

Decoupling the AGN outflow and star-forming disc kinematics in the nuclear region of NGC 7582 with *JWST* NIRSpec and MIRI/MRS

Oscar Veenema ¹,^{*} Niranjan Thatte ¹, Dimitra Rigopoulou ^{1,2}, Ismael García-Bernete,³ Almudena Alonso-Herrero,³ Miguel Pereira-Santaella,⁴ Anelise Audibert,^{5,6} Enrica Bellocchi,^{7,8} Andrew J. Bunker ¹, Steph Campbell ⁹, Francoise Combes,¹⁰ Richard I. Davies ¹¹, Fergus R. Donnan ¹², Santiago García-Burillo,¹³ Omaira Gonzalez Martin ¹⁴, Laura Hermosa Muñoz ³, Erin K. S. Hicks,^{15,16,17} Sebastian F. Hoenig,¹⁸ Alvaro Labiano,¹⁹ Nancy A. Levenson,²⁰ Chris Packham,^{17,21} Cristina Ramos Almeida ^{5,6}, Claudio Ricci,^{22,23} Rogemar A. Riffel ^{24,25}, David Rosario,⁹ Taro Shimizu ¹¹ and Lulu Zhang ¹⁷

Affiliations are listed at the end of the paper

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ABSTRACT

We present a detailed study of the inner regions of NGC 7582, a nearby Seyfert 2 galaxy, from the Galaxy Activity, Torus, and Outflow Survey (GATOS). The galaxy hosts a circumnuclear star-forming disc and an active galactic nucleus (AGN)-driven biconical ionized outflow. Using *James Webb Space Telescope* Near-Infrared Spectrograph (NIRSpec) and Mid-Infrared Instrument/Medium-Resolution Spectrometer (MIRI/MRS) integral-field spectroscopy, we analyse ionic emission lines spanning a wide range of ionization potentials (IPs, ~ 8 –126 eV). Gaussian line-profile fitting reveals kinematic stratification: low-IP species ($\lesssim 20$ eV; e.g. [Fe II], [Ar II], and [Ne II]) trace ordered disc rotation with PA $\sim -12 \pm 3^\circ$, while high-IP species ($\gtrsim 35$ eV; e.g. [O IV], [Mg IV], and [Ne V]) follow the outflow with PA $\sim 54 \pm 10^\circ$. Outflowing gas exhibits systematically higher velocity dispersions (119 ± 13 km s⁻¹) than the disc (78 ± 11 km s⁻¹), consistent with turbulent or bulk motions. Intermediate-IP lines, [S III], [Ar III], and [Ne III], show contributions from both components, with the outflow characterized by higher dispersion, lower amplitude, and higher velocities in double-Gaussian fits. For these lines, a thin inclined disc plus 1D outflow model enables robust separation and quantification of the disc and outflow velocity fields. The outflow is consistent with a hollow bicone capable of accelerating gas beyond the local escape velocity, implying most material is unlikely to be re-accreted. The ionization cone opening angle shows no dependence on IP, indicating the AGN torus polar regions are largely unobscured. Our study provides new insights into AGN-driven outflows and circumnuclear disc dynamics, offering a framework to disentangle overlapping interstellar medium kinematics in nearby active galaxies.

Key words: galaxies: active – galaxies: individual: NGC 7582 – galaxies: kinematics and dynamics – galaxies: nuclei – galaxies: Seyfert – infrared: galaxies.

1 INTRODUCTION

In the central regions of galaxies, the kinematics of the gas in the interstellar medium (ISM) can deviate markedly from the ordered disc-like rotation observed on galactic scales (e.g. M. Mingozi et al. 2019; I. García-Bernete et al. 2021; K. F. Heckler, T. V. Ricci & R. A. Riffel 2022; D. Esparza-Arredondo et al. 2025; L. Lin et al. 2025). These inner zones are dynamically complex, shaped by processes such as active galactic nuclei (AGNs) feedback, winds from nuclear and circumnuclear star formation, and gravitational interactions from mergers, to list but a few.

AGN can launch powerful winds and jets that ionize and heat the surrounding ISM, while nuclear or circumnuclear starbursts generate supernova-driven feedback. Together these mechanisms can produce gaseous streaming motions (inflows and outflows, e.g. C. Ciccone et al. 2014; C. Harrison et al. 2014), shocks (e.g. K.-Y. Huang et al. 2022; R. Davies et al. 2024), and enhanced turbulence (e.g. J. Mullaney et al. 2013; G. Venturi et al. 2021; L. Ulivi et al. 2024) within the ISM. This often creates a complex kinematic environment which can directly affect the wider structure and composition of the galaxy as a whole (P. Guillard et al. 2015; C. M. Harrison & C. Ramos Almeida 2024), highlighting the interplay between AGN feedback and galaxy evolution.

ISM kinematics can be studied through the analysis of emission lines, each tracing different phases, stratified by quanti-

* E-mail: oscar.veenema@physics.ox.ac.uk

ties such as gas temperature and level of ionization. Molecular gas, primarily composed of molecular hydrogen (H_2) and carbon monoxide (CO), is particularly well studied. CO traces the coldest molecular phase (M. Aravena et al. 2010; C. Ramos Almeida et al. 2022; G. Jones et al. 2023), while H_2 spans a broader and warmer temperature range (D. Rigopoulou et al. 2002; H. Roussel et al. 2007; A. Togi & J. D. T. Smith 2016). Owing to its abundance, H_2 is especially valuable for probing molecular gas on nuclear and circumnuclear scales (M. Rosenberg, P. van der Werf & F. Israel 2013; M. Pereira-Santaella et al. 2022; J. Costa-Souza et al. 2024; R. Davies et al. 2024; R. A. Riffel et al. 2025; D. Esparza-Arredondo et al. 2025). Beyond the molecular phase, atomic hydrogen emits a series of recombination lines which allows us to trace star formation (R. C. Jr Kennicutt 1998; D. Calzetti et al. 2007; M. S. Z. de Mellos et al. 2024). Also important are atomic (ionic) emission lines originating from heavy elements. These lines span a wide range of ionization potentials (IPs), allowing different photon energies to be probed. Since each feedback mechanism interacts with and ionizes the ISM differently, these metal lines originate from and serve as sensitive kinematic tracers across a wide range of ionization energies, making them invaluable diagnostics for understanding and disentangling the roles of AGN and star formation on the structure and dynamics at the centres of galaxies (L. J. Kewley et al. 2006; J. Ma et al. 2021).

Many atomic lines lie in the near- and mid-infrared (IR), a wavelength range ($\lambda_{\text{rest}} \sim 1 - 30 \mu\text{m}$) challenging to observe from the ground due to strong atmospheric absorption and high thermal background. Resolving the complex kinematics of the nuclear regions of galaxies requires high spatial resolution (typically tens of parsecs or better), making the *James Webb Space Telescope* (*JWST*) Near-Infrared Spectrograph (NIRSpec) (P. Jakobsen et al. 2022; T. Böker et al. 2023) and Mid-Infrared Instrument/Medium-Resolution Spectrometer (MIRI/MRS) (M. Wells et al. 2015; I. Argyriou et al. 2023) integral field spectrographs (IFS) exceptionally well suited for such work, especially for nearby galaxies.

In addition to being characterized by their IP, many emission lines in the IR have been shown to trace specific physical components or processes within galaxies, particularly on nuclear scales. For example, $[\text{Fe II}] 5.34 \mu\text{m}$ is widely used as a tracer of shocked gas, especially when its flux ratio is high relative to hydrogen recombination lines (K. Kawara, M. Nishida & Y. Taniguchi 1988; L. Colina et al. 2015; V. U et al. 2022; A. Alonso-Herrero et al. 2025). Similarly, $[\text{Mg IV}] 4.49 \mu\text{m}$ has been shown to trace direct AGN photoionization as well as faster shocks originating from more highly ionized regions (M. Pereira-Santaella et al. 2024). Several IR lines also serve as diagnostics of recent star-formation, including $[\text{Ar II}] 6.99 \mu\text{m}$ and $[\text{Ne II}] 12.81 \mu\text{m}$ (L. C. Ho & E. Keto 2007; M.-Y. Zhuang, L. C. Ho & J. Shangquan 2019; C. Whitcomb et al. 2020; J. Young et al. 2023), while higher ionization neon lines in the mid-IR, such as $[\text{Ne III}] 15.56 \mu\text{m}$, $[\text{Ne V}] 14.32 \mu\text{m}$, and $[\text{Ne VI}] 7.65 \mu\text{m}$ also provide a means to distinguish between different excitation mechanisms (A. Feltre et al. 2023; L. Zhang et al. 2025). Notably, $[\text{Ne II}]$ and $[\text{Ne III}]$ trace star-forming regions, with $[\text{Ne III}]$ also being enhanced in more energetic or turbulent environments, whereas $[\text{Ne V}]$ and $[\text{Ne VI}]$, with significantly higher IPs, are exclusively associated with AGN activity (E. Sturm et al. 2002; M. Pereira-Santaella et al. 2010b; L. Hermosa Muñoz et al. 2025). Additional mid-IR lines such as $[\text{S IV}] 10.51 \mu\text{m}$ and $[\text{O IV}] 25.89 \mu\text{m}$ are also strong tracers of AGN ionization and can highlight features such as ionized outflows (D. Dicken et al. 2014).

