt \\
& \quad + \int_0^T \int_{\Omega \times D} M \left[\varepsilon \nabla_x \hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,+} \beta^L(\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \right] \cdot \nabla_x \varphi \, dq \, dx \, dt \\
(4.5c) \quad & - \int_0^T \int_{\Omega \times D} M \sum_{i=1}^K \left[\sigma(u_{\kappa,\alpha,L}^{\Delta t,+}) q_i \right] \beta^L(\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt = 0 \quad \forall \varphi \in L^2(0, T; X);
\end{aligned}$$

subject to the initial conditions $\rho_{\kappa,\alpha,L}^{\Delta t}(0) = \rho^0 \in L_{\geq 0}^\infty(\Omega)$, $u_{\kappa,\alpha,L}^{\Delta t}(0) = u^0 \in \underline{H}_0^1(\Omega)$ and $\hat{\psi}_{\kappa,\alpha,L}^{\Delta t}(0) = \hat{\psi}^0 \in X \cap Z_2$, where we recall (3.19) and (3.21). We emphasize that (4.5a–c) is an equivalent restatement of problem $(P_{\kappa,\alpha,L}^{\Delta t})$, for which existence of a solution has been established (cf. Lemma 3.4).

We are now ready to embark on the derivation of the required bounds, uniform in the time step Δt and the cut-off parameter L , on norms of $\rho_{\kappa,\alpha,L}^{[\Delta t]}(t) \in H_0^1(\Omega) \cap L_{\geq 0}^{\frac{4\Gamma}{3}}(\Omega)$, $u_{\kappa,\alpha,L}^{\Delta t,+}(t) \in \underline{H}_0^1(\Omega)$, $\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}(t) \in X \cap Z_2$ and $\varrho_{\kappa,\alpha,L}^{\Delta t,+}(t) \in H^1(\Omega)$, $t \in (0, T]$.

4.1. L , Δt -independent bounds on the spatial derivatives of $u_{\kappa,\alpha,L}^{\Delta t}$ and $\hat{\psi}_{\kappa,\alpha,L}^{\Delta t}$. We note that it is *not* possible to pass to the limit $\delta \rightarrow 0_+$ in (3.86) to obtain strong enough L -independent bounds due to the fourth, seventh and eighth of the ten terms on the left-hand side. Similarly, it is *not* possible to pass to the limit in these terms even before we use the bound $[\mathcal{F}_\delta^L]''(\cdot) \geq \frac{1}{L}$; recall its use in (3.81) to obtain (3.82), and hence (3.86). However, it is a simple matter to pass to the limit $\delta \rightarrow 0_+$ in (3.85). Noting (3.94c,d), (3.95a), (3.96e,f) and the convexity of $P_\kappa(\cdot)$, we may pass to the limit $\delta \rightarrow 0_+$ in (3.85) to obtain for $n = 1, \dots, N$ that

$$\begin{aligned}
& \frac{1}{2} \int_\Omega \left[\rho_{\kappa,\alpha,L}^n |u_{\kappa,\alpha,L}^n|^2 + \rho_{\kappa,\alpha,L}^{n-1} |u_{\kappa,\alpha,L}^n - u_{\kappa,\alpha,L}^{n-1}|^2 \right] dx + \int_\Omega P_\kappa(\rho_{\kappa,\alpha,L}^n) dx \\
& \quad + \alpha \kappa \int_{t_{n-1}}^{t_n} \left[\|\nabla_x [(\rho_{\kappa,\alpha,L}^{[\Delta t],n})^2]\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x [(\rho_{\kappa,\alpha,L}^{[\Delta t],n})^{\frac{\Gamma}{2}}]\|_{L^2(\Omega)}^2 \right] dt + \Delta t \mu^S c_0 \|u_{\kappa,\alpha,L}^n\|_{H^1(\Omega)}^2 \\
& \quad + \mathfrak{z} \left[\|\varrho_{\kappa,\alpha,L}^n\|_{L^2(\Omega)}^2 + \|\varrho_{\kappa,\alpha,L}^n - \varrho_{\kappa,\alpha,L}^{n-1}\|_{L^2(\Omega)}^2 \right] + 2 \Delta t \mathfrak{z} \varepsilon \|\nabla_x \varrho_{\kappa,\alpha,L}^n\|_{L^2(\Omega)}^2 \\
& \leq \frac{1}{2} \int_\Omega \rho_{\kappa,\alpha,L}^{n-1} |u_{\kappa,\alpha,L}^{n-1}|^2 dx + \int_\Omega P_\kappa(\rho_{\kappa,\alpha,L}^{n-1}) dx + \Delta t \int_\Omega \rho_{\kappa,\alpha,L}^n f^n \cdot u_{\kappa,\alpha,L}^n dx \\
& \quad + \mathfrak{z} \|\varrho_{\kappa,\alpha,L}^{n-1}\|_{L^2(\Omega)}^2 + k \Delta t (K+1) \int_\Omega \varrho_{\kappa,\alpha,L}^n \nabla_x \cdot u_{\kappa,\alpha,L}^n dx \\
(4.6) \quad & - k \Delta t \sum_{i=1}^K \int_\Omega C_i(M \hat{\psi}_{\kappa,\alpha,L}^n) : \nabla_x u_{\kappa,\alpha,L}^n dx.
\end{aligned}$$

Summing the above over n , and adopting the notation (4.1a), (4.4a,b) and (3.24), we obtain for $n = 1, \dots, N$ that

$$\begin{aligned}
& \frac{1}{2} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,+}(t_n) |\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+}(t_n)|^2 dx + \frac{1}{2\Delta t} \int_0^{t_n} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,-} |\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} - \tilde{u}_{\kappa,\alpha,L}^{\Delta t,-}|^2 dx dt \\
& + \int_{\Omega} P_{\kappa}(\rho_{\kappa,\alpha,L}^{\Delta t,+}(t_n)) dx + \mu^S c_0 \int_0^{t_n} \|\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+}\|_{H^1(\Omega)}^2 dt \\
& + \alpha \kappa \int_0^{t_n} \left[\|\tilde{\nabla}_x [(\rho_{\kappa,\alpha,L}^{\Delta t,+})^2]\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\tilde{\nabla}_x [(\rho_{\kappa,\alpha,L}^{\Delta t,+})^{\frac{\Gamma}{2}}]\|_{L^2(\Omega)}^2 \right] dt + \mathfrak{z} \|\varrho_{\kappa,\alpha,L}^{\Delta t,+}(t_n)\|_{L^2(\Omega)}^2 \\
& + \mathfrak{z} \int_0^{t_n} \left[\|\varrho_{\kappa,\alpha,L}^{\Delta t,+} - \varrho_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^2(\Omega)}^2 + 2\varepsilon \|\tilde{\nabla}_x \varrho_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(\Omega)}^2 \right] dt \\
& \leq \frac{1}{2} \int_{\Omega} \rho^0 |u^0|^2 dx + \int_{\Omega} P_{\kappa}(\rho^0) dx + \mathfrak{z} \|\varrho^0\|_{L^2(\Omega)}^2 \\
& + \int_0^{t_n} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,+} \tilde{f} \cdot \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} dx dt + k(K+1) \int_0^{t_n} \int_{\Omega} \varrho_{\kappa,\alpha,L}^{\Delta t,+} \tilde{\nabla}_x \cdot \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} dx dt \\
& - k \sum_{i=1}^K \int_0^{t_n} \int_{\Omega} C_i(M \hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) : \tilde{\nabla}_x \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} dx dt.
\end{aligned} \tag{4.7}$$

We now require the appropriate $\hat{\psi}_{\kappa,\alpha,L}^n$ analogue of (3.81). The appropriate choice of test function in (3.29c) for this purpose is $\varphi = [\mathcal{F}^L]'(\hat{\psi}_{\kappa,\alpha,L}^n)$. While Lemma 3.4 guarantees that $\hat{\psi}_{\kappa,\alpha,L}^n$ belongs to Z_2 , and is therefore nonnegative a.e. on $\Omega \times D$, there is unfortunately no reason why $\hat{\psi}_{\kappa,\alpha,L}^n$ should be strictly positive on $\Omega \times D$, and therefore the expression $[\mathcal{F}^L]'(\hat{\psi}_{\kappa,\alpha,L}^n)$ may in general be undefined. Similarly to (3.76), we shall circumvent this problem by choosing $\varphi = [\mathcal{F}^L]'(\hat{\psi}_{\kappa,\alpha,L}^n + \varsigma)$ in (3.29c), which leads, for any fixed $\varsigma \in \mathbb{R}_{>0}$, to

$$\begin{aligned}
0 &= \int_{\Omega \times D} M \frac{\hat{\psi}_{\kappa,\alpha,L}^n - \hat{\psi}_{\kappa,\alpha,L}^{n-1}}{\Delta t} [[\mathcal{F}^L]'(\hat{\psi}_{\kappa,\alpha,L}^n + \varsigma)] dq dx \\
& - \int_{\Omega \times D} M \tilde{u}_{\kappa,\alpha,L}^n \cdot \left(\tilde{\nabla}_x [[\mathcal{F}^L]'(\hat{\psi}_{\kappa,\alpha,L}^n + \varsigma)] \right) \beta^L(\hat{\psi}_{\kappa,\alpha,L}^n) dq dx \\
& + \frac{1}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K \int_{\Omega \times D} A_{ij} M \tilde{\nabla}_{q_j} \hat{\psi}_{\kappa,\alpha,L}^n \cdot \tilde{\nabla}_{q_i} [[\mathcal{F}^L]'(\hat{\psi}_{\kappa,\alpha,L}^n + \varsigma)] dq dx \\
& + \varepsilon \int_{\Omega \times D} M \tilde{\nabla}_x \hat{\psi}_{\kappa,\alpha,L}^n \cdot \tilde{\nabla}_x [[\mathcal{F}^L]'(\hat{\psi}_{\kappa,\alpha,L}^n + \varsigma)] dq dx \\
& - \sum_{i=1}^K \int_{\Omega \times D} M \beta^L(\hat{\psi}_{\kappa,\alpha,L}^n) [\sigma(\tilde{u}_{\kappa,\alpha,L}^n) q_i] \cdot \tilde{\nabla}_{q_i} [[\mathcal{F}^L]'(\hat{\psi}_{\kappa,\alpha,L}^n + \varsigma)] dq dx \\
& =: \sum_{i=1}^5 \mathbf{T}_i.
\end{aligned} \tag{4.8}$$

It follows from (3.40a) that

$$\begin{aligned}
\mathbf{T}_1 &\geq \frac{1}{\Delta t} \int_{\Omega \times D} M \left[\mathcal{F}^L(\hat{\psi}_{\kappa,\alpha,L}^n + \varsigma) - \mathcal{F}^L(\hat{\psi}_{\kappa,\alpha,L}^{n-1} + \varsigma) \right] dq dx \\
& + \frac{1}{2\Delta t L} \int_{\Omega \times D} M (\hat{\psi}_{\kappa,\alpha,L}^n - \hat{\psi}_{\kappa,\alpha,L}^{n-1})^2 dq dx.
\end{aligned} \tag{4.9}$$

In addition, it follows from (3.40a) that

$$\begin{aligned}
T_2 &= - \int_{\Omega \times D} M \frac{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^n)}{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^n + \varsigma)} \underset{\sim}{u}_{\kappa,\alpha,L}^n \cdot \underset{\sim}{\nabla}_x \widehat{\psi}_{\kappa,\alpha,L}^n \underset{\sim}{dq} \underset{\sim}{dx} \\
&= \int_{\Omega} \left(\int_D M \widehat{\psi}_{\kappa,\alpha,L}^n \underset{\sim}{dq} \right) \underset{\sim}{\nabla}_x \cdot \underset{\sim}{u}_{\kappa,\alpha,L}^n \underset{\sim}{dx} \\
(4.10) \quad &+ \int_{\Omega \times D} M \left[1 - \frac{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^n)}{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^n + \varsigma)} \right] \underset{\sim}{u}_{\kappa,\alpha,L}^n \cdot \underset{\sim}{\nabla}_x \widehat{\psi}_{\kappa,\alpha,L}^n \underset{\sim}{dq} \underset{\sim}{dx}.
\end{aligned}$$

Thanks to (2.5), we have that

$$(4.11a) \quad T_3 \geq \frac{a_0}{4\lambda} \int_{\Omega \times D} M [(\mathcal{F}^L)''(\widehat{\psi}_{\kappa,\alpha,L}^n + \varsigma)] |\underset{\sim}{\nabla}_q \widehat{\psi}_{\kappa,\alpha,L}^n|^2 \underset{\sim}{dq} \underset{\sim}{dx},$$

$$(4.11b) \quad T_4 \geq \varepsilon \int_{\Omega \times D} M [(\mathcal{F}^L)''(\widehat{\psi}_{\kappa,\alpha,L}^n + \varsigma)] |\underset{\sim}{\nabla}_x \widehat{\psi}_{\kappa,\alpha,L}^n|^2 \underset{\sim}{dq} \underset{\sim}{dx}.$$

It is tempting to bound $(\mathcal{F}^L)''(\widehat{\psi}_{\kappa,\alpha,L}^n + \varsigma)$ below further by $(\widehat{\psi}_{\kappa,\alpha,L}^n + \varsigma)^{-1}$ using (3.40b). We have refrained from doing so as the precise form of (4.11b) will be required to absorb the extraneous term that the process of shifting $\widehat{\psi}_{\kappa,\alpha,L}^n$ by the addition of $\varsigma > 0$ generates in the last term in (4.10). Similarly, (4.11a) is required for the last line in (4.12) below. Finally, it follows from (3.40a) and (1.6a) that

$$\begin{aligned}
T_5 &= - \sum_{i=1}^K \int_{\Omega} \left[\int_D M [(\underset{\sim}{\nabla}_x \underset{\sim}{u}_{\kappa,\alpha,L}^n) \underset{\sim}{q}_i] \cdot \underset{\sim}{\nabla}_{q_i} \widehat{\psi}_{\kappa,\alpha,L}^n \underset{\sim}{dq} \right] \underset{\sim}{dx} \\
&+ \int_{\Omega \times D} M \left[1 - \frac{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^n)}{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^n + \varsigma)} \right] \sum_{i=1}^K [(\underset{\sim}{\nabla}_x \underset{\sim}{u}_{\kappa,\alpha,L}^n) \underset{\sim}{q}_i] \cdot \underset{\sim}{\nabla}_{q_i} \widehat{\psi}_{\kappa,\alpha,L}^n \underset{\sim}{dq} \underset{\sim}{dx} \\
&= - \int_{\Omega \times D} M \sum_{i=1}^K U'(\tfrac{1}{2}|\underset{\sim}{q}_i|^2) \widehat{\psi}_{\kappa,\alpha,L}^n (\underset{\sim}{q}_i \underset{\sim}{q}_i^T) : \underset{\sim}{\nabla}_x \underset{\sim}{u}_{\kappa,\alpha,L}^n \underset{\sim}{dq} \underset{\sim}{dx} \\
&+ K \int_{\Omega} \left(\int_D M \widehat{\psi}_{\kappa,\alpha,L}^n \underset{\sim}{dq} \right) \underset{\sim}{\nabla}_x \cdot \underset{\sim}{u}_{\kappa,\alpha,L}^n \underset{\sim}{dx} \\
(4.12) \quad &+ \int_{\Omega \times D} M \left[1 - \frac{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^n)}{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^n + \varsigma)} \right] \sum_{i=1}^K [(\underset{\sim}{\nabla}_x \underset{\sim}{u}_{\kappa,\alpha,L}^n) \underset{\sim}{q}_i] \cdot \underset{\sim}{\nabla}_{q_i} \widehat{\psi}_{\kappa,\alpha,L}^n \underset{\sim}{dq} \underset{\sim}{dx}.
\end{aligned}$$

Substituting (4.9)–(4.12) into (4.8), multiplying by Δt , summing over n and adopting the notation (4.4a,b) yields, for $n = 1, \dots, N$, that

$$\begin{aligned}
&\int_{\Omega \times D} M \mathcal{F}^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}(t_n) + \varsigma) \underset{\sim}{dq} \underset{\sim}{dx} + \frac{1}{2\Delta t L} \int_0^{t_n} \int_{\Omega \times D} M (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,-})^2 \underset{\sim}{dq} \underset{\sim}{dx} \underset{\sim}{dt} \\
&+ \int_0^{t_n} \int_{\Omega \times D} M [(\mathcal{F}^L)''(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)] \left[\frac{a_0}{4\lambda} |\underset{\sim}{\nabla}_q \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}|^2 + \varepsilon |\underset{\sim}{\nabla}_x \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}|^2 \right] \underset{\sim}{dq} \underset{\sim}{dx} \underset{\sim}{dt}
\end{aligned}$$

$$\begin{aligned}
(4.13) \quad & \leq \int_{\Omega \times D} M \mathcal{F}^L(\beta^L(\widehat{\psi}^0) + \varsigma) \, \mathrm{d}q \, \mathrm{d}x \\
& - (K+1) \int_0^{t_n} \int_{\Omega} \left(\int_D M \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} \, \mathrm{d}q \right) \nabla_x \cdot \widehat{u}_{\kappa,\alpha,L}^{\Delta t,+} \, \mathrm{d}x \, \mathrm{d}t \\
& + \int_0^{t_n} \int_{\Omega \times D} M \sum_{i=1}^K U'_i(\tfrac{1}{2}|q|^2) \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}(q_i q_i^T) : \nabla_x \widehat{u}_{\kappa,\alpha,L}^{\Delta t,+} \, \mathrm{d}q \, \mathrm{d}x \, \mathrm{d}t \\
& - \int_0^{t_n} \int_{\Omega \times D} M \left[1 - \frac{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+})}{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)} \right] \widehat{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \nabla_x \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} \, \mathrm{d}q \, \mathrm{d}x \, \mathrm{d}t \\
& - \int_0^{t_n} \int_{\Omega \times D} M \left[1 - \frac{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+})}{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)} \right] \sum_{i=1}^K \left[(\nabla_x \widehat{u}_{\kappa,\alpha,L}^{\Delta t,+}) q_i \right] \cdot \nabla_{q_i} \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} \, \mathrm{d}q \, \mathrm{d}x \, \mathrm{d}t,
\end{aligned}$$

where we have noted (3.23b). The denominator in the prefactor of the second integral on the left-hand side motivates us to link Δt to L so that $\Delta t L = o(1)$, as $\Delta t \rightarrow 0_+$ (or, equivalently, $\Delta t = o(L^{-1})$, as $L \rightarrow \infty$), in order to drive the integral multiplied by the prefactor to 0 in the limit of $\Delta t \rightarrow 0_+$, once the product of the two has been bounded above by a constant, independent of Δt and L .

Comparing (4.13) with (4.7), and noting (1.12), we see that after multiplying (4.13) by k and adding the resulting inequality to (4.7) the last two terms on the right-hand side of (4.7) are cancelled by k times the second and third terms on the right-hand side of (4.13). Hence, for $n = 1, \dots, N$, we deduce that

$$\begin{aligned}
(4.14) \quad & \frac{1}{2} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,+}(t_n) |\widehat{u}_{\kappa,\alpha,L}^{\Delta t,+}(t_n)|^2 \, \mathrm{d}x + \frac{1}{2\Delta t} \int_0^{t_n} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,-} |\widehat{u}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{u}_{\kappa,\alpha,L}^{\Delta t,-}|^2 \, \mathrm{d}x \, \mathrm{d}t \\
& + \int_{\Omega} P_{\kappa}(\rho_{\kappa,\alpha,L}^{\Delta t,+}(t_n)) \, \mathrm{d}x + k \int_{\Omega \times D} M \mathcal{F}^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}(t_n) + \varsigma) \, \mathrm{d}q \, \mathrm{d}x \\
& + \alpha \kappa \int_0^{t_n} \left[\|\nabla_x[(\rho_{\kappa,\alpha,L}^{\Delta t,+})^2]\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x[(\rho_{\kappa,\alpha,L}^{\Delta t,+})^{\frac{\Gamma}{2}}]\|_{L^2(\Omega)}^2 \right] \, \mathrm{d}t \\
& + \mu^S c_0 \int_0^{t_n} \|\widehat{u}_{\kappa,\alpha,L}^{\Delta t,+}\|_{H^1(\Omega)}^2 \, \mathrm{d}t + \frac{k}{2\Delta t L} \int_0^{t_n} \int_{\Omega \times D} M (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,-})^2 \, \mathrm{d}q \, \mathrm{d}x \, \mathrm{d}t \\
& + k \int_0^{t_n} \int_{\Omega \times D} M [(\mathcal{F}^L)''(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)] \left[\frac{a_0}{4\lambda} |\nabla_q \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}|^2 + \varepsilon |\nabla_x \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}|^2 \right] \, \mathrm{d}q \, \mathrm{d}x \, \mathrm{d}t \\
& + \mathfrak{z} \|\varrho_{\kappa,\alpha,L}^{\Delta t,+}(t_n)\|_{L^2(\Omega)}^2 + \mathfrak{z} \int_0^{t_n} \left[\|\varrho_{\kappa,\alpha,L}^{\Delta t,+} - \varrho_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^2(\Omega)}^2 + 2\varepsilon \|\nabla_x \varrho_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(\Omega)}^2 \right] \, \mathrm{d}t \\
& \leq \frac{1}{2} \int_{\Omega} \rho^0 |u^0|^2 \, \mathrm{d}x + \int_{\Omega} P_{\kappa}(\rho^0) \, \mathrm{d}x + \mathfrak{z} \|\varrho^0\|_{L^2(\Omega)}^2 \\
& + \int_0^{t_n} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,+} f \cdot \widehat{u}_{\kappa,\alpha,L}^{\Delta t,+} \, \mathrm{d}x \, \mathrm{d}t + k \int_{\Omega \times D} M \mathcal{F}^L(\beta^L(\widehat{\psi}^0) + \varsigma) \, \mathrm{d}q \, \mathrm{d}x \\
& - k \int_0^{t_n} \int_{\Omega \times D} M \left[1 - \frac{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+})}{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)} \right] \widehat{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \nabla_x \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} \, \mathrm{d}q \, \mathrm{d}x \, \mathrm{d}t \\
& - k \int_0^{t_n} \int_{\Omega \times D} M \left[1 - \frac{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+})}{\beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)} \right] \sum_{i=1}^K \left[(\nabla_x \widehat{u}_{\kappa,\alpha,L}^{\Delta t,+}) q_i \right] \cdot \nabla_{q_i} \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} \, \mathrm{d}q \, \mathrm{d}x \, \mathrm{d}t.
\end{aligned}$$

Similarly to (3.71), we have on choosing $\eta = 1$ in (3.29a) that

$$(4.15) \quad \int_{\Omega} \rho_{\kappa,\alpha,L}^n \, \mathrm{d}x = \int_{\Omega} \rho_{\kappa,\alpha,L}^{n-1} \, \mathrm{d}x = \int_{\Omega} \rho^0 \, \mathrm{d}x, \quad n = 1, \dots, N.$$

Noting (4.15) and (4.4b), we have that

$$(4.16) \quad \left| \int_0^{t_n} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,+} f \cdot u_{\kappa,\alpha,L}^{\Delta t,+} dx dt \right| \leq \frac{1}{2} \left[\int_0^{t_n} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,+} |u_{\kappa,\alpha,L}^{\Delta t,+}|^2 dx dt + \int_0^{t_n} \|f\|_{L^\infty(\Omega)}^2 dt \int_{\Omega} \rho^0 dx \right].$$

Next we recall from [11, (4.25)] the bound

$$(4.17) \quad \int_{\Omega \times D} M \mathcal{F}^L(\beta^L(\widehat{\psi}^0) + \varsigma) dq dx \leq \frac{3\varsigma}{2} |\Omega| + \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}^0 + \varsigma) dq dx.$$

Let $\underline{b} := (b_1, \dots, b_K)$, recall (3.4), and $b := |\underline{b}|_1 := b_1 + \dots + b_K$; then we can bound the magnitude of the last term on the right-hand side of (4.14), on noting (3.40a) and (1.6b), by

$$(4.18) \quad \begin{aligned} & \frac{k a_0}{8\lambda} \left(\int_0^{t_n} \int_{\Omega \times D} M[[\mathcal{F}^L]''(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)] |\nabla_q \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}|^2 dq dx dt \right) \\ & + \varsigma \frac{2k \lambda b}{a_0} \left(\int_0^{t_n} \int_{\Omega} |\nabla_x u_{\kappa,\alpha,L}^{\Delta t,+}|^2 dx dt \right), \end{aligned}$$

see [13, (4.20)] for the details. Similarly, the second to last term on the right-hand side of (4.14) can be bounded by

$$(4.19) \quad \frac{k \varepsilon}{2} \left(\int_0^{t_n} \int_{\Omega \times D} M[[\mathcal{F}^L]''(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)] |\nabla_x \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}|^2 dq dx dt \right) + \varsigma \frac{k}{2\varepsilon} \left(\int_0^{t_n} \int_{\Omega} |u_{\kappa,\alpha,L}^{\Delta t,+}|^2 dx dt \right).$$

Noting (4.14)–(4.19), and using (3.40b) to bound the expression $[\mathcal{F}^L]''(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)$ from below by $\mathcal{F}''(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma) = (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)^{-1}$ and (3.41) to bound $\mathcal{F}^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)$ by $\mathcal{F}(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma)$ from below yields, for $n = 1, \dots, N$, that

$$(4.20) \quad \begin{aligned} & \frac{1}{2} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,+}(t_n) |u_{\kappa,\alpha,L}^{\Delta t,+}(t_n)|^2 dx + \frac{1}{2\Delta t} \int_0^{t_n} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,-} |u_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,-}|^2 dx dt \\ & + \int_{\Omega} P_{\kappa}(\rho_{\kappa,\alpha,L}^{\Delta t,+}(t_n)) dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}(t_n) + \varsigma) dq dx \\ & + \alpha \kappa \int_0^{t_n} \left[\|\nabla_x[(\rho_{\kappa,\alpha,L}^{\Delta t,+})^2]\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x[(\rho_{\kappa,\alpha,L}^{\Delta t,+})^{\frac{\Gamma}{2}}]\|_{L^2(\Omega)}^2 \right] dt \\ & + \mu^S c_0 \int_0^{t_n} \|u_{\kappa,\alpha,L}^{\Delta t,+}\|_{H^1(\Omega)}^2 dt + \frac{k}{2\Delta t L} \int_0^{t_n} \int_{\Omega \times D} M (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,-})^2 dq dx dt \\ & + \frac{k}{2} \int_0^{t_n} \int_{\Omega \times D} \frac{M}{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} + \varsigma} \left[\frac{a_0}{4\lambda} |\nabla_q \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}|^2 + \varepsilon |\nabla_x \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}|^2 \right] dq dx dt \\ & + \mathfrak{z} \|\varrho_{\kappa,\alpha,L}^{\Delta t,+}(t_n)\|_{L^2(\Omega)}^2 + \mathfrak{z} \int_0^{t_n} \left[\|\varrho_{\kappa,\alpha,L}^{\Delta t,+} - \varrho_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^2(\Omega)}^2 + 2\varepsilon \|\nabla_x \varrho_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(\Omega)}^2 \right] dt \\ & \leq \frac{1}{2} \int_{\Omega} \rho^0 |u^0|^2 dx + \int_{\Omega} P_{\kappa}(\rho^0) dx + \mathfrak{z} \|\varrho^0\|_{L^2(\Omega)}^2 \\ & + \frac{1}{2} \int_0^{t_n} \|f\|_{L^\infty(\Omega)}^2 dt \int_{\Omega} \rho^0 dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}^0 + \varsigma) dq dx + \frac{3k\varsigma}{2} |\Omega| \\ & + 2k\varsigma \max \left\{ \frac{\lambda b}{a_0}, \frac{1}{4\varepsilon} \right\} \int_0^{t_n} \|u_{\kappa,\alpha,L}^{\Delta t,+}\|_{H^1(\Omega)}^2 dt + \frac{1}{2} \int_0^{t_n} \int_{\Omega} \rho_{\kappa,\alpha,L}^{\Delta t,+} |u_{\kappa,\alpha,L}^{\Delta t,+}|^2 dx dt. \end{aligned}$$

Passing to the limit $\varsigma \rightarrow 0_+$ in (4.20), and then applying a discrete Gronwall inequality yields, for $n = 1, \dots, N$, that

$$\begin{aligned}
& \frac{1}{2} \int_{\Omega} \rho_{\kappa, \alpha, L}^{\Delta t, +}(t_n) |u_{\kappa, \alpha, L}^{\Delta t, +}(t_n)|^2 \, dx + \frac{1}{2\Delta t} \int_0^{t_n} \int_{\Omega} \rho_{\kappa, \alpha, L}^{\Delta t, -} |u_{\kappa, \alpha, L}^{\Delta t, +} - u_{\kappa, \alpha, L}^{\Delta t, -}|^2 \, dx \, dt \\
& + \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L}^{\Delta t, +}(t_n)) \, dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}(t_n)) \, dq \, dx \\
& + \alpha \kappa \int_0^{t_n} \left[\|\nabla_x [(\rho_{\kappa, \alpha, L}^{\Delta t})^2]\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x [(\rho_{\kappa, \alpha, L}^{\Delta t})^{\frac{\Gamma}{2}}]\|_{L^2(\Omega)}^2 \right] \, dt \\
& + \mu^S c_0 \int_0^{t_n} \|u_{\kappa, \alpha, L}^{\Delta t, +}\|_{H^1(\Omega)}^2 \, dt + \frac{k}{2\Delta t L} \int_0^{t_n} \int_{\Omega \times D} M (\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +} - \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, -})^2 \, dq \, dx \, dt \\
& + k \int_0^{t_n} \int_{\Omega \times D} M \left[\frac{a_0}{2\lambda} \left| \nabla_q \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \right|^2 + 2\varepsilon \left| \nabla_x \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \right|^2 \right] \, dq \, dx \, dt \\
& + \mathfrak{z} \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}(t_n)\|_{L^2(\Omega)}^2 + \mathfrak{z} \int_0^{t_n} \left[\|\varrho_{\kappa, \alpha, L}^{\Delta t, +} - \varrho_{\kappa, \alpha, L}^{\Delta t, -}\|_{L^2(\Omega)}^2 + 2\varepsilon \|\nabla_x \varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^2(\Omega)}^2 \right] \, dt \\
& \leq e^{t_n} \left[\frac{1}{2} \int_{\Omega} \rho^0 |u^0|^2 \, dx + \int_{\Omega} P_{\kappa}(\rho^0) \, dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}^0) \, dq \, dx \right. \\
& \quad \left. + \mathfrak{z} \|\varrho^0\|_{L^2(\Omega)}^2 + \frac{1}{2} \int_0^{t_n} \|f\|_{L^\infty(\Omega)}^2 \, dt \int_{\Omega} \rho^0 \, dx \right] \\
& \leq e^{t_n} \left[\frac{1}{2} \int_{\Omega} \rho^0 |u_0|^2 \, dx + \int_{\Omega} P_{\kappa}(\rho^0) \, dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_0) \, dq \, dx \right. \\
& \quad \left. + \mathfrak{z} \int_{\Omega} \left(\int_D M \widehat{\psi}_0 \, dq \right)^2 \, dx + \frac{1}{2} \int_0^{t_n} \|f\|_{L^\infty(\Omega)}^2 \, dt \int_{\Omega} \rho^0 \, dx \right] \\
(4.21) \quad & \leq C,
\end{aligned}$$

where C is a positive constant, independent of the parameters Δt , L , α and κ . Here, we have noted (3.20a), (3.23a) and (3.24) for the penultimate inequality in (4.21), and (3.18a) for the final inequality.

Next we bound the extra stress term (1.12). As we do not have a bound on $\|M^{\frac{1}{2}} \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^2(\Omega \times D)}$ in (4.21), we will need a weaker bound than (3.15). First, we deduce from (1.12), (1.6a) and as $M = 0$ on ∂D that

$$(4.22) \quad C_i(M \varphi) \approx - \int_D (\nabla_{q_i} M) q_i^T \varphi \, dq = \int_D M (\nabla_{q_i} \varphi) q_i^T \, dq + \left(\int_D M \varphi \, dq \right) I.$$

Hence, for $r \in [1, 2)$, on noting that $\nabla_{q_i} \varphi = \nabla_{q_i} (\sqrt{\varphi})^2 = 2\sqrt{\varphi} \nabla_{q_i} \sqrt{\varphi}$ for any sufficiently smooth nonnegative function φ , we have that

$$\begin{aligned}
\|C_i(M \varphi)\|_{L^r(\Omega)} & \leq C \left[\int_{\Omega} \left(\int_D M \varphi \, dq \right)^{\frac{r}{2}} \left(\int_D M |\nabla_{q_i} \sqrt{\varphi}|^2 \, dq \right)^{\frac{r}{2}} \, dx + \int_{\Omega} \left(\int_D M \varphi \, dq \right)^r \, dx \right]^{\frac{1}{r}} \\
(4.23) \quad & \leq C \left[\|\nabla_{q_i} \sqrt{\varphi}\|_{L_M^2(\Omega \times D)} \left\| \int_D M \varphi \, dq \right\|_{L^{\frac{r}{2-r}}(\Omega)}^{\frac{1}{2}} + \left\| \int_D M \varphi \, dq \right\|_{L^r(\Omega)} \right].
\end{aligned}$$

Therefore, for $r \in [1, 2)$ and $s \in [1, 2]$, it follows that, for any such function φ ,

$$\begin{aligned}
& \|C_i(M \varphi)\|_{L^s(0, T; L^r(\Omega))} \\
& \leq C \left[\|\nabla_{q_i} \sqrt{\varphi}\|_{L^2(0, T; L_M^2(\Omega \times D))} \left\| \int_D M \varphi \, dq \right\|_{L^{\frac{r}{2-r}}(0, T; L^{\frac{r}{2-r}}(\Omega))}^{\frac{1}{2}} + \left\| \int_D M \varphi \, dq \right\|_{L^s(0, T; L^r(\Omega))} \right], \\
(4.24) \quad &
\end{aligned}$$

where $v = \frac{s}{2-s}$ if $s \in [1, 2)$ and $v = \infty$ if $s = 2$. We deduce from (4.24) and (4.21) that, for $i = 1, \dots, K$,

$$(4.25) \quad \|C_i(M \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +})\|_{L^s(0, T; L^r(\Omega))} \leq C \quad \text{if} \quad \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^v(0, T; L^{\frac{r}{2-r}}(\Omega))} \leq C,$$

where $r \in [1, 2)$, $s \in [1, 2]$ and $v = \frac{s}{2-s}$ if $s \in [1, 2)$ and $v = \infty$ if $s = 2$.

It follows from (4.21) and (3.2) that

$$(4.26) \quad \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^{\frac{2}{\vartheta}}(0, T; L^v(\Omega))} \leq C \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^\infty(0, T; L^2(\Omega))}^{1-\vartheta} \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^2(0, T; H^1(\Omega))}^{\vartheta} \leq C,$$

where $\vartheta = \frac{(v-2)d}{2v}$, and $v \in (2, \infty)$ if $d = 2$ and $v \in (2, 6]$ if $d = 3$. For example, we have that

$$(4.27) \quad \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^\infty(0, T; L^2(\Omega))} + \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^{\frac{2(d+2)}{d}}(\Omega_T)} + \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^4(0, T; L^{\frac{2d}{d-1}}(\Omega))} \leq C,$$

and hence we deduce from (4.25) and (1.11) that

$$(4.28a) \quad \|C_i(M \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +})\|_{L^2(0, T; L^{\frac{4}{3}}(\Omega))} + \|C_i(M \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +})\|_{L^{\frac{4(d+2)}{3d+4}}(\Omega_T)} \leq C, \quad i = 1, \dots, K,$$

$$(4.28b) \quad \|\tau_1(M \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +})\|_{L^2(0, T; L^{\frac{4}{3}}(\Omega))} + \|\tau_1(M \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +})\|_{L^{\frac{4(d+2)}{3d+4}}(\Omega_T)} \leq C,$$

where C is independent of Δt , L , α and κ .

Remark 4.1. We note from (4.21) and (4.26) that if $\mathfrak{z} = 0$, then the bounds (4.27) and (4.28a,b) no longer hold. In this case, we have only the following weaker bounds.

Similarly to (4.15), we have on choosing $\varphi = 1$ in (3.29c), and noting (4.4b) and (3.23b), that, for a.a. $t \in (0, T)$,

$$(4.29) \quad \int_{\Omega} \varrho_{\kappa, \alpha, L}^{\Delta t, +} dx \underset{\sim}{=} \int_{\Omega \times D} M \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +} dq dx \underset{\sim}{=} \int_{\Omega \times D} M \widehat{\psi}^0 dq dx \underset{\sim}{=} \leq C.$$

Next we deduce from (4.21) and (4.29) that

$$(4.30) \quad \begin{aligned} \|\nabla_x \varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^2(0, T; L^1(\Omega))}^2 &= 4 \left\| \int_D M \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \nabla_x \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} dq \right\|_{L^2(0, T; L^1(\Omega))}^2 \\ &\leq 4 \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^\infty(0, T; L^1(\Omega))} \left\| \nabla_x \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \right\|_{L^2(0, T; L_M^2(\Omega \times D))}^2 \leq C. \end{aligned}$$

It follows from Sobolev embedding, (4.30) and (4.29) that

$$(4.31) \quad \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^2(0, T; L^{\frac{d}{d-1}}(\Omega))} \leq C \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^2(0, T; W^{1,1}(\Omega))} \leq C.$$

Therefore, we obtain from (4.25), (4.31) and (1.11) that

$$(4.32) \quad \|\tau_1(M \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +})\|_{L^{\frac{4}{3}}(0, T; L^{\frac{2d}{2d-1}}(\Omega))} \leq C,$$

where C is independent of Δt , L , α and κ .

4.2. L , Δt -independent bounds on the spatial and temporal derivatives of $\rho_{\kappa, \alpha, L}^{[\Delta t]}$. In addition to the bounds on $\rho_{\kappa, \alpha, L}^{[\Delta t]}$ and $\rho_{\kappa, \alpha, L}^{\Delta t, +}$ in (4.21), we establish further relevant bounds here. Similarly to (3.87), on choosing, for any $s \in (0, T]$, $\eta(\cdot, t) = \chi_{[0, s]} [\rho_{\kappa, \alpha, L}^{[\Delta t]}(\cdot, t)]^{\vartheta-1}$, for $\vartheta = 2$ and $\frac{\vartheta}{2}$, in (4.5a), we obtain, on noting (4.21), that

$$(4.33) \quad \begin{aligned} &\frac{1}{\vartheta} \|\rho_{\kappa, \alpha, L}^{[\Delta t]}(\cdot, s)\|_{L^\vartheta(\Omega)}^\vartheta + \frac{4\alpha(\vartheta-1)}{\vartheta^2} \int_0^s \|\nabla_x [(\rho_{\kappa, \alpha, L}^{[\Delta t]})^{\frac{\vartheta}{2}}]\|_{L^2(\Omega)}^2 dt \\ &= \frac{1}{\vartheta} \left[\|\rho^0\|_{L^\vartheta(\Omega)}^\vartheta + (\vartheta-1) \int_0^s \int_{\Omega} u_{\kappa, \alpha, L}^{\Delta t, +} \cdot \nabla_x [(\rho_{\kappa, \alpha, L}^{[\Delta t]})^\vartheta] dx dt \right] \\ &\leq \frac{1}{\vartheta} \left[\|\rho^0\|_{L^\vartheta(\Omega)}^\vartheta + \int_0^s \|u_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^2(\Omega)}^2 dt + \frac{(\vartheta-1)^2}{4} \int_0^s \|\nabla_x [(\rho_{\kappa, \alpha, L}^{[\Delta t]})^\vartheta]\|_{L^2(\Omega)}^2 dt \right] \\ &\leq C, \end{aligned}$$

where C is independent of Δt and L . Similarly to (3.88), it follows from a Poincaré inequality, (4.21) and (4.33) for $\vartheta = \frac{\Gamma}{2}$ that

$$(4.34) \quad \begin{aligned} \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^\Gamma(0,T;L^\Gamma(\Omega))}^\Gamma &= \|(\rho_{\kappa,\alpha,L}^{[\Delta t]})^{\frac{\Gamma}{2}}\|_{L^2(0,T;L^2(\Omega))}^2 \\ &\leq C \left[\|\nabla_x [(\rho_{\kappa,\alpha,L}^{[\Delta t]})^{\frac{\Gamma}{2}}]\|_{L^2(0,T;L^2(\Omega))}^2 + \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^\infty(0,T;L^{\frac{\Gamma}{2}}(\Omega))}^\Gamma \right] \leq C. \end{aligned}$$

Similarly to (3.89), we obtain from (3.33), (3.2), (4.21) and (4.33) for $\vartheta = \frac{\Gamma}{2}$ that

$$(4.35) \quad \begin{aligned} \left| \int_0^T c(u_{\kappa,\alpha,L}^{\Delta t,+})(\rho_{\kappa,\alpha,L}^{[\Delta t]}, \eta) dt \right| &\leq C \|\eta\|_{L^2(0,T;H^1(\Omega))} + C \left| \int_0^T \|u_{\kappa,\alpha,L}^{\Delta t,+}\|_{H^1(\Omega)} \|\eta\|_{H^1(\Omega)} dt \right| \\ &\leq C \left[1 + \|u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(0,T;H^1(\Omega))} \right] \|\eta\|_{L^2(0,T;H^1(\Omega))} \\ &\leq C \|\eta\|_{L^2(0,T;H^1(\Omega))} \quad \forall \eta \in L^2(0,T;H^1(\Omega)). \end{aligned}$$

Hence, similarly to (3.90), we deduce from (4.33) for $\vartheta = 2$ and $\frac{\Gamma}{2}$, (4.35), on noting (4.5a) and (3.33), (4.21) and (4.34) that

$$(4.36) \quad \begin{aligned} \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^\infty(0,T;L^{\frac{\Gamma}{2}}(\Omega))}^\Gamma &+ \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^2(0,T;H^1(\Omega))}^2 \\ &+ \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{H^1(0,T;H^1(\Omega)')}^2 + \|(\rho_{\kappa,\alpha,L}^{[\Delta t]})^{\frac{\Gamma}{2}}\|_{L^2(0,T;H^1(\Omega))}^2 \leq C, \end{aligned}$$

where C is independent of Δt and L . Similarly to (3.91), we deduce from (3.2) and the last bound in (4.36) that

$$(4.37) \quad \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^\Gamma(0,T;L^{3\Gamma}(\Omega))}^\Gamma = \|(\rho_{\kappa,\alpha,L}^{[\Delta t]})^{\frac{\Gamma}{2}}\|_{L^2(0,T;L^6(\Omega))}^2 \leq C \|(\rho_{\kappa,\alpha,L}^{[\Delta t]})^{\frac{\Gamma}{2}}\|_{L^2(0,T;H^1(\Omega))}^2 \leq C.$$

Next, similarly to (3.92), it follows from (3.1), the first bound in (4.36) and (4.37) that

$$(4.38) \quad \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^{\frac{4\Gamma}{3}}(0,T;L^{\frac{4\Gamma}{3}}(\Omega))}^{\frac{4\Gamma}{3}} \leq \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^\infty(0,T;L^{\frac{\Gamma}{2}}(\Omega))}^{\frac{\Gamma}{2}} \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^\Gamma(0,T;L^{3\Gamma}(\Omega))}^\Gamma \leq C,$$

where C is independent of Δt and L . Finally, it follows from (4.2), (2.3) and (4.38) that

$$(4.39) \quad \|p_{\kappa,\alpha,L}^{\{\Delta t\}}\|_{L^{\frac{4}{3}}(\Omega_T)}^{\frac{4}{3}} \leq \|p_\kappa(\rho_{\kappa,\alpha,L}^{[\Delta t]})\|_{L^{\frac{4}{3}}(\Omega_T)}^{\frac{4}{3}} \leq C \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^{\frac{4\Gamma}{3}}(\Omega_T)}^{\frac{4\Gamma}{3}} \leq C,$$

where C is independent of Δt and L .

4.3. Passing to the limit $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$) in the continuity equation (4.5a). As noted after (4.13), we shall assume that

$$(4.40) \quad \Delta t = o(L^{-1}) \quad \text{as } \Delta t \rightarrow 0_+ \text{ } (L \rightarrow \infty).$$

Requiring, for example, that $0 < \Delta t \leq C_0/(L \log L)$, $L > 1$, with an arbitrary (but fixed) constant C_0 will suffice to ensure that (4.40) holds. We have the following convergence results.

Lemma 4.1. *There exists a subsequence (not indicated) of $\{(\rho_{\kappa,\alpha,L}^{[\Delta t]}, u_{\kappa,\alpha,L}^{\Delta t,+}, \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+})\}_{\Delta t>0}$, and functions*

$$(4.41) \quad \rho_{\kappa,\alpha} \in L^2(0,T;H^1(\Omega)) \cap H^1(0,T;H^1(\Omega)') \cap C_w([0,T];L^{\frac{\Gamma}{2}}(\Omega)) \cap L_{\geq 0}^{\frac{4\Gamma}{3}}(\Omega_T)$$

with $\rho_{\kappa,\alpha}(\cdot, 0) = \rho^0(\cdot)$ and $u_{\kappa,\alpha} \in L^2(0,T;H_0^1(\Omega))$ such that, as $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$),

$$(4.42a) \quad \rho_{\kappa,\alpha,L}^{[\Delta t]} \rightarrow \rho_{\kappa,\alpha} \quad \text{weakly in } L^2(0,T;H^1(\Omega)), \quad \text{weakly in } H^1(0,T;H^1(\Omega)'),$$

$$(4.42b) \quad \rho_{\kappa,\alpha,L}^{[\Delta t]} \rightarrow \rho_{\kappa,\alpha} \quad \text{in } C_w([0,T];L^{\frac{\Gamma}{2}}(\Omega)), \quad \text{weakly in } L^{\frac{4\Gamma}{3}}(\Omega_T),$$

$$(4.42c) \quad \left(\rho_{\kappa,\alpha,L}^{[\Delta t]} \right)^{\frac{\Gamma}{2}} \rightarrow \rho_{\kappa,\alpha}^{\frac{\Gamma}{2}} \quad \text{weakly in } L^2(0,T;H^1(\Omega)),$$

$$(4.42d) \quad \rho_{\kappa,\alpha,L}^{[\Delta t]} \rightarrow \rho_{\kappa,\alpha} \quad \text{strongly in } L^2(0,T;L^r(\Omega)), \quad \text{strongly in } L^v(\Omega_T),$$

where $r \in [1, \infty)$ if $d = 2$ and $r \in [1, 6)$ if $d = 3$, and $v \in [1, \frac{4\Gamma}{3})$;

$$(4.43a) \quad p_\kappa(\rho_{\kappa,\alpha,L}^{[\Delta t]}) \rightarrow p_\kappa(\rho_{\kappa,\alpha}) \quad \text{strongly in } L^s(\Omega_T),$$

$$(4.43b) \quad p_{\kappa,\alpha,L}^{\{\Delta t\}} \rightarrow p_\kappa(\rho_{\kappa,\alpha}) \quad \text{weakly in } L^s(\Omega_T),$$

where $s \in (1, \frac{4}{3})$; and

$$(4.44) \quad u_{\kappa,\alpha,L}^{\Delta t,+} \rightharpoonup u_{\kappa,\alpha} \quad \text{weakly in } L^2(0, T; H_0^1(\Omega)).$$

Moreover, we have that

$$(4.45) \quad \int_0^T \left\langle \frac{\partial \rho_{\kappa,\alpha}}{\partial t}, \eta \right\rangle_{H^1(\Omega)} dt + \int_0^T \int_\Omega \left(\alpha \nabla_x \rho_{\kappa,\alpha} - \rho_{\kappa,\alpha} u_{\kappa,\alpha} \right) \cdot \nabla_x \eta \, dx \, dt = 0$$

$$\forall \eta \in L^2(0, T; H^1(\Omega)).$$

Proof. The convergence results (4.42a,b) follow immediately from (4.36), (3.14a,b) and (4.38). The strong convergence results (4.42d) follow from (4.42a), (3.11), (4.38) and the interpolation result (3.1). The weak convergence result (4.42c) is then a consequence of (4.36) and (4.42d). Therefore, we have the desired result (4.41). As $L^2(0, T; H^1(\Omega)) \cap H^1(0, T; H^1(\Omega)') \hookrightarrow C([0, T]; L^2(\Omega))$, we obtain that $\rho_{\kappa,\alpha}(\cdot, 0) = \rho^0(\cdot)$.

Next, it follows from (2.3), for $s \in [1, \frac{4}{3}]$, that

$$(4.46) \quad \|p_\kappa(\rho_{\kappa,\alpha}) - p_\kappa(\rho_{\kappa,\alpha,L}^{[\Delta t]})\|_{L^s(\Omega_T)} \leq C \left[\|\rho_{\kappa,\alpha}\|_{L^{s\Gamma}(\Omega_T)}^{\Gamma-1} + \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^{s\Gamma}(\Omega_T)}^{\Gamma-1} \right] \|\rho_{\kappa,\alpha} - \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^{s\Gamma}(\Omega_T)}.$$

Hence, the result (4.43a) follows from (4.46) and (4.42d).

It follows from (4.2) and (4.43a) that, for $s \in (1, \frac{4}{3})$,

$$(4.47) \quad \int_{\Omega_T} p_{\kappa,\alpha,L}^{\{\Delta t\}} \eta \, dx \, dt = \int_{\Omega_T} p_\kappa(\rho_{\kappa,\alpha,L}^{[\Delta t]}) \eta^{\{\Delta t\}} \, dx \, dt \quad \forall \eta \in L^{\frac{s}{s-1}}(\Omega_T),$$

where

$$(4.48) \quad \eta^{\{\Delta t\}}(\cdot, t) := \frac{1}{\Delta t} \int_{t_{n-1}}^{t_n} \eta(\cdot, t') \, dt', \quad t \in (t_{n-1}, t_n], \quad n = 1, \dots, N.$$

We note that

$$(4.49) \quad \lim_{\Delta t \rightarrow 0_+} \|\eta - \eta^{\{\Delta t\}}\|_{L^r(\Omega_T)} = 0 \quad \forall \eta \in L^r(\Omega_T), \quad r \in [1, \infty).$$

Therefore, the desired result (4.43b) follows from (4.47), (4.43a) and (4.49). Finally, the weak convergence result (4.44) follows immediately from (4.21).

It is now a simple matter to pass to the limit $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$) for the subsequence in (4.5a), on noting (4.42a-c) and (4.44), to obtain (4.45). \square

In order to pass to the limit $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$) in the momentum equation (4.5b), we will need a strong convergence result for $\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}$. First, it follows from (4.4b), (3.20a), and (4.21) that

$$(4.50) \quad \begin{aligned} \|\nabla_x u_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^2(0,T;L^2(\Omega))}^2 &= \Delta t \|\nabla_x u_0\|^2 + \int_{\Delta t}^T \|\nabla_x u_{\kappa,\alpha,L}^{\Delta t,-}\|^2 \, dt \\ &\leq \int_\Omega \rho^0 |u_0|^2 \, dx + \int_0^{T-\Delta t} \|\nabla_x u_{\kappa,\alpha,L}^{\Delta t,+}\|^2 \, dt \leq C; \end{aligned}$$

hence, we obtain from (3.2), a Poincaré inequality, (4.21), (4.50) and (4.4a,b) that

$$(4.51) \quad \|u_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^2(0,T;L^6(\Omega))} \leq \|u_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^2(0,T;H^1(\Omega))} \leq \|\nabla_x u_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^2(0,T;L^2(\Omega))} \leq C,$$

where C is independent of Δt , L , α and κ .

Lemma 4.2. *There exists a $C \in \mathbb{R}_{>0}$, independent of Δt and L , such that*

$$(4.52a) \quad \begin{aligned} & \|\rho_{\kappa,\alpha,L}^{\Delta t, \pm}\|_{L^\infty(0,T;L^\Gamma(\Omega))} + \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t, \pm}} u_{\kappa,\alpha,L}^{\Delta t, \pm} \right\|_{L^\infty(0,T;L^2(\Omega))} \\ & + \|\rho_{\kappa,\alpha,L}^{\Delta t, \pm} u_{\kappa,\alpha,L}^{\Delta t, \pm}\|_{L^\infty(0,T;L^{\frac{2\Gamma}{\Gamma+1}}(\Omega))} + \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t, \pm}} u_{\kappa,\alpha,L}^{\Delta t, \pm} \right\|_{L^2(0,T;L^{\frac{6\Gamma}{\Gamma+3}}(\Omega))} \leq C, \end{aligned}$$

$$(4.52b) \quad \begin{aligned} & \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t, +}\|_{L^\infty(0,T;L^1(\Omega))} + \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t, +}\|_{L^2(0,T;L^{\frac{6\Gamma}{\Gamma+12}}(\Omega))} \\ & + \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t, +} \right\|_{L^2(0,T;L^{\frac{6\Gamma}{\Gamma+6}}(\Omega))} + \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t, +}\|_{L^v(\Omega_T)} \leq C, \end{aligned}$$

$$(4.52c) \quad \|\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^v(\Omega_T)} \leq C,$$

where $v = \frac{8\Gamma-12}{3\Gamma} \geq \frac{13}{6}$ as $\Gamma \geq 8$.

Hence, in addition to (4.41), $\rho_{\kappa,\alpha} \in L^\infty(0,T;L^\Gamma(\Omega))$ and for a further subsequence of the subsequence of Lemma 4.1, it follows that, as $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$),

$$(4.53a) \quad \nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]} \rightharpoonup \nabla_x \rho_{\kappa,\alpha} \quad \text{weakly in } L^v(\Omega_T), \quad \text{strongly in } L^2(\Omega_T),$$

$$(4.53b) \quad \rho_{\kappa,\alpha,L}^{\Delta t, +} \rightarrow \rho_{\kappa,\alpha} \quad \text{weakly-}\star \text{ in } L^\infty(0,T;L^\Gamma(\Omega)), \quad \text{strongly in } L^2(0,T;H^1(\Omega)'),$$

and, for any nonnegative $\eta \in C[0,T]$,

$$(4.53c) \quad \int_0^T \left(\int_\Omega P_\kappa(\rho_{\kappa,\alpha}) \, dx \right) \eta \, dt \leq \liminf_{\Delta t \rightarrow 0_+ (L \rightarrow \infty)} \int_0^T \left(\int_\Omega P_\kappa(\rho_{\kappa,\alpha,L}^{\Delta t, +}) \, dx \right) \eta \, dt.$$

Proof. The first two bounds in (4.52a) follow immediately from (4.4a,b), (4.21), (3.5), (3.18a) and (3.20a). The third bound in (4.52a) follows immediately from the first two on noting that

$$(4.54) \quad \|\rho_{\kappa,\alpha,L}^{\Delta t, \pm} u_{\kappa,\alpha,L}^{\Delta t, \pm}\|_{L^{\frac{2\Gamma}{\Gamma+1}}(\Omega)} \leq \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t, \pm}} \right\|_{L^{2\Gamma}(\Omega)} \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t, \pm}} u_{\kappa,\alpha,L}^{\Delta t, \pm} \right\|_{L^2(\Omega)}.$$

It follows from the first bound in (4.52a), (3.2) and (4.51) that

$$(4.55) \quad \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t, \pm}} u_{\kappa,\alpha,L}^{\Delta t, \pm} \right\|_{L^2(0,T;L^{\frac{6\Gamma}{\Gamma+3}}(\Omega))} \leq \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t, \pm}} \right\|_{L^\infty(0,T;L^{2\Gamma}(\Omega))} \|u_{\kappa,\alpha,L}^{\Delta t, \pm}\|_{L^2(0,T;L^6(\Omega))} \leq C,$$

and hence the fourth bound in (4.52a).

Next we note, for any $\eta \in L^2(0,T;H^1(\Omega))$ and for a.a. $s \in (t_{n-1}, t_n]$, that

$$(4.56) \quad \begin{aligned} & \left| \int_\Omega \left[\rho_{\kappa,\alpha,L}^{[\Delta t]}(x, s) - \rho_{\kappa,\alpha,L}^{\Delta t, +}(x, s) \right] \eta(x, s) \, dx \right| \\ & = \left| \int_s^{t_n} \int_\Omega \eta(x, s) \frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t]}}{\partial t}(x, t) \, dx \, dt \right| \leq \|\eta(\cdot, s)\|_{H^1(\Omega)} \int_{t_{n-1}}^{t_n} \left\| \frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t]}}{\partial t} \right\|_{H^1(\Omega)'} \, dt. \end{aligned}$$

It follows from (4.56) with $\eta = |u_{\kappa,\alpha,L}^{\Delta t, +}|$, (4.52a), (4.36) and (4.21) that

$$(4.57) \quad \begin{aligned} & \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t, +}\|_{L^\infty(0,T;L^1(\Omega))} \leq \|\rho_{\kappa,\alpha,L}^{\Delta t, +} u_{\kappa,\alpha,L}^{\Delta t, +}\|_{L^\infty(0,T;L^1(\Omega))} \\ & + \frac{1}{2} \left[\left\| \frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t]}}{\partial t} \right\|_{L^2(0,T;H^1(\Omega)')}^2 + \|u_{\kappa,\alpha,L}^{\Delta t, +}\|_{L^2(0,T;H^1(\Omega))}^2 \right] \leq C, \end{aligned}$$

and hence the first desired bound in (4.52b). Similarly to (4.55), it follows from (4.36), (3.2) and (4.51) that

$$(4.58a) \quad \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t, +}\|_{L^2(0,T;L^{\frac{6\Gamma}{\Gamma+12}}(\Omega))} \leq \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^\infty(0,T;L^{\frac{\Gamma}{2}}(\Omega))} \|u_{\kappa,\alpha,L}^{\Delta t, +}\|_{L^2(0,T;L^6(\Omega))},$$

$$(4.58b) \quad \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t, +} \right\|_{L^2(0,T;L^{\frac{6\Gamma}{\Gamma+6}}(\Omega))} \leq \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} \right\|_{L^\infty(0,T;L^\Gamma(\Omega))} \|u_{\kappa,\alpha,L}^{\Delta t, +}\|_{L^2(0,T;L^6(\Omega))},$$

and hence the second and third bounds in (4.52b). It follows from (3.1) with $v = \frac{8\Gamma-12}{3\Gamma}$, $r = 1$ and $s = \frac{6\Gamma}{\Gamma+12}$ that $v\vartheta = 2$ (with $\vartheta \in (0, 1)$ for $\Gamma \geq 8$) and so

$$(4.59) \quad \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^v(\Omega_T)} \leq \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^\infty(0,T;L^1(\Omega))}^{1-\vartheta} \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(0,T;L^s(\Omega))}^\vartheta.$$

Thus, (4.59) and the first two bounds in (4.52b) yield the fourth bound in (4.52b). On noting this bound and recalling from (3.4) that $\partial\Omega \in C^{2,\theta}$, $\theta \in (0, 1)$, and $\rho^0 \in L^\infty(\Omega)$ satisfying (3.18a), we can now apply the parabolic regularity result, Lemma 7.38 in Novotný & Straškraba [58] (or Lemma G.2 in Appendix G), to (4.5a) to obtain that the solution $\rho_{\kappa,\alpha,L}^{[\Delta t]}$ satisfies the bound (4.52c).

The first desired result in (4.53a) follows immediately from (4.52c). Next we obtain from (4.33) for $\vartheta = 2$ that, for any $s \in (0, T]$,

$$(4.60) \quad \begin{aligned} & \frac{1}{2} \|\rho_{\kappa,\alpha,L}^{[\Delta t]}(\cdot, s)\|_{L^2(\Omega)}^2 + \alpha \int_0^s \|\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^2(\Omega)}^2 dt \\ &= \frac{1}{2} \left[\|\rho^0\|_{L^2(\Omega)}^2 - \int_0^s \int_\Omega (\nabla_x \cdot u_{\kappa,\alpha,L}^{\Delta t,+}) (\rho_{\kappa,\alpha,L}^{[\Delta t]})^2 dx dt \right]. \end{aligned}$$

Integrating (4.60) over $s \in (0, T)$, and performing integration by parts, yields that

$$(4.61) \quad \begin{aligned} & \frac{1}{2} \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^2(\Omega_T)}^2 + \alpha \int_0^T (T-t) \|\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^2(\Omega)}^2 dt \\ &= \frac{1}{2} \left[T \|\rho^0\|_{L^2(\Omega)}^2 - \int_0^T (T-t) \int_\Omega (\nabla_x \cdot u_{\kappa,\alpha,L}^{\Delta t,+}) (\rho_{\kappa,\alpha,L}^{[\Delta t]})^2 dx dt \right]. \end{aligned}$$

Similarly, on choosing for any $s \in (0, T]$, $\eta(\cdot, t) = \chi_{[0,s]} \rho_{\kappa,\alpha}(\cdot, t)$ in (4.45), and integrating over $s \in (0, T)$ yields that

$$(4.62) \quad \begin{aligned} & \frac{1}{2} \|\rho_{\kappa,\alpha}\|_{L^2(\Omega_T)}^2 + \alpha \int_0^T (T-t) \|\nabla_x \rho_{\kappa,\alpha}\|_{L^2(\Omega)}^2 dt \\ &= \frac{1}{2} \left[T \|\rho^0\|_{L^2(\Omega)}^2 - \int_0^T (T-t) \int_\Omega (\nabla_x \cdot u_{\kappa,\alpha}) (\rho_{\kappa,\alpha})^2 dx dt \right]. \end{aligned}$$

We deduce from (4.61), (4.62), (4.42d) (with $v = 4$) and (4.44) that

$$(4.63) \quad \lim_{\Delta t \rightarrow 0_+ (L \rightarrow \infty)} \int_0^T (T-t) \|\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^2(\Omega)}^2 dt = \int_0^T (T-t) \|\nabla_x \rho_{\kappa,\alpha}\|_{L^2(\Omega)}^2 dt.$$

By applying the elementary identity $|a - b|^2 = |a|^2 - |b|^2 - 2(a - b) \cdot b$ with $a = \nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}$ and $b = \nabla_x \rho_{\kappa,\alpha,L}$, it follows from (4.63) and (4.42a) that

$$(4.64a) \quad \lim_{\Delta t \rightarrow 0_+ (L \rightarrow \infty)} \int_0^T (T-t) \|\nabla_x (\rho_{\kappa,\alpha,L} - \rho_{\kappa,\alpha,L}^{[\Delta t]})\|_{L^2(\Omega)}^2 dt = 0,$$

and hence, for a.a. $t \in (0, T)$,

$$(4.64b) \quad \|\nabla_x (\rho_{\kappa,\alpha,L} - \rho_{\kappa,\alpha,L}^{[\Delta t]})(\cdot, t)\|_{L^2(\Omega)}^2 \rightarrow 0, \quad \text{as } \Delta t \rightarrow 0_+ (L \rightarrow \infty).$$

Therefore, we obtain the second desired result (4.53a) from (4.64b), (4.52c) and Vitali's convergence theorem. The details of the argument are as follows. With $v \geq \frac{13}{6} > 2$, the bound (4.52c) implies that $|\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}|^2$ is equi-integrable in $L^1(\Omega_T)$, i.e., $\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}$ is 2-equi-integrable. Further, thanks to (4.64a), a subsequence of $\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}$ is a.e. convergent on Ω_T (cf. Theorem 2.20 (iii) in [33]), and thus by Egoroff's theorem (cf. Theorem 2.22 in [33]) it also converges in measure. Hence, by Vitali's convergence theorem (cf. Theorem 2.24 in [33], with $p = 2$), we have strong convergence of the subsequence (not indicated).

The first stated convergence result in (4.53b) follows directly from the first bound in (4.52a). Next, it follows from (4.56) and (4.36) that

$$(4.65) \quad \|\rho_{\kappa,\alpha,L}^{[\Delta t]} - \rho_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(0,T;H^1(\Omega)')}^2 \leq (\Delta t)^2 \left\| \frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t]}}{\partial t} \right\|_{L^2(0,T;H^1(\Omega)')}^2 \leq C (\Delta t)^2.$$

Hence the desired convergence results (4.53b) follow immediately from (4.65) and (4.42d) with $r > \frac{2d}{d+2}$ (to ensure that $L^r(\Omega) \subset H^1(\Omega)'$).

Finally, it follows for any nonnegative $\eta \in C[0, T]$, on noting the convexity of $P_\kappa(\cdot)$, that

$$(4.66) \quad \int_0^T \left(\int_\Omega P_\kappa(\rho_{\kappa,\alpha,L}^{\Delta t,+}) \, dx \right) \eta \, dt \geq \int_0^T \left(\int_\Omega \left[P_\kappa(\rho_{\kappa,\alpha}) + P'_\kappa(\rho_{\kappa,\alpha}) (\rho_{\kappa,\alpha,L}^{\Delta t,+} - \rho_{\kappa,\alpha}) \right] \, dx \right) \eta \, dt.$$

This yields the desired result (4.53c) on noting (4.53b) and that $P'_\kappa(\rho_{\kappa,\alpha}) \in L^1(0, T; L^{\frac{\Gamma}{\Gamma-1}}(\Omega))$ by (4.41). \square

The following result is also required to pass to the limit $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$) in the momentum equation (4.5b).

Lemma 4.3. *There exists a $C \in \mathbb{R}_{>0}$, independent of Δt and L , such that*

$$(4.67) \quad \|\nabla_x \cdot (\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+})\|_{L^s(\Omega_T)} + \left\| \frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t]}}{\partial t} \right\|_{L^s(\Omega_T)} + \|\Delta_x \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^s(\Omega_T)} \leq C,$$

where $s = \frac{8\Gamma-12}{7\Gamma-6} \geq \frac{26}{25}$ as $\Gamma \geq 8$. In addition, we have that

$$(4.68a) \quad \lim_{\Delta t \rightarrow 0_+} \lim_{(L \rightarrow \infty)} \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+} - \rho_{\kappa,\alpha,L}^{\Delta t,-} u_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^2(\Omega_T)} = 0,$$

$$(4.68b) \quad \lim_{\Delta t \rightarrow 0_+} \lim_{(L \rightarrow \infty)} \|\rho_{\kappa,\alpha,L}^{\Delta t,+} u_{\kappa,\alpha,L}^{\Delta t,+} - \rho_{\kappa,\alpha,L}^{\Delta t,-} u_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^2(\Omega_T)} = 0.$$

Proof. It follows from (4.34) and (4.52c) that

$$(4.69) \quad \|\rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^{\frac{8\Gamma-12}{3\Gamma}}(\Omega_T)} + \|\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^{\frac{8\Gamma-12}{3\Gamma}}(\Omega_T)} \leq C,$$

where we have noted that $\frac{8\Gamma-12}{3\Gamma} < \Gamma$, with $\Gamma \geq 8$. Hence, (4.69) and (4.51) yield that

$$(4.70) \quad \|\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^s(\Omega_T)} \|\nabla_x u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^s(\Omega_T)} \leq \|\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^{\frac{8\Gamma-12}{3\Gamma}}(\Omega_T)} \|\nabla_x u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(\Omega_T)} \leq C,$$

where $s = \frac{8\Gamma-12}{7\Gamma-6}$. From the first bound in (4.69) and (4.70), we obtain the first bound in (4.67). As $\partial\Omega \in C^{2,\theta}$, $\theta \in (0, 1)$, it follows from (3.17), (3.18a) and elliptic regularity that, for all $r \in [1, \infty)$,

$$(4.71) \quad \|\rho^0\|_{W^{2,r}(\Omega)} \leq C(\alpha) \|\rho_0 - \rho^0\|_{L^\infty(\Omega)} \leq C(\alpha) \quad \text{and} \quad \nabla_x \rho^0 \cdot \tilde{n} = 0 \text{ on } \partial\Omega.$$

On noting (4.71) and the first bound in (4.67), we can now apply to (4.5a) the parabolic regularity result, Lemma 7.37 in Novotný & Straškraba [58] (or Lemma G.1 in Appendix G), to obtain that the solution $\rho_{\kappa,\alpha,L}^{[\Delta t]}$ satisfies the last two bounds in (4.67).

Next, we note that

$$\begin{aligned}
& \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+} - \rho_{\kappa,\alpha,L}^{\Delta t,-} u_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^2(\Omega_T)}^2 = \int_0^T \int_{\Omega} |\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+} - \rho_{\kappa,\alpha,L}^{\Delta t,-} u_{\kappa,\alpha,L}^{\Delta t,-}|^2 dx dt \\
& = \int_0^T \int_{\Omega} \left| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]} \rho_{\kappa,\alpha,L}^{\Delta t,-}} \left(u_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,-} \right) \right. \\
& \quad \left. + \left(\sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} - \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} \right) \left(\sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t,+} + \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} u_{\kappa,\alpha,L}^{\Delta t,-} \right) \right|^2 dx dt \\
& \leq 2 \int_0^T \int_{\Omega} \rho_{\kappa,\alpha,L}^{[\Delta t]} \rho_{\kappa,\alpha,L}^{\Delta t,-} |u_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,-}|^2 dx dt \\
& \quad + 2 \int_0^T \int_{\Omega} \left| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} - \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} \right|^2 \left| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t,+} + \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} u_{\kappa,\alpha,L}^{\Delta t,-} \right|^2 dx dt \\
& \leq 2 \int_0^T \int_{\Omega} \rho_{\kappa,\alpha,L}^{[\Delta t]} \rho_{\kappa,\alpha,L}^{\Delta t,-} |u_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,-}|^2 dx dt \\
& \quad + 2 \int_0^T \int_{\Omega} \left| \rho_{\kappa,\alpha,L}^{[\Delta t]} - \rho_{\kappa,\alpha,L}^{\Delta t,-} \right| \left| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t,+} + \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} u_{\kappa,\alpha,L}^{\Delta t,-} \right|^2 dx dt.
\end{aligned} \tag{4.72}$$

It follows from (4.36) and (4.21) that

$$\begin{aligned}
& \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} \left(u_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,-} \right) \right\|_{L^2(0,T;L^{\frac{2\Gamma}{\Gamma+2}}(\Omega))} \\
& \leq \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} \right\|_{L^\infty(0,T;L^\Gamma(\Omega))} \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} \left(u_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,-} \right) \right\|_{L^2(\Omega_T)} \leq C (\Delta t)^{\frac{1}{2}}.
\end{aligned} \tag{4.73}$$

Similarly, it follows from (4.51), (4.36) and (4.52a) and that

$$\begin{aligned}
& \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} \left(u_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,-} \right) \right\|_{L^2(0,T;L^{\frac{6\Gamma}{\Gamma+9}}(\Omega))} \\
& \leq \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]} \rho_{\kappa,\alpha,L}^{\Delta t,-}} \right\|_{L^\infty(0,T;L^{\frac{2\Gamma}{3}}(\Omega))} \|u_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^2(0,T;L^6(\Omega))} \\
& \leq C \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} \right\|_{L^\infty(0,T;L^\Gamma(\Omega))} \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} \right\|_{L^\infty(0,T;L^{2\Gamma}(\Omega))} \leq C.
\end{aligned} \tag{4.74}$$

We deduce from (4.73), (4.74) and (3.1) as $\frac{6\Gamma}{\Gamma+9} > 2 > \frac{2\Gamma}{\Gamma+2}$ that

$$\lim_{\Delta t \rightarrow 0_+ (L \rightarrow \infty)} \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,-}} \left(u_{\kappa,\alpha,L}^{\Delta t,+} - u_{\kappa,\alpha,L}^{\Delta t,-} \right) \right\|_{L^2(\Omega_T)} = 0, \tag{4.75}$$

and so the first term on the right-hand side of (4.72) converges to zero, as $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$).

Next, we deal with the second term on the right-hand side of (4.72). It follows from (4.67) that

$$\|\rho_{\kappa,\alpha,L}^{[\Delta t]} - \rho_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^s(\Omega_T)} \leq \Delta t \left\| \frac{\partial \rho_{\kappa,\alpha,L}^{[\Delta t]}}{\partial t} \right\|_{L^s(\Omega_T)} \leq C \Delta t, \tag{4.76}$$

and from (4.36) and (4.52a) that

$$\|\rho_{\kappa,\alpha,L}^{[\Delta t]} - \rho_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^\infty(0,T;L^{\frac{\Gamma}{2}}(\Omega))} \leq C. \tag{4.77}$$

Hence, the bounds (4.76) and (4.77) yield, on noting (3.1) and as $\Gamma \geq 8$, that

$$\lim_{\Delta t \rightarrow 0_+ (L \rightarrow \infty)} \|\rho_{\kappa,\alpha,L}^{[\Delta t]} - \rho_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^v(0,T;L^r(\Omega))} = 0, \quad \text{for any } v \in [1, \infty), r \in [1, 4). \tag{4.78}$$

As $\frac{6\Gamma}{\Gamma+3} \geq \frac{48}{11} > 4$ for $\Gamma \geq 8$, it follows from (3.1) and (4.52a) that

$$(4.79) \quad \begin{aligned} & \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,\pm}} u_{\kappa,\alpha,L}^{\Delta t,\pm} \right\|_{L^3(0,T;L^3(\Omega))} \\ & \leq \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,\pm}} u_{\kappa,\alpha,L}^{\Delta t,\pm} \right\|_{L^\infty(0,T;L^2(\Omega))}^{\frac{1}{3}} \left\| \sqrt{\rho_{\kappa,\alpha,L}^{\Delta t,\pm}} u_{\kappa,\alpha,L}^{\Delta t,\pm} \right\|_{L^2(0,T;L^4(\Omega))}^{\frac{2}{3}} \leq C. \end{aligned}$$

In addition, (4.52b) and (4.51) yield that

$$(4.80) \quad \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t,+} \right\|_{L^4(0,T;L^{\frac{12}{7}}(\Omega))} \leq \left\| \rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+} \right\|_{L^\infty(0,T;L^1(\Omega))}^{\frac{1}{2}} \left\| u_{\kappa,\alpha,L}^{\Delta t,+} \right\|_{L^2(0,T;L^6(\Omega))}^{\frac{1}{2}} \leq C.$$

As $\frac{6\Gamma}{\Gamma+6} \geq \frac{24}{7}$ for $\Gamma \geq 8$, it then follows from (3.1), (4.80) and (4.52b) that

$$(4.81) \quad \begin{aligned} & \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t,+} \right\|_{L^{\frac{28}{13}}(0,T;L^3(\Omega))} \\ & \leq \left(\int_0^T \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t,+} \right\|_{L^{\frac{12}{7}}(\Omega)}^{\frac{4}{13}} \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t,+} \right\|_{L^{\frac{24}{7}}(\Omega)}^{\frac{24}{13}} dt \right)^{\frac{13}{28}} \\ & \leq \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t,+} \right\|_{L^4(0,T;L^{\frac{12}{7}}(\Omega))}^{\frac{1}{7}} \left\| \sqrt{\rho_{\kappa,\alpha,L}^{[\Delta t]}} u_{\kappa,\alpha,L}^{\Delta t,+} \right\|_{L^2(0,T;L^{\frac{24}{7}}(\Omega))}^{\frac{6}{7}} \leq C. \end{aligned}$$

Combining (4.78), (4.79) and (4.81) yields that the second term on the right-hand side of (4.72) converges to zero, as $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$). Therefore, we have the desired result (4.68a).