Numerous studies have utilized *JWST*/NIRSpec and MIRI/MRS IFS to investigate the nuclear regions of nearby galaxies (e.g. I. García-Bernete et al. 2022, 2024; L. Armus et al. 2023; F. Donnan et al. 2023; A. Alonso-Herrero et al. 2024; L. Hermosa Muñoz et al. 2024; M. Ceci et al. 2025; C. Marconcini et al. 2025b; L. M. Feuillet et al. 2025; R. A. Riffel et al. 2026; C. Ramos Almeida et al. 2025; O. Veenema et al. 2025).

Of particular relevance to this study, as they probe the kinematics of different ISM ionization phases as a function of IP, R. Davies et al. (2024) studied the nucleus of NGC 5728 with MIRI/MRS, demonstrating that molecular gas is being depleted by outflows along an AGN-driven ionization cone. They also showed that the kinematic major axis differs when traced by mid-IR $[\text{Ne II}]$, $[\text{Ne III}]$, $[\text{Ne V}]$, and $[\text{Ne VI}]$ lines, with the higher-IP lines exhibiting clear offsets relative to lower-IP lines due to the different position angles (PAs) of the ordered disc rotation and ionization cone (this was also shown in several other Seyfert galaxies by L. Zhang et al. 2024; L. Hermosa Muñoz et al. 2024).

Also of particular interest to this study, L. Ulivi et al. (2025) analysed the inner ~ 1 kpc of Arp 220 using *JWST* NIRSpec IFS. By examining multiple near-IR emission lines and fitting multi-Gaussian profiles, they disentangled the kinematics of the ISM, showing two counter-rotating discs, each associated with one of the nuclei, as well as two distinct outflows, each driven by an active nucleus.

An important factor in resolving and interpreting kinematic signatures in galaxies is the viewing geometry, specifically, the galaxy's inclination and our line of sight to its morphological components (such as nuclear rings, discs, jets, and outflows). The prevalence and morphology of these features across the galaxy population remain uncertain, motivating the need for detailed case-by-case studies. One such example is NGC 7582, a galaxy with a high inclination ($i \sim 58^\circ$) and a morphologically complex nuclear region. As we will show, NGC 7582 exhibits a pronounced kinematic dichotomy between low- and high-ionization circumnuclear gas, which is distinguishable in part due to our viewing angle.

The remainder of this paper is structured as follows. In Section 2, we review previous studies of NGC 7582. Section 3 describes the data collection, reduction, and analysis methods. Section 4 gives our combined results and discussion, with Section 4.1 presenting results from single-Gaussian fitting to many mid-IR emission lines across the nuclear and circumnuclear regions, showing strong variation between disc and outflow tracing lines. Section 4.2 extends this analysis to double-Gaussian fits and discusses the decoupling of the disc and outflow kinematics in three intermediate-IP lines. Then in Section 4.3, we apply a physically motivated thin inclined disc model to the intermediate-IP lines to isolate the outflow kinematic signature, while in Section 4.4 we instead apply a simple 1D outflow model to isolate the disc rotation. Building on these, Section 4.5 introduces a hybrid approach that simultaneously models both disc and outflow components. In Section 4.6, we consider the expulsion of gas in the outflow. Then in Section 4.7, we measure the opening angle of the front facing ionization cone in order to constrain the morphology of the AGN dusty torus in NGC 7582. Finally, Section 5 summarizes our findings.

2 NGC 7582

NGC 7582 is a nearby Seyfert 2 galaxy at redshift $z \sim 0.00525$, corresponding to a distance of $D \sim 22$ Mpc (assuming a flat Λ -

cold dark matter Universe with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_m = 0.3$), therefore $1 \text{ arcsec} \sim 100 \text{ pc}$. Previous studies have identified several key nuclear and circumnuclear features, including a star-forming disc/ring at a radius of $\sim 200 \text{ pc}$ (R. A. Riffel et al. 2009; A. Alonso-Herrero et al. 2020; S. García-Burillo et al. 2021; S. Juneau et al. 2022), composed of clumps with elevated star-formation containing roughly ~ 500 O-type stars each, with a star formation rate of $0.23 - 0.28 M_\odot \text{ yr}^{-1}$, R. A. Riffel et al. 2009) to the north and south of a Compton-thick central AGN (I. García-Bernete et al. 2016; T. Ricci et al. 2018), with a hydrogen column density of $N_{\text{H}} \sim 10^{23} - 10^{24} \text{ cm}^{-2}$ (E. Rivers et al. 2015). The galaxy hosts a prominent bar (R. A. Riffel et al. 2009), and the $\sim 200 \text{ pc}$ radius circumnuclear ring is likely located near the inner Lindblad resonance, where bar-driven gas inflows accumulate and trigger the observed enhanced star formation (M. C. Sormani, E. Sobacchi & J. L. Sanders 2024).

Within $\sim 50 \text{ pc}$ of the nucleus, a marked depletion of molecular gas has been observed, while a dusty, molecular-rich torus is found on scales of $\gtrsim 10 \text{ pc}$ surrounding the AGN (S. García-Burillo et al. 2021, 2024). Additionally, a pair of ionization cones originating from the AGN and extending over $\gtrsim 3 \text{ kpc}$ has been confirmed, containing highly ionized, outflowing gas (S. Morris et al. 1985; S. Juneau et al. 2022).

In Fig. 1, we present integrated flux maps (see Section 3.2 for details) of the mid-IR [Ar II] $15.8 \mu\text{m}$, [Ar III] $8.99 \mu\text{m}$, and [Ne V] $14.32 \mu\text{m}$ emission lines from *JWST* MIRI/MRS (see Section 4.1), with annotations highlighting the key morphological components of the nuclear and circumnuclear regions of NGC 7582. These include the circumnuclear star-forming disc/ring, the embedded southern star-forming clumps, the AGN position, and the edges of the biconical ionized outflow. The figure serves as a visual guide to the spatial distribution of these structures, which are traced to varying degrees across the full set of emission lines analysed in this work, and is intended to aid the interpretation of the flux and kinematic maps presented throughout the paper.

NGC 7582 has also been shown to have strong radio emission (D. A. Forbes & M. J. Ward 1993; M. Orienti & M. Prieto 2010) from its AGN out to $\sim 4 \text{ arcsec}$ north and $\sim 2 \text{ arcsec}$ south. This emission is rather diffuse, with current debate on whether it could be originated from intense starburst clumps in the circumnuclear disc, or a rather young radio jet (11 000–23 000 yr old, S. Juneau et al. 2022).

Of particular relevance is the work of T. Ricci et al. (2018), who analysed optical Gemini Multi-Object Spectrograph (GMOS) (J. Allington-Smith et al. 2002) and near-IR Spectrograph for Integral Field Observation in the Near-Infrared (SINFONI) (F. Eisenhauer et al. 2003) observations of NGC 7582. They showed that the low-ionization gas in the circumnuclear ring/disc, traced by $\text{H}\alpha$ and [N II] 6584\AA , follows the galactic disc rotation described by S. Morris et al. (1985), with a major axis at $\text{PA} \sim 0 \pm 5^\circ$. In contrast, the high-ionization [O III] 5007\AA line traces the biconical ionized outflow, oriented at a different PA of $\sim 47 \pm 5^\circ$. T. Ricci et al. (2018) also proposed a model of the circumnuclear structure of NGC 7582, comprising of the highly inclined circumnuclear ring surrounding the AGN torus, with the biconical outflow emerging from within the ring and oriented approximately orthogonal to it. They further incorporated the putative radio jet oriented nearly perpendicular to the ionization cones (see their fig. 17).

Also noteworthy is the study by S. Juneau et al. (2022), who used Very Large Telescope/Multi Unit Spectroscopic Explorer

(VLT/MUSE) (R. Bacon et al. 2010) observations of NGC 7582 to show that its circumnuclear star-forming ring/disc constitutes a kinematically distinct core (KDC), on a scale of $\sim 600 \text{ pc}$. The stars in the KDC corotate with the main disc of the galaxy along a similar kinematic axis (which they found had a $\text{PA} \sim -23^\circ$), but exhibit a much higher rotation velocity that rises steeply before peaking along the major axis. The low-ionization gas kinematics in the ring/disc, traced by $\text{H}\alpha$ and [N II] 6584\AA , were found to closely follow the stellar component, rotating at the same PA and with similar velocity.

S. Juneau et al. (2022) also examined the kinematics of the ionized conical outflow, again with [O III] 5007\AA . They found the outflow to have edges at $\text{PA} \sim 15^\circ$ and $\sim 115^\circ$, originating from the AGN torus and emerging through the circumnuclear ring/disc, consistent with Atacama Large Millimeter/submillimeter Array (ALMA) observations by S. García-Burillo et al. (2021). Their analysis supports a biconical outflow with hollow ionization cones, where enhanced emission along the cone edges arises from limb brightening, indicating that the cone interior contains more diffuse gas. They further identified evidence of gas at the cone edges moving more slowly as it interacts with the surrounding medium, possibly producing shocks. The redshifted side of the outflow shows weaker emission, attributed to obscuration by dust lanes.

These previous investigations of the gas kinematics in the nuclear regions of NGC 7582 (T. Ricci et al. 2018; S. Juneau et al. 2022) have primarily relied on optical and near-IR observations from ground-based facilities. While highly informative, such studies are limited by strong dust extinction at shorter wavelengths and by challenges inherent to ground-based observing, including atmospheric absorption and turbulence. By contrast, a mid-IR study is far less affected by dust obscuration.