We now adapt the argument above to prove the desired result (4.68b). The bounds (4.72)–(4.75) remain true with $\rho_{\kappa,\alpha,L}^{[\Delta t]}$ replaced by $\rho_{\kappa,\alpha,L}^{\Delta t,+}$. Similarly, (4.76) remains true with $\rho_{\kappa,\alpha,L}^{[\Delta t]}$ on the left-hand side of the inequality replaced by $\rho_{\kappa,\alpha,L}^{\Delta t,+}$. Hence, the bounds (4.77) and (4.78) remain true with $\rho_{\kappa,\alpha,L}^{[\Delta t]}$ replaced by $\rho_{\kappa,\alpha,L}^{\Delta t,+}$. Therefore, on combining all these modified bounds with (4.79), we obtain the desired result (4.68b). \square

4.4. L , Δt -independent bounds on the time-derivatives of $m_{\kappa,\alpha,L}^{\Delta t}$ and $\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t}$. On noting from (4.4a,b) that

$$(4.82) \quad m_{\kappa,\alpha,L}^{\Delta t} = \frac{t - t_{n-1}}{\Delta t} m_{\kappa,\alpha,L}^{\Delta t,+} + \frac{t_n - t}{\Delta t} m_{\kappa,\alpha,L}^{\Delta t,-}, \quad t \in (t_{n-1}, t_n], \quad n = 1, \dots, N,$$

an elementary calculation yields, for any $s \in [1, \infty]$, that

$$(4.83) \quad \int_0^T \|m_{\kappa,\alpha,L}^{\Delta t}\|_{L^s(\Omega)}^2 dt \leq \frac{1}{2} \int_0^T \left(\|m_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^s(\Omega)}^2 + \|m_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^s(\Omega)}^2 \right) dt.$$

In order to pass to the limit $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$) in the momentum equation (4.5b), we require the following result.

Lemma 4.4. *There exists a $C \in \mathbb{R}_{>0}$, independent of Δt and L , such that*

$$(4.84a) \quad \|m_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^\infty(0,T;L^{\frac{2\Gamma}{\Gamma+1}}(\Omega))} + \|m_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^2(0,T;L^{\frac{6\Gamma}{\Gamma+6}}(\Omega))} + \|m_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^v(\Omega_T)} \leq C,$$

$$(4.84b) \quad \|m_{\kappa,\alpha,L}^{\Delta t(\pm)} \otimes u_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^2(0,T;L^{\frac{6\Gamma}{4\Gamma+3}}(\Omega))} \leq C,$$

where $v = \frac{10\Gamma-6}{3(\Gamma+1)} \geq \frac{74}{27}$ as $\Gamma \geq 8$.

Proof. The first bound in (4.84a) for $m_{\kappa,\alpha,L}^{\Delta t,\pm}$ follows immediately from the third bound in (4.52a). The corresponding bound for $m_{\kappa,\alpha,L}^{\Delta t}$ is a direct consequence of (4.82). Similarly to (4.58a), it follows from (4.52a) and (4.51) that

$$(4.85) \quad \|\rho_{\kappa,\alpha,L}^{\Delta t,\pm} u_{\kappa,\alpha,L}^{\Delta t,\pm}\|_{L^2(0,T;L^{\frac{6\Gamma}{\Gamma+6}}(\Omega))} \leq \|\rho_{\kappa,\alpha,L}^{\Delta t,\pm}\|_{L^\infty(0,T;L^\Gamma(\Omega))} \|u_{\kappa,\alpha,L}^{\Delta t,\pm}\|_{L^2(0,T;L^6(\Omega))} \leq C,$$

and hence the second bound in (4.84a) for $m_{\kappa,\alpha,L}^{\Delta t,\pm}$. The corresponding bound for $\tilde{m}_{\kappa,\alpha,L}^{\Delta t}$ is then a direct consequence of (4.83). Similarly to (4.59), it follows from (3.1) with $v = \frac{10\Gamma-6}{3(\Gamma+1)}$, $r = \frac{2\Gamma}{\Gamma+1}$ and $s = \frac{6\Gamma}{\Gamma+6}$ that $v\vartheta = 2$ (with $\vartheta \in (0,1)$ thanks to $\Gamma \geq 8$), and so

$$(4.86) \quad \|\tilde{m}_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^v(\Omega_T)} \leq \|\tilde{m}_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^\infty(0,T;L^r(\Omega))}^{1-\vartheta} \|\tilde{m}_{\kappa,\alpha,L}^{\Delta t(\pm)}\|_{L^2(0,T;L^s(\Omega))}^\vartheta.$$

Hence (4.86) and the first two bounds in (4.84a) yield the third bound in (4.84a). Finally, combining the first bound in (4.84a) and (4.51) yields the bound (4.84b). \square

Next, we need to bound the time-derivative of $\tilde{m}_{\kappa,\alpha,L}^{\Delta t}$ independently of Δt and L . It follows from (4.39) that we will need to choose at least $w \in L^4(0,T;W_0^{1,4}(\Omega))$ in (4.5b). We now rewrite the time-derivative of $\rho_{\kappa,\alpha,L}^{\Delta t}$ in (4.5b) using (4.5a). Adopting the notation (4.48), we have for any $w \in L^4(0,T;W_0^{1,4}(\Omega))$ that $\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}^{\{\Delta t\}} \in L^2(0,T;H^1(\Omega))$, and so (4.5b) yields that

$$(4.87) \quad \begin{aligned} - \int_0^T \int_\Omega \frac{\partial \rho_{\kappa,\alpha,L}^{\Delta t}}{\partial t} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w} \, dx \, dt &= - \int_0^T \left\langle \frac{\partial \rho_{\kappa,\alpha,L}^{\Delta t}}{\partial t}, \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}^{\{\Delta t\}} \right\rangle_{H^1(\Omega)} dt \\ &= \int_0^T \int_\Omega \left(\alpha \nabla_x \rho_{\kappa,\alpha,L}^{\Delta t} - \rho_{\kappa,\alpha,L}^{\Delta t} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \right) \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}^{\{\Delta t\}}) \, dx \, dt \\ &= \alpha \int_0^T \int_\Omega \nabla_x \rho_{\kappa,\alpha,L}^{\Delta t} \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}^{\{\Delta t\}}) \, dx \, dt \\ &\quad - \int_0^T \int_\Omega \rho_{\kappa,\alpha,L}^{\Delta t} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}) \, dx \, dt \\ &\quad + \int_0^T \int_\Omega \rho_{\kappa,\alpha,L}^{\Delta t} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot (\tilde{w} - \tilde{w}^{\{\Delta t\}})) \, dx \, dt. \end{aligned}$$

Next we note that, for all $w \in L^4(0,T;W_0^{1,4}(\Omega))$,

$$(4.88) \quad \begin{aligned} \int_0^T \int_\Omega \rho_{\kappa,\alpha,L}^{\Delta t} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}) \, dx \, dt \\ = \int_0^T \int_\Omega m_{\kappa,\alpha,L}^{\Delta t,-} \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}) \, dx \, dt \\ + \int_0^T \int_\Omega \left(\rho_{\kappa,\alpha,L}^{\Delta t} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} - \rho_{\kappa,\alpha,L}^{\Delta t,-} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,-} \right) \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}) \, dx \, dt. \end{aligned}$$

Therefore, on combining (4.5b), (4.87) and (4.88), one can rewrite (4.5b) as

$$(4.89) \quad \begin{aligned} \int_0^T \int_\Omega \left[\frac{\partial m_{\kappa,\alpha,L}^{\Delta t}}{\partial t} \cdot \tilde{w} + \frac{\alpha}{2} \nabla_x \rho_{\kappa,\alpha,L}^{\Delta t} \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}^{\{\Delta t\}}) - \left[(m_{\kappa,\alpha,L}^{\Delta t,-} \cdot \nabla_x) \tilde{w} \right] \cdot \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \right] dx \, dt \\ + \int_0^T \int_\Omega S(\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+}) : \nabla_x \tilde{w} \, dx \, dt - \int_0^T \int_\Omega p_{\kappa,\alpha,L}^{\{\Delta t\}} \nabla_x \cdot \tilde{w} \, dx \, dt \\ - \frac{1}{2} \int_0^T \int_\Omega \left(\rho_{\kappa,\alpha,L}^{\Delta t} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} - \rho_{\kappa,\alpha,L}^{\Delta t,-} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,-} \right) \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w}) \, dx \, dt \\ + \frac{1}{2} \int_0^T \int_\Omega \rho_{\kappa,\alpha,L}^{\Delta t} \tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot \nabla_x (\tilde{u}_{\kappa,\alpha,L}^{\Delta t,+} \cdot (\tilde{w} - \tilde{w}^{\{\Delta t\}})) \, dx \, dt \\ = \int_0^T \int_\Omega \left[\rho_{\kappa,\alpha,L}^{\Delta t,+} \tilde{f}^{\{\Delta t\}} \cdot \tilde{w} - \tau_1(M \hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) : \nabla_x \tilde{w} \right] dx \, dt \\ - 2\mathfrak{z} \int_0^T \int_\Omega \left(\int_D M \beta^L(\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \, dq \right) \nabla_x \varrho_{\kappa,\alpha,L}^{\Delta t,+} \cdot \tilde{w} \, dx \, dt \quad \forall w \in L^4(0,T;W_0^{1,4}(\Omega)). \end{aligned}$$

We have the following result.

Lemma 4.5. *There exists a $C \in \mathbb{R}_{>0}$, independent of Δt and L , such that*

$$(4.90) \quad \left\| \frac{\partial m_{\kappa,\alpha,L}^{\Delta t}}{\partial t} \right\|_{L^s(0,T;W_0^{1,4}(\Omega)')} \leq C,$$

where $s = \frac{8\Gamma-12}{7\Gamma-6} \geq \frac{26}{25}$ as $\Gamma \geq 8$.

Proof. Let $v = \frac{8\Gamma-12}{3\Gamma} \geq \frac{13}{6}$ as in Lemma 4.2. Then for $s' = \frac{s}{s-1} = \frac{8\Gamma-12}{\Gamma-6} > 8$, we have that

$$(4.91) \quad \frac{1}{s'} + \frac{1}{v} + \frac{1}{2} = 1.$$

It follows from (4.89), (4.91), $W^{1,4}(\Omega) \hookrightarrow L^\infty(\Omega)$, $H^1(\Omega) \hookrightarrow L^6(\Omega)$, (4.52a–c), (4.51), (4.84a), on noting that $\frac{10\Gamma-6}{3(\Gamma+1)} \geq \frac{8\Gamma-12}{7\Gamma-6}$, (4.39), (4.28b), (1.25), (4.21) and (3.31a) that, for any $w \in L^{s'}(0,T;W_0^{1,4}(\Omega))$,

$$(4.92) \quad \begin{aligned} & \left| \int_0^T \int_\Omega \frac{\partial m_{\kappa,\alpha,L}^{\Delta t}}{\partial t} \cdot w \, dx \, dt \right| \\ & \leq C \|\nabla_x \rho_{\kappa,\alpha,L}^{[\Delta t]}\|_{L^v(\Omega_T)} \|u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(0,T;H^1(\Omega))} \|w^{\{\Delta t\}}\|_{L^{s'}(0,T;W^{1,4}(\Omega))} \\ & \quad + C \left[\|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^v(\Omega_T)} + \|m_{\kappa,\alpha,L}^{\Delta t,-}\|_{L^v(\Omega_T)} + 1 \right] \|u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(0,T;H^1(\Omega))} \|w\|_{L^{s'}(0,T;W^{1,4}(\Omega))} \\ & \quad + C \|\rho_{\kappa,\alpha,L}^{[\Delta t]} u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^v(\Omega_T)} \|u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(0,T;H^1(\Omega))} \|w^{\{\Delta t\}}\|_{L^{s'}(0,T;W^{1,4}(\Omega))} \\ & \quad + \|\rho_{\kappa,\alpha,L}^{\{\Delta t\}}\|_{L^{\frac{4}{3}}(\Omega_T)} \|w\|_{L^4(0,T;W^{1,4}(\Omega))} + C \|\tau_1(M \hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+})\|_{L^2(0,T;L^{\frac{4}{3}}(\Omega))} \|w\|_{L^2(0,T;W^{1,4}(\Omega))} \\ & \quad + \|\varrho_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^\infty(0,T;L^2(\Omega))} \|\varrho_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(0,T;H^1(\Omega))} \|w\|_{L^2(0,T;L^\infty(\Omega))} \\ & \quad + \|\rho_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^\infty(0,T;L^2(\Omega))} \|f^{\{\Delta t\}}\|_{L^2(0,T;L^\infty(\Omega))} \|w\|_{L^2(\Omega_T)} \\ & \leq C \left[\|w\|_{L^{s'}(0,T;W^{1,4}(\Omega))} + \|w^{\{\Delta t\}}\|_{L^{s'}(0,T;W^{1,4}(\Omega))} \right]. \end{aligned}$$

The desired result (4.90) then follows immediately from (4.92) on noting, as $s' \geq 4$, that, for any $w \in L^{s'}(0,T;W_0^{1,4}(\Omega))$,

$$(4.93) \quad \begin{aligned} \|w^{\{\Delta t\}}\|_{L^{s'}(0,T;W^{1,4}(\Omega))}^{s'} &= \Delta t \sum_{n=1}^N \left\| \frac{1}{\Delta t} \int_{t_{n-1}}^{t_n} w \, dt \right\|_{W^{1,4}(\Omega)}^{s'} \\ &\leq C \Delta t \sum_{n=1}^N \left(\frac{1}{\Delta t} \int_{t_{n-1}}^{t_n} \|w\|_{W^{1,4}(\Omega)}^4 \, dt \right)^{\frac{s'}{4}} \leq C \|w\|_{L^{s'}(0,T;W^{1,4}(\Omega))}^{s'}. \end{aligned}$$

□

Next, we bound the time derivative of $\hat{\psi}_{\kappa,\alpha,L}^{\Delta t}$.

Lemma 4.6. *There exists a $C \in \mathbb{R}_{>0}$, independent of Δt and L , such that*

$$(4.94) \quad \left\| M \frac{\partial \hat{\psi}_{\kappa,\alpha,L}^{\Delta t}}{\partial t} \right\|_{L^2(0,T;H^s(\Omega \times D)')} \leq C,$$

where $s > 1 + \frac{1}{2}(K+1)d$.

Proof. It follows from (4.5c), (4.21), (1.25) and (4.27) that, for any $\varphi \in L^2(0, T; W^{1,\infty}(\Omega \times D))$,

$$\begin{aligned}
& \left| \int_0^T \int_{\Omega \times D} M \frac{\partial \hat{\psi}_{\kappa,\alpha,L}^{\Delta t}}{\partial t} \varphi \, dq \, dx \, dt \right| \\
& \leq 2\varepsilon \left| \int_0^T \int_{\Omega \times D} M \sqrt{\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}} \nabla_x \sqrt{\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}} \cdot \nabla_x \varphi \, dq \, dx \, dt \right| \\
& \quad + \frac{1}{2\lambda} \left| \sum_{i=1}^K \sum_{j=1}^K A_{ij} \int_0^T \int_{\Omega \times D} M \sqrt{\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}} \nabla_{q_j} \sqrt{\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}} \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt \right| \\
& \quad + \left| \int_0^T \int_{\Omega \times D} M u_{\kappa,\alpha,L}^{\Delta t,+} \beta^L(\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \cdot \nabla_x \varphi \, dq \, dx \, dt \right| \\
& \quad + \left| \int_0^T \int_{\Omega \times D} M \sum_{i=1}^K \left[\sigma(u_{\kappa,\alpha,L}^{\Delta t,+}) q_i \right] \beta^L(\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt \right| \\
& \leq C \max \left\{ 1, \|\varrho_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^\infty(0,T;L^2(\Omega))} \right\} \left[\|\nabla_x \sqrt{\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}}\|_{L^2(0,T;L_M^2(\Omega \times D))} \right. \\
& \quad \left. + \|\nabla_q \sqrt{\hat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}}\|_{L^2(0,T;L_M^2(\Omega \times D))} + \|u_{\kappa,\alpha,L}^{\Delta t,+}\|_{L^2(0,T;H^1(\Omega))} \right] \|\varphi\|_{L^2(0,T;W^{1,\infty}(\Omega \times D))} \\
(4.95) \quad & \leq C \|\varphi\|_{L^2(0,T;W^{1,\infty}(\Omega \times D))}.
\end{aligned}$$

The desired result (4.94) then follows on noting that $H^s(\Omega \times D) \hookrightarrow W^{1,\infty}(\Omega \times D)$ for the stated bound on s . \square

Remark 4.2. We note that allowing $\mathfrak{z} = 0$ would impact on the proof of Lemma 4.5. As is already clear from the formal energy inequality (1.23), by setting $\mathfrak{z} = 0$ one loses control over the $L^\infty(0, T; L^2(\Omega)) \cap L^2(0, T; H^1(\Omega))$ norm of $\varrho_{\kappa,\alpha,L}^{\Delta t,+}$; instead, one can only control weaker norms of $\varrho_{\kappa,\alpha,L}^{\Delta t,+}$, leading to (4.32) in place of (4.28b). While this weaker control is not sufficient to prove (4.90) as stated, one can prove a weaker result by replacing the $W^{1,4}(\Omega)$ norm on the test function w by the $W^{1,2d}(\Omega)$ norm throughout the proof. Admitting $\mathfrak{z} = 0$ in Lemma 4.6, on the other hand, results in unsurmountable difficulties: the proof of the lemma cannot be completed without an $L^r(0, T; L^2(\Omega))$ norm bound on $\varrho_{\kappa,\alpha,L}^{\Delta t,+}$, with $r > 2$ at least; in particular, we are unable to prove (4.94), or a weaker result on the time-difference of $\hat{\psi}_{\kappa,\alpha,L}^{\Delta t}$, when $\mathfrak{z} = 0$. Remark 4.1, and equation (4.31) in particular, indicate that an $L^r(0, T; L^2(\Omega))$ norm bound on $\varrho_{\kappa,\alpha,L}^{\Delta t,+}$, with $r > 2$, is unlikely to hold without requiring $\mathfrak{z} > 0$, regardless of the choice of \mathfrak{k} in (1.15).

4.5. Passing to the limit $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$) in the momentum equation (4.5b) and the Fokker–Planck equation (4.5c). We have the following convergence results.

Lemma 4.7. *We have that*

$$(4.96) \quad m_{\kappa,\alpha} := \rho_{\kappa,\alpha} u_{\kappa,\alpha} \in L^v(\Omega_T) \cap W^{1,s}(0, T; W_0^{1,4}(\Omega)') \cap C_w([0, T]; L^{\frac{2\Gamma}{\Gamma+1}}(\Omega)),$$

where $v = \frac{10\Gamma-6}{3(\Gamma+1)}$ and $s = \frac{8\Gamma-12}{7\Gamma-6}$, with $\Gamma \geq 8$. In addition, for a further subsequence of the subsequences of Lemmas 4.1 and 4.2 it follows that, as $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$),

$$(4.97a) \quad m_{\kappa,\alpha,L}^{\Delta t(\pm)} \rightarrow m_{\kappa,\alpha} \quad \text{weakly in } L^v(\Omega_T),$$

$$(4.97b) \quad m_{\kappa,\alpha,L}^{\Delta t} \rightarrow m_{\kappa,\alpha} \quad \text{weakly in } W^{1,s}(0, T; W_0^{1,4}(\Omega)'),$$

$$(4.97c) \quad m_{\kappa,\alpha,L}^{\Delta t(\pm)} \rightarrow m_{\kappa,\alpha} \quad \text{strongly in } L^2(0, T; H^1(\Omega)'),$$

$$(4.97d) \quad m_{\kappa,\alpha,L}^{\Delta t(\pm)} \otimes u_{\kappa,\alpha,L}^{\Delta t,+} \rightarrow m_{\kappa,\alpha} \otimes u_{\kappa,\alpha} \quad \text{weakly in } L^2(0, T; L^{\frac{6\Gamma}{4\Gamma+3}}(\Omega)),$$

$$(4.97e) \quad m_{\kappa,\alpha,L}^{\Delta t(\pm)} \rightarrow m_{\kappa,\alpha} \quad \text{in } C_w([0, T]; L^{\frac{2\Gamma}{\Gamma+1}}(\Omega)).$$

Proof. The weak convergence result (4.97a) for some limit function $\underline{m}_{\kappa,\alpha} \in \underline{L}^v(\Omega_T)$, which is the common limit of $\underline{m}_{\kappa,\alpha,L}^{\Delta t}$, $\underline{m}_{\kappa,\alpha,L}^{\Delta t,+}$ and $\underline{m}_{\kappa,\alpha,L}^{\Delta t,-}$, follows immediately from (4.84a), (4.68b), (4.3) and (4.82). The weak convergence result (4.97b) and the strong convergence result (4.97c) for $\underline{m}_{\kappa,\alpha,L}^{\Delta t}$ follow immediately for $\underline{m}_{\kappa,\alpha} \in \underline{L}^v(\Omega_T) \cap W^{1,s}(0,T;W_0^{1,4}(\Omega)')$ from (4.84a), (4.90) and (3.11), on noting that $\underline{L}^v(\Omega)$ is compactly embedded in $H^1(\Omega)'$, which is in turn continuously embedded in $W^{1,4}(\Omega)'$. The corresponding result (4.97c) for $\underline{m}_{\kappa,\alpha,L}^{\Delta t,\pm}$ then follows from (4.97c) for $\underline{m}_{\kappa,\alpha,L}^{\Delta t}$, and (4.68b). It follows from (4.53b), (4.44) and (4.97c) for $\underline{m}_{\kappa,\alpha,L}^{\Delta t,+}$ that $\underline{m}_{\kappa,\alpha} = \rho_{\kappa,\alpha} \underline{y}_{\kappa,\alpha}$, and hence (4.96). The result (4.97d) follows immediately from (4.84b), (4.97c) and (4.44). Finally (4.97e) follows from the bound on the first term in (4.84a), (4.90) and (3.14a,b). \square

Next, noting (4.4a,b), a simple calculation yields that [see (6.32)–(6.34) in [10] for details]:

$$(4.98) \quad \int_0^T \int_{\Omega \times D} M |\nabla_x \sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t}}|^2 \, dq \, dx \, dt \leq 2 \int_0^T \int_{\Omega \times D} M \left[|\nabla_x \sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}}|^2 + |\nabla_x \sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,-}}|^2 \right] \, dq \, dx \, dt,$$

and an analogous result with ∇_x replaced by ∇_q . Then, the bound (4.21), on noting (3.23a), (4.98) and the convexity of \mathcal{F} , imply the existence of a $C \in \mathbb{R}_{>0}$, independent of Δt and L , such that:

$$(4.99) \quad \begin{aligned} & \text{ess.sup}_{t \in [0,T]} \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,(\pm)}(t)) \, dq \, dx + \frac{1}{\Delta t L} \int_0^T \int_{\Omega \times D} M (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,-})^2 \, dq \, dx \, dt \\ & + \int_0^T \int_{\Omega \times D} M |\nabla_x \sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,(\pm)}}|^2 \, dq \, dx \, dt + \int_0^T \int_{\Omega \times D} M |\nabla_q \sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,(\pm)}}|^2 \, dq \, dx \, dt \leq C. \end{aligned}$$

Lemma 4.8. *For a further subsequence of the subsequences of Lemmas 4.1, 4.2 and 4.7, there exists a function*

$$(4.100a) \quad \widehat{\psi}_{\kappa,\alpha} \in L^v(0,T;Z_1) \cap H^1(0,T;M^{-1}(H^s(\Omega \times D))'),$$

where $v \in [1, \infty)$ and $s > 1 + \frac{1}{2}(K+1)d$, with finite relative entropy and Fisher information,

$$(4.100b) \quad \mathcal{F}(\widehat{\psi}_{\kappa,\alpha}) \in L^\infty(0,T;L_M^1(\Omega \times D)) \quad \text{and} \quad \sqrt{\widehat{\psi}_{\kappa,\alpha}} \in L^2(0,T;H_M^1(\Omega \times D)),$$

such that, as $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$),

$$(4.101a) \quad M^{\frac{1}{2}} \nabla_x \sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,(\pm)}} \rightarrow M^{\frac{1}{2}} \nabla_x \sqrt{\widehat{\psi}_{\kappa,\alpha}} \quad \text{weakly in } L^2(0,T;L^2(\Omega \times D)),$$

$$(4.101b) \quad M^{\frac{1}{2}} \nabla_q \sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,(\pm)}} \rightarrow M^{\frac{1}{2}} \nabla_q \sqrt{\widehat{\psi}_{\kappa,\alpha}} \quad \text{weakly in } L^2(0,T;L^2(\Omega \times D)),$$

$$(4.101c) \quad M \frac{\partial \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t}}{\partial t} \rightarrow M \frac{\partial \widehat{\psi}_{\kappa,\alpha}}{\partial t} \quad \text{weakly in } L^2(0,T;H^s(\Omega \times D)'),$$

$$(4.101d) \quad \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,(\pm)} \rightarrow \widehat{\psi}_{\kappa,\alpha} \quad \text{strongly in } L^v(0,T;L_M^1(\Omega \times D)),$$

$$(4.101e) \quad \beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,(\pm)}) \rightarrow \widehat{\psi}_{\kappa,\alpha} \quad \text{strongly in } L^v(0,T;L_M^1(\Omega \times D)),$$

$$(4.101f) \quad \tau(M \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \rightarrow \tau(M \widehat{\psi}_{\kappa,\alpha}) \quad \text{strongly in } L^r(\Omega_T),$$

where $r \in [1, \frac{4(d+2)}{3d+4})$; and, for a.a. $t \in (0,T)$,

$$(4.101g) \quad \int_{\Omega \times D} M(q) \mathcal{F}(\widehat{\psi}_{\kappa,\alpha}(x, q, t)) \, dq \, dx \leq \liminf_{\Delta t \rightarrow 0_+ (L \rightarrow \infty)} \int_{\Omega \times D} M(q) \mathcal{F}(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}(x, q, t)) \, dq \, dx.$$

In addition, we have that

$$(4.102) \quad \varrho_{\kappa,\alpha} := \int_D M \widehat{\psi}_{\kappa,\alpha} \, dq \in L^\infty(0,T;L^2(\Omega)) \cap L^2(0,T;H^1(\Omega)),$$

and, as $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$),

$$(4.103a) \quad \varrho_{\kappa,\alpha,L}^{\Delta t,+} \rightarrow \varrho_{\kappa,\alpha} \quad \text{weakly-}\star \text{ in } L^\infty(0,T;L^2(\Omega)), \quad \text{weakly in } L^2(0,T;H^1(\Omega)),$$

$$(4.103b) \quad \varrho_{\kappa,\alpha,L}^{\Delta t,+} \int_D M \beta^L(\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \, dq \xrightarrow{\sim} \varrho_{\kappa,\alpha} \quad \text{strongly in } L^{\frac{5\varsigma}{3(\varsigma-1)}}(0,T;L^\varsigma(\Omega)),$$

for any $\varsigma \in (1,6)$.

Proof. In order to prove the strong convergence result (4.101d), we will apply Dubinskii's compactness result (3.13) with $\mathfrak{X} = L_M^1(\Omega \times D)$, $\mathfrak{X}_1 = M^{-1}H^s(\Omega \times D)'$ and

$$(4.104) \quad \mathfrak{M} = \{\varphi \in Z_1 : \int_{\Omega \times D} M \left[|\nabla_q \sqrt{\varphi}|^2 + |\nabla_x \sqrt{\varphi}|^2 \right] \, dq \, dx < \infty\}.$$

See Section 5 in [11] for the proof of the compactness of the embedding $\mathfrak{M} \hookrightarrow \mathfrak{X}$, and the continuity of the embedding $\mathfrak{X} \hookrightarrow \mathfrak{X}_1$. Hence, the desired result (4.101d) for $\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t}$ and $v = 1$ follows from (3.13) with $\varsigma_0 = \varsigma_1 = 2$, and the stated choices of \mathfrak{M} , \mathfrak{X} and \mathfrak{X}_1 above, on noting (4.99) and (4.94). The desired result (4.101d) for $\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,\pm}$ and $v = 1$ then follows from (4.101d) for $\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t}$ and $v = 1$, (4.82) with $\widehat{m}_{\kappa,\alpha,L}^{\Delta t(\pm)}$ replaced by $\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t(\pm)}$, the second bound in (4.99) and (4.40). The desired result (4.101d) for $v \in (1,\infty)$ then follows from (4.101d) for $v = 1$, the first bound in (4.99) (note that $\mathcal{F}(s) \geq [s - e + 1]_+$ for all $s \geq 0$, since $\mathcal{F}(s) \geq 0$ and, by convexity of \mathcal{F} , $\mathcal{F}(s) \geq \mathcal{F}(e) + (s - e)\mathcal{F}'(e) = s - e + 1$) and an interpolation result, see Lemma 5.1 in [11]. The weak convergence result (4.101c) follows immediately from (4.94). The weak convergence results (4.101a,b) follow immediately from the last two terms in (4.99), on noting an argument similar to that in the proof of Lemma 3.3 in [11] in order to identify the limit. The result (4.101e) follows from (4.101d) and the Lipschitz continuity of β^L , see (5.8) in [14] for details. The result (4.101g) follows from (4.101d) and Fatou's lemma, see (6.46) in [11] for details. In addition, the convergence results (4.101a–d,g) yield the desired results (4.100a,b).

The results (4.103a) for some limit function $\varrho_{\kappa,\alpha}$ follow immediately from the bounds on $\varrho_{\kappa,\alpha,L}^{\Delta t,+}$ in (4.21). The fact that $\varrho_{\kappa,\alpha} = \int_D M \widehat{\psi}_{\kappa,\alpha} \, dq$ follows from (3.93), (4.4b) and (4.101d), and hence the desired result (4.102). The strong convergence results (4.103b) follow from noting the embedding $H^1(\Omega) \hookrightarrow L^6(\Omega)$, (1.25), (4.101d,e) and (3.1).

Finally, we need to prove (4.101f). Similarly to (4.25)–(4.28a,b), we deduce from (4.100b), (4.102), (3.2) and (4.24) that

$$(4.105) \quad \|\rho_{\kappa,\alpha}\|_{L^{\frac{2(d+2)}{d}}(\Omega_T)} + \|\tau_1(M \widehat{\psi}_{\kappa,\alpha})\|_{L^{\frac{4(d+2)}{3d+4}}(\Omega_T)} \leq C.$$

On recalling (1.11) and (4.22), we have that

$$(4.106) \quad \tau_1(M \varphi) \underset{\approx}{=} k \int_D M \left(\sum_{i=1}^K \nabla_{q_i} \varphi \otimes q_i \right) \, dq - k \left(\int_D M \varphi \, dq \right) I \underset{\approx}{=}.$$

Let $D_0 \subset \overline{D_0} \subset D$ be an arbitrary Lipschitz subdomain of D , then (4.106) yields that

$$(4.107) \quad \begin{aligned} & \int_{\Omega_T} |\tau_1(M \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) - \tau_1(M \widehat{\psi}_{\kappa,\alpha})| \, dx \, dt \\ & \leq k |b|_{\frac{1}{2}}^{\frac{1}{2}} \int_0^T \int_{\Omega} \int_{D \setminus D_0} M \sum_{i=1}^K \left(|\nabla_q \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}| + |\nabla_q \widehat{\psi}_{\kappa,\alpha}| \right) \, dq \, dx \, dt \\ & \quad + k \int_0^T \int_{\Omega} \left| \int_{D_0} M \sum_{i=1}^K q_i \otimes \nabla_{q_i} (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}) \, dq \right| \, dx \, dt \\ & \quad + k d^{\frac{1}{2}} \int_0^T \int_{\Omega \times D} M |\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}| \, dq \, dx \, dt =: T_1 + T_2 + T_3, \end{aligned}$$

where we have recalled (3.4). Further, we deduce from (4.101b,d) that

$$M \nabla_{q_i} \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} = 2M \sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}} \nabla_{q_i} \sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}} \rightarrow 2M \sqrt{\widehat{\psi}_{\kappa,\alpha}} \nabla_{q_i} \sqrt{\widehat{\psi}_{\kappa,\alpha}} = M \nabla_{q_i} \widehat{\psi}_{\kappa,\alpha}, \quad i = 1, \dots, K,$$

weakly in $L^1(0, T; \widetilde{L}^1(\Omega \times D)) = \widetilde{L}^1(\Omega_T \times D)$ as $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$). By the Dunford–Pettis theorem the sequence $\{M \nabla_q \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}\}_{\Delta t > 0}$ is therefore equi-integrable in $\widetilde{L}^1(\Omega_T \times D)$; hence, for any $\delta > 0$ there exists a $\delta_0 = \delta_0(\delta)$ such that for any set $D_0 \subset D$ with $T |\Omega| |D \setminus D_0| < \delta_0$,

$$k |b|_1^{\frac{1}{2}} \int_0^T \int_{\Omega} \int_{D \setminus D_0} M \sum_{i=1}^K \left(|\nabla_q \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}| + |\nabla_q \widehat{\psi}_{\kappa,\alpha}| \right) dq dx dt < \frac{\delta}{3}.$$

We therefore select $D_0 \subset \overline{D_0} \subset D$ to be a Lipschitz subdomain of D such that $T |\Omega| |D \setminus D_0| < \delta_0$, which implies that $0 < T_1 < \frac{\delta}{3}$; that, now, fixes our choice of D_0 .

Next, we bound T_2 . By performing partial integration over D_0 , we have that

$$\begin{aligned} T_2 &= k \int_0^T \int_{\Omega} \left| \int_{D_0} M \sum_{i=1}^K q_i \otimes \nabla_{q_i} (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}) dq \right| dx dt \\ &\leq k \int_0^T \int_{\Omega} \left| - \int_{D_0} \sum_{i=1}^K \left(\nabla_{q_i} M \otimes q_i \right) (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}) dq \right| dx dt \\ &\quad + k \int_0^T \int_{\Omega} \left| -K \left[\int_{D_0} M (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}) dq \right] I \right| dx dt \\ &\quad + k \int_0^T \int_{\Omega} \left| \int_{\partial D_0} \sum_{i=1}^K M (n_i \otimes q_i) (\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}) d\sigma(q) \right| dx dt, \end{aligned}$$

where the d -component column vector n_i is the i th component of the Kd -component unit outward (column) normal vector $\underline{n} = (n_1^T, \dots, n_K^T)^T$ to the boundary ∂D_0 of D_0 . As the closure of the Lipschitz subdomain D_0 is a strict subset of the open set D , we have, on noting (1.6a) and (1.7a,b), that $\sup_{q \in D_0} \left(\frac{1}{M(q)} |\nabla_q M(q)| \right) \leq C(\delta_0) < \infty$. Hence,

$$\begin{aligned} T_2 &\leq k \int_0^T \int_{\Omega} \int_{D_0} \left[|b|_1^{\frac{1}{2}} |\nabla_q M| + K d^{\frac{1}{2}} M \right] |\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}| dq dx dt \\ &\quad + k |b|_1^{\frac{1}{2}} \int_0^T \int_{\Omega} \int_{\partial D_0} M |\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}| d\sigma(q) dx dt \\ &\leq k \left(|b|_1^{\frac{1}{2}} C(\delta_0) + K d^{\frac{1}{2}} \right) \int_0^T \int_{\Omega} \int_{D_0} M |\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}| dq dx dt \\ &\quad + k |b|_1^{\frac{1}{2}} \|M\|_{L^\infty(D)} \int_0^T \int_{\Omega} \int_{\partial D_0} |\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}| d\sigma(q) dx dt \\ (4.108) \quad &=: T_{21} + T_{22}. \end{aligned}$$

Thus, thanks to (4.101d) with $v = 1$, there exists a Δt_0 such that for all $\Delta t \leq \Delta t_0$, we have that $0 < T_{21} < \frac{\delta}{6}$ and $0 < T_3 < \frac{\delta}{3}$.