In a previous study (O. Veenema et al. 2025), we investigated the mid-IR molecular H_2 emission within the central $\sim 200 \text{ pc}$ of NGC 7582, finding that it follows the rotation of the circumnuclear disc/ring. We also showed evidence of heating by slow shocks ($v_s \sim 10 \text{ km s}^{-1}$) in the southern clumps of the circumnuclear star-forming disc, likely driven by intense star formation. Altogether, these earlier results highlight that the kinematics of the central $\sim \text{kpc}$ of NGC 7582 are highly intricate and merit dedicated, detailed study.

3 METHOD, DATA COLLECTION, AND REDUCTION

3.1 Data collection and reduction

The data used here are part of the Galaxy Activity, Torus, and Outflow Survey (GATOS; S. García-Burillo et al. 2021; A. Alonso-Herrero et al. 2021). These data are part of the *JWST* Cycle 2 GO proposal ID 3535 (PIs: I. García-Bernete and D. Rigopoulou), which were originally presented and the data reduction discussed in O. Veenema et al. (2025). This study uses near-IR to mid-IR ($2.87\text{--}28.1 \mu\text{m}$) observations using integral-field spectrographs MIRI/MRS ($4.9\text{--}27.9 \mu\text{m}$) with a spectral resolution of $R \sim 1300\text{--}3700$ (A. Labiano et al. 2021) ($\sigma_{\text{inst}} \sim 35 - 97 \text{ km s}^{-1}$), acquired on 2023 October 31, and NIRSpec with the grating-filter pairs *G395H/F290LP* ($2.87\text{--}5.27 \mu\text{m}$) with $R \sim 2700$ (P. Jakobsen et al. 2022; T. Böker et al. 2022) ($\sigma_{\text{inst}} \sim 47 \text{ km s}^{-1}$), acquired on 2024 April 7.

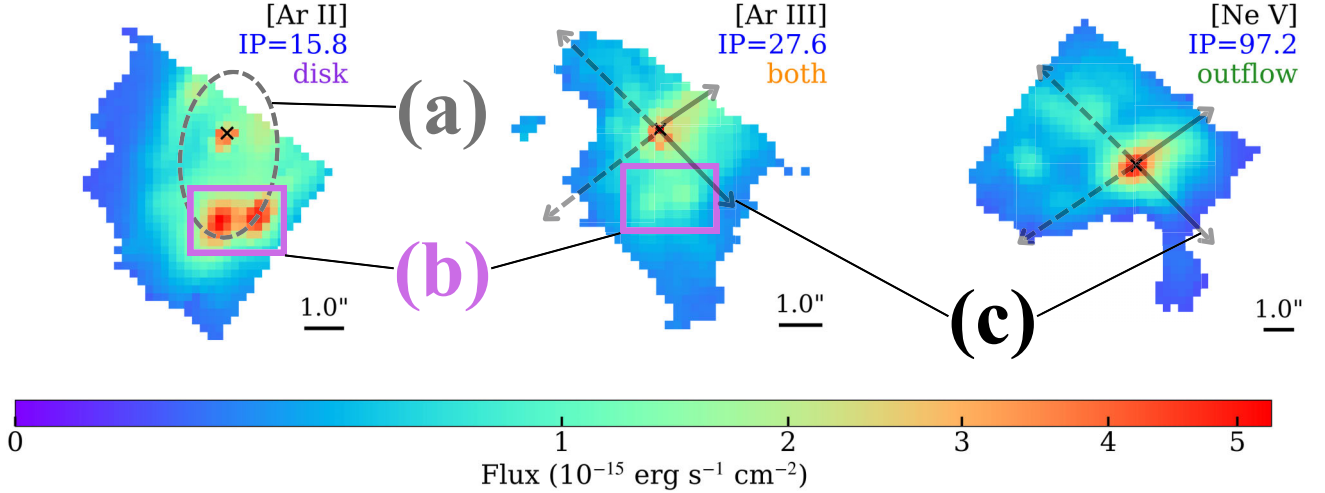


Figure 1. Integrated flux maps of the [Ar II] 15.8 μm , [Ar III] 8.99 μm , and [Ne V] 14.32 μm emission lines, annotated to highlight the main nuclear and circumnuclear structures in NGC 7582. We give the line name, IP, and associated kinematic trace (see Table 1) at the top right of each flux map. North is up, and east is to the left. The AGN position (defined as the peak continuum emission in each channel) is marked by a cross. (a) Outline of the circumnuclear disc/ring: a highly inclined ($i \sim 58^\circ$; M. Wold et al. 2006; S. Garcia-Burillo et al. 2021) structure rich in molecular gas and hosting multiple star-forming regions. This component is most clearly traced by the low-IP emission lines. (b) Southern circumnuclear star-forming clumps: prominent regions embedded within the disc/ring, identified as SF 1 and SF 2 by O. Veenema et al. (2025, also referred to as M1 and M2 by M. Wold & E. Galliano 2006; T. Ricci et al. 2018). These clumps are bright in low- to intermediate-IP tracers. (c) Ionized biconical outflow: the AGN-driven, limb-brightened cone structure, with solid arrows indicating the front-facing (blueshifted, western) cone edges and dashed arrows marking the receding (eastern) side. The outflow is primarily traced by mid- and high-IP lines, appearing as diffuse extensions along the arrowed directions from the nucleus, with the western cone more prominently detected. As with Fig. 3, this figure illustrates how different emission lines preferentially trace distinct physical components of the galaxy, depending on their IP, and provides a visual guide to interpreting the kinematic and flux maps presented throughout this work.

3.2 Calculating emission-line flux and kinematics

In this study, we focus on the kinematics of various emission lines in the mid-IR. For each line, per spaxel, we first masked the spectral emission region and fitted a third-order polynomial out to 0.02 μm of the continuum on either side. This continuum fit was then subtracted from all wavelength bins to isolate the emission line.

We then masked out spaxels which had a line peak signal-to-noise (S/N) less than 3, and those with emission-line fluxes that were lower than the 60th percentile of the flux of that specific emission line across every spaxel. This approach retains the brightest 40 per cent of spaxels for each line, preferentially selecting regions with sufficient flux for reliable Gaussian kinematic fitting. This gives a variable S/N cut for each emission line, with the minimum S/N (i.e. the S/N at the 60th flux percentile) being > 6 , but typically much higher, for all emission lines. This additional flux threshold is important because the velocity and velocity dispersion maps are not flux-weighted, and without such a cut, low-flux, or more noisy spaxels would contribute equally to the kinematic analyses presented here. This threshold therefore balances statistical robustness and spatial coverage: lower cuts admit more noise-dominated spaxels, while higher cuts overly restrict the field of view (FoV) and reduce the number of spaxels available for analysis.

For each emission line in every remaining (unmasked) spaxel, we fit a single-Gaussian profile, with the best-fit central wavelength (λ) used to derive the gas velocity, $v = c \times (\lambda - \lambda_0)/\lambda_0$, where c is the speed of light and λ_0 is the rest-frame wavelength of the emission line. We also attempted double-Gaussian fits across all lines. However, we found that for many emission lines, double-Gaussian functions did not statistically improve the fit compared

to a single-Gaussian function. These double-Gaussian fits were implemented under various model assumptions, discussed in detail later. All single- and double-Gaussian fitting was performed using Levenberg–Marquardt (LM) optimization via the CAPFIT PYTHON library (M. Cappellari 2023). We show several example individual spaxel fits in Appendix A in Figs A1–A3.

We generated continuum subtracted integrated flux emission and velocity maps for each emission line using the native resolution and FoV of the NIRSspec or MIRI/MRS datacubes, without reprojection or convolution. In this study, we are analysing many different emission lines that originate from different instruments (NIRSspec or MIRI/MRS) or MIRI/MRS channels, each with slightly different spaxel scales, resolutions, and FoVs. Hence, the spatial sampling and exact coverage vary slightly between lines. For clarity, and to prevent overcrowding in multipanel figures, we omit explicit RA and Dec. axes in most figures, however we do mark the central AGN position (as the position of the peak continuum flux in each channel) on every map as a point of reference and include scalebars on our main reference figures showing the single-Gaussian velocity, flux, and single-Gaussian velocity dispersion maps (Figs 2–4).