Finally, we shall show that, for Δt_0 sufficiently small, also $0 < T_{22} < \frac{\delta}{6}$. In the process of doing so we shall repeatedly use the following result. As the closure of D_0 is a compact subset of D , we have from (1.7a) that

$$(4.109) \quad \sup_{q \in D_0} [M(q)]^{-1} \leq C(D_0) < \infty.$$

We begin by noting that (4.99) and (4.109) imply that $\{\sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}}\}_{\Delta t > 0}$ is a bounded sequence in $L^2(0, T; H^1(\Omega \times D_0))$; hence, by Sobolev embedding, it is also a bounded sequence in $L^2(0, T; L^{\frac{2(K+1)d}{(K+1)d-2}}(\Omega \times D_0))$. Further, by (4.99) and (4.109), $\{\sqrt{\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}}\}_{\Delta t > 0}$ is a bounded sequence in

$L^\infty(0, T; L^2(\Omega \times D_0))$. It then follows from (3.1) that $\{\sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}}\}_{\Delta t > 0}$ is a bounded sequence in $L^{\frac{2((K+1)d+2)}{(K+1)d}}(0, T; L^{\frac{2((K+1)d+2)}{(K+1)d}}(\Omega \times D_0))$; thus,

$$(4.110) \quad \int_0^T \int_{\Omega \times D_0} |\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}|^{\frac{(K+1)d+2}{(K+1)d}} \, dq \, dx \, dt \leq C(D_0),$$

where the constant $C(D_0)$ is independent of Δt and L . Now, for any $s \in (1, 2)$, we have by Hölder's inequality, (4.99), (4.109) and the inequality $(a^{\frac{s}{2}} + b^{\frac{s}{2}}) \leq 2^{1-\frac{s}{2}}(a+b)^{\frac{s}{2}}$ with $a, b \geq 0$, which follows from the concavity of the function $x \in [0, \infty) \mapsto x^{\frac{s}{2}} \in [0, \infty)$, that

$$\begin{aligned} & \int_0^T \int_{\Omega \times D_0} \left(|\nabla_x \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}|^s + |\nabla_q \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}|^s \right) \, dq \, dx \, dt \\ &= 2^s \int_0^T \int_{\Omega \times D_0} |\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}|^{\frac{s}{2}} \left(\left| \nabla_x \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \right|^s + \left| \nabla_q \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \right|^s \right) \, dq \, dx \, dt \\ &\leq 2^{\frac{s}{2}+1} \left(\int_0^T \int_{\Omega \times D_0} |\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}|^{\frac{s}{2-s}} \, dq \, dx \, dt \right)^{\frac{2-s}{2}} \\ &\quad \times \left(\int_0^T \int_{\Omega \times D_0} \left| \nabla_x \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \right|^2 + \left| \nabla_q \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \right|^2 \, dq \, dx \, dt \right)^{\frac{s}{2}} \\ &\leq C \left(\int_0^T \int_{\Omega \times D_0} |\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}|^{\frac{s}{2-s}} \, dq \, dx \, dt \right)^{\frac{2-s}{2}}. \end{aligned}$$

Comparing this with (4.110) indicates that $s \in (1, 2)$ should be chosen so that

$$\frac{s}{2-s} \leq \frac{(K+1)d+2}{(K+1)d}.$$

The largest such s is $s = \frac{(K+1)d+2}{(K+1)d+1}$; using this value of s , we then deduce on noting (4.110) that

$$(4.111) \quad \|\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^s(0, T; W^{1, s}(\Omega \times D_0))} \leq C(D_0).$$

Note further that, thanks to (4.109) and (4.101d), we have for any $v \in [1, \infty)$ that

$$(4.112) \quad \|\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +} - \widehat{\psi}_{\kappa, \alpha}\|_{L^v(0, T; L^1(\Omega \times D_0))} \rightarrow 0, \quad \text{as } \Delta t \rightarrow 0_+ \ (L \rightarrow \infty).$$

We shall now use (4.111) and (4.112) to show that T_{22} converges to 0 as $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$). To this end, we shall make use of the following sharp trace inequality, established recently by Auchmuty (cf. Theorem 6.3 inequality (6.3) in [3]): suppose that \mathcal{O} is a bounded Lipschitz domain and let $r \in (1, 2)$; then, the following inequality holds for all $\varphi \in W^{1, r}(\mathcal{O})$:

$$(4.113) \quad \int_{\partial \mathcal{O}} |\varphi|^{2-\frac{1}{r}} \, d\sigma \leq \frac{|\partial \mathcal{O}|}{|\mathcal{O}|} \|\varphi\|_{L^{2-\frac{1}{r}}(\mathcal{O})}^{2-\frac{1}{r}} + \left(2 - \frac{1}{r}\right) k_{\mathcal{O}} \|\varphi\|_{L^1(\mathcal{O})}^{1-\frac{1}{r}} \|\nabla \varphi\|_{L^r(\mathcal{O})},$$

where $k_{\mathcal{O}}$ is a positive constant, which depends on \mathcal{O} only. We deduce from (4.113) and (3.1) for any $r \in (1, 2)$ that

$$(4.114) \quad \int_{\partial \mathcal{O}} |\varphi|^{2-\frac{1}{r}} \, d\sigma \leq C(\mathcal{O}, r) \|\varphi\|_{L^1(\mathcal{O})}^{1-\frac{1}{r}} \|\varphi\|_{W^{1, r}(\mathcal{O})} \quad \forall \varphi \in W^{1, r}(\mathcal{O}).$$

We apply (4.114) with $\mathcal{O} = D_0$, integrate the resulting inequality over $(0, T) \times \Omega$ and apply Hölder's inequality; this yields for any $r \in (1, 2)$ and for all $\varphi \in L^r(0, T; W^{1, r}(\Omega \times D_0))$ that

$$(4.115) \quad \int_0^T \int_{\Omega \times \partial D_0} |\varphi|^{2-\frac{1}{r}} \, d\sigma(q) \, dx \, dt \leq C(D_0, r) \|\varphi\|_{L^1(0, T; L^1(\Omega \times D_0))}^{1-\frac{1}{r}} \|\varphi\|_{L^r(0, T; W^{1, r}(\Omega \times D_0))}.$$

Motivated by the bound (4.111), we fix

$$r = s = \frac{(K+1)d+2}{(K+1)d+1} \in (1, 2)$$

in (4.115). It follows from (4.115), (4.111) and (4.112) that

$$(4.116) \quad \int_0^T \int_{\Omega \times \partial D_0} |\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+} - \widehat{\psi}_{\kappa,\alpha}|^{2-\frac{1}{s}} d\sigma(q) dx dt \rightarrow 0, \quad \text{as } \Delta t \rightarrow 0_+ \ (L \rightarrow \infty).$$

Since $2 - \frac{1}{s} > 1$, it follows from (4.116) that T_{22} converges to 0, as $\Delta t \rightarrow 0_+ \ (L \rightarrow \infty)$. We thus deduce that there exists a Δt_0 such that for all $\Delta t \leq \Delta t_0$, we have that $0 < T_{22} < \frac{\delta}{6}$. Finally, by recalling the inequalities (4.107) and (4.108) and the bounds on T_1 , T_{21} , T_{22} and T_3 , it then follows that for each $\delta > 0$ there exists a Δt_0 such that for all $\Delta t \leq \Delta t_0$, we have that

$$\int_{\Omega_T} |\tau_1(M\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) - \tau_1(M\widehat{\psi}_{\kappa,\alpha})| dx dt < \delta.$$

Thus we have proved that

$$\tau_1(M\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \rightarrow \tau_1(M\widehat{\psi}_{\kappa,\alpha}) \quad \text{strongly in } L^1(\Omega_T), \quad \text{as } \Delta t \rightarrow 0_+ \ (L \rightarrow \infty).$$

This, together with (4.28b), (4.105) and (3.1), implies that, as $\Delta t \rightarrow 0_+ \ (L \rightarrow \infty)$,

$$(4.117) \quad \tau_1(M\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+}) \rightarrow \tau_1(M\widehat{\psi}_{\kappa,\alpha}) \quad \text{strongly in } L^r(\Omega_T) \text{ for all } r \in \left[1, \frac{4(d+2)}{3d+4}\right).$$

We note further that, according to (4.103b) with $\varsigma = 2$, $\varrho_{\kappa,\alpha,L}^{\Delta t,+} \rightarrow \varrho_{\kappa,\alpha}$ strongly in $L^{\frac{10}{3}}(0, T; L^2(\Omega))$ as $\Delta t \rightarrow 0_+ \ (L \rightarrow \infty)$; therefore $(\varrho_{\kappa,\alpha,L}^{\Delta t,+})^2 \rightarrow \varrho_{\kappa,\alpha}^2$ strongly in $L^{\frac{5}{3}}(0, T; L^1(\Omega))$, and thus strongly in $L^1(0, T; L^1(\Omega)) = L^1(\Omega_T)$, as $\Delta t \rightarrow 0_+ \ (L \rightarrow \infty)$. Also, by (4.27), since $\frac{8(d+2)}{3d+4} < \frac{2(d+2)}{d}$ for $d \in \{2, 3\}$, we have that $\{(\varrho_{\kappa,\alpha,L}^{\Delta t,+})^2\}_{\Delta t > 0}$ is a bounded sequence in $L^{\frac{4(d+2)}{3d+4}}(\Omega_T)$; consequently from (4.105) and (3.1), $(\varrho_{\kappa,\alpha,L}^{\Delta t,+})^2 \rightarrow (\varrho_{\kappa,\alpha})^2$ strongly in $L^r(\Omega_T)$ for all $r \in \left[1, \frac{4(d+2)}{3d+4}\right)$. Combining this with (4.117) we deduce (4.101f) thanks to (1.10). \square

We are now ready to pass to the limit with $\Delta t \rightarrow 0_+ \ (L \rightarrow \infty)$ in (4.5a–c) to prove the existence of a weak solution to the regularized problem $(P_{\kappa,\alpha})$.

Theorem 4.1. *The triple $(\rho_{\kappa,\alpha}, \underline{u}_{\kappa,\alpha}, \widehat{\psi}_{\kappa,\alpha})$, defined as in Lemmas 4.1 and 4.8, is a global weak solution to problem $(P_{\kappa,\alpha})$, in the sense that*

$$(4.118a) \quad \int_0^T \left\langle \frac{\partial \rho_{\kappa,\alpha}}{\partial t}, \eta \right\rangle_{H^1(\Omega)} dt + \int_0^T \int_{\Omega} \left(\alpha \nabla_x \rho_{\kappa,\alpha} - \rho_{\kappa,\alpha} \underline{u}_{\kappa,\alpha} \right) \cdot \nabla_x \eta dx dt = 0$$

$$\forall \eta \in L^2(0, T; H^1(\Omega)),$$

with $\rho_{\kappa,\alpha}(\cdot, 0) = \rho^0(\cdot)$,

$$(4.118b) \quad \int_0^T \left\langle \frac{\partial(\rho_{\kappa,\alpha} \underline{u}_{\kappa,\alpha})}{\partial t}, w \right\rangle_{W_0^{1,4}(\Omega)} dt + \frac{\alpha}{2} \int_0^T \int_{\Omega} \nabla_x \rho_{\kappa,\alpha} \cdot \nabla_x (\underline{u}_{\kappa,\alpha} \cdot w) dx dt$$

$$+ \int_0^T \int_{\Omega} \left[S(\underline{u}_{\kappa,\alpha}) - \rho_{\kappa,\alpha} \underline{u}_{\kappa,\alpha} \otimes \underline{u}_{\kappa,\alpha} - p_{\kappa}(\rho_{\kappa,\alpha}) I \right] : \nabla_x w dx dt$$

$$= \int_0^T \int_{\Omega} \left[\rho_{\kappa,\alpha} f \cdot w - \left(\tau_1(M\widehat{\psi}_{\kappa,\alpha}) - \mathfrak{z} \varrho_{\kappa,\alpha}^2 I \right) : \nabla_x w \right] dx dt \quad \forall w \in L^r(0, T; W_0^{1,4}(\Omega)),$$

with $(\rho_{\kappa,\alpha} \underline{u}_{\kappa,\alpha})(\cdot, 0) = (\rho^0 \underline{u}_0)(\cdot)$ and $r = \frac{8\Gamma-12}{\Gamma-6}$, $\Gamma \geq 8$, and

$$\begin{aligned}
& \int_0^T \left\langle M \frac{\partial \widehat{\psi}_{\kappa,\alpha}}{\partial t}, \varphi \right\rangle_{H^s(\Omega \times D)} dt + \frac{1}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \int_0^T \int_{\Omega \times D} M \nabla_{q_j} \widehat{\psi}_{\kappa,\alpha} \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt \\
& + \int_0^T \int_{\Omega \times D} M \left[\varepsilon \nabla_x \widehat{\psi}_{\kappa,\alpha} - \underline{u}_{\kappa,\alpha} \widehat{\psi}_{\kappa,\alpha} \right] \cdot \nabla_x \varphi \, dq \, dx \, dt \\
& - \int_0^T \int_{\Omega \times D} M \sum_{i=1}^K \left[\sigma(\underline{u}_{\kappa,\alpha}) q_i \right] \widehat{\psi}_{\kappa,\alpha} \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt = 0 \quad \forall \varphi \in L^2(0, T; H^s(\Omega \times D)),
\end{aligned}
\tag{4.118c}$$

with $\widehat{\psi}_{\kappa,\alpha}(\cdot, 0) = \widehat{\psi}_0(\cdot)$ and $s > 1 + \frac{1}{2}(K+1)d$.

In addition, the weak solution $(\rho_{\kappa,\alpha}, \underline{u}_{\kappa,\alpha}, \widehat{\psi}_{\kappa,\alpha})$ satisfies, for a.a. $t' \in (0, T)$, the following energy inequality:

$$\begin{aligned}
& \frac{1}{2} \int_{\Omega} \rho_{\kappa,\alpha}(t') |\underline{u}_{\kappa,\alpha}(t')|^2 \, dx + \int_{\Omega} P_{\kappa}(\rho_{\kappa,\alpha}(t')) \, dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_{\kappa,\alpha}(t')) \, dq \, dx \\
& + \alpha \kappa \int_0^{t'} \left[\|\nabla_x(\rho_{\kappa,\alpha}^2)\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x(\rho_{\kappa,\alpha}^{\frac{\Gamma}{2}})\|_{L^2(\Omega)}^2 \right] dt + \mu^S c_0 \int_0^{t'} \|\underline{u}_{\kappa,\alpha}\|_{H^1(\Omega)}^2 dt \\
& + k \int_0^{t'} \int_{\Omega \times D} M \left[\frac{a_0}{2\lambda} \left| \nabla_q \sqrt{\widehat{\psi}_{\kappa,\alpha}} \right|^2 + 2\varepsilon \left| \nabla_x \sqrt{\widehat{\psi}_{\kappa,\alpha}} \right|^2 \right] \, dq \, dx \, dt \\
& + \mathfrak{z} \|\varrho_{\kappa,\alpha}(t')\|_{L^2(\Omega)}^2 + 2\mathfrak{z} \varepsilon \int_0^{t'} \|\nabla_x \varrho_{\kappa,\alpha}\|_{L^2(\Omega)}^2 dt \\
& \leq e^{t'} \left[\frac{1}{2} \int_{\Omega} \rho^0 |\underline{u}_0|^2 \, dx + \int_{\Omega} P_{\kappa}(\rho^0) \, dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_0) \, dq \, dx \right. \\
& \quad \left. + \mathfrak{z} \int_{\Omega} \left(\int_D M \widehat{\psi}_0 \, dq \right)^2 \, dx + \frac{1}{2} \int_0^{t'} \|f\|_{L^\infty(\Omega)}^2 dt \int_{\Omega} \rho^0 \, dx \right]
\end{aligned}
\tag{4.119}$$

$\leq C$,

where $C \in \mathbb{R}_{>0}$ is independent of α and κ .

Proof. The limit equation (4.118a) has already been established in Lemma 4.1, see (4.45).

We now pass to the limit $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$), subject to (4.40), for the subsequence of $\{(\rho_{\kappa,\alpha,L}^{[\Delta t]}, \underline{u}_{\kappa,\alpha,L}^{\Delta t,+}, \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t,+})\}_{\Delta t>0}$ of Lemma 4.8 in (4.89) initially for a fixed test function $w \in \mathcal{C}_0^\infty(\Omega_T)$. We consider first the five terms on the left-hand side of (4.89). On noting (4.97b,d), (4.53a), (4.44), (4.52c) and (4.49), we obtain the first two terms on the left-hand side of (4.118b) and the $\rho_{\kappa,\alpha} \underline{u}_{\kappa,\alpha} \otimes \underline{u}_{\kappa,\alpha}$ term from the first term on the left-hand side of (4.89). The second and third terms on the left-hand side of (4.89) give rise to the remaining terms on the left-hand side of (4.118b), on noting (4.44) and (4.43b). The fourth and fifth terms on the left-hand side of (4.89) converge to zero, on noting (4.68a), (4.51), (4.52b) and (4.49).

We now consider the two terms on the right-hand side of (4.89). The first term gives rise to the f and $\underline{\tau}_1$ contributions on the right-hand side of (4.118b), on noting (4.53b), (3.31b) and (4.101f). The second term on the right-hand side of (4.89) converges to the $\varrho_{\kappa,\alpha}^2$ term on the right-hand side of (4.118b), on noting (4.103a,b) and performing integration by parts. Therefore, we have obtained (4.118b) for any $w \in \mathcal{C}_0^\infty(\Omega_T)$. The desired result (4.118b) for any $w \in L^r(0, T; W_0^{1,4}(\Omega))$ then follows from the denseness of $\mathcal{C}_0^\infty(\Omega_T)$ in $L^r(0, T; W_0^{1,4}(\Omega))$, (4.96), (4.53a,b), (4.44), (4.91), (4.97d), (4.41), (3.5), (3.4), (4.101f), and finally (4.102), which with (3.2) yields, similarly to (4.26) and (4.27), that $\varrho_{\kappa,\alpha} \in L^4(0, T; L^{\frac{2d}{d-1}}(\Omega))$.

Similarly, we now pass to the limit $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$) for the subsequence in (4.5c) initially for a fixed test function $\varphi \in C([0, T]; C^\infty(\overline{\Omega \times D}))$. The first term of (4.5c) converges to the first

term of (4.118c), on noting (4.101c). For the second term of (4.5c), we note that

$$\begin{aligned}
 & \int_0^T \int_{\Omega \times D} M \nabla_{q_i} \widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +} \cdot \nabla_{q_j} \varphi \, dq \, dx \, dt \\
 &= 2 \int_0^T \int_{\Omega \times D} M \left(\sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} - \sqrt{\widehat{\psi}_{\kappa, \alpha}} \right) \nabla_{q_i} \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \cdot \nabla_{q_j} \varphi \, dq \, dx \, dt \\
 &+ 2 \int_0^T \int_{\Omega \times D} M \sqrt{\widehat{\psi}_{\kappa, \alpha}} \nabla_{q_i} \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \cdot \nabla_{q_j} \varphi \, dq \, dx \, dt =: T_1 + T_2.
 \end{aligned}
 \tag{4.120}$$

Next, on noting (4.99) and that $|\sqrt{c_1} - \sqrt{c_2}| \leq \sqrt{|c_1 - c_2|}$ for all $c_1, c_2 \in \mathbb{R}_{\geq 0}$, we have that

$$\begin{aligned}
 |T_1| &\leq C \left\| \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} - \sqrt{\widehat{\psi}_{\kappa, \alpha}} \right\|_{L^2(0, T; L_M^2(\Omega \times D))} \|\nabla_{q_j} \varphi\|_{L^\infty(0, T; L^\infty(\Omega \times D))} \\
 &\leq C \|\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +} - \widehat{\psi}_{\kappa, \alpha}\|_{L^1(0, T; L_M^1(\Omega \times D))}^{\frac{1}{2}} \|\nabla_{q_j} \varphi\|_{L^\infty(0, T; L^\infty(\Omega \times D))},
 \end{aligned}
 \tag{4.121}$$

and so (4.101d) yields that T_1 converges to zero as $\Delta t \rightarrow 0_+$. Similarly, as $M^{\frac{1}{2}} \sqrt{\widehat{\psi}_{\kappa, \alpha}} \nabla_{q_j} \varphi \in L^2(0, T; L^2(\Omega \times D))$, it follows from (4.101b) that, as $\Delta t \rightarrow 0_+$,

$$T_2 \rightarrow 2 \int_0^T \int_{\Omega \times D} M \sqrt{\widehat{\psi}_{\kappa, \alpha}} \nabla_{q_i} \sqrt{\widehat{\psi}_{\kappa, \alpha}} \cdot \nabla_{q_j} \varphi \, dq \, dx \, dt = \int_0^T \int_{\Omega \times D} M \nabla_{q_i} \widehat{\psi}_{\kappa, \alpha} \cdot \nabla_{q_j} \varphi \, dq \, dx \, dt.$$

Hence the second term in (4.5c) converges to the second term in (4.118c). For the fourth term in (4.5c), we note that

$$\begin{aligned}
 & \int_0^T \int_{\Omega \times D} M \left[\sigma(\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}) q_i \right] \beta^L(\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}) \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt \\
 &= \int_0^T \int_{\Omega \times D} M \left[\nabla_x u_{\kappa, \alpha, L}^{\Delta t, +} q_i \right] \left(\beta^L(\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}) - \widehat{\psi}_{\kappa, \alpha} \right) \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt \\
 &+ \int_0^T \int_{\Omega \times D} M \left[\nabla_x u_{\kappa, \alpha, L}^{\Delta t, +} q_i \right] \widehat{\psi}_{\kappa, \alpha} \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt =: T_3 + T_4.
 \end{aligned}
 \tag{4.123}$$

Next, on noting (4.51), (3.1), (1.25), (4.29) and (4.102), we have that

$$\begin{aligned}
 |T_3| &\leq C \left\| \int_D M |\beta^L(\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}) - \widehat{\psi}_{\kappa, \alpha}| \, dq \right\|_{L^2(0, T; L^2(\Omega))} \|\nabla_{q_i} \varphi\|_{L^\infty(0, T; L^\infty(\Omega \times D))} \\
 &\leq C \|\beta^L(\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}) - \widehat{\psi}_{\kappa, \alpha}\|_{L^2(0, T; L_M^1(\Omega \times D))}^{\frac{2}{5}} \|\varrho_{\kappa, \alpha}^{\Delta t, +} + \varrho_{\kappa, \alpha}\|_{L^2(0, T; L^6(\Omega))}^{\frac{3}{5}} \|\nabla_{q_i} \varphi\|_{L^\infty(0, T; L^\infty(\Omega))},
 \end{aligned}
 \tag{4.124}$$

and so (4.101e) and (4.26) (with $v = 6$ and $\vartheta = \frac{d}{3}$, $d = 2, 3$) yield that T_3 converges to zero as $\Delta t \rightarrow 0_+$. Similarly, as (4.102) yields that $\int_D M \widehat{\psi}_{\kappa, \alpha} q_i \otimes \nabla_{q_i} \varphi \, dq \in L^2(\Omega_T)$, it follows from (4.44) that, as $\Delta t \rightarrow 0_+$,

$$T_4 \rightarrow \int_0^T \int_{\Omega \times D} M [(\nabla_x u_{\kappa, \alpha}) q_i] \widehat{\psi}_{\kappa, \alpha} \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt.
 \tag{4.125}$$

Hence the last term in (4.5c) converges to the last term in (4.118c). Similarly to the second and last terms, the third term in (4.5c) converges to the third term in (4.118c). Therefore, we have obtained (4.118c) for any $\varphi \in C([0, T]; C^\infty(\overline{\Omega \times D}))$. The desired result (4.118c), for any $\varphi \in L^2(0, T; H^s(\Omega \times D))$, then follows from the denseness of the function space $C([0, T]; C^\infty(\overline{\Omega \times D}))$ in $L^2(0, T; H^s(\Omega \times D))$, $H^s(\Omega \times D) \hookrightarrow W^{1, \infty}(\Omega \times D)$, (4.100a,b), (4.102) and (4.44).

Next we shall verify the attainment of the respective initial data by $\rho_{\kappa, \alpha} u_{\kappa, \alpha}$ and $\widehat{\psi}_{\kappa, \alpha}$. We have already established that $\rho_{\kappa, \alpha} u_{\kappa, \alpha} \in C_w([0, T]; L^{\frac{2\Gamma}{\Gamma+1}}(\Omega))$, see (4.96). That $\widehat{\psi}_{\kappa, \alpha} \in C_w([0, T];$

$L_M^1(\Omega \times D)$) follows from $\mathcal{F}(\widehat{\psi}_{\kappa,\alpha}) \in L^\infty(0, T; L_M^1(\Omega \times D))$ and $\widehat{\psi}_{\kappa,\alpha} \in H^1(0, T; M^{-1}(H^s(\Omega \times D)))'$ (cf. (4.100b) and (4.94)) with $s > 1 + \frac{1}{2}(K+1)d$, by Lemma 3.1(b) on taking $\mathfrak{X} := L_M^\Phi(\Omega \times D)$, the Maxwellian weighted Orlicz space with Young's function $\Phi(r) = \mathcal{F}(1 + |r|)$ (cf. Kufner, John & Fučík [41], Sec. 3.18.2) whose separable predual $\mathfrak{E} := E_M^\Psi(\Omega \times D)$ has Young's function $\Psi(r) = \exp |r| - |r| - 1$, and $\mathfrak{Y} := M^{-1}(H^s(\Omega \times D))'$ whose predual with respect to the duality pairing $\langle M \cdot, \cdot \rangle_{H^s(\Omega \times D)}$ is $\mathfrak{F} := H^s(\Omega \times D)$, with $s > 1 + \frac{1}{2}(K+1)d$, and noting the embedding $C_{w*}([0, T]; L_M^\Phi(\Omega \times D)) \hookrightarrow C_w([0, T]; L_M^1(\Omega \times D))$. The last embedding and that $\mathfrak{F} \hookrightarrow \mathfrak{E}$ are proved by adapting Def. 3.6.1. and Thm. 3.2.3 in Kufner, John & Fučík [41] to the measure $M(q) dq d\mathfrak{x}$ to show that $L^\infty(\Omega \times D) \hookrightarrow L_M^\Xi(\Omega \times D)$ for any Young's function Ξ , and then adapting Theorem 3.17.7 *ibid.* to deduce that $\mathfrak{F} \hookrightarrow L^\infty(\Omega \times D) \hookrightarrow E_M^\Psi(\Omega \times D) = \mathfrak{E}$.

We are now ready to prove that $\rho_{\kappa,\alpha} u_{\kappa,\alpha}$ and $\widehat{\psi}_{\kappa,\alpha}$ satisfy the initial conditions $(\rho_{\kappa,\alpha} u_{\kappa,\alpha})(\cdot, 0) = (\rho^0 u_0)(\cdot)$ and $\widehat{\psi}_{\kappa,\alpha}(\cdot, \cdot, 0) = \widehat{\psi}_0(\cdot, \cdot)$ in the sense of $C_w([0, T]; \mathcal{L}^{\frac{2F}{F+1}}(\Omega))$ and $C_w([0, T]; L_M^1(\Omega \times D))$, respectively. The desired result for $\rho_{\kappa,\alpha} u_{\kappa,\alpha}$ follows immediately from (4.97e) and (3.20b). We now consider $\widehat{\psi}_{\kappa,\alpha}$. According to (4.94), there exists a $C \in \mathbb{R}_{>0}$, independent of Δt and L , such that

$$\left| \int_0^T \int_{\Omega \times D} M \frac{\partial \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t}}{\partial t} \varphi d\mathfrak{q} d\mathfrak{x} dt \right| \leq C \|\varphi\|_{L^2(0,T;H^s(\Omega \times D))} \quad \forall \varphi \in L^2(0,T;H^s(\Omega \times D)),$$

where $s > 1 + \frac{1}{2}(K+1)d$. Choosing, in particular

$$\varphi(\mathfrak{x}, t) = \phi(\mathfrak{x}) \left(1 - \frac{t}{\delta}\right)_+, \quad \phi \in H^s(\Omega \times D), \quad 0 < \delta < T,$$

integrating by parts with respect to t and using that $\widehat{\psi}_{\kappa,\alpha,L}^{\Delta t}(\cdot, \cdot, 0) = \widehat{\psi}^0(\cdot, \cdot)$, we have that

$$\left| \frac{1}{\delta} \int_0^\delta \int_{\Omega \times D} M \widehat{\psi}_{\kappa,\alpha,L}^{\Delta t} \phi d\mathfrak{q} d\mathfrak{x} dt - \int_{\Omega \times D} M \widehat{\psi}^0 \phi d\mathfrak{q} d\mathfrak{x} \right| \leq C \delta^{\frac{1}{2}} \|\phi\|_{H^s(\Omega \times D)} \quad \forall \phi \in H^s(\Omega \times D).$$

For $\delta \in (0, T)$ and ϕ fixed, we now pass to the limit $\Delta t \rightarrow 0_+$ ($L \rightarrow \infty$) in this inequality using (4.101d) and (3.23b) to deduce that

$$\left| \frac{1}{\delta} \int_0^\delta \int_{\Omega \times D} M \widehat{\psi}_{\kappa,\alpha} \phi d\mathfrak{q} d\mathfrak{x} dt - \int_{\Omega \times D} M \widehat{\psi}_0 \phi d\mathfrak{q} d\mathfrak{x} \right| \leq C \delta^{\frac{1}{2}} \|\phi\|_{H^s(\Omega \times D)} \quad \forall \phi \in H^s(\Omega \times D),$$

where we have recalled that $H^s(\Omega \times D) \hookrightarrow W^{1,\infty}(\Omega \times D)$. Thus, noting the weak continuity result $\widehat{\psi}_{\kappa,\alpha} \in C_w([0, T]; L_M^1(\Omega \times D))$ established above, it follows on passing to the limit $\delta \rightarrow 0_+$ that

$$\int_{\Omega \times D} M \widehat{\psi}_{\kappa,\alpha}(0) \phi d\mathfrak{q} d\mathfrak{x} = \int_{\Omega \times D} M \widehat{\psi}_0 \phi d\mathfrak{q} d\mathfrak{x} \quad \forall \phi \in H^s(\Omega \times D).$$

Hence, we have $\widehat{\psi}_{\kappa,\alpha}(\cdot, \cdot, 0) = \widehat{\psi}_0$ in $L_M^1(\Omega \times D)$.

It remains to prove the inequality (4.119). For $t' \in (0, T]$ fixed, let $n = n(t', \Delta t)$ be a positive integer such that $0 \leq (n-1)\Delta t < t' \leq n\Delta t \leq T$. It follows from (4.21) and (4.4b), on noting that

the interval $(0, t']$ is contained in $(0, t_n]$, that

$$\begin{aligned}
& \frac{1}{2} \int_{\Omega} \rho_{\kappa, \alpha, L}^{\Delta t, +}(t') |u_{\kappa, \alpha, L}^{\Delta t, +}(t')|^2 dx + \int_{\Omega} P_{\kappa}(\rho_{\kappa, \alpha, L}^{\Delta t, +}(t')) dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}(t')) dq dx \\
& + \alpha \kappa \int_0^{t'} \left[\|\nabla_x [(\rho_{\kappa, \alpha, L}^{\Delta t, +})^2]\|_{L^2(\Omega)}^2 + \frac{4}{\Gamma} \|\nabla_x [(\rho_{\kappa, \alpha, L}^{\Delta t, +})^{\frac{\Gamma}{2}}]\|_{L^2(\Omega)}^2 \right] dt + \mu^S c_0 \int_0^{t'} \|u_{\kappa, \alpha, L}^{\Delta t, +}\|_{H^1(\Omega)}^2 dt \\
& + k \int_0^{t'} \int_{\Omega \times D} M \left[\frac{a_0}{2\lambda} \left| \nabla_q \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \right|^2 + 2\varepsilon \left| \nabla_x \sqrt{\widehat{\psi}_{\kappa, \alpha, L}^{\Delta t, +}} \right|^2 \right] dq dx dt \\
& + 3 \|\varrho_{\kappa, \alpha, L}^{\Delta t, +}(t')\|_{L^2(\Omega)}^2 + 23\varepsilon \int_0^{t'} \|\nabla_x \varrho_{\kappa, \alpha, L}^{\Delta t, +}\|_{L^2(\Omega)}^2 dt \\
& \leq e^{t_n} \left[\frac{1}{2} \int_{\Omega} \rho^0 |u_0|^2 dx + \int_{\Omega} P_{\kappa}(\rho^0) dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_0) dq dx \right. \\
& \quad \left. + 3 \int_D \left(\int_{\Omega} M \widehat{\psi}_0 dx \right)^2 dx + \frac{1}{2} \int_0^{t_n} \|f\|_{L^\infty(\Omega)}^2 dt \int_{\Omega} \rho^0 dx \right] \\
& \leq C,
\end{aligned}
\tag{4.126}$$

where $C \in \mathbb{R}_{>0}$ is independent of α and κ . Clearly $n = n(t', \Delta t) \geq \frac{t'}{\Delta t} \rightarrow \infty$ as $\Delta t \rightarrow 0_+$. Since $t' \in (t_{n-1}, t_n]$ and $t_n - t_{n-1} = \Delta t$, we deduce that as $\Delta t \rightarrow 0_+$ (and, hence, $n = n(t', \Delta t) \rightarrow \infty$) both t_{n-1} and t_n converge to t' ; hence

$$(4.127) \quad e^{t_n} \rightarrow e^{t'} \quad \text{and} \quad \int_0^{t_n} \|f\|_{L^\infty(\Omega)}^2 dt \rightarrow \int_0^{t'} \|f\|_{L^\infty(\Omega)}^2 dt, \quad \text{as } \Delta t \rightarrow 0_+.$$

We multiply (4.126) by any nonnegative $\eta \in C_0^\infty(0, T)$, integrate over $(0, T)$, and pass to the limit $\Delta t \rightarrow 0_+$ (and $L \rightarrow \infty$) in the resulting inequality. It then follows from (4.97d), (4.53c), (4.101g); weak lower-semicontinuity, via the weak convergence results (4.42a,c), (4.44), (4.101a,b) and (4.103a); and (4.127), that we obtain the inequality (4.119) multiplied by η and integrated over $(0, T)$. The desired result (4.119) then follows from the well-known variant of du Bois-Reymond's lemma according to which, if $\phi \in L^1(0, T)$, then

$$(4.128) \quad \int_0^T \phi \eta dt \geq 0 \quad \forall \eta \in C_0^\infty(0, T) \text{ with } \eta \geq 0 \text{ on } (0, T) \quad \Rightarrow \quad \phi \geq 0 \quad \text{a.e. in } (0, T).$$

□

5. EXISTENCE OF A SOLUTION TO (P_κ)

It follows from the bounds on $\varrho_{\kappa, \alpha}$ in (4.119), similarly to (4.26) and (4.27), that

$$(5.1) \quad \|\varrho_{\kappa, \alpha}\|_{L^\infty(0, T; L^2(\Omega))} + \|\varrho_{\kappa, \alpha}\|_{L^{\frac{2(d+2)}{d}}(0, T; L^{\frac{2d}{d-1}}(\Omega))} + \|\varrho_{\kappa, \alpha}\|_{L^4(0, T; L^{\frac{2d}{d-1}}(\Omega))} \leq C,$$

where throughout this section C is a generic positive constant, independent of α . Hence, we deduce from (5.1), (4.24) and (4.119), similarly to (4.28b), that

$$(5.2) \quad \|\tau_1(M \widehat{\psi}_{\kappa, \alpha})\|_{L^2(0, T; L^{\frac{4}{3}}(\Omega))} + \|\tau_1(M \widehat{\psi}_{\kappa, \alpha})\|_{L^{\frac{4(d+2)}{3d+4}}(0, T; L^{\frac{4}{3}}(\Omega))} \leq C.$$

Similarly to (4.94), it follows from (5.1) and (4.119) that

$$(5.3) \quad \left\| M \frac{\partial \widehat{\psi}_{\kappa, \alpha}}{\partial t} \right\|_{L^2(0, T; H^s(\Omega \times D)')} \leq C,$$

where $s > 1 + \frac{1}{2}(K+1)d$. We have the following analogue of Lemma 4.8.

Lemma 5.1. *There exist functions*

$$(5.4a) \quad u_{\kappa} \in L^2(0, T; H_0^1(\Omega)) \quad \text{and} \quad \widehat{\psi}_{\kappa} \in L^v(0, T; Z_1) \cap H^1(0, T; M^{-1}(H^s(\Omega \times D))'),$$

where $v \in [1, \infty)$ and $s > 1 + \frac{1}{2}(K+1)d$, with finite relative entropy and Fisher information,

$$(5.4b) \quad \mathcal{F}(\widehat{\psi}_{\kappa}) \in L^\infty(0, T; L_M^1(\Omega \times D)) \quad \text{and} \quad \sqrt{\widehat{\psi}_{\kappa}} \in L^2(0, T; H_M^1(\Omega \times D)),$$

and a subsequence of $\{(\rho_{\kappa, \alpha}, \underline{u}_{\kappa, \alpha}, \widehat{\psi}_{\kappa, \alpha})\}_{\alpha > 0}$ such that, as $\alpha \rightarrow 0_+$,

$$(5.5) \quad u_{\kappa, \alpha} \rightarrow u_{\kappa} \quad \text{weakly in } L^2(0, T; H_0^1(\Omega)),$$

and

$$(5.6a) \quad M^{\frac{1}{2}} \nabla_x \sqrt{\widehat{\psi}_{\kappa, \alpha}} \rightarrow M^{\frac{1}{2}} \nabla_x \sqrt{\widehat{\psi}_{\kappa}} \quad \text{weakly in } L^2(0, T; L^2(\Omega \times D)),$$

$$(5.6b) \quad M^{\frac{1}{2}} \nabla_q \sqrt{\widehat{\psi}_{\kappa, \alpha}} \rightarrow M^{\frac{1}{2}} \nabla_q \sqrt{\widehat{\psi}_{\kappa}} \quad \text{weakly in } L^2(0, T; L^2(\Omega \times D)),$$

$$(5.6c) \quad M \frac{\partial \widehat{\psi}_{\kappa, \alpha}}{\partial t} \rightarrow M \frac{\partial \widehat{\psi}_{\kappa}}{\partial t} \quad \text{weakly in } L^2(0, T; H^s(\Omega \times D)'),$$

$$(5.6d) \quad \widehat{\psi}_{\kappa, \alpha} \rightarrow \widehat{\psi}_{\kappa} \quad \text{strongly in } L^v(0, T; L_M^1(\Omega \times D)),$$

$$(5.6e) \quad \tau(M \widehat{\psi}_{\kappa, \alpha}) \rightarrow \tau(M \widehat{\psi}_{\kappa}) \quad \text{strongly in } L^r(\Omega_T),$$

where $r \in [1, \frac{4(d+2)}{3d+4})$, and, for a.a. $t \in (0, T)$,

$$(5.6f) \quad \int_{\Omega \times D} M(q) \mathcal{F}(\widehat{\psi}_{\kappa}(x, q, t)) \, dq \, dx \leq \liminf_{\alpha \rightarrow 0_+} \int_{\Omega \times D} M(q) \mathcal{F}(\widehat{\psi}_{\kappa, \alpha}(x, q, t)) \, dq \, dx.$$

In addition, we have that

$$(5.7) \quad \varrho_{\kappa} := \int_D M \widehat{\psi}_{\kappa} \, dq \in L^\infty(0, T; L^2(\Omega)) \cap L^2(0, T; H^1(\Omega)),$$

and, as $\alpha \rightarrow 0_+$,

$$(5.8a) \quad \varrho_{\kappa, \alpha} \rightarrow \varrho_{\kappa} \quad \text{weakly-}\star \text{ in } L^\infty(0, T; L^2(\Omega)), \quad \text{weakly in } L^2(0, T; H^1(\Omega)),$$

$$(5.8b) \quad \varrho_{\kappa, \alpha} \rightarrow \varrho_{\kappa} \quad \text{strongly in } L^{\frac{5\varsigma}{3(\varsigma-1)}}(0, T; L^5(\Omega)),$$

for any $\varsigma \in (1, 6)$.

Proof. The convergence result (5.5) and the first result in (5.4a) follow immediately from the bound on $\underline{u}_{\kappa, \alpha}$ in (4.119). The remainder of the results follow from the bounds on $\widehat{\psi}_{\kappa, \alpha}$ and $\varrho_{\kappa, \alpha}$ in (4.119) in the same way as the results of Lemma 4.8. \square

We have the following analogue of Lemmas 4.2 and 4.4.

Lemma 5.2. *Let $\Gamma \geq 8$; then, there exists a $C \in \mathbb{R}_{>0}$, independent of α , such that*

$$(5.9a) \quad \begin{aligned} & \|\rho_{\kappa, \alpha}\|_{L^\infty(0, T; L^\Gamma(\Omega))} + \|\underline{u}_{\kappa, \alpha}\|_{L^2(0, T; H^1(\Omega))} + \left\| \sqrt{\rho_{\kappa, \alpha}} \underline{u}_{\kappa, \alpha} \right\|_{L^\infty(0, T; L^2(\Omega))} \\ & + \|\rho_{\kappa, \alpha} \underline{u}_{\kappa, \alpha}\|_{L^\infty(0, T; L^{\frac{2\Gamma}{\Gamma+1}}(\Omega))} + \|\rho_{\kappa, \alpha} \underline{u}_{\kappa, \alpha}\|_{L^2(0, T; L^{\frac{6\Gamma}{\Gamma+6}}(\Omega))} \\ & + \|\rho_{\kappa, \alpha} \underline{u}_{\kappa, \alpha}\|_{L^{\frac{10\Gamma-6}{3(\Gamma+1)}}(\Omega_T)} + \left\| \rho_{\kappa, \alpha} |\underline{u}_{\kappa, \alpha}|^2 \right\|_{L^2(0, T; L^{\frac{6\Gamma}{4\Gamma+3}}(\Omega))} \leq C, \end{aligned}$$

$$(5.9b) \quad \begin{aligned} & \sqrt{\alpha} \|\nabla_x \rho_{\kappa, \alpha}\|_{L^2(\Omega_T)} + \alpha \|\nabla_x \rho_{\kappa, \alpha}\|_{L^{\frac{10\Gamma-6}{3(\Gamma+1)}}(\Omega_T)} + \alpha \|(\nabla_x \rho_{\kappa, \alpha} \cdot \nabla_x) \underline{u}_{\kappa, \alpha}\|_{L^{\frac{5\Gamma-3}{4\Gamma}}(\Omega_T)} \\ & + \alpha \|\nabla_x \rho_{\kappa, \alpha} \otimes \underline{u}_{\kappa, \alpha}\|_{L^{\frac{5\Gamma-3}{4\Gamma}}(\Omega_T)} \leq C, \end{aligned}$$

$$(5.9c) \quad \left\| \frac{\partial \rho_{\kappa, \alpha}}{\partial t} \right\|_{L^2(0, T; H^1(\Omega)')} \leq C.$$

Hence, there exists a function $\rho_\kappa \in C_w([0, T]; L_{\geq 0}^\Gamma(\Omega)) \cap H^1(0, T; H^1(\Omega)')$, and for a further subsequence of the subsequence of Lemma 5.1, it follows that, as $\alpha \rightarrow 0_+$,

$$(5.10a) \quad \rho_{\kappa, \alpha} \rightarrow \rho_\kappa \quad \text{in } C_w([0, T]; L^\Gamma(\Omega)), \quad \text{weakly in } H^1(0, T; H^1(\Omega)'),$$

$$(5.10b) \quad \rho_{\kappa, \alpha} \rightarrow \rho_\kappa \quad \text{strongly in } L^2(0, T; H^1(\Omega)'),$$

$$(5.10c) \quad \alpha \nabla_x \rho_{\kappa, \alpha} \rightarrow 0 \quad \text{strongly in } L^r(\Omega_T), \quad r \in [1, \frac{10\Gamma-6}{3(\Gamma+1)}),$$

$$(5.10d) \quad \alpha (\nabla_x \rho_{\kappa, \alpha} \cdot \nabla_x) u_{\kappa, \alpha} \rightarrow 0 \quad \text{weakly in } L^{\frac{5\Gamma-3}{4\Gamma}}(\Omega_T),$$

$$(5.10e) \quad \alpha \nabla_x \rho_{\kappa, \alpha} \otimes u_{\kappa, \alpha} \rightarrow 0 \quad \text{weakly in } L^{\frac{5\Gamma-3}{4\Gamma}}(\Omega_T), \quad \text{strongly in } L^1(0, T; L^{\frac{3}{2}}(\Omega)),$$

and, for any nonnegative $\eta \in C[0, T]$,

$$(5.10f) \quad \int_0^T \left(\int_\Omega P_\kappa(\rho_\kappa) \, dx \right) \eta \, dt \leq \liminf_{\alpha \rightarrow 0_+} \int_0^T \left(\int_\Omega P_\kappa(\rho_{\kappa, \alpha}) \, dx \right) \eta \, dt.$$

Proof. The first three bounds in (5.9a) follow immediately from (4.119). The last four bounds in (5.9a) follow, similarly to (4.84a,b), from the first two bounds in (5.9a).

On recalling (4.41), we choose $\eta = \rho_{\kappa, \alpha}$ in (4.45) to obtain, on noting (5.9a), that

$$(5.11) \quad \begin{aligned} \frac{1}{2} \|\rho_{\kappa, \alpha}(\cdot, T)\|_{L^2(\Omega)}^2 + \alpha \|\nabla_x \rho_{\kappa, \alpha}\|_{L^2(\Omega_T)}^2 &= \frac{1}{2} \left[\|\rho^0\|_{L^2(\Omega)}^2 - \int_0^T \int_\Omega (\nabla_x \cdot u_{\kappa, \alpha}) \rho_{\kappa, \alpha}^2 \, dx \right] \\ &\leq C \left[1 + \|u_{\kappa, \alpha}\|_{L^2(0, T; H^1(\Omega))} \|\rho_{\kappa, \alpha}\|_{L^4(0, T; L^4(\Omega))}^2 \right] \leq C. \end{aligned}$$

Hence the first bound (5.9b). On noting the sixth bound in (5.9a) and recalling from (3.4) that $\partial\Omega \in C^{2, \theta}$, $\theta \in (0, 1)$, and $\rho^0 \in L^\infty(\Omega)$ satisfying (3.18a), we can now apply the parabolic regularity result, Lemma 7.38 in Novotný & Straškraba [58] (or Lemma G.2 in Appendix G), to (4.118a) to obtain that the solution $\rho_{\kappa, \alpha}$ satisfies the second bound in (5.9b). The third and fourth bounds in (5.9b) follow from the second bounds in (5.9a,b), similarly to (4.70). The bound (5.9c) follows immediately from (4.118a), the fifth bound in (5.9a) and the first bound in (5.9b).

The convergence results (5.10a,b) follow immediately from the first bound in (5.9a,c), (3.14a,b) and (3.11). The first two bounds in (5.9b) and (3.1) yield the desired result (5.10c). The convergence result (5.10d) and the first result in (5.10e) follow from the final two bounds in (5.9b), the second bound in (5.9a) and (5.10c). The second result in (5.10e) also follows from the second bound in (5.9a) and (5.10c) with $r = 2$ on noting that, for all $\eta_1 \in L^2(\Omega_T)$ and $\eta_2 \in L^2(0, T; H^1(\Omega))$,

$$(5.12) \quad \|\eta_1 \eta_2\|_{L^1(0, T; L^{\frac{3}{2}}(\Omega))} \leq \|\eta_1\|_{L^2(\Omega_T)} \|\eta_2\|_{L^2(0, T; L^6(\Omega))} \leq C \|\eta_1\|_{L^2(\Omega_T)} \|\eta_2\|_{L^2(0, T; H^1(\Omega))}.$$

Finally, the result (5.10f) follows, similarly to (4.66), from (5.10a) and the convexity of P_κ . \square

Next, we set $v = \frac{10\Gamma-6}{3(\Gamma+1)} \geq \frac{74}{27}$, which appears in Lemma 5.2, and $s' = \frac{5\Gamma-3}{\Gamma-3} \geq 5$ so that (4.91) holds. Then, similarly to (4.92), it follows from (4.118b), (4.91), $W^{1,4}(\Omega) \hookrightarrow L^\infty(\Omega)$, $H^1(\Omega) \hookrightarrow L^6(\Omega)$, (5.9a,b), (5.1), on noting that $\frac{2d}{d-1} > \frac{8}{3}$, (5.2), and (3.4) that, for any $w \in L^{s'}(0, T; W_0^{1,4}(\Omega))$,

$$(5.13) \quad \begin{aligned} &\left| \int_0^T \left\langle \frac{\partial(\rho_{\kappa, \alpha} u_{\kappa, \alpha})}{\partial t}, w \right\rangle_{W_0^{1,4}(\Omega)} \, dt - \int_0^T \int_\Omega p_\kappa(\rho_{\kappa, \alpha}) \nabla_x \cdot w \, dx \, dt \right| \\ &\leq C \left[\alpha \|\nabla_x \rho_{\kappa, \alpha}\|_{L^v(\Omega_T)} + \|\rho_{\kappa, \alpha} u_{\kappa, \alpha}\|_{L^v(\Omega_T)} + 1 \right] \|u_{\kappa, \alpha}\|_{L^2(0, T; H^1(\Omega))} \|w\|_{L^{s'}(0, T; W^{1,4}(\Omega))} \\ &\quad + C \left[\|\tau_1(M \hat{\psi}_{\kappa, \alpha})\|_{L^2(0, T; L^{\frac{4}{3}}(\Omega))} + \|\varrho_{\kappa, \alpha}\|_{L^4(0, T; L^{\frac{2d}{d-1}}(\Omega))}^2 \right] \|w\|_{L^2(0, T; W^{1,4}(\Omega))} \\ &\quad + \|\rho_{\kappa, \alpha}\|_{L^\infty(0, T; L^2(\Omega))} \|f\|_{L^2(0, T; L^\infty(\Omega))} \|w\|_{L^2(\Omega_T)} \\ &\leq C \|w\|_{L^{s'}(0, T; W^{1,4}(\Omega))}. \end{aligned}$$

For $r, s \in (1, \infty)$, let

$$(5.14) \quad L_0^r(\Omega) := \{\zeta \in L^r(\Omega) : \int_{\Omega} \zeta \, d\tilde{x} = 0\}, \quad \tilde{E}^{r,s}(\Omega) := \{w \in L^r(\Omega) : \nabla_x \cdot w \in L^s(\Omega)\}$$

$$\text{and } \tilde{E}_0^{r,s}(\Omega) := \{w \in \tilde{E}^{r,s}(\Omega) : w \cdot \tilde{n} = 0 \text{ on } \partial\Omega\}.$$

The equality $w \cdot \tilde{n} = 0$ on $\partial\Omega$ should be understood in the sense of traces of Sobolev functions, with equality in $W^{1-\frac{v}{v'}, v'}(\partial\Omega)'$, where $\frac{1}{v} + \frac{1}{v'} = 1$ and $v = \min\{r, s\}$; cf. Lemma 3.10 in [58].

We now introduce the Bogovskiĭ operator $\mathcal{B} : L_0^r(\Omega) \rightarrow \tilde{W}_0^{1,r}(\Omega)$, $r \in (1, \infty)$, such that

$$(5.15) \quad \int_{\Omega} (\nabla_x \cdot \mathcal{B}(\zeta) - \zeta) \eta \, d\tilde{x} = 0 \quad \forall \eta \in L^{\frac{r}{r-1}}(\Omega);$$

which satisfies

$$(5.16a) \quad \|\mathcal{B}(\zeta)\|_{W^{1,r}(\Omega)} \leq C \|\zeta\|_{L^r(\Omega)} \quad \forall \zeta \in L_0^r(\Omega),$$

$$(5.16b) \quad \|\mathcal{B}(\nabla_x \cdot w)\|_{L^r(\Omega)} \leq C \|w\|_{L^r(\Omega)} \quad \forall w \in \tilde{E}_0^{r,s}(\Omega),$$

see Lemma 3.17 in Novotný & Straškraba [58] (or Lemma C.1 in Appendix C).

Lemma 5.3. *There exists a $C(\alpha) \in \mathbb{R}_{>0}$ such that*

$$(5.17) \quad \|\nabla_x \cdot (\rho_{\kappa,\alpha} u_{\kappa,\alpha})\|_{L^s(\Omega_T)} + \left\| \frac{\partial \rho_{\kappa,\alpha}}{\partial t} \right\|_{L^s(\Omega_T)} + \|\Delta_x \rho_{\kappa,\alpha}\|_{L^s(\Omega_T)} \leq C(\alpha),$$

where $s = \frac{5\Gamma-3}{4\Gamma}$. In addition, there exists a $C \in \mathbb{R}_{>0}$, independent of α , such that

$$(5.18) \quad \|\rho_{\kappa,\alpha}\|_{L^{\Gamma+1}(\Omega_T)} \leq C.$$

Proof. We prove (5.17), similarly to (4.67). The first bound in (5.17) follows from (5.9b). On noting (4.71) and the first bound in (5.17), we can now apply the parabolic regularity result, Lemma 7.37 in Novotný & Straškraba [58] (or Lemma G.1 in Appendix G), to (4.118a) to obtain that the solution $\rho_{\kappa,\alpha}$ satisfies the last two bounds in (4.67). It follows from (5.17), (5.9a,b) and (4.118a) that

$$(5.19) \quad \frac{\partial \rho_{\kappa,\alpha}}{\partial t} = \nabla_x \cdot \left(\alpha \nabla_x \rho_{\kappa,\alpha} - \rho_{\kappa,\alpha} u_{\kappa,\alpha} \right) \in L^s(\Omega_T)$$

$$\text{with } \left(\alpha \nabla_x \rho_{\kappa,\alpha} - \rho_{\kappa,\alpha} u_{\kappa,\alpha} \right) \cdot \tilde{n} = 0 \quad \text{on } \partial\Omega \times (0, T),$$

and hence, on recalling (5.14) and that $s = \frac{5\Gamma-3}{4\Gamma} < \frac{10\Gamma-6}{3(\Gamma+1)} = r$, we have that

$$(5.20) \quad \alpha \nabla_x \rho_{\kappa,\alpha} - \rho_{\kappa,\alpha} u_{\kappa,\alpha} \in L^s(0, T; \tilde{E}_0^{r,s}(\Omega)).$$

Therefore (5.19), (5.20), (5.16b) and (5.9a,b) yield that

$$(5.21) \quad \left\| \mathcal{B} \left(\frac{\partial \rho_{\kappa,\alpha}}{\partial t} \right) \right\|_{L^s(0,T;L^r(\Omega))} = \|\mathcal{B}(\nabla_x \cdot (\alpha \nabla_x \rho_{\kappa,\alpha} - \rho_{\kappa,\alpha} u_{\kappa,\alpha}))\|_{L^s(0,T;L^r(\Omega))}$$

$$\leq C \|\alpha \nabla_x \rho_{\kappa,\alpha} - \rho_{\kappa,\alpha} u_{\kappa,\alpha}\|_{L^s(0,T;L^r(\Omega))} \leq C.$$

On recalling the notation used in (3.88), then, similarly to (4.15), we obtain, on choosing $\eta = 1$ in (4.118a) and noting (3.18a), that, for all $t \in [0, T]$,

$$(5.22) \quad 0 \leq f \rho_{\kappa,\alpha}(t) = f \rho^0 \leq \|\rho_0\|_{L^\infty(\Omega)}.$$

We now choose $w = \eta \mathcal{B}((I - f)\rho_{\kappa,\alpha})$ in (5.13), where $\eta \in C_0^\infty(0, T)$, to obtain, on noting (2.3), (1.3), (5.22), (5.16a), (5.21) and (5.9a) that, for $s' = \frac{5\Gamma-3}{\Gamma-3}$,

$$\begin{aligned}
& \left| \int_0^T \eta \int_\Omega (c_p \rho_{\kappa,\alpha}^{\gamma+1} + \kappa (\rho_{\kappa,\alpha}^5 + \rho_{\kappa,\alpha}^{\Gamma+1})) \, dx \, dt \right| \\
& \leq C \|\eta\|_{L^\infty(0,T)} \left[\|f \rho_{\kappa,\alpha}\|_{L^\infty(0,T)} \|p_\kappa(\rho_{\kappa,\alpha})\|_{L^1(\Omega_T)} + \|\mathcal{B}((I - f)\rho_{\kappa,\alpha})\|_{L^{s'}(0,T;W^{1,4}(\Omega))} \right] \\
& \quad + \left| \int_0^T \int_\Omega \rho_{\kappa,\alpha} u_{\kappa,\alpha} \cdot \left[\frac{d\eta}{dt} \mathcal{B}((I - f)\rho_{\kappa,\alpha}) + \eta \mathcal{B}\left(\frac{\partial \rho_{\kappa,\alpha}}{\partial t}\right) \right] \, dx \, dt \right| \\
& \leq C \|\eta\|_{L^\infty(0,T)} \left[\|\rho_{\kappa,\alpha}\|_{L^\Gamma(\Omega_T)}^\Gamma + \|\rho_{\kappa,\alpha}\|_{L^{s'}(0,T;L^4(\Omega))} \right] \\
& \quad + \|\rho_{\kappa,\alpha} u_{\kappa,\alpha}\|_{L^\infty(0,T;L^{\frac{2\Gamma}{\Gamma+1}}(\Omega))} \left\| \frac{d\eta}{dt} \right\|_{L^1(0,T)} \|\mathcal{B}((I - f)\rho_{\kappa,\alpha})\|_{L^\infty(0,T;L^{\frac{2\Gamma}{\Gamma+1}}(\Omega))} \\
& \quad + \|\rho_{\kappa,\alpha} u_{\kappa,\alpha}\|_{L^\infty(0,T;L^{\frac{2\Gamma}{\Gamma+1}}(\Omega))} \|\eta\|_{L^\infty(0,T)} \left\| \mathcal{B}\left(\frac{\partial \rho_{\kappa,\alpha}}{\partial t}\right) \right\|_{L^1(0,T;L^{\frac{2\Gamma}{\Gamma+1}}(\Omega))} \\
& \leq C \left[\|\eta\|_{L^\infty(0,T)} + \left\| \frac{d\eta}{dt} \right\|_{L^1(0,T)} \right].
\end{aligned} \tag{5.23}$$

We now consider (5.23) with $\eta = \eta_m \in C_0^\infty(0, T)$, $m \in \mathbb{N}$, where $\eta_m \in [0, 1]$ with $\eta_m(t) = 1$ for $t \in [\frac{1}{m}, T - \frac{1}{m}]$ and $\|\frac{d\eta_m}{dt}\|_{L^\infty(0,T)} \leq 2m$ yielding $\|\frac{d\eta_m}{dt}\|_{L^1(0,T)} \leq 4$. As $\eta_m \rightarrow 1$ pointwise in $(0, T)$, as $m \rightarrow \infty$, we obtain the desired result (5.18). \square

We have the following analogue of Lemmas 4.5 and 4.7.

Lemma 5.4. *Let $\Gamma \geq 8$; then, there exists a $C \in \mathbb{R}_{>0}$, independent of α , such that*

$$(5.24) \quad \left\| \frac{\partial(\rho_{\kappa,\alpha} u_{\kappa,\alpha})}{\partial t} \right\|_{L^{\frac{\Gamma+1}{\Gamma}}(0,T;W_0^{1,\Gamma+1}(\Omega)')} \leq C.$$

Hence, for a further subsequence of the subsequence of Lemma 5.2, it follows that, as $\alpha \rightarrow 0_+$,

$$(5.25a) \quad \rho_{\kappa,\alpha} u_{\kappa,\alpha} \rightharpoonup \rho_\kappa u_\kappa \quad \text{weakly in } L^{\frac{10\Gamma-6}{3(\Gamma+1)}}(\Omega_T), \quad \text{weakly in } W^{1,\frac{\Gamma+1}{\Gamma}}(0,T;W_0^{1,\Gamma+1}(\Omega)'),$$

$$(5.25b) \quad \rho_{\kappa,\alpha} u_{\kappa,\alpha} \rightharpoonup \rho_\kappa u_\kappa \quad \text{in } C_w([0,T];L^{\frac{2\Gamma}{\Gamma+1}}(\Omega)), \quad \text{strongly in } L^2(0,T;H^1(\Omega)'),$$

$$(5.25c) \quad \rho_{\kappa,\alpha} u_{\kappa,\alpha} \otimes u_{\kappa,\alpha} \rightharpoonup \rho_\kappa u_\kappa \otimes u_\kappa \quad \text{weakly in } L^2(0,T;L^{\frac{6\Gamma}{4\Gamma+3}}(\Omega)),$$

$$(5.25d) \quad \rho_{\kappa,\alpha} \rightarrow \rho_\kappa \quad \text{weakly in } L^{\Gamma+1}(\Omega_T),$$

$$(5.25e) \quad p_\kappa(\rho_{\kappa,\alpha}) \rightarrow \overline{p_\kappa(\rho_\kappa)} \quad \text{weakly in } L^{\frac{\Gamma+1}{\Gamma}}(\Omega_T),$$

where $\overline{p_\kappa(\rho_\kappa)} \in L^{\frac{\Gamma+1}{\Gamma}}_{\geq 0}(\Omega_T)$ remains to be identified.

Proof. We deduce from (5.13), (5.18) and as $4 < s' < \Gamma + 1$ that

$$(5.26) \quad \left| \int_0^T \left\langle \frac{\partial(\rho_{\kappa,\alpha} u_{\kappa,\alpha})}{\partial t}, w \right\rangle_{W_0^{1,\Gamma+1}(\Omega)} \, dt \right| \leq C \|w\|_{L^{\Gamma+1}(0,T;W^{1,\Gamma+1}(\Omega))} \quad \forall w \in L^{\Gamma+1}(0,T;W_0^{1,\Gamma+1}(\Omega)),$$

and hence the desired result (5.24).

The results (5.25a–c) follow similarly to (4.97a–e) from (5.9a), (5.24), (3.14a,b), (3.11), (5.10b) and (5.5). The results (5.25d–f) follow immediately from (5.18), (2.3) and (1.3). \square

We have the following analogue of Theorem 4.1.

Lemma 5.5. *The triple $(\rho_\kappa, \underline{u}_\kappa, \widehat{\psi}_\kappa)$, defined as in Lemmas 5.1 and 5.2, satisfies*

$$(5.27a) \quad \int_0^T \left\langle \frac{\partial \rho_\kappa}{\partial t}, \eta \right\rangle_{H^1(\Omega)} dt - \int_0^T \int_\Omega \rho_\kappa \underline{u}_\kappa \cdot \nabla_x \eta \, dx \, dt = 0 \quad \forall \eta \in L^2(0, T; H^1(\Omega)),$$

with $\rho_\kappa(\cdot, 0) = \rho_0(\cdot)$,

$$(5.27b) \quad \begin{aligned} & \int_0^T \left\langle \frac{\partial(\rho_\kappa \underline{u}_\kappa)}{\partial t}, w \right\rangle_{W_0^{1, \Gamma+1}(\Omega)} dt + \int_0^T \int_\Omega \left[S(\underline{u}_\kappa) - \rho_\kappa \underline{u}_\kappa \otimes \underline{u}_\kappa - \overline{p_\kappa(\rho_\kappa)} I \right] : \nabla_x w \, dx \, dt \\ & = \int_0^T \int_\Omega \left[\rho_\kappa f \cdot w - \left(\tau_1(M \widehat{\psi}_\kappa) - \mathfrak{z} \varrho_\kappa^2 I \right) : \nabla_x w \right] dx \, dt \\ & \quad \forall w \in L^{\Gamma+1}(0, T; W_0^{1, \Gamma+1}(\Omega)), \end{aligned}$$

with $(\rho_\kappa \underline{u}_\kappa)(\cdot, 0) = (\rho_0 \underline{u}_0)(\cdot)$, and

$$(5.27c) \quad \begin{aligned} & \int_0^T \left\langle M \frac{\partial \widehat{\psi}_\kappa}{\partial t}, \varphi \right\rangle_{H^s(\Omega \times D)} dt + \frac{1}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \int_0^T \int_{\Omega \times D} M \nabla_{q_j} \widehat{\psi}_\kappa \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt \\ & + \int_0^T \int_{\Omega \times D} M \left[\varepsilon \nabla_x \widehat{\psi}_\kappa - \underline{u}_\kappa \widehat{\psi}_\kappa \right] \cdot \nabla_x \varphi \, dq \, dx \, dt \\ & - \int_0^T \int_{\Omega \times D} M \sum_{i=1}^K \left[\sigma(\underline{u}_\kappa) q_i \right] \widehat{\psi}_\kappa \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt = 0 \quad \forall \varphi \in L^2(0, T; H^s(\Omega \times D)), \end{aligned}$$

with $\widehat{\psi}_\kappa(\cdot, 0) = \widehat{\psi}_0(\cdot)$ and $s > 1 + \frac{1}{2}(K+1)d$.

In addition, the triple $(\rho_\kappa, \underline{u}_\kappa, \widehat{\psi}_\kappa)$ satisfies, for a.a. $t' \in (0, T)$,

$$(5.28) \quad \begin{aligned} & \frac{1}{2} \int_\Omega \rho_\kappa(t') |\underline{u}_\kappa(t')|^2 \, dx + \int_\Omega P_\kappa(\rho_\kappa(t')) \, dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_\kappa(t')) \, dq \, dx \\ & + \mu^S c_0 \int_0^{t'} \|\underline{u}_\kappa\|_{H^1(\Omega)}^2 \, dt + k \int_0^{t'} \int_{\Omega \times D} M \left[\frac{a_0}{2\lambda} \left| \nabla_q \sqrt{\widehat{\psi}_\kappa} \right|^2 + 2\varepsilon \left| \nabla_x \sqrt{\widehat{\psi}_\kappa} \right|^2 \right] \, dq \, dx \, dt \\ & + \mathfrak{z} \|\varrho_\kappa(t')\|_{L^2(\Omega)}^2 + 2\mathfrak{z} \varepsilon \int_0^{t'} \|\nabla_x \varrho_\kappa\|_{L^2(\Omega)}^2 \, dt \\ & \leq e^{t'} \left[\frac{1}{2} \int_\Omega \rho_0 |\underline{u}_0|^2 \, dx + \int_\Omega P_\kappa(\rho_0) \, dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_0) \, dq \, dx \right. \\ & \quad \left. + \mathfrak{z} \int_\Omega \left(\int_D M \widehat{\psi}_0 \, dq \right)^2 \, dx + \frac{1}{2} \int_0^{t'} \|f\|_{L^\infty(\Omega)}^2 \, dt \int_\Omega \rho_0 \, dx \right] \\ & \leq C, \end{aligned}$$

where $C \in \mathbb{R}_{>0}$ is independent of κ .

Proof. Passing to the limit $\alpha \rightarrow 0_+$ for the subsequence of Lemma 5.4 in (4.118a) yields (5.27a) subject to the stated initial condition, on noting (5.10a,c), (5.25a) and (3.18b).

Similarly to the proof of (4.118b), passing to the limit $\alpha \rightarrow 0_+$ for the subsequence of Lemma 5.4 in (4.118b) yields (5.27b) subject to the stated initial condition, on noting (5.5), (5.10a,d,e), (5.25a-c,e), (5.6e), (5.8b) and (3.18b). Similarly to the proof of (4.118c), passing to the limit $\alpha \rightarrow 0_+$ for the subsequence of Lemma 5.4 in (4.118c) yields (5.27c) subject to the stated initial condition, on noting (5.5), (5.6a-d), (5.7) and (3.2). Similarly to the proof of (4.119), we deduce (5.28) from (4.119) using the results (5.25c), (5.10f), (5.6a,b,f), (5.5), (5.8a) and (3.18c). \square

Finally, to obtain the complete analogue of Theorem 4.1, we have to identify $\overline{p_\kappa(\rho_\kappa)}$, which appears in (5.27b) and (5.25e), by establishing that $\overline{p_\kappa(\rho_\kappa)} = p_\kappa(\rho_\kappa)$. Due to the presence of the extra stress term in the momentum equation, we require a modification of the effective viscous

flux compactness result, Proposition 7.36 in Novotný & Straškraba [58]. Such results require pseudodifferential operators identified via the Fourier transform \mathfrak{F} . We briefly recall the key ideas, and refer to Section 4.4.1 in [58] for the details.

With

$$(5.29) \quad \mathfrak{S}(\mathbb{R}^d) := \left\{ \eta \in C^\infty(\mathbb{R}^d) : \sup_{\tilde{x} \in \mathbb{R}^d} \left| x_1^{\lambda_1} \cdots x_d^{\lambda_d} \frac{\partial^{|\lambda|} \eta}{\partial x_1^{\lambda_1} \cdots \partial x_d^{\lambda_d}} \right| \leq C(|\zeta|, |\lambda|) \quad \forall \zeta, \lambda \in \mathbb{N}^d \right\},$$

the space of smooth rapidly decreasing (complex-valued) functions, we introduce the Fourier transform $\mathfrak{F} : \mathfrak{S}(\mathbb{R}^d) \rightarrow \mathfrak{S}(\mathbb{R}^d)$, and its inverse $\mathfrak{F}^{-1} : \mathfrak{S}(\mathbb{R}^d) \rightarrow \mathfrak{S}(\mathbb{R}^d)$, defined by

$$(5.30) \quad [\mathfrak{F}(\eta)](\tilde{y}) = \frac{1}{(2\pi)^{\frac{d}{2}}} \int_{\mathbb{R}^d} e^{-i\tilde{x} \cdot \tilde{y}} \eta(\tilde{x}) \, d\tilde{x} \quad \text{and} \quad [\mathfrak{F}^{-1}(\eta)](\tilde{x}) = \frac{1}{(2\pi)^{\frac{d}{2}}} \int_{\mathbb{R}^d} e^{i\tilde{x} \cdot \tilde{y}} \eta(\tilde{y}) \, d\tilde{y}.$$

These are extended to $\mathfrak{F}, \mathfrak{F}^{-1} : \mathfrak{S}(\mathbb{R}^d)' \rightarrow \mathfrak{S}(\mathbb{R}^d)'$, where $\mathfrak{S}(\mathbb{R}^d)'$, the dual of $\mathfrak{S}(\mathbb{R}^d)$, is the space of tempered distributions, via

$$(5.31) \quad \langle \mathfrak{F}(\eta), \xi \rangle_{\mathfrak{S}(\mathbb{R}^d)} = \langle \eta, \mathfrak{F}(\xi) \rangle_{\mathfrak{S}(\mathbb{R}^d)} \quad \text{and} \quad \langle \mathfrak{F}^{-1}(\eta), \xi \rangle_{\mathfrak{S}(\mathbb{R}^d)} = \langle \eta, \mathfrak{F}^{-1}(\xi) \rangle_{\mathfrak{S}(\mathbb{R}^d)} \quad \forall \xi \in \mathfrak{S}(\mathbb{R}^d).$$

We now introduce the inverse divergence operator $\mathcal{A}_j : \mathfrak{S}(\mathbb{R}^d) \rightarrow \mathfrak{S}(\mathbb{R}^d)'$, $j = 1, \dots, d$, such that

$$(5.32) \quad \mathcal{A}_j(\eta) = -\mathfrak{F}^{-1} \left(\frac{i y_j}{|y|^2} [\mathfrak{F}(\eta)](\tilde{y}) \right).$$

It follows from Theorems 1.55 and 1.57 in [58] (or Lemmas B.1 and B.2 in Appendix B) and Sobolev embedding that, for $j = 1, \dots, d$,

$$(5.33a) \quad \|\nabla_x \mathcal{A}_j(\eta)\|_{L^r(\mathbb{R}^d)} \leq C(r) \|\eta\|_{L^r(\mathbb{R}^d)} \quad \forall \eta \in \mathfrak{S}(\mathbb{R}^d), \quad r \in (1, \infty),$$

$$(5.33b) \quad \|\mathcal{A}_j(\eta)\|_{L^{\frac{dr}{d-r}}(\mathbb{R}^d)} \leq C(r) \|\eta\|_{L^r(\mathbb{R}^d)} \quad \forall \eta \in \mathfrak{S}(\mathbb{R}^d), \quad r \in (1, d).$$

Hence, we deduce from (5.33a,b) that \mathcal{A}_j can be extended to $\mathcal{A}_j : L^r(\mathbb{R}^d) \rightarrow D^{1,r}(\mathbb{R}^d)$ for $r \in (1, \infty)$, $j = 1, \dots, d$, where $D^{1,r}(\mathbb{R}^d)$ is a homogeneous Sobolev space; see Section 1.3.6 in [58] (or Appendix A here). In addition, by duality, \mathcal{A}_j can be extended to $\mathcal{A}_j : D^{1,r}(\mathbb{R}^d)' \rightarrow D^{1,r}(\mathbb{R}^d)$ for $r \in (1, \infty)$, $j = 1, \dots, d$, see (4.4.4) in [58]. Moreover, as $\mathcal{A}_j(\eta)$ is real, for a real-valued function η , and from the Parseval–Plancherel formula we have, for all $\eta \in L^r(\mathbb{R}^d)$ and $\xi \in L^{\frac{r}{r-1}}(\mathbb{R}^d)$, $r \in (1, \infty)$, having compact support that

$$(5.34) \quad \int_{\mathbb{R}^d} \mathcal{A}_j(\eta) \xi \, d\tilde{x} = - \int_{\mathbb{R}^d} \eta \mathcal{A}_j(\xi) \, d\tilde{x}, \quad j = 1, \dots, d.$$

Finally, we introduce the so-called Riesz operator $\mathcal{R}_{kj} : L^r(\mathbb{R}^d) \rightarrow L^r(\mathbb{R}^d)$, $r \in (1, \infty)$, defined by

$$(5.35) \quad \mathcal{R}_{kj}(\eta) = \frac{\partial}{\partial x_k} \mathcal{A}_j(\eta), \quad j, k = 1, \dots, d.$$

We note for all $\eta \in L^r(\mathbb{R}^d)$ and $\xi \in L^{\frac{r}{r-1}}(\mathbb{R}^d)$, $r \in (1, \infty)$, that

$$(5.36a) \quad \sum_{j=1}^d \mathcal{R}_{jj}(\eta) = \sum_{j=1}^d \frac{\partial}{\partial x_j} \mathcal{A}_j(\eta) = \eta,$$

$$(5.36b) \quad \mathcal{R}_{kj}(\eta) = \mathcal{R}_{jk}(\eta) \quad \text{and} \quad \int_{\mathbb{R}^d} \mathcal{R}_{jk}(\eta) \xi \, d\tilde{x} = \int_{\mathbb{R}^d} \eta \mathcal{R}_{jk}(\xi) \, d\tilde{x}, \quad j, k = 1, \dots, d.$$

Below we use the notation $\mathcal{A}(\cdot)$ and $\mathcal{R}(\cdot)$ with components $\mathcal{A}_i(\cdot)$ and $\mathcal{R}_{ij}(\cdot)$, $i, j = 1, \dots, d$, respectively. We shall adopt the convention that whenever any of these operators is applied to a function or a distribution that has been defined on Ω only, it is tacitly understood that the function or distribution in question has been extended by 0 from Ω to the whole of \mathbb{R}^d .