4 RESULTS AND DISCUSSION

4.1 Single-Gaussian fitting

We begin by broadly examining the kinematics of several emission lines detected across the combined NIRSspec and MIRI/MRS wavelength range for NGC 7582. These lines, along with their wavelengths, IPs, interpreted kinematic trace, major velocity axis position angle, and average velocity dispersion (σ) from single-Gaussian fitting are listed in Table 1. Fig. 2 presents the velocity

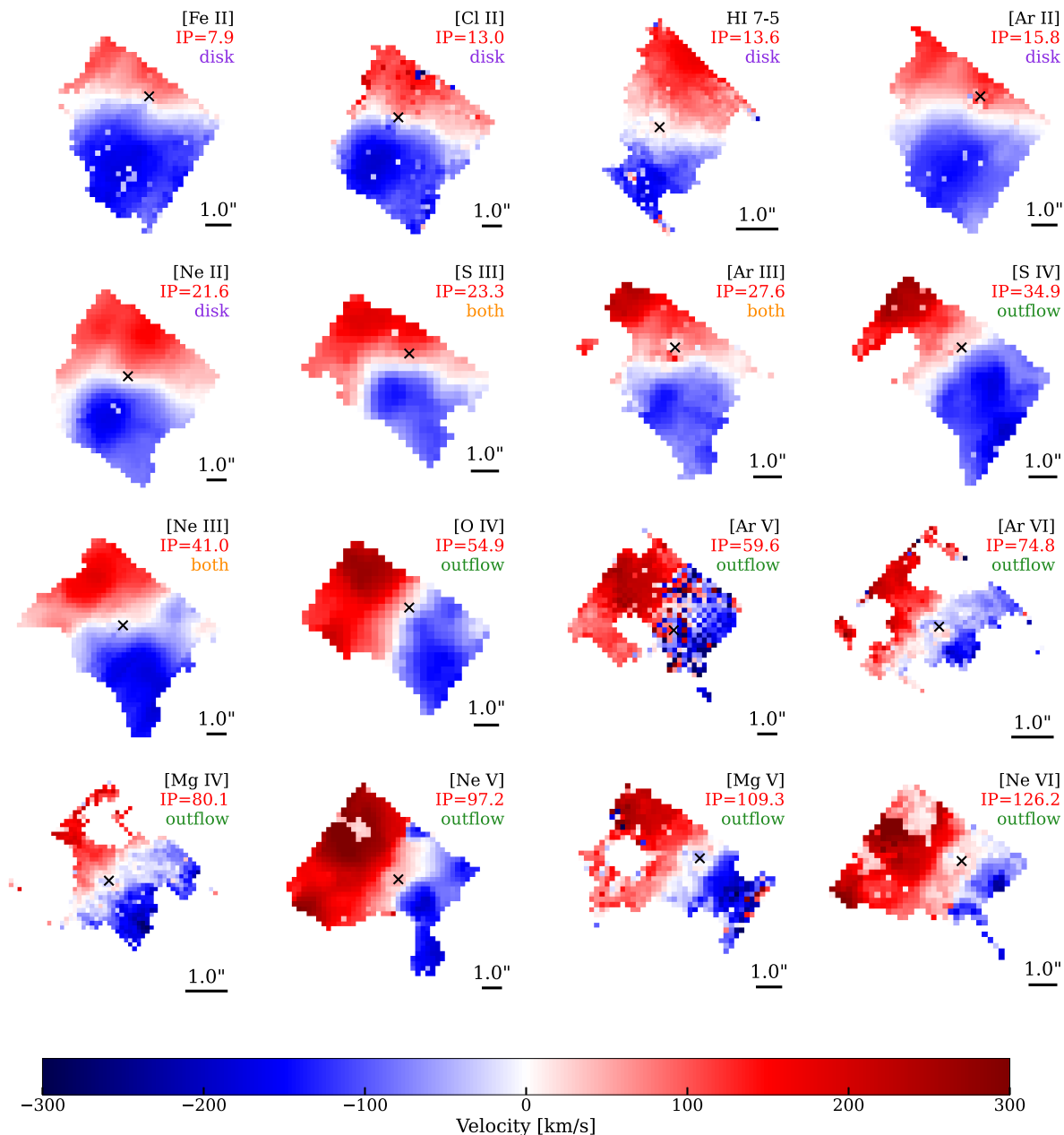


Figure 2. Velocity maps in the central ~ 200 pc region for various atomic emission lines in NGC 7582 fit using single-Gaussians ordered by IP (in eV) from left to right, top to bottom. Labels at the top right give the line name, IP, and kinematic trace for each velocity map. The lower IP lines typically trace the ordered circumnuclear star-forming ring/disc rotation, whereas the higher IP lines typically trace the biconical ionized outflow motion. North is up, and east is to the left. Each black cross denotes the AGN position, i.e. the photometric centre of the corresponding continuum. A 1 arcsec ~ 100 pc scale bar is included at the bottom right of each subplot.

maps from the single-Gaussian fits to the continuum-subtracted flux for each line, ordered by increasing IP. We also show the continuum subtracted line flux emission for each emission line in Fig. 3, as well as the velocity dispersion in Fig. 4.

Notably, when fitting single-Gaussian profiles to the emission lines and deriving the corresponding velocity maps, we observe a striking kinematic dichotomy between low-IP ($\lesssim 20$ eV) and high-IP ($\gtrsim 35$ eV) lines, as shown in Fig. 2. The low-IP lines exhibit a rotation pattern with blueshifted emission to the south of the AGN and redshifted emission to the north. In contrast, high-IP lines display a different kinematic axis, with blueshifted emission to the southwest and redshifted to the northeast.

Additionally, we analysed the kinematics of the pure rotational H_2 lines from S(1) to S(7), which trace the warm molecular gas. These lines all exhibit the same strongly disc-like rotation as the low-IP atomic lines. Several of these line kinematics are presented and discussed in detail in O. Veenema et al. (2025), and F. R. Donnan et al. (2026).

4.1.1 Low- and high-IP lines: single Gaussian

The kinematics of low-IP lines in the circumnuclear region in the optical and near-IR have been discussed in previous studies (S.

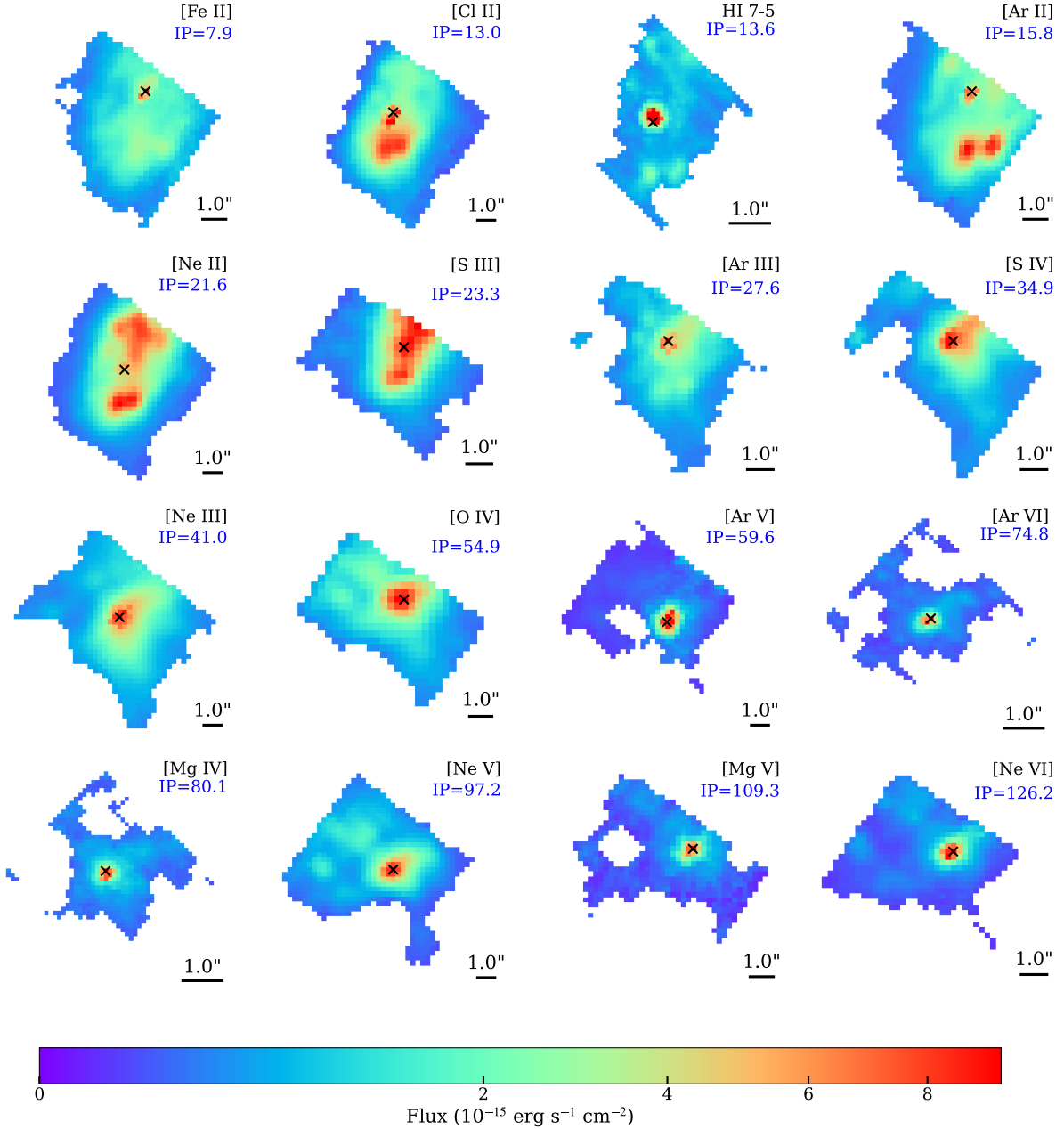


Figure 3. Continuum-subtracted flux maps in the central ~ 200 pc region for various atomic emission lines in NGC 7582 ordered by IP (in eV) from left to right, top to bottom. The lower IP lines typically trace the circumnuclear star-forming ring/disc, with multiple prominent clumps of star formation visible, whereas the higher IP lines typically trace gas excited by the biconical ionized outflow. North is up, and east is to the left. Each black cross denotes the AGN position, i.e. the photometric centre of the corresponding continuum. A 1 arcsec ~ 100 pc scale bar is included at the bottom right of each subplot.