We now have the following modification of Proposition 7.36 in Novotný & Straškraba [58], which is adequate for our purposes.

Lemma 5.6. *Given $\{(g_n, u_n, m_n, p_n, \tau_n, f_n, F_n)\}_{n \in \mathbb{N}}$, we assume for any $\zeta \in C_0^\infty(\Omega)$ that, as $n \rightarrow \infty$,*

$$(5.37a) \quad g_n \rightarrow g \quad \text{in } C_w([0, T]; L^q(\Omega)), \quad \text{weakly } (-\star) \text{ in } L^\omega(\Omega_T),$$

$$(5.37b) \quad u_n \rightarrow u \quad \text{weakly in } L^2(0, T; H_0^1(\Omega)),$$

$$(5.37c) \quad m_n \rightarrow m \quad \text{in } C_w([0, T]; L^z(\Omega)),$$

$$(5.37d) \quad p_n \rightarrow p \quad \text{weakly in } L^r(\Omega_T),$$

$$(5.37e) \quad \tau_n \rightarrow \tau \quad \text{strongly in } L^1(0, T; L^{\frac{q}{q-1}}(\Omega)),$$

$$(5.37f) \quad f_n \rightarrow f \quad \text{weakly in } L^2(0, T; H^1(\Omega)'),$$

$$(5.37g) \quad \mathcal{A}(\zeta f_n) \rightarrow \mathcal{A}(\zeta f) \quad \text{strongly in } L^2(0, T; L^{\frac{z}{z-1}}(\Omega)),$$

$$(5.37h) \quad F_n \rightarrow F \quad \text{weakly in } L^s(\Omega_T),$$

where $q \in (d, \infty)$, $r, s \in (1, \infty)$, $\omega \in [\max\{2, \frac{r}{r-1}\}, \infty]$ and $z \in (\frac{6q}{5q-6}, \infty)$.

In addition, suppose that

$$(5.38a) \quad \frac{\partial g_n}{\partial t} + \nabla_x \cdot (u_n g_n) = f_n \quad \text{in } C_0^\infty(\Omega_T)',$$

$$(5.38b) \quad \begin{aligned} \frac{\partial m_n}{\partial t} + \nabla_x \cdot (m_n \otimes u_n) - \mu \Delta_x u_n - (\mu + \lambda) \nabla_x (\nabla_x \cdot u_n) + \nabla_x p_n \\ = F_n + \nabla_x \cdot \tau_n \end{aligned} \quad \text{in } C_0^\infty(\Omega_T)'.$$

Then it follows that, for any $\zeta \in C_0^\infty(\Omega)$ and $\eta \in C_0^\infty(0, T)$,

$$(5.39) \quad \lim_{n \rightarrow \infty} \int_0^T \eta \left(\int_\Omega \zeta g_n [p_n - (2\mu + \lambda) \nabla_x \cdot u_n] dx \right) dt = \int_0^T \eta \left(\int_\Omega \zeta g [p - (2\mu + \lambda) \nabla_x \cdot u] dx \right) dt.$$

Proof. We adapt the proof of Proposition 7.36 in [58], by just pointing out the key differences. As $q > d$, then q^* , the Sobolev conjugate of q in the notation (1.3.64) of [58], is such that $q^* = \infty$. Hence our restrictions on r, s, ω and z satisfy the restrictions of Proposition 7.36 in [58]. With any $\tilde{\zeta} \in C_0^\infty(\Omega)$, it follows from (5.38a) and properties (5.34)–(5.36a,b) of \mathcal{A}_j and \mathcal{R}_{kj} that, for $i = 1, \dots, d$,

$$(5.40) \quad \frac{\partial}{\partial t} \mathcal{A}_i(\tilde{\zeta} g_n) + \sum_{j=1}^d \mathcal{R}_{ij}(\tilde{\zeta} g_n u_n^j) = \mathcal{A}_i(\tilde{\zeta} f_n) + \mathcal{A}_i(g_n u_n \cdot \nabla_x \tilde{\zeta}) \quad \text{in } C_0^\infty(\Omega_T)',$$

where we adopt the notation u_n^j for the j^{th} component of u_n . With any $\zeta, \tilde{\zeta} \in C_0^\infty(\Omega)$ and $\eta \in C_0^\infty(0, T)$, we now consider $\eta \zeta \mathcal{A}(\tilde{\zeta} g_n)$ as a test function for (5.38b). It follows from (5.37a) and (5.33a,b) that $\mathcal{A}(\tilde{\zeta} g_n) \in L^\infty(0, T; W^{1,q}(\Omega)) \cap L^\omega(0, T; W^{1,\omega}(\Omega))$, and hence $\mathcal{A}(\tilde{\zeta} g_n) \in L^\infty(\Omega_T)$ as $q > d$. Similarly to (5.9a),

$$g_n u_n \in L^2(0, T; L^{\frac{6q}{q+6}}(\Omega))$$

and

$$m_n \otimes u_n \in L^2(0, T; L^{\frac{6z}{z+6}}(\Omega)).$$

As $z \in (\frac{6q}{5q-6}, \infty)$, and therefore $\frac{z}{z-1} \in (1, \frac{6q}{q+6})$ and $\frac{6z}{z+6} \in (\frac{q}{q-1}, 6)$, it follows from (5.40), (5.33a,b), (5.35) and (5.37g) that

$$\frac{\partial}{\partial t} \mathcal{A}(\tilde{\zeta} g_n) \in L^2(0, T; L^{\frac{z}{z-1}}(\Omega)).$$

Noting the above and (5.37b–e,h), we see that $\eta \zeta \mathcal{A}(\tilde{\zeta} g_n)$ is a valid test function for (5.38b), and we obtain, on using integration by parts several times and properties (5.34)–(5.36a,b) of \mathcal{A}_j and

\mathcal{R}_{kj} , that

$$\begin{aligned}
& \int_0^T \eta \left(\int_{\Omega} \zeta \tilde{\zeta} g_n [p_n - (2\mu + \lambda) \nabla_x \cdot \tilde{u}_n] \tilde{\zeta} \, dx \right) dt \\
&= \mu \int_0^T \eta \left(\int_{\Omega} \left(\nabla_x \tilde{u}_n : \mathcal{A}(\tilde{\zeta} g_n) \otimes \nabla_x \zeta - \tilde{u}_n \otimes \nabla_x \zeta : \mathcal{R}(\tilde{\zeta} g_n) + \tilde{\zeta} g_n \tilde{u}_n \cdot \nabla_x \zeta \right) \tilde{\zeta} \, dx \right) dt \\
&+ \int_0^T \eta \left(\int_{\Omega} \left((\tau_n - m_n \otimes \tilde{u}_n) : \mathcal{A}(\tilde{\zeta} g_n) \otimes \nabla_x \zeta - [p_n - (\mu + \lambda) \nabla_x \cdot \tilde{u}_n] \mathcal{A}(\tilde{\zeta} g_n) \cdot \nabla_x \zeta \right) \tilde{\zeta} \, dx \right) dt \\
&+ \int_0^T \eta \left(\int_{\Omega} \zeta \left((\tau_n - m_n \otimes \tilde{u}_n) : \mathcal{R}(\tilde{\zeta} g_n) - F_n \cdot \mathcal{A}(\tilde{\zeta} g_n) \right) \tilde{\zeta} \, dx \right) dt \\
&- \int_0^T \frac{d\eta}{dt} \left(\int_{\Omega} \zeta m_n \cdot \mathcal{A}(\tilde{\zeta} g_n) \tilde{\zeta} \, dx \right) dt - \int_0^T \eta \left(\int_{\Omega} \zeta m_n \cdot [\mathcal{A}(\tilde{\zeta} f_n) + \mathcal{A}(g_n \tilde{u}_n \cdot \nabla_x \tilde{\zeta})] \tilde{\zeta} \, dx \right) dt \\
&+ \int_0^T \eta \left(\int_{\Omega} \tilde{\zeta} \sum_{i=1}^d \sum_{j=1}^d g_n u_n^j \mathcal{R}_{ij}(\zeta m_n^i) \tilde{\zeta} \, dx \right) dt.
\end{aligned}
\tag{5.41}$$

The equation (5.41) is exactly the same as (7.5.12) in [58], except for the extra τ_n terms and the change of notation.

We will just concentrate on the terms involving τ_n , as the other terms are dealt with as in [58]. It follows from (5.37a), (5.33a,b), see (7.5.18)–(7.5.20) in [58] for the details, that

$$(5.42a) \quad \mathcal{A}(\tilde{\zeta} g_n) \rightarrow \mathcal{A}(\tilde{\zeta} g) \quad \text{weakly in } L^\infty(0, T; W^{1,q}(\Omega)), \quad \text{strongly in } L^v(\Omega_T),$$

$$(5.42b) \quad \mathcal{R}(\tilde{\zeta} g_n) \rightarrow \mathcal{R}(\tilde{\zeta} g) \quad \text{weakly in } L^\infty(0, T; L^q(\Omega)), \quad \text{strongly in } L^v(0, T; H^{-1}(\Omega)),$$

where $v \in [1, \infty)$. It follows from (5.37e) and (5.42a,b) that, as $n \rightarrow \infty$,

$$\begin{aligned}
& \int_0^T \eta \left(\int_{\Omega} \tau_n : [\mathcal{A}(\tilde{\zeta} g_n) \otimes \nabla_x \zeta + \mathcal{R}(\tilde{\zeta} g_n) \zeta] \tilde{\zeta} \, dx \right) dt \\
& \rightarrow \int_0^T \eta \left(\int_{\Omega} \tau : [\mathcal{A}(\tilde{\zeta} g) \otimes \nabla_x \zeta + \mathcal{R}(\tilde{\zeta} g) \zeta] \tilde{\zeta} \, dx \right) dt.
\end{aligned}
\tag{5.43}$$

Combining (5.43) with the convergence, as $n \rightarrow \infty$, of other terms in (5.41) as in the proof of Proposition 7.36 in [58], which involves the use of the crucial Commutator lemma (Lemma 4.25 in [58], or Lemma D.3 in Appendix D), we obtain that

$$\begin{aligned}
& \lim_{n \rightarrow \infty} \int_0^T \eta \left(\int_{\Omega} \zeta \tilde{\zeta} g_n [p_n - (2\mu + \lambda) \nabla_x \cdot \tilde{u}_n] \tilde{\zeta} \, dx \right) dt \\
&= \mu \int_0^T \eta \left(\int_{\Omega} \left(\nabla_x \tilde{u} : \mathcal{A}(\tilde{\zeta} g) \otimes \nabla_x \zeta - \tilde{u} \otimes \nabla_x \zeta : \mathcal{R}(\tilde{\zeta} g) + \tilde{\zeta} g \tilde{u} \cdot \nabla_x \zeta \right) \tilde{\zeta} \, dx \right) dt \\
&+ \int_0^T \eta \left(\int_{\Omega} \left((\tau - m \otimes \tilde{u}) : \mathcal{A}(\tilde{\zeta} g) \otimes \nabla_x \zeta - [p - (\mu + \lambda) \nabla_x \cdot \tilde{u}] \mathcal{A}(\tilde{\zeta} g) \cdot \nabla_x \zeta \right) \tilde{\zeta} \, dx \right) dt \\
&+ \int_0^T \eta \left(\int_{\Omega} \zeta \left((\tau - m \otimes \tilde{u}) : \mathcal{R}(\tilde{\zeta} g) - F \cdot \mathcal{A}(\tilde{\zeta} g) \right) \tilde{\zeta} \, dx \right) dt \\
&- \int_0^T \frac{d\eta}{dt} \left(\int_{\Omega} \zeta m \cdot \mathcal{A}(\tilde{\zeta} g) \tilde{\zeta} \, dx \right) dt - \int_0^T \eta \left(\int_{\Omega} \zeta m \cdot [\mathcal{A}(\tilde{\zeta} f) + \mathcal{A}(g \tilde{u} \cdot \nabla_x \tilde{\zeta})] \tilde{\zeta} \, dx \right) dt \\
&+ \int_0^T \eta \left(\int_{\Omega} \tilde{\zeta} \sum_{i=1}^d \sum_{j=1}^d g u^j \mathcal{R}_{ij}(\zeta m^i) \tilde{\zeta} \, dx \right) dt,
\end{aligned}
\tag{5.44}$$

which is exactly the same as (7.5.25) in [58], except for the extra τ terms and the change of notation.

In addition, the equations (5.38a,b) are exactly the same as in (7.5.7)–(7.5.8) in [58] except for the extra τ_n term. One can use (5.37a–h) to pass to the limit $n \rightarrow \infty$ in (5.38a,b) to obtain

$$(5.45a) \quad \frac{\partial g}{\partial t} + \nabla_x \cdot (\tilde{u} g) = f \quad \text{in } C_0^\infty(\Omega_T)',$$

$$(5.45b) \quad \frac{\partial m}{\partial t} + \nabla_x \cdot (\tilde{m} \otimes \tilde{u}) - \mu \Delta_x \tilde{u} - (\mu + \lambda) \nabla_x (\nabla_x \cdot \tilde{u}) + \nabla_x p = \tilde{F} + \nabla_x \cdot \tilde{\tau} \quad \text{in } C_0^\infty(\Omega_T)',$$

see [58] for details. Clearly, the τ_n term in (5.38b) is easily dealt with using (5.37e). Similarly to (5.40), we deduce that $\eta \zeta \mathcal{A}(\tilde{\zeta} g)$ is a valid test function for (5.45b), for any $\zeta, \tilde{\zeta} \in C_0^\infty(\Omega)$ and $\eta \in C_0^\infty(0, T)$, and we obtain (5.41) without the subscript n . Combining this with (5.44), we deduce that, for any $\zeta, \tilde{\zeta} \in C_0^\infty(\Omega)$ and $\eta \in C_0^\infty(0, T)$,

$$(5.46) \quad \lim_{n \rightarrow \infty} \int_0^T \eta \left(\int_\Omega \zeta \tilde{\zeta} g_n [p_n - (2\mu + \lambda) \nabla_x \cdot \tilde{u}_n] dx \right) dt = \int_0^T \eta \left(\int_\Omega \zeta \tilde{\zeta} g [p - (2\mu + \lambda) \nabla_x \cdot \tilde{u}] dx \right) dt.$$

Hence we arrive at (5.39) by taking $\tilde{\zeta} \in C_0^\infty(\Omega)$ such that $\tilde{\zeta} \equiv 1$ on the support of ζ . \square

We need also the following variation of Lemma 5.6 for later use in Section 6.

Corollary 5.1. *The results of Lemma 5.6 hold with the assumptions (5.37f,g) replaced by*

$$(5.47) \quad f_n \rightarrow f \quad \text{weakly in } L^2(\Omega_T), \quad \text{as } n \rightarrow \infty.$$

Proof. One can still pass to the limit $n \rightarrow \infty$ in (5.38a) to obtain (5.45a) using (5.47) in place of (5.37f,g). We deduce from (5.47), (5.33b) and (5.34) that

$$(5.48) \quad \mathcal{A}(\zeta f_n) \rightarrow \mathcal{A}(\zeta f) \quad \text{weakly in } L^2(0, T; L^6(\Omega)), \quad \text{as } n \rightarrow \infty.$$

As $\frac{z}{z-1} < \frac{6q}{q+6} < 6$, (5.48) ensures that one can still conclude from (5.40) that $\eta \zeta \mathcal{A}(\tilde{\zeta} g_n)$ is a valid test function for (5.38b). Similarly, one can deduce that $\eta \zeta \mathcal{A}(\tilde{\zeta} g)$ is a valid test function for (5.45b). The only other place where (5.37f,g) are used in the proof of Lemma 5.6 is in dealing with the term involving f_n on the right-hand side of (5.41); that is, the term

$$(5.49) \quad - \int_0^T \eta \left(\int_\Omega \zeta \tilde{m}_n \cdot \mathcal{A}(\tilde{\zeta} f_n) dx \right) dt = \int_0^T \eta \left(\int_\Omega \tilde{\zeta} f_n \sum_{i=1}^d \mathcal{A}_i(\zeta \tilde{m}_n^i) dx \right) dt,$$

where we have noted (5.34). Similarly to (5.42a), it follows from (5.37c), (5.33a,b) and Sobolev embedding, as $z > \frac{6}{5}$, that

$$(5.50) \quad \mathcal{A}(\tilde{\zeta} \tilde{m}_n) \rightarrow \mathcal{A}(\tilde{\zeta} \tilde{m}) \quad \text{weakly in } L^\infty(0, T; W^{1,z}(\Omega)), \quad \text{strongly in } L^v(0, T; L^3(\Omega)),$$

where $v \in [1, \infty)$. Therefore, (5.50), (5.47) and (5.34) imply that we can pass to the limit $n \rightarrow \infty$ in (5.49) to obtain the term involving f on the right-hand side of (5.44). \square

In order to identify $\overline{p_\kappa(\rho_\kappa)}$ in (5.27b) and (5.25e), we now apply Lemma 5.6 with (5.38a,b) being (4.118a,b) so that $\mu = \frac{\mu^S}{2}$ and $\lambda = \mu^B - \frac{\mu^S}{d}$, $g_n = \rho_{\kappa,\alpha}$, $\tilde{u}_n = \underline{u}_{\kappa,\alpha}$, $\tilde{m}_n = \rho_{\kappa,\alpha} \underline{u}_{\kappa,\alpha}$, $p_n = p_\kappa(\rho_{\kappa,\alpha})$, $\tau_n = \tau(M \hat{\psi}_{\kappa,\alpha}) + \frac{\alpha}{2} (\underline{u}_{\kappa,\alpha} \otimes \nabla_x \rho_{\kappa,\alpha})$, $f_n = \alpha \Delta_x \rho_{\kappa,\alpha}$ and $\tilde{F}_n = \rho_{\kappa,\alpha} \tilde{f} - \frac{\alpha}{2} (\nabla_x \rho_{\kappa,\alpha} \cdot \nabla_x) \underline{u}_{\kappa,\alpha}$. With $\{(\rho_{\kappa,\alpha}, \underline{u}_{\kappa,\alpha}, \hat{\psi}_{\kappa,\alpha})\}_{\alpha>0}$ being the subsequence (not indicated) of Lemma 5.4, we have that (5.37a–d) hold with $g = \rho_\kappa$, $\underline{u} = \underline{u}_\kappa$, $\tilde{m} = \rho_\kappa \underline{u}_\kappa$ and $p = \overline{p_\kappa(\rho_\kappa)}$, and $q = \Gamma$, $\omega = \Gamma + 1$, $z = \frac{2\Gamma}{\Gamma+1}$ and $r = \frac{\Gamma+1}{\Gamma}$ on recalling (5.10a), (5.25b,d,e) and (5.5). We note that $\omega = \Gamma + 1 = \frac{r}{r-1} > 2$ and, as $\Gamma \geq 8$, $z = \frac{2\Gamma}{\Gamma+1} \geq \frac{16}{9} > \frac{24}{17} \geq \frac{6\Gamma}{5\Gamma-6} = \frac{6q}{5q-6}$. Hence, the constraints on q, r, ω and z hold. The results (5.37e,h) hold with $\tau = \tau(M \hat{\psi}_\kappa)$, $\tilde{F} = \rho_\kappa \tilde{f}$ and $s = \frac{5\Gamma-3}{4\Gamma}$, on recalling (5.6e), (5.10d,e) and (5.25d), and noting that $\frac{4(d+2)}{3d+4} \geq \frac{20}{13} > \frac{8}{7} \geq \frac{q}{q-1}$. Finally, the results (5.37f,g) hold with $f = 0$ on recalling (5.10c) and the properties of \mathcal{A} and \mathcal{R} , on noting that $\frac{z}{z-1} = \frac{2\Gamma}{\Gamma-1} < \frac{10\Gamma-6}{3(\Gamma+1)}$, see (7.9.21)

in [58] for details. Hence, we obtain from (5.39) for the subsequence of Lemma 5.4 that, for all $\zeta \in C_0^\infty(\Omega)$ and $\eta \in C_0^\infty(0, T)$,

$$(5.51) \quad \lim_{\alpha \rightarrow 0+} \int_0^T \eta \left(\int_\Omega \zeta \rho_{\kappa, \alpha} [p_\kappa(\rho_{\kappa, \alpha}) - \mu^* \nabla_x \cdot u_{\kappa, \alpha}] dx \right) dt = \int_0^T \eta \left(\int_\Omega \zeta \rho_\kappa [\overline{p_\kappa(\rho_\kappa)} - \mu^* \nabla_x \cdot u_\kappa] dx \right) dt,$$

where $\mu^* := \frac{(d-1)}{d} \mu^S + \mu^B$. The first two bounds in (5.9a) yield that

$$(5.52) \quad \|\rho_{\kappa, \alpha} \nabla_x \cdot u_{\kappa, \alpha}\|_{L^2(0, T; L^{\frac{2\Gamma}{\Gamma+2}}(\Omega))} \leq C,$$

and hence there exists a $\overline{\rho_\kappa \nabla_x \cdot u_\kappa} \in L^2(0, T; L^{\frac{2\Gamma}{\Gamma+2}}(\Omega))$ such that for a subsequence (not indicated)

$$(5.53) \quad \rho_{\kappa, \alpha} \nabla_x \cdot u_{\kappa, \alpha} \rightharpoonup \overline{\rho_\kappa \nabla_x \cdot u_\kappa} \quad \text{weakly in } L^2(0, T; L^{\frac{2\Gamma}{\Gamma+2}}(\Omega)), \quad \text{as } \alpha \rightarrow 0+.$$

It follows from the monotonicity of $p_\kappa(\cdot)$ that

$$(5.54) \quad \begin{aligned} \rho_{\kappa, \alpha} p_\kappa(\rho_{\kappa, \alpha}) &= (\rho_{\kappa, \alpha} - \rho_\kappa) (p_\kappa(\rho_{\kappa, \alpha}) - p_\kappa(\rho_\kappa)) + (\rho_{\kappa, \alpha} - \rho_\kappa) p_\kappa(\rho_\kappa) + \rho_\kappa p_\kappa(\rho_{\kappa, \alpha}) \\ &\geq (\rho_{\kappa, \alpha} - \rho_\kappa) p_\kappa(\rho_\kappa) + \rho_\kappa p_\kappa(\rho_{\kappa, \alpha}) \end{aligned} \quad \text{a.e. in } \Omega_T.$$

We deduce from (5.51), (5.54), (5.53) and (5.25d,e) that for all nonnegative $\zeta \in C_0^\infty(\Omega)$ and $\eta \in C_0^\infty(0, T)$,

$$(5.55) \quad \int_0^T \eta \left(\int_\Omega \zeta [\overline{\rho_\kappa \nabla_x \cdot u_\kappa} - \rho_\kappa \nabla_x \cdot u_\kappa] dx \right) dt \geq 0 \quad \Rightarrow \quad \overline{\rho_\kappa \nabla_x \cdot u_\kappa} \geq \rho_\kappa \nabla_x \cdot u_\kappa \quad \text{a.e. in } \Omega_T,$$

where we have noted (4.128) with $(0, T)$ replaced by Ω_T for the final implication.

Next, we introduce $\mathfrak{L}(s) = s \log s$ for $s \in [0, \infty)$. On recalling (3.18a,b), we have for a subsequence (not indicated), via Lebesgue's dominated convergence theorem, that

$$(5.56) \quad \lim_{\alpha \rightarrow 0+} \int_\Omega \mathfrak{L}(\rho^0) dx = \int_\Omega \mathfrak{L}(\rho_0) dx.$$

We can now follow the discussion in Section 7.9.3 in Novotný & Straškraba [58] to deduce that $p_\kappa(\rho_\kappa) = p_\kappa(\rho_\kappa)$. For the benefit of the reader, we briefly outline the argument. One deduces from (5.27a) as $\rho_\kappa \in C_w([0, T]; L_{\geq 0}^\Gamma(\Omega))$, via renormalization and noting that $s \mathfrak{L}'(s) - \mathfrak{L}(s) = s$ for $s \in [0, \infty)$, that, for any $t' \in (0, T]$,

$$(5.57) \quad \int_\Omega [\mathfrak{L}(\rho_\kappa)(t') - \mathfrak{L}(\rho_0)] dx = - \int_0^{t'} \int_\Omega \rho_\kappa \nabla_x \cdot u_\kappa dx dt.$$

On noting (5.17), one can choose, similarly to (4.8), $\eta = \chi_{[0, t']} [\log(\rho_{\kappa, \alpha} + \varsigma) - 1]$, where $\varsigma \in \mathbb{R}_{>0}$, in (4.118a), and on passing to the limit $\varsigma \rightarrow 0+$ obtain that, for any $t' \in (0, T]$,

$$(5.58) \quad \int_\Omega [\mathfrak{L}(\rho_{\kappa, \alpha})(t') - \mathfrak{L}(\rho^0)] dx \leq - \int_0^{t'} \int_\Omega \rho_{\kappa, \alpha} \nabla_x \cdot u_{\kappa, \alpha} dx dt.$$

Subtracting (5.57) from (5.58), and passing to the limit $\alpha \rightarrow 0+$, one deduces from (5.56), (5.53) and (5.55) that, for any $t' \in (0, T]$,

$$(5.59) \quad \int_\Omega [\overline{\mathfrak{L}(\rho_\kappa)(t')} - \mathfrak{L}(\rho_\kappa)(t')] dx \leq \int_0^{t'} \int_\Omega [\rho_\kappa \nabla_x \cdot u_\kappa - \overline{\rho_\kappa \nabla_x \cdot u_\kappa}] dx dt \leq 0,$$

where, on noting (5.10a),

$$(5.60) \quad \mathfrak{L}(\rho_{\kappa, \alpha})(t') \rightarrow \overline{\mathfrak{L}(\rho_\kappa)(t')} \quad \text{weakly in } L^r(\Omega), \quad \text{for any } r \in [1, \Gamma), \quad \text{as } \alpha \rightarrow 0+.$$

As $\mathfrak{L}(s)$ is continuous and convex for $s \in [0, \infty)$, it follows from (5.10a) and (5.60), see e.g. Corollary 3.33 in [58] (or Lemma D.1 in Appendix D), that $\overline{\mathfrak{L}(\rho_\kappa)(t')} \geq \mathfrak{L}(\rho_\kappa)(t')$ a.e. in Ω for any $t' \in (0, T]$. Hence, we deduce from (5.59) that $\overline{\mathfrak{L}(\rho_\kappa)(t')} = \mathfrak{L}(\rho_\kappa)(t')$ a.e. in Ω for any $t' \in (0, T]$. Therefore, on applying Lemma 3.34 in [58] (or Lemma D.2 in Appendix D), we conclude from the above, (5.60)

and (5.10a) that $\rho_{\kappa,\alpha}(t) \rightarrow \rho_{\kappa}(t)$ strongly in $L^1(\Omega)$ for any $t \in (0, T]$, as $\alpha \rightarrow 0_+$. It immediately follows from this, (5.18), (3.1) and (5.25e), on possibly extracting a further subsequence (not indicated), that, as $\alpha \rightarrow 0_+$,

$$(5.61a) \quad \rho_{\kappa,\alpha} \rightarrow \rho_{\kappa} \quad \text{strongly in } L^r(\Omega_T), \quad \text{for any } r \in [1, \Gamma + 1),$$

$$(5.61b) \quad p_{\kappa}(\rho_{\kappa,\alpha}) \rightarrow p_{\kappa}(\rho_{\kappa}) \quad \text{weakly in } L^{\frac{\Gamma+1}{\Gamma}}(\Omega_T), \quad \text{that is, } \overline{p_{\kappa}(\rho_{\kappa,\alpha})} = p_{\kappa}(\rho_{\kappa}).$$

Finally, we have the following complete analogue of Theorem 4.1.

Theorem 5.1. *The triple $(\rho_{\kappa}, \underline{u}_{\kappa}, \widehat{\psi}_{\kappa})$, defined as in Lemmas 5.1 and 5.2, is a global weak solution to problem (P_{κ}) , in the sense that (5.27a,c), with their initial conditions, hold and*

$$(5.62) \quad \begin{aligned} & \int_0^T \left\langle \frac{\partial(\rho_{\kappa} \underline{u}_{\kappa})}{\partial t}, w \right\rangle_{W_0^{1,\Gamma+1}(\Omega)} dt + \int_0^T \int_{\Omega} \left[S(\underline{u}_{\kappa}) - \rho_{\kappa} \underline{u}_{\kappa} \otimes \underline{u}_{\kappa} - p_{\kappa}(\rho_{\kappa}) I \right] : \nabla_x w \, dx \, dt \\ & = \int_0^T \int_{\Omega} \left[\rho_{\kappa} f \cdot w - \left(\tau_1(M \widehat{\psi}_{\kappa}) - \mathfrak{z} \varrho_{\kappa}^2 I \right) : \nabla_x w \right] dx \, dt \\ & \quad \forall w \in L^{\Gamma+1}(0, T; W_0^{1,\Gamma+1}(\Omega)), \end{aligned}$$

with $(\rho_{\kappa} \underline{u}_{\kappa})(\cdot, 0) = (\rho_0 \underline{u}_0)(\cdot)$. In addition, the weak solution $(\rho_{\kappa}, \underline{u}_{\kappa}, \widehat{\psi}_{\kappa})$ satisfies (5.28).

Proof. The results (5.27a,c) and (5.28) have already been established in Lemma 5.5. Equation (5.62) was established in Lemma 5.5 with $p_{\kappa}(\rho_{\kappa})$ replaced by $\overline{p_{\kappa}(\rho_{\kappa})}$, see (5.27b). The desired result (5.62) then follows immediately from (5.27b) and (5.61b). \square

6. EXISTENCE OF A SOLUTION TO (P)

It follows from the bounds on ϱ_{κ} in (5.28), similarly to (4.26) and (4.27), that

$$(6.1) \quad \|\varrho_{\kappa}\|_{L^{\infty}(0,T;L^2(\Omega))} + \|\varrho_{\kappa}\|_{L^{\frac{2(d+2)}{d}}(0,T;L^6(\Omega))} + \|\varrho_{\kappa}\|_{L^4(0,T;L^{\frac{2d}{d-1}}(\Omega))} \leq C,$$

where throughout this section C is a generic positive constant, independent of κ . Hence, we deduce from (6.1), (4.24) and (5.28), similarly to (4.28b), that

$$(6.2) \quad \|\tau_1(M \widehat{\psi}_{\kappa})\|_{L^2(0,T;L^{\frac{4}{3}}(\Omega))} + \|\tau_1(M \widehat{\psi}_{\kappa})\|_{L^{\frac{4(d+2)}{3d+4}}(0,T;L^6(\Omega))} + \|\tau_1(M \widehat{\psi}_{\kappa})\|_{L^{\frac{4}{3}}(0,T;L^{\frac{12}{7}}(\Omega))} \leq C.$$

Similarly to (4.94), it follows from (6.1) and (5.28) that

$$(6.3) \quad \left\| M \frac{\partial \widehat{\psi}_{\kappa}}{\partial t} \right\|_{L^2(0,T;H^s(\Omega \times D)')} \leq C,$$

where $s > 1 + \frac{1}{2}(K+1)d$. We have the following analogue of Lemma 5.1.

Lemma 6.1. *There exist functions*

$$(6.4a) \quad \underline{u} \in L^2(0,T;H_0^1(\Omega)) \quad \text{and} \quad \widehat{\psi} \in L^v(0,T;Z_1) \cap H^1(0,T;M^{-1}(H^s(\Omega \times D))'),$$

where $v \in [1, \infty)$ and $s > 1 + \frac{1}{2}(K+1)d$, with finite relative entropy and Fisher information,

$$(6.4b) \quad \mathcal{F}(\widehat{\psi}) \in L^{\infty}(0,T;L_M^1(\Omega \times D)) \quad \text{and} \quad \sqrt{\widehat{\psi}} \in L^2(0,T;H_M^1(\Omega \times D)),$$

and a subsequence of $\{(\rho_{\kappa}, \underline{u}_{\kappa}, \widehat{\psi}_{\kappa})\}_{\kappa>0}$ such that, as $\kappa \rightarrow 0_+$,

$$(6.5) \quad \underline{u}_{\kappa} \rightarrow \underline{u} \quad \text{weakly in } L^2(0,T;H_0^1(\Omega)),$$

and

$$(6.6a) \quad M^{\frac{1}{2}} \nabla_x \sqrt{\widehat{\psi}_\kappa} \rightarrow M^{\frac{1}{2}} \nabla_x \sqrt{\widehat{\psi}} \quad \text{weakly in } L^2(0, T; L^2(\Omega \times D)),$$

$$(6.6b) \quad M^{\frac{1}{2}} \nabla_q \sqrt{\widehat{\psi}_\kappa} \rightarrow M^{\frac{1}{2}} \nabla_q \sqrt{\widehat{\psi}} \quad \text{weakly in } L^2(0, T; L^2(\Omega \times D)),$$

$$(6.6c) \quad M \frac{\partial \widehat{\psi}_\kappa}{\partial t} \rightarrow M \frac{\partial \widehat{\psi}}{\partial t} \quad \text{weakly in } L^2(0, T; H^s(\Omega \times D)'),$$

$$(6.6d) \quad \widehat{\psi}_\kappa \rightarrow \widehat{\psi} \quad \text{strongly in } L^v(0, T; L_M^1(\Omega \times D)),$$

$$(6.6e) \quad \tau(M \widehat{\psi}_\kappa) \rightarrow \tau(M \widehat{\psi}) \quad \text{strongly in } L^r(\Omega_T),$$

where $r \in [1, \frac{4(d+2)}{3d+4})$, and, for a.a. $t \in (0, T)$,

$$(6.6f) \quad \int_{\Omega \times D} M(q) \mathcal{F}(\widehat{\psi}_\kappa(x, q, t)) \, dq \, dx \leq \liminf_{\kappa \rightarrow 0+} \int_{\Omega \times D} M(q) \mathcal{F}(\widehat{\psi}(x, q, t)) \, dq \, dx.$$

In addition, we have that

$$(6.7) \quad \varrho := \int_D M \widehat{\psi} \, dq \in L^\infty(0, T; L^2(\Omega)) \cap L^2(0, T; H^1(\Omega)),$$

and, as $\kappa \rightarrow 0_+$,

$$(6.8a) \quad \varrho_\kappa \rightarrow \varrho \quad \text{weakly-}\star \text{ in } L^\infty(0, T; L^2(\Omega)), \quad \text{weakly in } L^2(0, T; H^1(\Omega)),$$

$$(6.8b) \quad \varrho_\kappa \rightarrow \varrho \quad \text{strongly in } L^{\frac{5\varsigma}{3(\varsigma-1)}}(0, T; L^\varsigma(\Omega)),$$

for any $\varsigma \in (1, 6)$.

Proof. The convergence result (6.5) and the first result in (6.4a) follow immediately from the bound on u_κ in (5.28). The remainder of the results follow from the bounds on $\widehat{\psi}_\kappa$ and ϱ_κ in (5.28) in the same way as the results of Lemma 4.8. \square

We have the following analogue of Lemma 5.2.

Lemma 6.2. *Let $\Gamma \geq 8$; then, there exists a $C \in \mathbb{R}_{>0}$, independent of κ , such that, for any $\gamma > \frac{3}{2}$ as in (1.3),*

$$(6.9a) \quad \begin{aligned} & \|\rho_\kappa\|_{L^\infty(0, T; L^\gamma(\Omega))} + \|u_\kappa\|_{L^2(0, T; H^1(\Omega))} + \kappa^{\frac{1}{\Gamma}} \|\rho_\kappa\|_{L^\infty(0, T; L^\Gamma(\Omega))} + \left\| \sqrt{\rho_\kappa} u_\kappa \right\|_{L^\infty(0, T; L^2(\Omega))} \\ & + \|\rho_\kappa u_\kappa\|_{L^\infty(0, T; L^{\frac{2\gamma}{\gamma+1}}(\Omega))} + \|\rho_\kappa u_\kappa\|_{L^2(0, T; L^{\frac{6\gamma}{\gamma+6}}(\Omega))} \\ & + \|\rho_\kappa u_\kappa\|_{L^{\frac{10\gamma-6}{3(\gamma+1)}}(\Omega_T)} + \left\| \rho_\kappa |u_\kappa|^2 \right\|_{L^2(0, T; L^{\frac{6\gamma}{4\gamma+3}}(\Omega))} \leq C, \end{aligned}$$

$$(6.9b) \quad \left\| \frac{\partial \rho_\kappa}{\partial t} \right\|_{L^2(0, T; W^{1,6}(\Omega)')} \leq C.$$

Hence, there exists a function $\rho \in C_w([0, T]; L_{\geq 0}^\gamma(\Omega)) \cap H^1(0, T; W^{1,6}(\Omega)'),$ and for a further subsequence of the subsequence of Lemma 6.1, it follows that, as $\kappa \rightarrow 0_+$,

$$(6.10a) \quad \rho_\kappa \rightarrow \rho \quad \text{in } C_w([0, T]; L^\gamma(\Omega)) \quad \text{weakly in } H^1(0, T; W^{1,6}(\Omega)'),$$

$$(6.10b) \quad \rho_\kappa \rightarrow \rho \quad \text{strongly in } L^2(0, T; H^1(\Omega)'),$$

and, for any nonnegative $\eta \in C[0, T]$,

$$(6.10c) \quad \int_0^T \left(\int_\Omega P(\rho) \, dx \right) \eta \, dt \leq \liminf_{\kappa \rightarrow 0+} \int_0^T \left(\int_\Omega P(\rho_\kappa) \, dx \right) \eta \, dt.$$

Proof. The first four bounds in (6.9a) follow immediately from (5.28). The last four bounds in (6.9a) follow, similarly to (4.84a,b), from the first two bounds in (6.9a). The bound (6.9b) follows immediately from (5.27a), the sixth bound in (6.9a), on noting that $\frac{6\gamma}{\gamma+6} > \frac{6}{5}$ as $\gamma > \frac{3}{2}$.

The convergence results (6.10a,b) follow immediately from (6.9a,b), (3.14a,b) and (3.11). The result (6.10c) follows, similarly to (4.66), from (6.10a) and the convexity of P . \square

Similarly to (5.13), it follows from (5.27b), (6.9a), (6.1), (6.2), (3.4), on noting that $\gamma > \frac{3}{2}$, and (3.2) that, for any $w \in L^\infty(0, T; W_0^{1,r}(\Omega))$ with $r = \max\{\Gamma + 1, v\}$ and $v = \max\{\frac{3\gamma}{2\gamma-3}, \frac{12}{5}\}$,

$$\begin{aligned}
& \left| \int_0^T \left\langle \frac{\partial(\rho_\kappa u_\kappa)}{\partial t}, w \right\rangle_{W_0^{1,\Gamma+1}(\Omega)} dt - \int_0^T \int_\Omega p_\kappa(\rho_\kappa) \nabla_x \cdot w \, dx \, dt \right| \\
& \leq C \|\rho_\kappa\|_{L^\infty(0,T;L^\gamma(\Omega))} \|u_\kappa\|_{L^2(0,T;L^6(\Omega))}^2 \|w\|_{L^\infty(0,T;W^{1,\frac{3\gamma}{2\gamma-3}}(\Omega))} \\
& \quad + C \|u_\kappa\|_{L^2(0,T;H^1(\Omega))} \|w\|_{L^2(0,T;H^1(\Omega))} \\
& \quad + C \left[\|\tau_1(M\widehat{\psi}_\kappa)\|_{L^{\frac{4}{3}}(0,T;L^{\frac{12}{7}}(\Omega))} + \|\varrho_\kappa\|_{L^2(0,T;L^6(\Omega))}^2 \right] \|\nabla_x w\|_{L^\infty(0,T;L^{\frac{12}{5}}(\Omega))} \\
& \quad + \|\rho_\kappa\|_{L^\infty(0,T;L^\gamma(\Omega))} \|f\|_{L^2(0,T;L^\infty(\Omega))} \|w\|_{L^2(0,T;L^3(\Omega))} \\
(6.11) \quad & \leq C \|w\|_{L^\infty(0,T;W^{1,v}(\Omega))}.
\end{aligned}$$

We deduce from (6.11) with $w = \eta v$, where $\eta \in C_0^\infty(0, T)$ and $v \in L^\infty(0, T; W_0^{1,v}(\Omega)) \cap H^1(0, T; L^{\frac{6\gamma}{5\gamma-6}}(\Omega))$ with $r = \max\{\Gamma + 1, v\}$ and $v = \max\{\frac{3\gamma}{2\gamma-3}, \frac{12}{5}\}$, on noting (6.9a) and (3.2) as $\frac{2\gamma}{\gamma-1} < 6$ for $\gamma > \frac{3}{2}$, that

$$\begin{aligned}
& \left| \int_0^T \eta \int_\Omega p_\kappa(\rho_\kappa) \nabla_x \cdot v \, dx \, dt \right| \\
& \leq \left| \int_0^T \int_\Omega \rho_\kappa u_\kappa \cdot \frac{\partial(\eta v)}{\partial t} \, dx \, dt \right| + C \|\eta\|_{L^\infty(0,T)} \|v\|_{L^\infty(I;W^{1,v}(\Omega))} \\
& \leq C \|\rho_\kappa u_\kappa\|_{L^\infty(0,T;L^{\frac{2\gamma}{\gamma+1}}(\Omega))} \left\| \frac{d\eta}{dt} \right\|_{L^1(0,T)} \|v\|_{L^\infty(I;L^{\frac{2\gamma}{\gamma-1}}(\Omega))} \\
& \quad + C \|\eta\|_{L^\infty(0,T)} \left[\|\rho_\kappa u_\kappa\|_{L^2(0,T;L^{\frac{6\gamma}{\gamma+6}}(\Omega))} \left\| \frac{\partial v}{\partial t} \right\|_{L^2(I;L^{\frac{6\gamma}{5\gamma-6}}(\Omega))} + \|v\|_{L^\infty(I;W^{1,v}(\Omega))} \right] \\
(6.12) \quad & \leq C \left[\left\| \frac{d\eta}{dt} \right\|_{L^1(0,T)} + \|\eta\|_{L^\infty(0,T)} \right] \left[\|v\|_{L^\infty(I;W^{1,v}(\Omega))} + \left\| \frac{\partial v}{\partial t} \right\|_{L^2(I;L^{\frac{6\gamma}{5\gamma-6}}(\Omega))} \right],
\end{aligned}$$

where $I = \text{supp}(\eta) \subset (0, T)$. With $v = v(\gamma)$ thus defined, let

$$(6.13) \quad \vartheta(\gamma) := \frac{\gamma}{v(\gamma)} = \begin{cases} \frac{2\gamma-3}{3} & \text{for } \frac{3}{2} < \gamma \leq 4, \\ \frac{5}{12}\gamma & \text{for } 4 \leq \gamma. \end{cases}$$

With $\vartheta(\gamma) \in \mathbb{R}_{>0}$ defined as above and $\ell \in \mathbb{N}$, we now introduce $\mathbf{b} : \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$ and $\mathbf{b}_\ell : \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$ such that

$$(6.14) \quad \mathbf{b}(s) := s^\vartheta \quad \text{and} \quad \mathbf{b}_\ell(s) := \begin{cases} \mathbf{b}(s) & \text{for } 0 \leq s \leq \ell, \\ \mathbf{b}(\ell) & \text{for } \ell \leq s. \end{cases}$$

We note from (6.14), (6.9a) and (3.2) that, for $v(\gamma) := \max\{\frac{3\gamma}{2\gamma-3}, \frac{12}{5}\}$ as in (6.12) and $\vartheta(\gamma)$ as in (6.13), we have that

$$(6.15a) \quad \|\mathbf{b}_\ell(\rho_\kappa)\|_{L^\infty(0,T;L^v(\Omega))} \leq \|\rho_\kappa^\vartheta\|_{L^\infty(0,T;L^v(\Omega))} \leq \|\rho_\kappa\|_{L^\infty(0,T;L^\gamma(\Omega))}^\vartheta \leq C,$$

$$(6.15b) \quad \|\mathbf{b}_\ell(\rho_\kappa) u_\kappa\|_{L^2(0,T;L^{\frac{6\gamma}{\gamma+6}}(\Omega))} \leq \|\rho_\kappa\|_{L^\infty(0,T;L^\gamma(\Omega))}^\vartheta \|u_\kappa\|_{L^2(0,T;L^6(\Omega))} \leq C,$$

$$(6.15c) \quad \|\mathbf{b}_\ell(\rho_\kappa) \nabla_x \cdot u_\kappa\|_{L^2(0,T;L^{\frac{2\gamma}{\gamma+2\vartheta}}(\Omega))} \leq \|\rho_\kappa\|_{L^\infty(0,T;L^\gamma(\Omega))}^\vartheta \|u_\kappa\|_{L^2(0,T;H^1(\Omega))} \leq C,$$

where $C \in \mathbb{R}_{>0}$ is independent of κ, ϑ and ℓ .

As $\Gamma > 2$, it follows from (5.27a), on extending ρ_κ and \underline{u}_κ from Ω to \mathbb{R}^d by zero, that

$$(6.16) \quad \frac{\partial \rho_\kappa}{\partial t} + \nabla_x \cdot (\rho_\kappa \underline{u}_\kappa) = 0 \quad \text{in } C_0^\infty(\mathbb{R}^d \times (0, T))',$$

see Lemmas 6.8 in Novotný & Straškraba [58] (or Lemma F.1 in Appendix F). Applying Lemma 6.11 in [58] (or Lemma F.3 in Appendix F) to (6.16), we have the renormalised equation, for any $\ell \in \mathbb{N}$,

$$(6.17) \quad \frac{\partial \mathbf{b}_\ell(\rho_\kappa)}{\partial t} + \nabla_x \cdot (\mathbf{b}_\ell(\rho_\kappa) \underline{u}_\kappa) + (\rho_\kappa (\mathbf{b}_\ell)_+'(\rho_\kappa) - \mathbf{b}_\ell(\rho_\kappa)) \nabla_x \cdot \underline{u}_\kappa = 0 \quad \text{in } C_0^\infty(\mathbb{R}^d \times (0, T))',$$

where $(\mathbf{b}_\ell)_+'(\cdot)$ is the right-derivative of $\mathbf{b}_\ell(\cdot)$ satisfying

$$(6.18) \quad (\mathbf{b}_\ell)_+'(s) = \begin{cases} \mathbf{b}'_\ell(s) & \text{for } 0 \leq s < \ell, \\ 0 & \text{for } \ell \leq s. \end{cases}$$

For any $\delta \in (0, \frac{T}{2})$, we now introduce the Friedrichs mollifier, with respect to the time variable, $\mathcal{S}_\delta : L^1(0, T; L^q(\Omega)) \rightarrow C^\infty(\delta, T - \delta; L^q(\Omega))$, $q \in [1, \infty]$,

$$(6.19) \quad \mathcal{S}_\delta(\eta)(\underline{x}, t) = \frac{1}{\delta} \int_0^T \omega\left(\frac{t-s}{\delta}\right) \eta(\underline{x}, s) \, ds \quad \text{a.e. in } \Omega \times (\delta, T - \delta),$$

where $\omega \in C_0^\infty(\mathbb{R})$, $\omega \geq 0$, $\text{supp}(\omega) \subset (-1, 1)$ and $\int_{\mathbb{R}} \omega \, ds = 1$. It follows from (6.17) and (6.19) that

$$(6.20) \quad \frac{\partial \mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa))}{\partial t} + \nabla_x \cdot \mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa) \underline{u}_\kappa) + \mathcal{S}_\delta([\rho_\kappa (\mathbf{b}_\ell)_+'(\rho_\kappa) - \mathbf{b}_\ell(\rho_\kappa)] \nabla_x \cdot \underline{u}_\kappa) = 0 \quad \text{in } C_0^\infty(\mathbb{R}^d \times (\delta, T - \delta))'.$$

In addition, it follows from (6.19), (6.14), (6.9a), (3.2), (6.18) and (6.20) that

$$(6.21) \quad \begin{aligned} \mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa)) &\in C^\infty(\delta, T - \delta; L^\infty(\mathbb{R}^d)), & \mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa) \underline{u}_\kappa) &\in C^\infty(\delta, T - \delta; L^6(\mathbb{R}^d)), \\ \mathcal{S}_\delta([\rho_\kappa (\mathbf{b}_\ell)_+'(\rho_\kappa) - \mathbf{b}_\ell(\rho_\kappa)] \nabla_x \cdot \underline{u}_\kappa) &\in C^\infty(\delta, T - \delta; L^2(\mathbb{R}^d)), & \nabla_x \cdot [\mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa) \underline{u}_\kappa)] &\in C^\infty(\delta, T - \delta; L^2(\mathbb{R}^d)). \end{aligned}$$

One can deduce from $\underline{u}_\kappa \in L^2(0, T; H_0^1(\Omega))$ and (6.21) that

$$(6.22) \quad \mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa) \underline{u}_\kappa) \in C^\infty(\delta, T - \delta; E_0^{6,2}(\Omega)),$$

where we recall (5.14). We note from (5.16a), (6.19), (6.14) and (6.15a) that $\mathcal{B}([\mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa))]) \in L^\infty(\delta, T - \delta; W_0^{1,r}(\Omega))$, $r \in [1, \infty)$, and, for $v(\gamma) := \max\{\frac{3\gamma}{2\gamma-3}, \frac{12}{5}\}$ as in (6.12), that

$$(6.23a) \quad \begin{aligned} \|\mathcal{B}((I - f)[\mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa))])\|_{L^\infty(\delta, T - \delta; W^{1,v}(\Omega))} &\leq C \|\mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa))\|_{L^\infty(\delta, T - \delta; L^v(\Omega))} \\ &\leq C \|\mathbf{b}_\ell(\rho_\kappa)\|_{L^\infty(0, T; L^v(\Omega))} \leq C, \end{aligned}$$

and from (6.20), Sobolev embedding, (6.22), (5.16a,b), (6.19), (6.14), (6.18) and (6.15b,c) with ϑ as in (6.13) that

$$\begin{aligned}
& \left\| \frac{\partial}{\partial t} \mathcal{B}((I - f)[\mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa))]) \right\|_{L^2(\delta, T-\delta; L^{\frac{6\gamma}{\gamma+6\vartheta}}(\Omega))} \\
& \leq \left\| \mathcal{B}(\nabla_x \cdot [\mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa)) u_\kappa]) \right\|_{L^2(\delta, T-\delta; L^{\frac{6\gamma}{\gamma+6\vartheta}}(\Omega))} \\
& \quad + \left\| \mathcal{B}((I - f)[\mathcal{S}_\delta([\rho_\kappa(\mathbf{b}_\ell)'_+(\rho_\kappa) - \mathbf{b}_\ell(\rho_\kappa)] \nabla_x \cdot u_\kappa]) \right\|_{L^2(\delta, T-\delta; L^{\frac{6\gamma}{\gamma+6\vartheta}}(\Omega))} \\
& \leq \left\| \mathcal{B}(\nabla_x \cdot [\mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa)) u_\kappa]) \right\|_{L^2(\delta, T-\delta; L^{\frac{6\gamma}{\gamma+6\vartheta}}(\Omega))} \\
& \quad + \left\| \mathcal{B}((I - f)[\mathcal{S}_\delta([\rho_\kappa(\mathbf{b}_\ell)'_+(\rho_\kappa) - \mathbf{b}_\ell(\rho_\kappa)] \nabla_x \cdot u_\kappa]) \right\|_{L^2(\delta, T-\delta; W^{1, \frac{2\gamma}{\gamma+2\vartheta}}(\Omega))} \\
& \leq C \left[\left\| \mathbf{b}_\ell(\rho_\kappa) u_\kappa \right\|_{L^2(0, T; L^{\frac{6\gamma}{\gamma+6\vartheta}}(\Omega))} + \left\| [\rho_\kappa(\mathbf{b}_\ell)'_+(\rho_\kappa) - \mathbf{b}_\ell(\rho_\kappa)] \nabla_x \cdot u_\kappa \right\|_{L^2(0, T; L^{\frac{2\gamma}{\gamma+2\vartheta}}(\Omega))} \right] \\
& \leq C \left[\left\| \mathbf{b}_\ell(\rho_\kappa) u_\kappa \right\|_{L^2(0, T; L^{\frac{6\gamma}{\gamma+6\vartheta}}(\Omega))} + \left\| \mathbf{b}_\ell(\rho_\kappa) \nabla_x \cdot u_\kappa \right\|_{L^2(0, T; L^{\frac{2\gamma}{\gamma+2\vartheta}}(\Omega))} \right] \leq C,
\end{aligned}
\tag{6.23b}$$

where $C \in \mathbb{R}_{>0}$ in (6.23a,b) is independent of κ, ϑ, ℓ and δ . We now have the following analogue of Lemma 5.3.

Lemma 6.3. *With $\vartheta(\gamma)$ as defined in (6.13), we have that*

$$\left\| \rho_\kappa \right\|_{L^{\gamma+\vartheta}(\Omega_T)} + \kappa^{\frac{1}{4+\vartheta}} \left\| \rho_\kappa \right\|_{L^{4+\vartheta}(\Omega_T)} + \kappa^{\frac{1}{\Gamma+\vartheta}} \left\| \rho_\kappa \right\|_{L^{\Gamma+\vartheta}(\Omega_T)} \leq C.
\tag{6.24}$$

Proof. For any $\ell \in \mathbb{N}$ and $\delta \in (0, \frac{T}{2})$, we choose $\vartheta = \mathcal{B}((I - f)[\mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa))]) \in L^\infty(\delta, T - \delta, W_0^{1,r}(\Omega))$, any $r \in [1, \infty)$, and $\eta \in C_0^\infty(0, T)$, with $\text{supp}(\eta) \subset (\delta, T - \delta)$, in (6.12) to obtain, on noting (6.23a,b), $\frac{6\gamma}{5\gamma-6} \leq \frac{6\gamma}{\gamma+6\vartheta}$ as $\vartheta \leq \frac{2\gamma}{3} - 1$, (6.9a) and (2.3), that

$$\begin{aligned}
& \left| \int_0^T \eta \int_\Omega p_\kappa(\rho_\kappa) \mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa)) \, dx \, dt \right| \\
& \leq C \left[\left\| \frac{d\eta}{dt} \right\|_{L^1(0, T)} + \|\eta\|_{L^\infty(0, T)} \right] \left[1 + \|p_\kappa(\rho_\kappa)\|_{L^1(\Omega_T)} \|\mathcal{S}_\delta(\mathbf{b}_\ell(\rho_\kappa))\|_{L^\infty(\delta, T-\delta)} \right] \\
& \leq C \left[\left\| \frac{d\eta}{dt} \right\|_{L^1(0, T)} + \|\eta\|_{L^\infty(0, T)} \right].
\end{aligned}
\tag{6.25}$$

We now consider (6.25) with $\eta = \eta_\ell \in C_0^\infty(0, T)$ with $\text{supp}(\eta_\ell) \subset (\frac{1}{\ell}, T - \frac{1}{\ell})$, $\ell \in \mathbb{N}$ with $\ell > \frac{4}{T}$, where $\eta_\ell \in [0, 1]$ with $\eta_\ell(t) = 1$ for $t \in [\frac{2}{\ell}, T - \frac{2}{\ell}]$ and $\|\frac{d\eta_\ell}{dt}\|_{L^\infty(0, T)} \leq 2\ell$ yielding $\|\frac{d\eta_\ell}{dt}\|_{L^1(0, T)} \leq 4$. For a fixed ℓ , we now let $\delta \rightarrow 0$ in (6.25) and using the standard convergence properties of mollifiers we obtain that

$$\left| \int_0^T \eta_\ell \int_\Omega p_\kappa(\rho_\kappa) \mathbf{b}_\ell(\rho_\kappa) \, dx \, dt \right| \leq C,
\tag{6.26}$$

where $C \in \mathbb{R}$ is independent of ℓ and κ . Letting $\ell \rightarrow \infty$ in (6.26), and noting that $\eta_\ell \rightarrow 1$ pointwise in $(0, T)$, $\mathbf{b}_\ell(\rho_\kappa) \rightarrow \mathbf{b}(\rho_\kappa) = \rho_\kappa^\vartheta$ pointwise in Ω_T and Fatou's lemma, we obtain that

$$\int_0^T \int_\Omega p_\kappa(\rho_\kappa) \rho_\kappa^\vartheta \, dx \, dt \leq C,
\tag{6.27}$$

where $C \in \mathbb{R}$ is independent of κ . Hence the desired result (6.24) follows from (6.27), (2.3) and (1.3). \square

Similarly to (6.11), it follows from (5.27b), (6.9a), (6.1), (6.2) and (3.4), on noting $\gamma > \frac{3}{2}$, that, for any $w \in L^{\Gamma+1}(0, T; W_0^{1,v}(\Omega))$ with $v = \max\{\Gamma + 1, \frac{6\gamma}{2\gamma-3}\}$,

$$\begin{aligned}
 & \left| \int_0^T \left\langle \frac{\partial(\rho_\kappa u_\kappa)}{\partial t}, w \right\rangle_{W_0^{1,\Gamma+1}(\Omega)} dt - \int_0^T \int_\Omega p_\kappa(\rho_\kappa) \nabla_x \cdot w \, dx \, dt \right| \\
 & \leq C \left[\|\rho_\kappa |u_\kappa|^2\|_{L^2(0,T;L^{\frac{6\gamma}{4\gamma+3}}(\Omega))} + \|u_\kappa\|_{L^2(0,T;H^1(\Omega))} \right] \|w\|_{L^2(0,T;W^{1,\frac{6\gamma}{2\gamma-3}}(\Omega))} \\
 & \quad + C \left[\|\tau_1(M\hat{\psi}_\kappa)\|_{L^2(0,T;L^{\frac{4}{3}}(\Omega))} + \|\varrho_\kappa\|_{L^4(0,T;L^{\frac{2d}{d-1}}(\Omega))}^2 \right] \|\nabla_x w\|_{L^2(0,T;L^4(\Omega))} \\
 & \quad + \|\rho_\kappa\|_{L^\infty(0,T;L^\gamma(\Omega))} \|f\|_{L^2(0,T;L^\infty(\Omega))} \|w\|_{L^2(0,T;L^3(\Omega))} \\
 (6.28) \quad & \leq C \|w\|_{L^2(0,T;W^{1,s}(\Omega))},
 \end{aligned}$$

where $s = \max\{4, \frac{6\gamma}{2\gamma-3}\}$. We now have the following analogue of Lemma 5.4.

Lemma 6.4. *There exists a $C \in \mathbb{R}_{>0}$, independent of κ , such that*

$$(6.29) \quad \left\| \frac{\partial(\rho_\kappa u_\kappa)}{\partial t} \right\|_{L^{\frac{\Gamma+\vartheta}{\Gamma}}(0,T;W_0^{1,r}(\Omega)')} \leq C,$$

where $\vartheta(\gamma)$ is defined as in (6.13), $r = \max\{s, \frac{\Gamma+\vartheta}{\vartheta}\}$ and $s = \max\{4, \frac{6\gamma}{2\gamma-3}\}$.

Hence, for a further subsequence of the subsequence of Lemma 6.2, it follows that, as $\kappa \rightarrow 0_+$,

$$(6.30a) \quad \rho_\kappa u_\kappa \rightarrow \rho u \quad \text{weakly in } L^{\frac{10\gamma-6}{3(\gamma+1)}}(\Omega_T), \quad \text{weakly in } W^{1,\frac{\Gamma+\vartheta}{\Gamma}}(0,T;W_0^{1,r}(\Omega)'),$$

$$(6.30b) \quad \rho_\kappa u_\kappa \rightarrow \rho u \quad \text{in } C_w([0,T];L^{\frac{2\gamma}{\gamma+1}}(\Omega)), \quad \text{strongly in } L^2(0,T;H^1(\Omega)'),$$

$$(6.30c) \quad \rho_\kappa u_\kappa \rightarrow \rho u \quad \text{weakly in } L^2(0,T;L^{\frac{6\gamma}{\gamma+6}}(\Omega)),$$

$$(6.30d) \quad \rho_\kappa u_\kappa \otimes u_\kappa \rightarrow \rho u \otimes u \quad \text{weakly in } L^2(0,T;L^{\frac{6\gamma}{4\gamma+3}}(\Omega)),$$

$$(6.30e) \quad \rho_\kappa \rightarrow \rho \quad \text{weakly in } L^{\gamma+\vartheta}(\Omega_T),$$

$$(6.30f) \quad \rho_\kappa^\gamma \rightarrow \overline{\rho^\gamma} \quad \text{weakly in } L^{\frac{\gamma+\vartheta}{\gamma}}(\Omega_T),$$

$$(6.30g) \quad \kappa(\rho_\kappa^4 + \rho_\kappa^\Gamma) \rightarrow 0 \quad \text{weakly in } L^{\frac{\Gamma+\vartheta}{\Gamma}}(\Omega_T),$$

where $\overline{\rho^\gamma} \in L_{\geq 0}^{\frac{\gamma+\vartheta}{\gamma}}(\Omega_T)$ remains to be identified.

Proof. It follows from (6.28), (2.3) and (6.24) that, for all $w \in L^{\frac{\Gamma+\vartheta}{\vartheta}}(0,T;W_0^{1,r}(\Omega))$,

$$\begin{aligned}
 & \left| \int_0^T \left\langle \frac{\partial(\rho_\kappa u_\kappa)}{\partial t}, w \right\rangle_{W_0^{1,r}(\Omega)} dt \right| \\
 & \leq C \|w\|_{L^2(0,T;W^{1,s}(\Omega))} + \|p_\kappa(\rho_\kappa)\|_{L^{\frac{\Gamma+\vartheta}{\Gamma}}(\Omega_T)} \|w\|_{L^{\frac{\Gamma+\vartheta}{\vartheta}}(0,T;W^{1,\frac{\Gamma+\vartheta}{\vartheta}}(\Omega))} \\
 (6.31) \quad & \leq C \|w\|_{L^{\frac{\Gamma+\vartheta}{\vartheta}}(0,T;W^{1,r}(\Omega))},
 \end{aligned}$$

where we have noted from (6.13) that $\frac{\Gamma+\vartheta}{\vartheta} \geq \frac{\gamma+\vartheta}{\vartheta} \geq 2$. The desired result (6.29) then follows from (6.31).

The results (6.30a–d) follow similarly to (4.97a–e) from (6.9a), (6.29), (3.14a,b), (3.11), (6.10b), and (6.5). The results (6.30e–g) follow immediately from (6.24) and (2.3). \square

We now have the following analogue of Lemma 5.5.

Lemma 6.5. *The triple $(\rho, \underline{u}, \widehat{\psi})$, defined as in Lemmas 6.1 and 6.2, satisfies*

$$(6.32a) \quad \int_0^T \left\langle \frac{\partial \rho}{\partial t}, \eta \right\rangle_{W^{1,6}(\Omega)} dt - \int_0^T \int_{\Omega} \rho \underline{u} \cdot \nabla_x \eta \, dx \, dt = 0 \quad \forall \eta \in L^2(0, T; W^{1,6}(\Omega)),$$

with $\rho(\cdot, 0) = \rho_0(\cdot)$,

$$(6.32b) \quad \int_0^T \left\langle \frac{\partial(\rho \underline{u})}{\partial t}, w \right\rangle_{W_0^{1,r}(\Omega)} dt + \int_0^T \int_{\Omega} \left[S(\underline{u}) - \rho \underline{u} \otimes \underline{u} - c_p \overline{\rho^\gamma} I \right] : \nabla_x w \, dx \, dt \\ = \int_0^T \int_{\Omega} \left[\rho f \cdot w - \left(\tau_1(M \widehat{\psi}) - \mathfrak{z} \varrho^2 I \right) : \nabla_x w \right] \, dx \, dt \quad \forall w \in L^{\frac{\gamma+\vartheta}{\vartheta}}(0, T; W_0^{1,r}(\Omega)),$$

with $(\rho \underline{u})(\cdot, 0) = (\rho_0 \underline{u}_0)(\cdot)$, $\vartheta(\gamma)$ defined as in (6.13) and $r = \max\{4, \frac{6\gamma}{2\gamma-3}\}$, and

$$(6.32c) \quad \int_0^T \left\langle M \frac{\partial \widehat{\psi}}{\partial t}, \varphi \right\rangle_{H^s(\Omega \times D)} dt + \frac{1}{4\lambda} \sum_{i=1}^K \sum_{j=1}^K A_{ij} \int_0^T \int_{\Omega \times D} M \nabla_{q_j} \widehat{\psi} \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt \\ + \int_0^T \int_{\Omega \times D} M \left[\varepsilon \nabla_x \widehat{\psi} - u \widehat{\psi} \right] \cdot \nabla_x \varphi \, dq \, dx \, dt \\ - \int_0^T \int_{\Omega \times D} M \sum_{i=1}^K \left[\sigma(\underline{u}) q_i \right] \widehat{\psi} \cdot \nabla_{q_i} \varphi \, dq \, dx \, dt = 0 \quad \forall \varphi \in L^2(0, T; H^s(\Omega \times D)),$$

with $\widehat{\psi}(\cdot, 0) = \widehat{\psi}_0(\cdot)$ and $s > 1 + \frac{1}{2}(K+1)d$.

In addition, the triple $(\rho, \underline{u}, \widehat{\psi})$ satisfies, for a.a. $t' \in (0, T)$,

$$(6.33) \quad \frac{1}{2} \int_{\Omega} \rho(t') |\underline{u}(t')|^2 \, dx + \int_{\Omega} P(\rho(t')) \, dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}(t')) \, dq \, dx \\ + \mu^S c_0 \int_0^{t'} \|\underline{u}\|_{H^1(\Omega)}^2 \, dt + k \int_0^{t'} \int_{\Omega \times D} M \left[\frac{a_0}{2\lambda} \left| \nabla_q \sqrt{\widehat{\psi}} \right|^2 + 2\varepsilon \left| \nabla_x \sqrt{\widehat{\psi}} \right|^2 \right] \, dq \, dx \, dt \\ + \mathfrak{z} \|\varrho(t')\|_{L^2(\Omega)}^2 + 2\mathfrak{z} \varepsilon \int_0^{t'} \|\nabla_x \varrho\|_{L^2(\Omega)}^2 \, dt \\ \leq e^{t'} \left[\frac{1}{2} \int_{\Omega} \rho_0 |\underline{u}_0|^2 \, dx + \int_{\Omega} P(\rho_0) \, dx + k \int_{\Omega \times D} M \mathcal{F}(\widehat{\psi}_0) \, dq \, dx \right. \\ \left. + \mathfrak{z} \int_{\Omega} \left(\int_D M \widehat{\psi}_0 \, dq \right)^2 \, dx + \frac{1}{2} \int_0^{t'} \|f\|_{L^\infty(\Omega)}^2 \, dt \int_{\Omega} \rho_0 \, dx \right].$$

Proof. Passing to the limit $\kappa \rightarrow 0_+$ for the subsequence of Lemma 6.4 in (5.27a) yields (6.32a) subject to the stated initial condition, on noting (6.10a), (6.30c) and that $\frac{6\gamma}{\gamma+6} > \frac{6}{5}$ as $\gamma > \frac{3}{2}$.

Similarly to the proof of (4.118b), passing to the limit $\kappa \rightarrow 0_+$ for the subsequence of Lemma 6.4 in (5.62) for any $\underline{w} \in C_0^\infty(\Omega_T)$ yields (6.32b) for any $\underline{w} \in C_0^\infty(\Omega_T)$ subject to the stated initial condition, on noting (6.5), (6.10a), (6.30a,b,d,f,g), (6.6e) and (6.8b). The desired result (6.32b) for any $\underline{w} \in L^{\frac{\gamma+\vartheta}{\vartheta}}(0, T; W_0^{1,r}(\Omega))$ then follows from (6.28), (6.30f) and noting from (6.13) that $r \geq \frac{\gamma+\vartheta}{\vartheta} \geq 2$. Similarly to the proof of (4.118c), passing to the limit $\kappa \rightarrow 0_+$ for the subsequence of Lemma 6.4 in (5.27c) yields (6.32c) subject to the stated initial condition, on noting (6.5), (6.6a-d), (6.7) and (3.2). Similarly to the proof of (4.119), we deduce (6.33) from (5.28) using the results (6.30d), (6.10c), (6.6a,b,f), (6.5) and (6.8a). \square

We need to identify $\overline{\rho^\gamma}$ in (6.32b) and (6.30f). Similarly to (6.14), with $\ell \in \mathbb{N}$, we now introduce $\mathfrak{t} : \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$ and $\mathfrak{t}_\ell : \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$ such that

$$(6.34) \quad \mathfrak{t}(s) := s \quad \text{and} \quad \mathfrak{t}_\ell(s) := \begin{cases} \mathfrak{t}(s) & \text{for } 0 \leq s \leq \ell, \\ \mathfrak{t}(\ell) & \text{for } \ell \leq s. \end{cases}$$

Then, similarly to (6.17), we have the renormalised equation, for any $\ell \in \mathbb{N}$,

$$(6.35) \quad \frac{\partial \mathbf{t}_\ell(\rho_\kappa)}{\partial t} + \nabla_x \cdot (\mathbf{t}_\ell(\rho_\kappa) \mathbf{u}_\kappa) + (\rho_\kappa (\mathbf{t}_\ell)_+'(\rho_\kappa) - \mathbf{t}_\ell(\rho_\kappa)) \nabla_x \cdot \mathbf{u}_\kappa = 0 \quad \text{in } C_0^\infty(\mathbb{R}^d \times (0, T))',$$

where $(\mathbf{t}_\ell)_+'(\cdot)$ is defined similarly to (6.18). It follows from (6.34), (5.28) and (6.35) that, for any fixed $\ell \in \mathbb{N}$,

$$(6.36) \quad \|\mathbf{t}_\ell(\rho_\kappa)\|_{L^\infty(\Omega_T)} + \|(\rho_\kappa (\mathbf{t}_\ell)_+'(\rho_\kappa) - \mathbf{t}_\ell(\rho_\kappa)) \nabla_x \cdot \mathbf{u}_\kappa\|_{L^2(\Omega_T)} + \left\| \frac{\partial \mathbf{t}_\ell(\rho_\kappa)}{\partial t} \right\|_{L^2(0, T; H^1(\Omega)')} \leq C(\ell).$$

In order to identify $\overline{\rho^\gamma}$ in (6.32b) and (6.30f), we now apply Corollary 5.1 with (5.38a,b) being (6.35) and (5.62) so that $\mu = \frac{\mu^S}{2}$ and $\lambda = \mu^B - \frac{\mu^S}{d}$, $g_n = \mathbf{t}_\ell(\rho_\kappa)$ for a fixed $\ell \in \mathbb{N}$, $\mathbf{u}_n = \mathbf{u}_\kappa$, $\mathbf{m}_n = \rho_\kappa \mathbf{u}_\kappa$, $p_n = p_\kappa(\rho_\kappa)$, $\mathbf{z}_n = \mathbf{z}_\kappa(M \hat{\psi}_\kappa)$, $f_n = -(\rho_\kappa (\mathbf{t}_\ell)_+'(\rho_\kappa) - \mathbf{t}_\ell(\rho_\kappa)) \nabla_x \cdot \mathbf{u}_\kappa$ and $F_n = \rho_\kappa \hat{f}_\kappa$. With $\{(\rho_\kappa, \mathbf{u}_\kappa, \hat{\psi}_\kappa)\}_{\kappa > 0}$ being the subsequence (not indicated) of Lemma 6.4, we have that (5.37a–d) hold with $g = \overline{\mathbf{t}_\ell(\rho)}$, $\mathbf{u} = \mathbf{u}$, $\mathbf{m} = \rho \mathbf{u}$ and $p = c_p \overline{\rho^\gamma}$, and $q < \infty$, $\omega = \infty$, $z = \frac{2\gamma}{\gamma+1}$ and $r = \frac{\gamma+\vartheta}{\Gamma}$ on recalling (6.36), (3.14a,b), (6.5), (6.30b) and (6.30f,g). We note that $z = \frac{2\gamma}{\gamma+1} > \frac{6}{5}$ as $\gamma > \frac{3}{2}$. Hence, the constraints on q, r, ω and z hold. The results (5.37e,h) hold with $\mathbf{z}_\kappa = \mathbf{z}_\kappa(M \hat{\psi}_\kappa)$, $F = \rho \hat{f}$ and $s = \gamma + \vartheta$, on recalling (6.6e) and (6.30e), and noting that $\frac{4(d+2)}{3d+4} \geq \frac{q}{q-1}$. Finally, the result (5.47) holds with $f = -(\rho (\mathbf{t}_\ell)_+'(\rho) - \mathbf{t}_\ell(\rho)) \nabla_x \cdot \mathbf{u}$ on recalling (6.36). Hence, we obtain from (5.39) for the subsequence of Lemma 6.4 that, for any fixed $\ell \in \mathbb{N}$ and for all $\zeta \in C_0^\infty(\Omega)$ and $\eta \in C_0^\infty(0, T)$,

$$(6.37) \quad \lim_{\kappa \rightarrow 0_+} \int_0^T \eta \left(\int_\Omega \zeta \mathbf{t}_\ell(\rho_\kappa) [p_\kappa(\rho_\kappa) - \mu^* \nabla_x \cdot \mathbf{u}_\kappa] \, dx \right) dt = \int_0^T \eta \left(\int_\Omega \zeta \overline{\mathbf{t}_\ell(\rho)} [c_p \overline{\rho^\gamma} - \mu^* \nabla_x \cdot \mathbf{u}] \, dx \right) dt,$$

where $\mu^* := \frac{(d-1)}{d} \mu^S + \mu^B$.