Morris et al. 1985; R. A. Riffel et al. 2009; T. Ricci et al. 2018), which attributed the observed kinematics to rotation within the actively star-forming circumnuclear ring/disc. We too see evidence for this in our flux maps (Fig. 3), with many showing strong flux from the southern two prominent star-forming clumps (seen as the bright spots roughly 1 arcsec south of the AGN, particularly visible for [Ar II]), and some of the northern clumps (when resolved and within the FoV).

The high-IP line flux maps (Fig. 3) are dominated by compact emission from the nuclear region, consistent with photoionization by the AGN as the primary source of these lines. This is

particularly evident for the highest IP species ($IP \gtrsim 55$ eV), which are largely confined to the central spaxels. However, we also identify visual evidence for extended emission in many of these high-IP line flux maps. This is most notable in [Ne V], [S IV], [Ar VI], and [Mg IV], which exhibit faint but coherent structure extending westward from the nucleus. This extension spatially coincides with the blueshifted side of the ionized outflow (Fig. 1) and displays a more diffuse morphology than the compact nuclear emission, and is broadly consistent with a hollow-cone geometry. We therefore interpret the bulk of the high-IP emission as tracing AGN photoionization in the nucleus, while also noting evidence

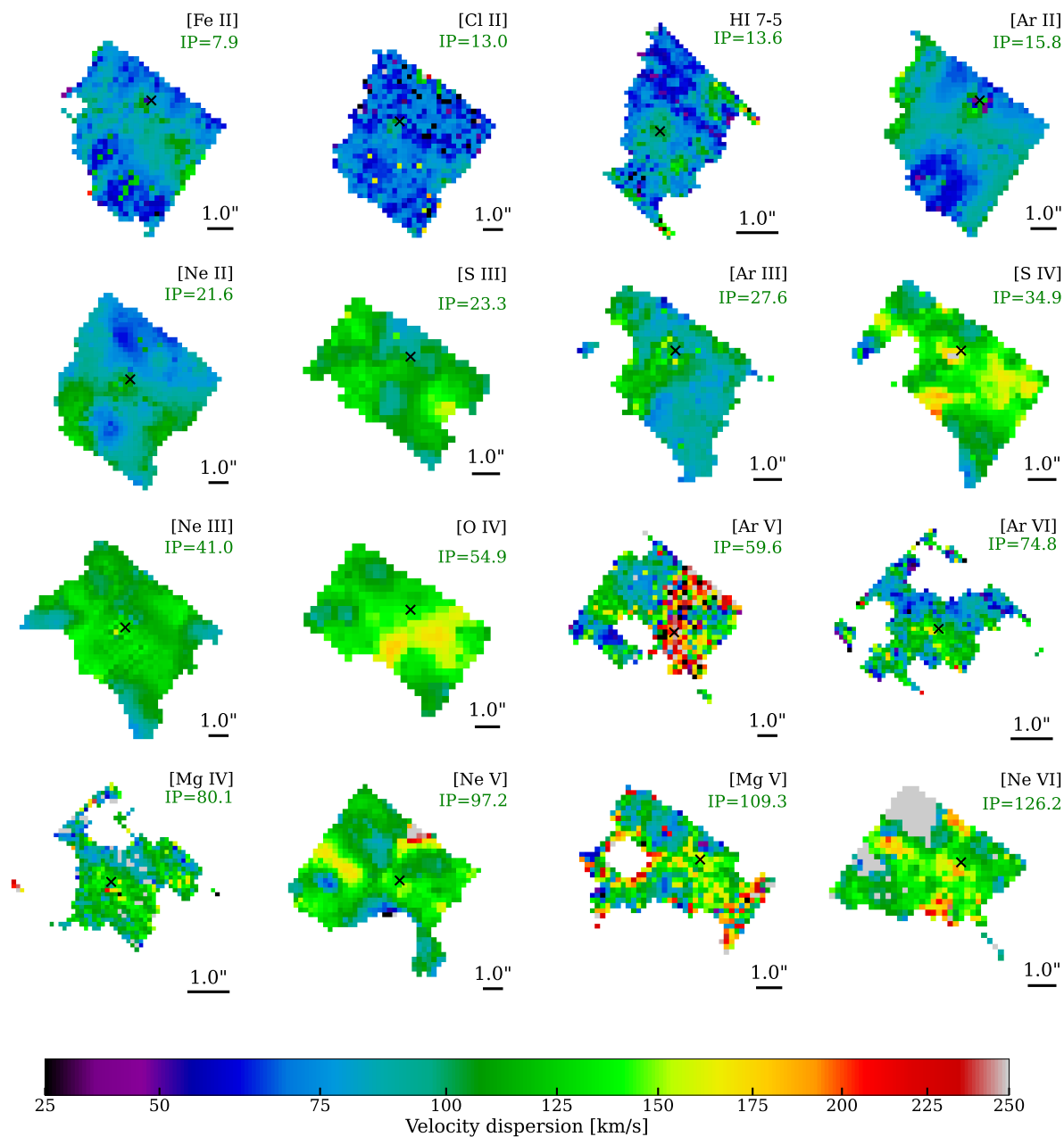


Figure 4. Velocity dispersion maps in the central ~ 200 pc region for various atomic emission lines in NGC 7582 ordered by IP (in eV) from left to right, top to bottom. These maps include the contribution from instrumental dispersion. The lower IP lines show lower velocity dispersions throughout the FoV, especially over the star-forming ring. The higher IP lines tend to have much larger velocity dispersions, particularly in the biconical outflow. Each black cross denotes the AGN position, i.e. the photometric centre of the corresponding continuum. A 1 arcsec ~ 100 pc scale bar is included at the bottom right of each subplot.

for an additional, spatially extended component likely associated with the outflow. The ionization cone is also slightly traced in the higher IP line velocity dispersion maps (Fig. 4) as regions of significantly enhanced dispersion (i.e. the more yellow and red regions at a PA $\sim 45^\circ$).

Moreover, S. Juneau et al. (2022) found the low-IP gas rotates with a kinematic major axis PA $\sim -23^\circ$ (in our notation going from red to blue shifting), while the high-IP [O III] line exhibited a PA $\sim 57^\circ$, with T. Ricci et al. (2018) reporting similar PAs ($\sim 0^\circ$ for the low-IP and $\sim 47^\circ$ for the high-IP). Additionally, S. García-Burillo et al. (2021) found that the cold molecular gas traced through CO 3–2 kinematics followed the disc rotation with a

PA $\sim -16^\circ$. We show this is consistent within uncertainties with the results we observe in the mid-IR by fitting for the kinematic major axis PA using PAFIT (D. Krajnović et al. 2006) on all single-Gaussian velocity maps in Fig. 2. We show this as a function of IP in Fig. 5 (left) and tabulate in Table 1. We estimate the kinematic PA uncertainty from the χ^2 minimization curve (from PAFIT) as the angular range over which χ^2 increases by one from its minimum (i.e. 1σ confidence), weighted by the velocity uncertainties.

Furthermore, S. Juneau et al. (2022) measured the ionized gas kinematics in the inner region using optical emission lines, finding that low-IP tracers such as H α closely follow the stellar disc rotation, with maximum rotational velocities of ~ 150

Table 1. Key fine structure transitions analysed in this study, including their IP (defined as the energy needed to reach the ionization stage producing that transition, A. Kramida, Y. Ralchenko & J. Reader 2022), rest-frame wavelength, the associated kinematic component traced, major velocity axis position angle (from PAFit), and median velocity dispersion (from single-Gaussian fitting).

Emission line	IP (eV)	Wavelength (μm)	Kinematic trace	Kinematic major axis PA (deg)	Average velocity dispersion (σ) (km s^{-1})
[Fe II]	7.9	5.34	Disc	-12 ± 13	76 ± 10
[Cl II]	13.0	14.37	Disc	-11 ± 9	64 ± 17
HI 7-5	13.6	4.65	Disc	-13 ± 8	71 ± 17
[Ar II]	15.8	6.99	Disc	-17 ± 19	85 ± 9
[Ne II]	21.6	12.81	Disc	-9 ± 11	92 ± 13
[S III]	23.3	18.71	Disc/outflow	13 ± 11	115 ± 14
[Ar III]	27.6	8.99	Disc/outflow	13 ± 13	94 ± 13
[S IV]	34.9	10.51	Outflow	58 ± 14	129 ± 17
[Ne III]	41.0	15.56	Disc/outflow	17 ± 9	115 ± 11
[O IV]	54.9	25.89	Outflow	42 ± 12	131 ± 10
[Ar V]	59.6	13.10	Outflow	62 ± 2	115 ± 31
[Ar VI]	74.8	4.53	Outflow	56 ± 8	92 ± 23
[Mg IV]	80.1	4.49	Outflow	62 ± 2	115 ± 25
[Ne V]	97.2	14.32	Outflow	65 ± 10	115 ± 22
[Mg V]	109.3	5.61	Outflow	43 ± 9	124 ± 32
[Ne VI]	126.2	7.65	Outflow	44 ± 12	131 ± 26