We deduce from (6.37), (6.34), (2.3), (1.3), (6.30f,g) and (6.5) that, for any fixed $\ell \in \mathbb{N}$,

$$(6.38) \quad c_p \left[\overline{\mathbf{t}_\ell(\rho) \rho^\gamma} - \overline{\mathbf{t}_\ell(\rho)} \overline{\rho^\gamma} \right] = \mu^* \left[\overline{\mathbf{t}_\ell(\rho) \nabla_x \cdot \mathbf{u}} - \overline{\mathbf{t}_\ell(\rho)} \nabla_x \cdot \mathbf{u} \right] \quad \text{a.e. in } \Omega_T,$$

where, as $\kappa \rightarrow 0_+$,

$$(6.39a) \quad \mathbf{t}_\ell(\rho_\kappa) \rho_\kappa^\gamma \rightarrow \overline{\mathbf{t}_\ell(\rho) \rho^\gamma} \quad \text{weakly in } L^{\frac{\gamma+\vartheta}{\gamma}}(\Omega_T),$$

$$(6.39b) \quad \text{and} \quad \mathbf{t}_\ell(\rho_\kappa) \nabla_x \cdot \mathbf{u}_\kappa \rightarrow \overline{\mathbf{t}_\ell(\rho) \nabla_x \cdot \mathbf{u}} \quad \text{weakly in } L^2(\Omega_T).$$

We can now follow the discussion in Sections 7.10.2–7.10.5 in Novotný & Straškraba [58] to deduce that $\overline{p(\rho)} = p(\rho)$. For the benefit of the reader, we briefly outline the argument. First, it follows from (6.30f) and (6.39a) that, for any fixed $\ell \in \mathbb{N}$,

$$(6.40) \quad \begin{aligned} & \int_{\Omega_T} \left[\overline{\mathbf{t}_\ell(\rho) \rho^\gamma} - \overline{\mathbf{t}_\ell(\rho)} \overline{\rho^\gamma} \right] \, dx \, dt \\ &= \lim_{\kappa \rightarrow 0_+} \left[\int_{\Omega_T} (\mathbf{t}_\ell(\rho_\kappa) - \mathbf{t}_\ell(\rho)) (\rho_\kappa^\gamma - \rho^\gamma) \, dx \, dt + \int_{\Omega_T} (\mathbf{t}_\ell(\rho) - \overline{\mathbf{t}_\ell(\rho)}) (\overline{\rho^\gamma} - \rho^\gamma) \, dx \, dt \right] \\ &\geq \limsup_{\kappa \rightarrow 0_+} \int_{\Omega_T} |\mathbf{t}_\ell(\rho_\kappa) - \mathbf{t}_\ell(\rho)|^{\gamma+1} \, dx \, dt, \end{aligned}$$

where we have noted that the second term on the second line is nonnegative as $\mathbf{t}_\ell(s)$ is concave and s^γ is convex for $s \in [0, \infty)$.

We now introduce $\mathfrak{L}_\ell : \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$, $\ell \in \mathbb{N}$, such that $s \mathfrak{L}'_\ell(s) - \mathfrak{L}_\ell(s) = \mathbf{t}_\ell(s)$ for all $s \in [0, \infty)$, so that

$$(6.41) \quad \mathfrak{L}_\ell(s) := \begin{cases} \mathfrak{L}(s) := s \log s & \text{for } 0 \leq s \leq \ell, \\ s \log \ell + s - \ell & \text{for } \ell \leq s. \end{cases}$$

Similarly to (5.57) and (6.17), one deduces, via renormalization, from (6.32a) and (5.27a) that, for any fixed $\ell \in \mathbb{N}$ and for any $t' \in (0, T]$,

$$(6.42a) \quad \int_{\Omega} [\mathfrak{L}_{\ell}(\rho)(t') - \mathfrak{L}_{\ell}(\rho_0)] \, d\tilde{x} = - \int_0^{t'} \int_{\Omega} \mathfrak{t}_{\ell}(\rho) \nabla_x \cdot u \, d\tilde{x} \, dt,$$

$$(6.42b) \quad \int_{\Omega} [\mathfrak{L}_{\ell}(\rho_{\kappa})(t') - \mathfrak{L}_{\ell}(\rho_0)] \, d\tilde{x} = - \int_0^{t'} \int_{\Omega} \mathfrak{t}_{\ell}(\rho_{\kappa}) \nabla_x \cdot u_{\kappa} \, d\tilde{x} \, dt.$$

Although establishing (6.42b) is straightforward as $\rho_{\kappa} \in C_w([0, T; L_{\geq 0}^{\Gamma}(\Omega)))$, establishing (6.42a) is not, since $\rho \in C_w([0, T; L_{\geq 0}^{\gamma}(\Omega)))$, and so ρ may not be in $L^2(\Omega_T)$ as $\gamma > \frac{3}{2}$. Nevertheless, (6.42a) can still be established, see Lemma 7.57 in [58]. We note that our $\vartheta(\gamma)$, recall (6.13), differs from the $\vartheta(\gamma)$ in [58] for $\gamma \geq 4$, due to the presence of the extra stress term in the momentum equation for our polymer model. However, as $\rho \in L^2(\Omega_T)$ for $\gamma \geq 4$, Lemma 7.57 in [58] is not required for such γ . Subtracting (6.42a) from (6.42b), and passing to the limit $\kappa \rightarrow 0_+$, one deduces from (6.39b) that, for any fixed $\ell \in \mathbb{N}$ and for any $t' \in (0, T]$,

$$(6.43) \quad \int_{\Omega} [\overline{\mathfrak{L}_{\ell}(\rho)(t')} - \mathfrak{L}_{\ell}(\rho)(t')] \, d\tilde{x} = \int_0^{t'} \int_{\Omega} [\mathfrak{t}_{\ell}(\rho) \nabla_x \cdot u - \overline{\mathfrak{t}_{\ell}(\rho) \nabla_x \cdot u}] \, d\tilde{x} \, dt,$$

where, on noting (6.10a) and the convexity of \mathfrak{L}_{ℓ} ,

$$(6.44) \quad \mathfrak{L}_{\ell}(\rho_{\kappa})(t') \rightarrow \overline{\mathfrak{L}_{\ell}(\rho)(t')} \geq \mathfrak{L}_{\ell}(\rho)(t') \quad \text{weakly in } L^{\gamma}(\Omega), \quad \text{as } \kappa \rightarrow 0_+.$$

It follows from (6.40), (6.38), (6.43), (6.44) and (6.4a) that, for any fixed $\ell \in \mathbb{N}$,

$$\begin{aligned} \limsup_{\kappa \rightarrow 0_+} \|\mathfrak{t}_{\ell}(\rho) - \mathfrak{t}_{\ell}(\rho_{\kappa})\|_{L^{\gamma+1}(\Omega_T)}^{\gamma+1} &\leq \frac{\mu^{\star}}{c_p} \int_{\Omega_T} [\overline{\mathfrak{t}_{\ell}(\rho) \nabla_x \cdot u} - \overline{\mathfrak{t}_{\ell}(\rho)} \nabla_x \cdot u] \, d\tilde{x} \, dt \\ &\leq \frac{\mu^{\star}}{c_p} \int_{\Omega_T} [\mathfrak{t}_{\ell}(\rho) - \overline{\mathfrak{t}_{\ell}(\rho)}] \nabla_x \cdot u \, d\tilde{x} \, dt \\ &\leq C \|\mathfrak{t}_{\ell}(\rho) - \overline{\mathfrak{t}_{\ell}(\rho)}\|_{L^2(\Omega_T)} \\ &\leq C \limsup_{\kappa \rightarrow 0_+} \|\mathfrak{t}_{\ell}(\rho) - \mathfrak{t}_{\ell}(\rho_{\kappa})\|_{L^2(\Omega_T)} \\ (6.45) \quad &\leq C \limsup_{\kappa \rightarrow 0_+} \|\mathfrak{t}_{\ell}(\rho) - \mathfrak{t}_{\ell}(\rho_{\kappa})\|_{L^{\gamma+1}(\Omega_T)}, \end{aligned}$$

where $C \in \mathbb{R}_{>0}$ is independent of ℓ and κ . It is easily deduced from (6.24), (6.30e) and (6.34) that, for all $\ell \in \mathbb{N}$, $\kappa > 0$ and $r \in [1, \gamma + \vartheta]$,

$$(6.46) \quad \|\rho_{\kappa} - \mathfrak{t}_{\ell}(\rho_{\kappa})\|_{L^r(\Omega_T)} + \|\rho - \mathfrak{t}_{\ell}(\rho)\|_{L^r(\Omega_T)} + \|\rho - \overline{\mathfrak{t}_{\ell}(\rho)}\|_{L^r(\Omega_T)} \leq C \ell^{1 - \frac{\gamma + \vartheta}{r}},$$

where $C \in \mathbb{R}_{>0}$ is independent of ℓ and κ . It follows from (6.45), (6.46) and (3.1) that

$$(6.47) \quad \lim_{\ell \rightarrow \infty} \limsup_{\kappa \rightarrow 0_+} \|\mathfrak{t}_{\ell}(\rho) - \mathfrak{t}_{\ell}(\rho_{\kappa})\|_{L^{\gamma+1}(\Omega_T)} = 0.$$

It immediately follows from (6.46), (6.47), (6.24), (3.1) and (6.30f), on possibly extracting a further subsequence (not indicated), that, as $\kappa \rightarrow 0_+$,

$$(6.48a) \quad \rho_{\kappa} \rightarrow \rho \quad \text{strongly in } L^s(\Omega_T), \quad \text{for any } s \in [1, \gamma + \vartheta(\gamma)),$$

$$(6.48b) \quad \rho_{\kappa}^{\gamma} \rightarrow \rho^{\gamma} \quad \text{weakly in } L^{\frac{\gamma + \vartheta}{\gamma}}(\Omega_T), \quad \text{that is, } \overline{\rho^{\gamma}} = \rho^{\gamma}.$$

Finally, we have the analogue of Theorem 5.1.

Theorem 6.1. *The triple $(\rho, \underline{u}, \widehat{\psi})$, defined as in Lemmas 6.1 and 6.2, is a global weak solution to problem (P), in the sense that (6.32a,c), with their initial conditions, hold and*

$$\begin{aligned} & \int_0^T \left\langle \frac{\partial(\rho \underline{u})}{\partial t}, w \right\rangle_{W_0^{1,r}(\Omega)} dt + \int_0^T \int_{\Omega} \left[S(\underline{u}) - \rho \underline{u} \otimes \underline{u} - c_p \rho^\gamma I \right] : \nabla_x w \, dx \, dt \\ &= \int_0^T \int_{\Omega} \left[\rho f \cdot w - \left(\tau_1(M \widehat{\psi}) - \mathfrak{z} \varrho^2 I \right) : \nabla_x w \right] \, dx \, dt \\ (6.49) \quad & \forall w \in L^{\frac{\gamma+\vartheta}{\vartheta}}(0, T; W_0^{1,r}(\Omega)), \end{aligned}$$

with $(\rho \underline{u})(\cdot, 0) = (\rho_0 \underline{u}_0)(\cdot)$, $\vartheta(\gamma)$ defined as in (6.13) and $r = \max\{4, \frac{6\gamma}{2\gamma-3}\}$. In addition, the weak solution $(\rho, \underline{u}, \widehat{\psi})$ satisfies (6.33).

Proof. The results (6.32a,c) and (6.33) have already been established in Lemma 6.5. Equation (6.49) was established in Lemma 6.5 with ρ^γ replaced by $\overline{\rho^\gamma}$, see (6.32b). The desired result (6.49) then follows immediately from (6.32b) and (6.48b). \square

7. CONCLUSIONS

We proved the existence of global-in-time weak solutions to a general class of models that arise from the kinetic theory of dilute solutions of nonhomogeneous polymeric liquids, where the polymer molecules are idealized as bead-spring chains with finitely extensible nonlinear elastic (FENE) type spring potentials. The class of models under consideration involved the unsteady, compressible, isentropic, isothermal Navier–Stokes system, with constant shear and bulk viscosity coefficients, in a bounded domain Ω in \mathbb{R}^d , $d = 2$ or 3 , for the density ρ , the velocity \underline{u} and the pressure p of the fluid, with an equation of state of the form $p(\rho) = c_p \rho^\gamma$, where c_p is a positive constant and $\gamma > \frac{3}{2}$. Our analysis applies, without alterations, to some other familiar *monotone* equations of state, such as the (Kirkwood-modified) Tait equation of state (cf. Remark 1.1). As a starting point for the extension of this work to nonmonotone equations of state we refer the reader to the work of Feireisl [30] in the context of the compressible Navier–Stokes equations; see also Section 7.12.3 in [58], which explains how, in the case of the compressible Navier–Stokes equations at least, the existence proof for monotone equations of state can be extended, with minor modifications, to the case of nonmonotone equations of state.

The right-hand side of the Navier–Stokes momentum equation included an elastic extra-stress tensor, defined as the sum of the classical Kramers expression and a quadratic interaction term. The elastic extra-stress tensor stems from the random movement of the polymer chains and involves the associated probability density function that satisfies a Fokker–Planck-type parabolic equation, a crucial feature of which is the presence of a centre-of-mass diffusion term. Our analysis required no structural assumptions on the drag term in the Fokker–Planck equation; in particular, the drag term need not be corotational. With a nonnegative initial density $\rho_0 \in L^\infty(\Omega)$ for the continuity equation; a square-integrable initial velocity datum \underline{u}_0 for the Navier–Stokes momentum equation; and a nonnegative initial probability density function ψ_0 for the Fokker–Planck equation, which has finite relative entropy with respect to the Maxwellian M associated with the spring potential in the model, we proved, *via* a limiting procedure on certain discretization and regularization parameters, the existence of a global-in-time bounded-energy weak solution $t \mapsto (\rho(t), \underline{u}(t), \psi(t))$ to the coupled Navier–Stokes–Fokker–Planck system, satisfying the initial condition $(\rho(0), \underline{u}(0), \psi(0)) = (\rho_0, \underline{u}_0, \psi_0)$.

The representation $p_p = k(K+1)\varrho + \mathfrak{z}\varrho^2$ of the polymeric pressure p_p used in the paper (cf. Remark 1.3) can be viewed as a quadratic truncation of a virial expansion of the polymeric pressure in terms of the polymer number density ϱ . The presence of the quadratic interaction term $\mathfrak{z}\varrho^2$, with $\mathfrak{z} > 0$, was found to be essential for the analysis, pointing at the necessity of accounting for effects that are not captured by the classical Kramers expression, such as excluded volume effects.

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APPENDIX

A. HOMOGENEOUS SOBOLEV SPACES

The linear space $C_0^\infty(\mathbb{R}^d)$, when equipped with the functional $|\cdot|_{1,r}$ defined by $|\eta|_{1,r} := \|\nabla_x \eta\|_{L^r(\mathbb{R}^d)}$, is a normed linear space. We denote by $D^{1,r}(\mathbb{R}^d)$ the closure of $C_0^\infty(\mathbb{R}^d)$ in this norm, and will refer to $D^{1,r}(\mathbb{R}^d)$ as a *homogeneous Sobolev space*. The space $D^{1,r}(\mathbb{R}^d)$ is separable for $r \in [1, \infty)$ and reflexive for $r \in (1, \infty)$. For $r \in [1, d)$ the space $D^{1,r}(\mathbb{R}^d)$ can be characterized as follows:

$$D^{1,r}(\mathbb{R}^d) = \{\eta \in C_0^\infty(\mathbb{R}^d)' : \eta \in L^{\frac{dr}{d-r}}(\mathbb{R}^d), \nabla_x \eta \in L^r(\mathbb{R}^d)\}.$$

More specifically, for $r \in [1, d)$, the following *Sobolev inequality* holds:

$$\|\eta\|_{L^{\frac{dr}{d-r}}(\mathbb{R}^d)} \leq C(r, d) \|\nabla_x \eta\|_{L^r(\mathbb{R}^d)} \quad \forall \eta \in D^{1,r}(\mathbb{R}^d),$$

where $C(r, d)$ is a positive constant depending on r and d but independent of η . For $r \in [d, \infty)$ the elements of $D^{1,r}(\mathbb{R}^d)$ are not distributions and can only be viewed as sets of equivalence classes.

B. FOURIER MULTIPLIERS

Let, as in (5.29), $\mathfrak{S}(\mathbb{R}^d)$ signify the linear space of smooth *rapidly decreasing (complex-valued) functions*, equipped with the topology of locally uniform convergence, and denote by $\mathfrak{S}(\mathbb{R}^d)'$ its dual space, referred to as the space of *tempered distributions*. On the space of rapidly decreasing functions the Fourier transform $\mathfrak{F} : \mathfrak{S}(\mathbb{R}^d) \rightarrow \mathfrak{S}(\mathbb{R}^d)$ and its inverse $\mathfrak{F}^{-1} : \mathfrak{S}(\mathbb{R}^d) \rightarrow \mathfrak{S}(\mathbb{R}^d)$ are defined by the formulae (5.30), which we extend to $\mathfrak{F}, \mathfrak{F}^{-1} : \mathfrak{S}(\mathbb{R}^d)' \rightarrow \mathfrak{S}(\mathbb{R}^d)'$ via the identities (5.31).

Lemma B.1 (Theorem 1.55 in [58]).

- (i) *The operator \mathfrak{F} defined by (5.30) can be extended by continuity to a bounded linear operator from $L^2(\mathbb{R}^d)$ to $L^2(\mathbb{R}^d)$. Furthermore, the following Parseval–Plancherel identity holds:*

$$\int_{\mathbb{R}^d} \overline{\eta(\underline{x})} \zeta(\underline{x}) \, d\underline{x} = \int_{\mathbb{R}^d} \overline{[\mathfrak{F}\eta](\underline{y})} [\mathfrak{F}(\zeta)](\underline{y}) \, d\underline{y} \quad \forall \eta, \zeta \in L^2(\mathbb{R}^d).$$

- (ii) *For any $\eta \in \mathfrak{S}(\mathbb{R}^d)'$ and any multiindex $\alpha \in \mathbb{N}^d$, one has the following identities:*

$$\frac{\partial^{|\alpha|}}{\partial y_1^{\alpha_1} \dots \partial y_d^{\alpha_d}} \mathfrak{F}\eta = \mathfrak{F}[(-i\underline{x})^\alpha \eta], \quad \mathfrak{F}\left(\frac{\partial^{|\alpha|}}{\partial x_1^{\alpha_1} \dots \partial x_d^{\alpha_d}} \eta\right) = (i\underline{y})^\alpha \mathfrak{F}(\eta).$$

- (iii) *For any $\eta \in \mathfrak{S}(\mathbb{R}^d)'$ and $\zeta \in \mathfrak{S}(\mathbb{R}^d)$, the following convolution identity holds:*

$$\mathfrak{F}(\eta * \zeta) = \mathfrak{F}(\eta) \mathfrak{F}(\zeta).$$

A bounded measurable function $m : \mathbb{R}^d \rightarrow \mathbb{R}$ is called a *Fourier multiplier* of type (r, s) , with $r, s \in [1, \infty)$, if there exists a positive constant $C(r, s)$ such that

$$\|\mathfrak{F}^{-1}(m \mathfrak{F}(\eta))\|_{L^s(\mathbb{R}^d)} \leq C(r, s) \|\eta\|_{L^r(\mathbb{R}^d)} \quad \forall \eta \in \mathfrak{S}(\mathbb{R}^d).$$

The linear operator

$$T : \mathfrak{S}(\mathbb{R}^d) \subset L^r(\mathbb{R}^d) \rightarrow L^s(\mathbb{R}^d), \quad T\eta = \mathfrak{F}^{-1}(m \mathfrak{F}(\eta)) = \mathfrak{F}^{-1}(m) * \eta,$$

with domain $\mathcal{D}(T) = \mathfrak{S}(\mathbb{R}^d)$ is a densely defined continuous linear operator from $L^r(\mathbb{R}^d)$ to $L^s(\mathbb{R}^d)$. Therefore, its closure, which is still denoted by T , is a continuous linear operator from $L^r(\mathbb{R}^d)$ into $L^s(\mathbb{R}^d)$; i.e.,

$$\|\mathfrak{F}^{-1}(m \mathfrak{F}(\eta))\|_{L^s(\mathbb{R}^d)} \leq C(r, s) \|\eta\|_{L^r(\mathbb{R}^d)} \quad \forall \eta \in L^r(\mathbb{R}^d).$$

Lemma B.2 (Lizorkin's multiplier theorem; Theorem 1.57 in [58]). *Let $r \in (1, \infty)$, $\beta \in [0, \frac{1}{r})$, and let $m : \underline{y} \in \mathbb{R}^d \mapsto m(\underline{y}) \in \mathbb{R}$ be such that $m \in L^\infty(\mathbb{R}^d)$, and m has the derivative*

$$\frac{\partial^d m}{\partial y_1 \dots \partial y_d}$$

as well as all lower-order partial derivatives, continuous on $\mathbb{R}^d \setminus \{0\}$. Suppose further that there exists a $B > 0$ such that

$$|y_1|^{\alpha_1+\beta} \dots |y_d|^{\alpha_d+\beta} \left| \frac{\partial^{|\alpha|} m}{\partial y_1^{\alpha_1} \dots \partial y_d^{\alpha_d}}(\underline{y}) \right| \leq B \quad \forall \underline{y} \in \mathbb{R}^d \setminus \{0\},$$

where, for each $i \in \{1, \dots, d\}$, α_i is 0 or 1. Then, m is a Fourier multiplier of type (r, s) with $\frac{1}{s} = \frac{1}{r} - \beta$. In particular if $\beta = 0$, then m is a Fourier multiplier of type (r, r) .

A particularly simple case is when $r = 2$: it follows from part (i) of Theorem B.1 that any $m \in L^\infty(\mathbb{R}^d)$ is a Fourier multiplier of type $(2, 2)$.

C. THE BOGOVSKIĬ OPERATOR

Lemma C.1 (Lemma 3.17 in [58]). *Suppose that Ω is a bounded Lipschitz domain in \mathbb{R}^d . Then, there exists a linear operator $\mathfrak{B} = (\mathfrak{B}^1, \dots, \mathfrak{B}^d)$, called the Bogovskiĭ operator, with the following properties:*

$$\mathfrak{B} : L_0^r(\Omega) \rightarrow \tilde{W}_0^{1,r}(\Omega), \quad 1 < r < \infty;$$

$$\nabla_x \cdot \mathfrak{B}(\eta) = \eta \quad \text{a.e. in } \Omega \quad \forall \eta \in L_0^r(\Omega), \quad 1 < r < \infty;$$

$$\|\nabla_x \mathfrak{B}(\eta)\|_{L^r(\Omega)} \leq C(r, \Omega) \|\eta\|_{L^r(\Omega)} \quad \forall \eta \in L_0^r(\Omega), \quad 1 < r < \infty.$$

For $1 < r, s < \infty$, let

$$\tilde{E}^{r,s}(\Omega) := \{\zeta \in \tilde{L}^r(\Omega) : \nabla_x \cdot \zeta \in L^s(\Omega)\},$$

equipped with the norm $\|\zeta\|_{\tilde{E}^{r,s}(\Omega)} = \|\zeta\|_{\tilde{L}^r(\Omega)} + \|\nabla_x \cdot \zeta\|_{L^s(\Omega)}$, and define $\tilde{E}_0^{r,s}(\Omega)$ as the closure of $\tilde{C}_0^\infty(\Omega)$ in the norm of $\tilde{E}^{r,s}(\Omega)$. If $\eta = \nabla_x \cdot \zeta$, where $\zeta \in \tilde{E}_0^{r,s}(\Omega)$ with some $r \in (1, \infty)$, then

$$\|\mathfrak{B}(\eta)\|_{L^r(\Omega)} \leq C(r, \Omega) \|\zeta\|_{\tilde{L}^r(\Omega)}.$$

Finally, if $\eta \in C_0^\infty(\Omega)$ and $\int \eta = 0$, then $\mathfrak{B}(\eta) \in \tilde{C}_0^\infty(\Omega)$.

D. CONVEXITY AND WEAK CONVERGENCE

Lemma D.1 (Corollary 3.33 in [58]).

- (i) *Let $r \in [1, \infty)$ and suppose that Ω is a domain in \mathbb{R}^d . Let further $\mathcal{H}(\eta) := \int_K |\eta|^r d\underline{x}$, $\eta \in L^r(K)$, where K is a measurable subset of Ω . Then,*

$$\mathcal{H}(\eta) \leq \liminf_{n \rightarrow \infty} \mathcal{H}(\eta_n) \quad \text{whenever } \eta_n \rightarrow \eta \text{ weakly in } L^r(\Omega).$$

- (ii) Let $r \in [1, \infty)$ and suppose that Ω is a domain in \mathbb{R}^d , I is an interval in \mathbb{R} and f is a convex lower-semicontinuous (respectively, concave upper-semicontinuous) function defined on I . Suppose further that η_n is a sequence of nonnegative functions from $L^r(\Omega)$ with values in I such that

$$\eta_n \rightarrow \eta \quad \text{weakly in } L^r(\Omega)$$

and

$$f(\eta_n) \rightarrow \overline{f(\eta)} \quad \text{weakly in } L^1(\Omega).$$

Then,

$$f(\eta) \leq \overline{f(\eta)} \quad (\text{respectively, } f(\eta) \geq \overline{f(\eta)}) \quad \text{a.e. in } \Omega.$$

Lemma D.2 (Lemma 3.34 in [58]). Let $r \in (1, \infty)$, suppose that Ω is a bounded domain in \mathbb{R}^d and I is an interval in \mathbb{R} . Suppose further that $f : I \rightarrow \mathbb{R}$ is a strictly convex function. Let $\{\eta_n\}_{n \in \mathbb{N}}$ be a sequence of functions in $L^r(\Omega)$ with values in I . If $\eta_n \rightarrow \eta$ weakly in $L^r(\Omega)$ and $f(\eta_n) \rightarrow f(\eta)$ weakly in $L^1(\Omega)$, then $\eta_n \rightarrow \eta$ strongly in $L^1(\Omega)$.

The following lemma, due to Feireisl, plays an important role in the proof of weak compactness of the effective viscous flux.

Lemma D.3 (Commutator lemma; Lemma 4.25 in [58]). Suppose that $d \in \{2, 3\}$, $r, s \in (1, \infty)$, $\frac{1}{r} + \frac{1}{s} = \frac{1}{v} < 1$, and

$$f_n \rightarrow f \quad \text{weakly in } L^r(\mathbb{R}^d),$$

$$g_n \rightarrow g \quad \text{weakly in } L^s(\mathbb{R}^d).$$

Consider the Riesz operators \mathcal{R}_{kj} , with $k, j = 1, \dots, d$, defined by (5.35). Then,

$$f_n \mathcal{R}_{kj}(g_n) - g_n \mathcal{R}_{kj}(f_n) \rightarrow f \mathcal{R}_{kj}(g) - g \mathcal{R}_{kj}(f) \quad \text{weakly in } L^v(\mathbb{R}^d).$$

E. AN ARZELA–ASCOLI TYPE THEOREM IN $C_w([0, T]; L^s(\Omega))$

Lemma E.1 (Lemma 6.2 in [58]). Let $r, s \in (1, \infty)$ and suppose that Ω is a bounded Lipschitz domain in \mathbb{R}^d , $d \geq 2$. Suppose further that $\{g_n\}_{n \in \mathbb{N}}$ is a sequence of functions defined on $[0, T]$ with values in $L^s(\Omega)$ such that $g_n \in C_w([0, T]; L^s(\Omega))$, and g_n is bounded in $C([0, T]; W^{-1, r}(\Omega)) \cap L^\infty(0, T; L^s(\Omega))$ independent of n . Then, there exists a subsequence (not indicated) such that

(i) $g_n \rightarrow g$ in $C_w([0, T]; L^s(\Omega))$;

(ii) If, in addition, $1 < r \leq \frac{d}{d-1}$ and $1 < s < \infty$, or $\frac{d}{d-1} < r < \infty$ and $\frac{dr}{d+r} < s < \infty$, then $g_n \rightarrow g$ strongly in $C([0, T]; W^{-1, r}(\Omega))$.

F. THE CONTINUITY EQUATION: EXTENSION TO \mathbb{R}^d AND RENORMALIZATION

Lemma F.1 (Lemma 6.8 in [58]). Let Ω be a bounded Lipschitz domain in \mathbb{R}^d , $d \geq 2$, and define $\Omega_T := \Omega \times (0, T)$. Suppose that $\rho \in L^2(\Omega_T)$, $\underline{u} \in L^2(0, T; H_0^1(\Omega))$ and $f \in L^1(\Omega_T)$ satisfy

$$\frac{\partial \rho}{\partial t} + \nabla_x \cdot (\underline{u} \rho) = f \quad \text{in } C_0^\infty(\Omega_T)'.$$

Then, upon extending (ρ, \underline{u}, f) to \mathbb{R}^d by $(0, 0, 0)$ outside Ω , one has that

$$\frac{\partial \rho}{\partial t} + \nabla_x \cdot (\underline{u} \rho) = f \quad \text{in } C_0^\infty(\mathbb{R}^d \times (0, T))'.$$

Next we collect some results of relevance for our analysis that concern the renormalization of the continuity equation. Consider a function b defined on $[0, \infty)$ such that

$$(F.1) \quad b \in C[0, \infty) \cap C^1(0, \infty), \quad |b'(s)| \leq Cs^{-\lambda_0} \quad \text{for } s \in (0, 1], \text{ where } C > 0 \text{ and } \lambda_0 < 1,$$

and

$$(F.2) \quad |b'(s)| \leq Cs^{\lambda_1} \quad \text{for } s \geq 1, \text{ where } C > 0 \text{ and } \lambda_1 > -1.$$

Lemma F.2 (Lemma 6.9 in [58]). *Let $d \geq 2$, $\beta \in [2, \infty)$, and let λ_1 be a real number such that*

$$(F.3) \quad \lambda_1 \leq \frac{\beta}{2} - 1.$$

Suppose further that $\rho \in L^\beta(0, T; L^\beta_{\text{loc}}(\mathbb{R}^d))$, $\rho \geq 0$ a.e. on $\mathbb{R}^d \times (0, T)$, $\underline{u} \in L^2(0, T; \underline{H}^1_{\text{loc}}(\mathbb{R}^d))$ and $f \in L^z(0, T; L^z_{\text{loc}}(\mathbb{R}^d))$, where $z = \left[\frac{\beta}{\lambda_1}\right]'$ is the conjugate of $\frac{\beta}{\lambda_1}$ if $\lambda_1 > 0$ and $z = 1$ if $\lambda_1 \leq 0$, are such that

$$\frac{\partial \rho}{\partial t} + \nabla_x \cdot (\underline{u} \rho) = f \quad \text{in } C_0^\infty(\mathbb{R}^d \times (0, T))'.$$

(i) *For any function $b \in C^1[0, \infty)$ satisfying (F.2) and (F.3) one has*

$$(F.4) \quad \frac{\partial b(\rho)}{\partial t} + \nabla_x \cdot (\underline{u} b(\rho)) + \{\rho b'(\rho) - b(\rho)\} \nabla_x \cdot \underline{u} = f b'(\rho) \quad \text{in } C_0^\infty(\mathbb{R}^d \times (0, T))'.$$

(ii) *If $f = 0$, then (F.4) holds for any b satisfying (F.1)–(F.3).*

Next we state an extension of the previous result to a class of piecewise C^1 functions b . Specifically, for $k > 0$, let

$$b_k(s) = \begin{cases} b(s) & \text{if } s \in [0, k), \\ b(k) & \text{if } s \in [k, \infty), \end{cases} \quad \text{where } b \in C^1[0, \infty).$$

Let $(b_k)'_+$ denote the right-derivative of b_k .

Lemma F.3 (Lemma 6.11 in [58]). *Let β , ρ , \underline{u} satisfy the assumptions of Lemma F.2 and let $f \in L^1_{\text{loc}}(\mathbb{R}^d \times (0, T))$ be such that*

$$\frac{\partial \rho}{\partial t} + \nabla_x \cdot (\underline{u} \rho) = f \quad \text{in } C_0^\infty(\mathbb{R}^d \times (0, T))'.$$

Then, part (i) of Lemma F.2 holds with b replaced with b_k and b' replaced with $(b_k)'_+$, for every $k > 0$. Furthermore, if $f = 0$, then part (ii) of Lemma F.2 holds, with the regularity assumption on $b \in C^1[0, \infty)$ appearing in the definition of b_k replaced with the relaxed assumption (F.1), for every $k > 0$.

G. CONTINUITY EQUATION WITH DISSIPATION

Suppose that Ω is a bounded domain in \mathbb{R}^d and $T > 0$. We shall consider the following parabolic initial-boundary-value problem, with homogeneous Neumann boundary condition:

$$(G.1a) \quad \frac{\partial z}{\partial t} - \alpha \Delta_x z = h \quad \text{in } \Omega_T := \Omega \times (0, T),$$

$$(G.1b) \quad z(\tilde{x}, 0) = z_0(\tilde{x}) \quad \text{for } \tilde{x} \in \Omega,$$

$$(G.1c) \quad \nabla_x z \cdot \tilde{n} = 0 \quad \text{on } \partial\Omega \times (0, T),$$

where $\alpha > 0$, z_0 and h are given functions defined on Ω and Ω_T , respectively, and z is an unknown function defined on Ω_T .

Lemma G.1 (Lemma 7.37 in [58]). *Let $\theta \in (0, 1]$, $r, s \in (1, \infty)$, and suppose that Ω is a bounded domain in \mathbb{R}^d . Suppose further that*

$$\Omega \in C^{2, \theta}, \quad \rho_0 \in \widetilde{W}^{2-\frac{2}{r}, s}(\Omega), \quad h \in L^r(0, T; L^s(\Omega)),$$

where $\widetilde{W}^{2-\frac{2}{r}, s}(\Omega)$ is the completion of the space $\{\zeta \in C^\infty(\overline{\Omega}) : \nabla_x \zeta \cdot \tilde{n}|_{\partial\Omega} = 0\}$ in $W^{2-\frac{2}{r}, s}(\Omega)$. Then, there exists a unique $z \in L^r(0, T; W^{2, s}(\Omega)) \cap C([0, T]; W^{2-\frac{2}{r}, s}(\Omega))$, $\frac{\partial z}{\partial t} \in L^r([0, T]; L^s(\Omega))$ satisfying equation (G.1a) a.e. on Ω_T , equation (G.1b) a.e. on Ω , equation (G.1c) in the sense of normal traces a.e. on $[0, T]$, and which satisfies the estimate

$$\begin{aligned} & \alpha^{1-\frac{1}{r}} \|z\|_{L^\infty(0, T; W^{2-\frac{2}{r}, s}(\Omega))} + \left\| \frac{\partial z}{\partial t} \right\|_{L^r(0, T; L^s(\Omega))} + \alpha \|z\|_{L^r(0, T; W^{2, s}(\Omega))} \\ & \leq C(r, s, \Omega) \left[\alpha^{1-\frac{1}{r}} \|z_0\|_{W^{2-\frac{2}{r}, s}(\Omega)} + \|h\|_{L^r(0, T; L^s(\Omega))} \right]. \end{aligned}$$

We shall also require the following regularity result concerning the case when the right-hand side h of (G.1a) is in divergence form, i.e., $h = \nabla_x \cdot \underline{b}$.

Lemma G.2 (Lemma 7.38 in [58]). *Let $\theta \in (0, 1]$, $r, s \in (1, \infty)$, and suppose that Ω is a bounded domain in \mathbb{R}^d . Suppose further that*

$$\Omega \in C^{2,\theta}, \quad z_0 \in L^s(\Omega), \quad \underline{b} \in L^r(0, T; L^s(\Omega)).$$

Then, there exists a unique $z \in L^r(0, T; W^{1,s}(\Omega)) \cap C([0, T]; L^s(\Omega))$, such that

$$\frac{d}{dt} \int_{\Omega} z \eta \, d\mathbf{x} + \alpha \int_{\Omega} \nabla_x z \cdot \nabla_x \eta \, d\mathbf{x} = - \int_{\Omega} \underline{b} \cdot \nabla_x \eta \, d\mathbf{x} \quad \text{in } (C_0^\infty[0, T])' \quad \forall \eta \in C^\infty(\overline{\Omega}),$$

$z(\underline{x}, 0) = z_0(\underline{x})$ for a.e. $\underline{x} \in \Omega$, and

$$\alpha^{1-\frac{1}{r}} \|z\|_{L^\infty(0, T; L^s(\Omega))} + \alpha \|\nabla_x z\|_{L^r(0, T; L^s(\Omega))} \leq C(r, s, \Omega) [\alpha^{1-\frac{1}{r}} \|z_0\|_{L^s(\Omega)} + \|\underline{b}\|_{L^r(0, T; L^s(\Omega))}].$$

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DEPARTMENT OF MATHEMATICS, IMPERIAL COLLEGE LONDON, LONDON SW7 2AZ, UK, jwb@imperial.ac.uk

MATHEMATICAL INSTITUTE, UNIVERSITY OF OXFORD, OXFORD OX2 6GG, UK, endre.suli@maths.ox.ac.uk