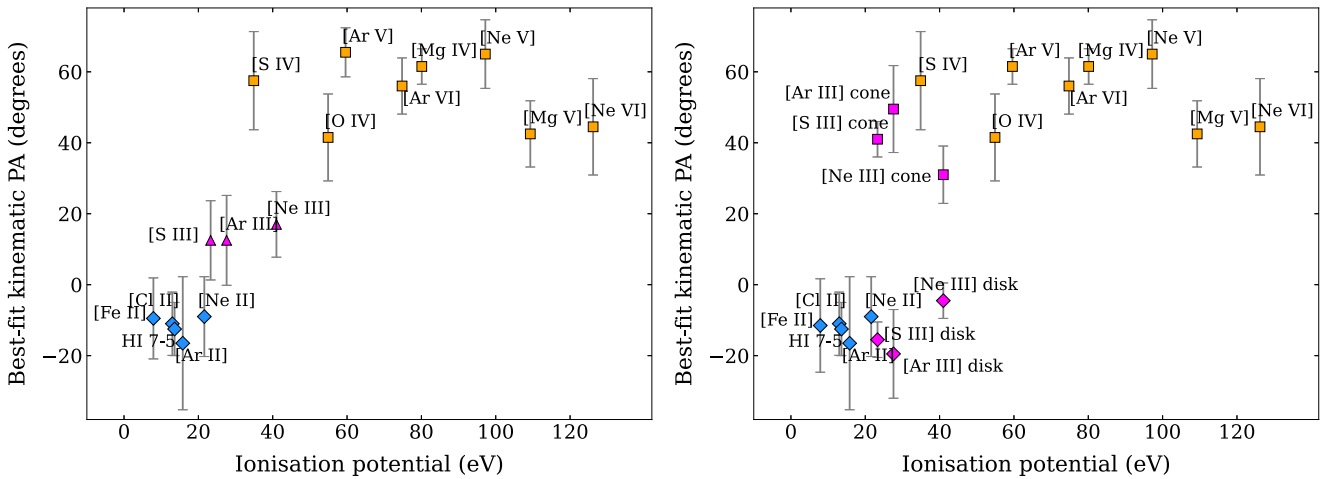


Figure 5. Left: best-fitting major kinematic axis PA versus IP for each of the lines from the single-Gaussian velocity maps fit using PAFIT (D. Krajnović et al. 2006). Errors are the 1σ errors. Blue diamond points are emission lines we denote as disc-tracing low-IP, magenta triangle points denote the mixed kinematics tracing ‘intermediate-IP’ lines, and orange square points denote outflow-tracing high-IP lines. The blue diamond points have an average PA $\sim -12 \pm 3^\circ$, whereas the orange square points have an average PA $\sim 54 \pm 10^\circ$. The magenta triangle points lie in the transition region between the other two groups and have an average PA $\sim 14 \pm 2^\circ$. Right: same as the left but with the intermediate-IP line PAs split into their disc component and outflow component velocity map PAs from double-Gaussian fitting (Fig. 7).

km s^{-1} . In contrast, the high-IP [O III] emission associated with the outflow reaches projected velocities up to $\sim 300 \text{ km s}^{-1}$. Our mid-IR measurements show broad agreement: the low-IP lines analysed here exhibit maximum velocities of $\sim 150\text{--}180 \text{ km s}^{-1}$ (e.g. see Table 2), while the high-IP lines reach comparable velocities along the ionization cone. Together with the agreement in kinematic major axis PAs described above, this comparison supports the interpretation that the low- and high-IP mid-IR lines trace the same components previously identified in the optical, providing a consistent picture across wavelength regimes.

Therefore, based on previous findings and physical interpretation, we propose that the low-IP lines are tracing gas rotating

within the circumnuclear star-forming disc, whereas the high-IP lines are tracing gas in the ionized biconical outflow. This classification is both physically motivated based on the emission and kinematic morphology, and supported by prior literature. Henceforth, we adopt the terms ‘disc-tracing’ and ‘cone/outflow-tracing’ as shorthand to describe the kinematic behaviour of respective emission lines.

We see that the disc-tracing lines, [Fe II], [Cl II], HI 7–5, [Ar II], and [Ne II], all exhibit nearly identical PAs, consistent with an average major axis PA $\sim -12 \pm 3^\circ$. The outflow-tracing lines, [S IV], [O IV], [Ar V], [Ar VI], [Mg IV], [Ne V], [Mg V], and [Ne VI] also exhibit similar PAs, albeit a different PA to the disc tracing lines, consistent with an average major axis PA $\sim 54 \pm 10^\circ$.

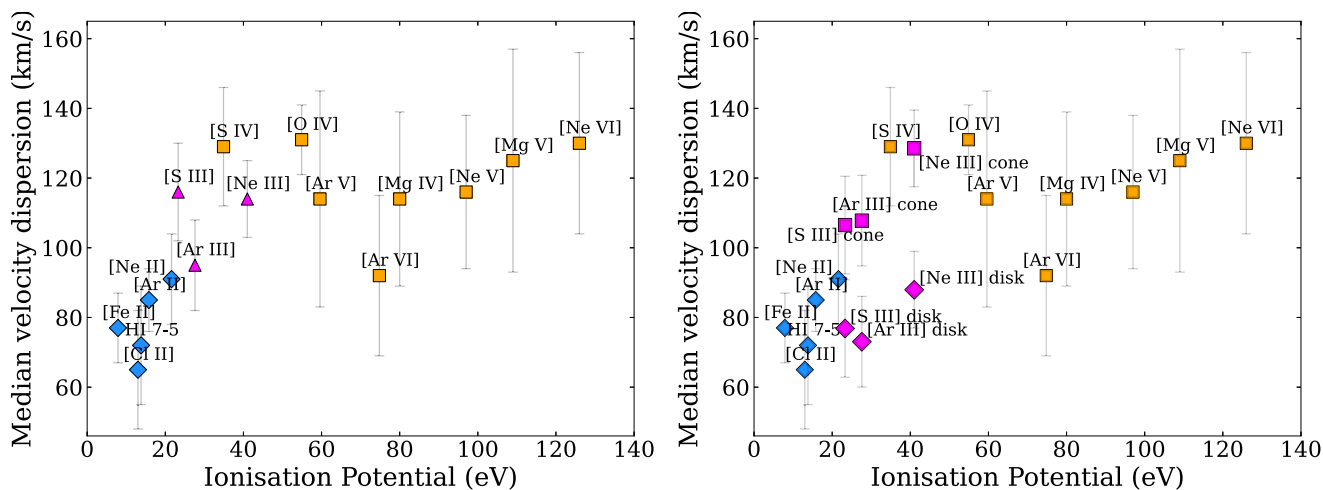


Figure 6. Left: median velocity dispersion of each emission line as a function of IP. We postulate that the blue diamond points lie roughly on a positive trend, whereas the orange square points display an approximately positive correlation within their errors. The magenta triangle points are in the transition region between the two distinct groups. Right: median velocity dispersion of each emission line, but with the intermediate-IP lines having both their median disc and outflow velocity dispersion found from double-Gaussian fitting (i.e. the dotted lines for each distribution shown in Fig. 8) plotted instead.

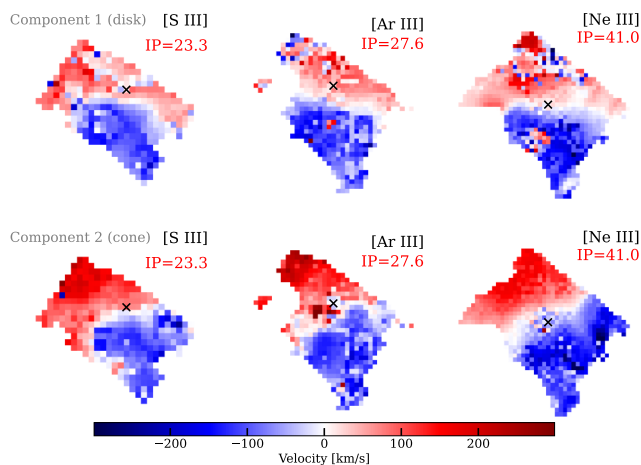


Figure 7. Velocity maps for the intermediate-IP emission lines in NGC 7582 fit using double-Gaussians ordered by IP (in eV). Top: disc Gaussian component (average major axis PA $\sim -13^\circ$), and bottom: cone (outflow) Gaussian component (average major axis PA $\sim 41^\circ$). The kinematics trace (disc or cone) was set based on the criteria $\sigma_{v\text{-disc}} < \sigma_{v\text{-cone}}$ for each spaxel. Crosses mark the AGN position. North is up, and east is to the left.

The masked flux maps also show the morphology of the disc and ionized outflow. Because only the brightest 40 percent of spaxels are retained (Figs 2–4), the spatial extent of each map directly traces the distribution of the corresponding emission line. Disc-tracing lines exhibit a roughly symmetric extension around the AGN, consistent with rotation about a central kinematic axis, while outflow-tracing lines extend predominantly to the north-east and south-west, in alignment along the ionization cone axis.

4.1.2 Intermediate-IP lines: single Gaussian

Of particular interest are what we will refer to as ‘intermediate-IP’ lines, namely [S III], [Ar III], and [Ne III], which exhibit mixed kinematic signatures in the single-Gaussian velocity maps (Fig.

2). Some of these lines display regions which align more closely with the disc and others more with the outflow kinematics. We interpret this mixed behaviour as a natural consequence of their intermediate IPs: these lines can be produced both in the lower ionization, star-forming disc and in the more highly ionized outflowing gas. This behaviour likely reflects overlapping gas elements along the line of sight, where gas excited by both components (the disc and the outflow) is capable of producing these emission lines. Thus, it is unsurprising that they exhibit features of both the disc and outflow kinematic signatures. We also see evidence for this mixed behaviour in the flux maps (Fig. 3), where these lines appear to show traces of the disc, and the edges of the blueshifted ionization cone. This is particularly visible in the [Ar III] emission map, where the cone edges appear as diffuse extensions westwards of the AGN, while we can also still resolve the two southern star-forming clumps, confirming that this line is able to trace both the outflow and the disc. This is also reflected in the best-fitting major axis PAs (Fig. 5, left). The [S III], [Ar III], and [Ne III] lines do not align with either the disc-tracing lines at $\sim -12^\circ$, $IP \lesssim 20$ eV, or the outflow-tracing lines at $\sim 54^\circ$, $IP \gtrsim 35$ eV, but instead fall between with an average PA $\sim 14 \pm 2^\circ$.

Also noteworthy is the contrasting behaviour of [S IV] and [Ne III]. While [S IV] ($IP = 34.8$ eV) appears to exclusively trace the outflow in both flux and velocity, [Ne III] ($IP = 41.0$ eV) shows contributions from both disc and outflow components. This demonstrates that IP alone does not fully determine which component (disc or outflow) the line traces. Rather, IP acts as a strong, but not exclusive determinant. The clean outflow signature in [S IV], despite its moderate IP, suggests it is preferentially excited by ionization mechanisms active in the outflow rather than in the disc. Previous studies have shown that the [S IV] line traces both star formation and AGN activity across many active galaxies (M. Pereira-Santaella et al. 2010a; D. Esparza-Arredondo et al. 2018). In some systems, it primarily follows star-forming regions, in others the AGN, and in some it reflects a combination of both. In NGC 7582, we find that [S IV] traces mostly the AGN and its ionized outflow, with no indication of contribution from the star-forming disc in either its velocity or flux maps.

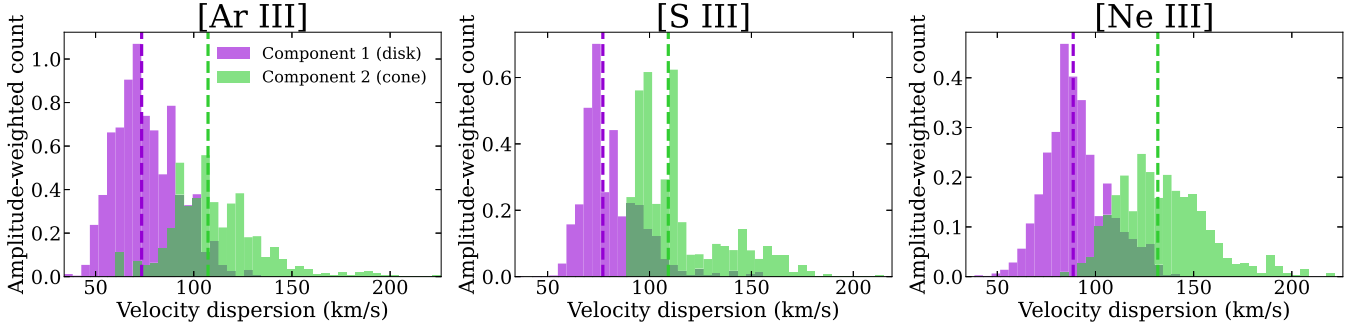


Figure 8. Gaussian amplitude-weighted histograms of the distribution of the velocity dispersion of each Gaussian component in each spaxel for the intermediate-IP lines from double-Gaussian fitting. The disc dispersion distribution is shown in purple and the outflow dispersion distribution is shown in green. Dashed lines show the median of each histogram.

Table 2. Best-fitting rotation curve parameters; maximum rotational velocity (V_{\max}) and the turnover radius (R_{turn}), derived from thin disc modelling for each disc-tracing emission line. Uncertainties are the 3σ errors.

Emission line	V_{\max} (km s $^{-1}$)	R_{turn} (pc)
[Fe II]	178 ± 9	86 ± 10
[Cl II]	177 ± 9	128 ± 20
HI 7-5	147 ± 9	75 ± 10
[Ar II]	146 ± 9	107 ± 10
[Ne II]	166 ± 12	118 ± 20

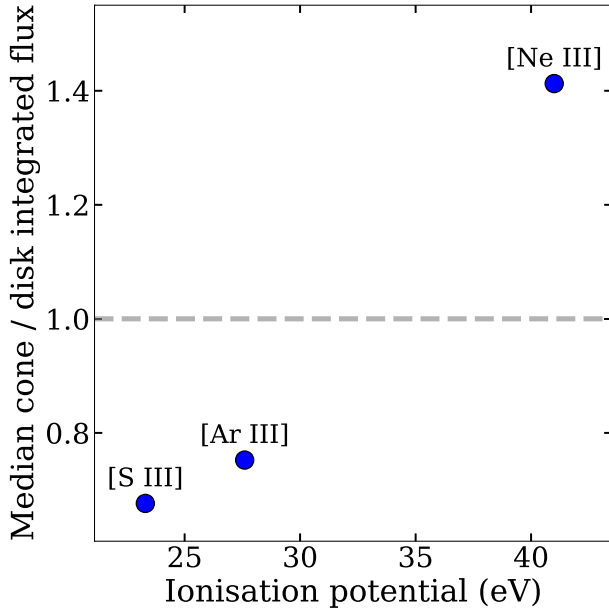


Figure 9. Median cone/disk integrated flux ratio across all unmasked spaxels for the intermediate-IP lines, showing that [S III] and [Ar III] are primarily excited in the disc, while [Ne III] is mostly excited in the cone. However, cone and disc fluxes are comparable for all three lines, differing by only up to ~ 50 per cent on average, indicating that each line can trace both components since neither dominates.

A similar trend was reported in other Seyfert galaxies by L. Zhang et al. (2024), who attributed the uniquely weaker detection of disc rotation in the [S IV] kinematics to obscuration, as this line lies near the centre of the $9.7 \mu\text{m}$ silicate absorption fea-

ture. Consequently, [S IV] emission from the star-forming disc is preferentially attenuated by silicate absorption, whereas emission from the ionized outflow, being located above and below the disc plane and thus subject to lower dust columns, remains largely unobscured. This naturally explains the weak or absent disc rotation signature in the [S IV] kinematics, despite the line’s intrinsic sensitivity to both star formation and AGN excitation based on IP.

It is also interesting to note that the separation in IP between the lines we classify as low-IP lies close to the IP of neutral He (24.6 eV), while the division between the high-IP lines, from [O IV] upwards, occurs near the IP of He II (54.4 eV). This is not a coincidence, with studies such as L. M. Feuillet et al. (2025) likewise showing that ionic species capable of tracing star formation predominantly have IPs below the He threshold, whereas those that trace AGN activity require IPs above the He II threshold. This cut-off is physically well motivated. Helium is the second most abundant element in the ISM, and its large photoionization cross-sections at 24.6 and 54.4 eV act as efficient photon sinks. In star-forming regions, the UV spectra of massive stars already decline steeply with increasing energy; any photons above 24.6 eV are quickly photoabsorbed by He, and the scarcity of stellar photons beyond 54.4 eV means that He II effectively absorbs almost all of them. Consequently, star-forming regions, such as the circumnuclear disc, produce very few photons capable of creating ions with IPs above the He II edge. By contrast, AGN generate a hard, power-law continuum extending well into the extreme-UV and X-ray, easily supplying photons energetic enough to ionize species with IPs $\gtrsim 50$ –100 eV. Thus, ions with IPs below 24.6 eV can be maintained in stellar photoionized gas, whereas species with IPs exceeding 54.4 eV necessarily trace the much harder radiation fields associated with AGN.

4.1.3 Velocity dispersion: single Gaussian

Similar IP-based behaviour is observed across the velocity dispersion maps of the low-, intermediate-, and high-IP lines (Fig. 4). These velocity dispersion maps look ‘noisy’, not due to low S/N, but because of the limitations of fitting a single-Gaussian to a line profile that arises from two distinct kinematic components (see Section 4.2 for physically motivated double-Gaussian fits). The low-IP lines exhibit the lowest velocity dispersions, with only modest enhancements near the AGN. The star-forming clumps north and south of the AGN, previously identified as part of the circumnuclear star-forming disc (R. A. Riffel et al. 2009; T. Ricci et al. 2018; O. Veenema et al. 2025), roughly coincide with

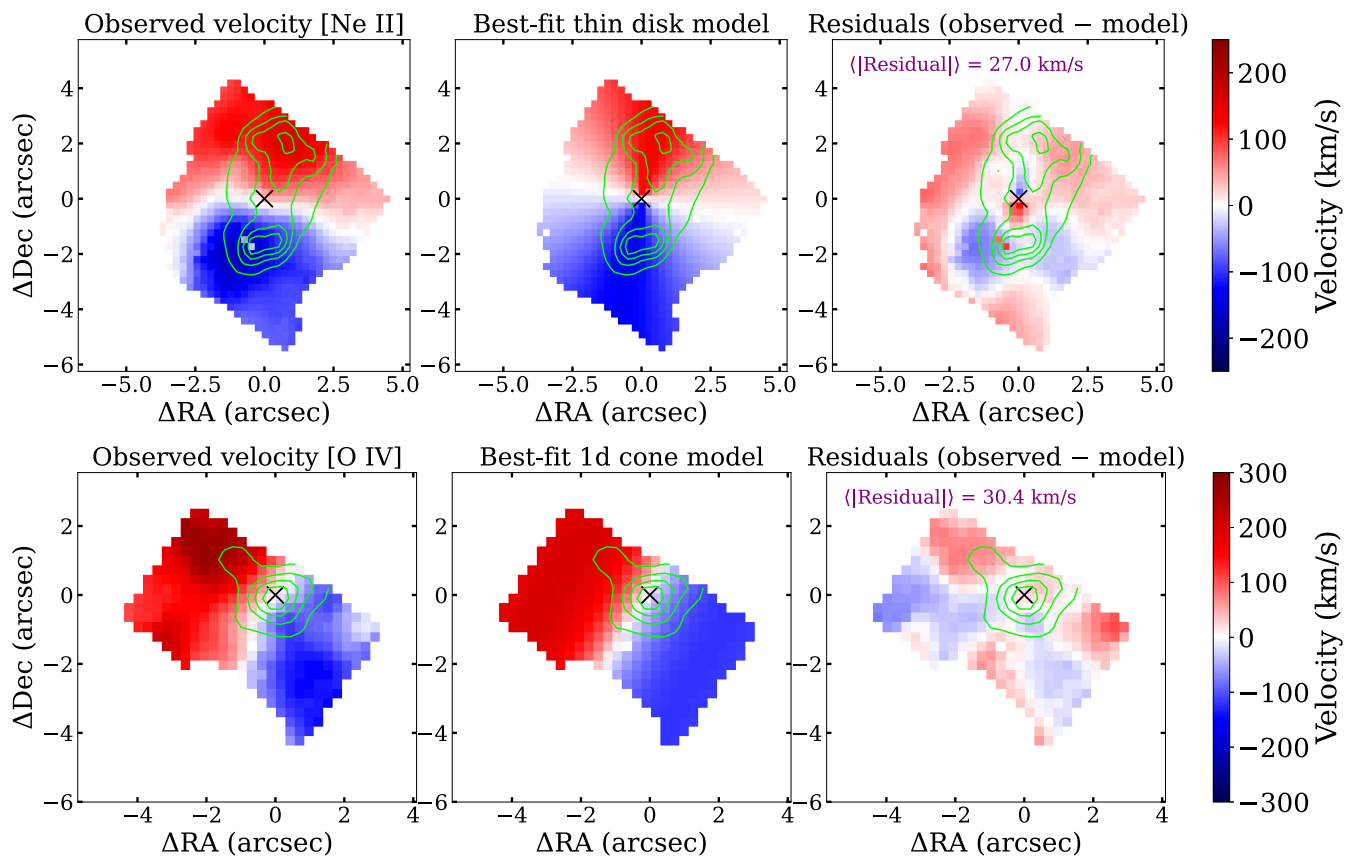


Figure 10. Top left: [Ne II] single-Gaussian velocity map, tracing the circumnuclear disc rotation. Top middle: inclined thin rotating disc model velocity field from the [Ne II] velocity map with best-fitting parameters inclination, $i = 58^\circ$, major velocity axis PA = -12° , and maximum rotational velocity (V_{max}) and turnover radius (R_{turn}) as given in Table 2. Top right: residuals of the model taken away from the observations (top left–top middle) with the mean absolute residual given, which is lower than the instrumental velocity dispersion in this MIRI/MRS channel. Green contours on the top plots show the [Ne II] integrated flux emission, which strongly traces the star-forming disc. Bottom left: [O IV] single-Gaussian velocity map. Bottom middle: best-fitting outflow model from [O IV]. Bottom right: residual (bottom left–bottom middle) map with the mean absolute residual shown, which is lower than the instrumental resolution in this MIRI/MRS channel. The green contours show the [O IV] integrated flux emission which is mostly PSF dominated from the AGN emission, with some elongation along the ionization cone axis direction. The cross denotes the AGN position.

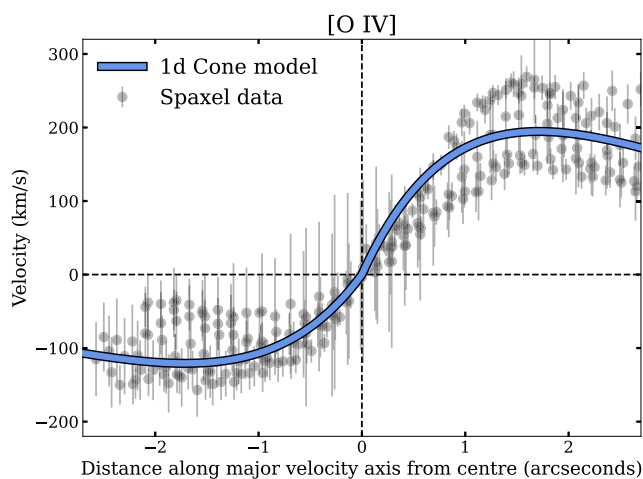


Figure 11. [O IV] single-Gaussian fit velocity for each spaxel as a function of distance along the outflow kinematic major axis (PA $\sim 42^\circ$ from PAFIT) from the AGN. The blue line shows the best-fitting asymmetrical exponential decay model which we use to then construct our analytical 1D outflow velocity field shown in Fig. 10 (bottom).

regions of reduced velocity dispersion in many low-IP lines, most notable in [Ne II], when compared with their flux distributions (Fig. 3). While our data alone do not directly demonstrate that these clumps are rotationally supported, the systematically lower dispersion observed in these more strongly emitting regions is consistent with dynamically colder gas, as expected for material embedded within a rotating disc where ordered motions dominate over local velocity dispersion.

The high-IP lines in Fig. 4 show broadly coherent behaviour, with systematically higher velocity dispersions across the FoV compared to the low-IP lines, and indications of modestly enhanced dispersion aligned with the ionization cone axis. Hence, the elevated dispersion, particularly along PA $\sim 45^\circ$ from the AGN, may be consistent with kinematic signatures of outflowing gas. This interpretation is supported by the results of S. Juneau et al. (2022), who analysed the velocity dispersion of ionized gas in the bicone (their fig. 6e) and found enhanced values within, and especially along the edges of, the front-facing cone. The reported dispersions of $\sim 140\text{--}200\text{ km s}^{-1}$ are comparable to the elevated values observed here in the high-IP lines. This agreement suggests consistency between the optical and mid-IR tracers, implying that the front facing ionization cone is not significantly af-

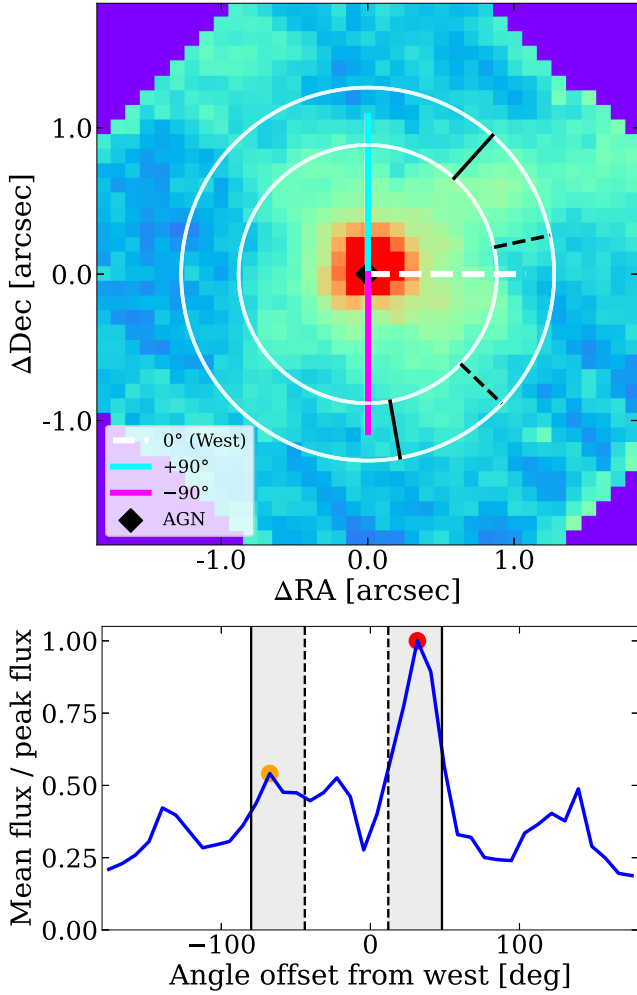


Figure 12. Top: [Mg IV] flux map showing the annulus considered and directionality of the angles. The colour scaling has been cut to three times below the maximum value to highlight the edges of the ionization cone. Bottom: average flux as a function of angle throughout the annulus for the [Mg IV] flux with the cone edge regions shaded and the positive peak identified as the red circle, and the negative peak identified as the orange circle. The dashed and solid black lines enclose the angular regions where we expect the cone edge to lie on both plots, with the dashed showing the inside of the edge and solid showing the outer extent of the edge. The two unmarked peaks beyond $\pm 100^\circ$ correspond to the edges of the eastern cone which is obscured by the disc.

affected by dust extinction, and supports the interpretation that the regions of increased dispersion are associated with the ionized outflow.

The intermediate-IP lines exhibit a distinct mixed behaviour in the velocity dispersion maps. Specifically, they show reduced dispersion in the north–south direction associated with the star-forming clumps, similar to the low-IP lines, alongside moderately enhanced dispersion along the north-east to south-west axis coincident with the ionization cone, as seen in the high-IP lines. This behaviour is particularly evident in [S III], where regions of elevated dispersion (yellow) to the top left and bottom right align with the ionization cone axis, while lower dispersion regions (darker green) north and south of the AGN trace the star-forming disc. This indicates that the intermediate-IP lines effectively capture the transition between the kinematic regimes traced by the

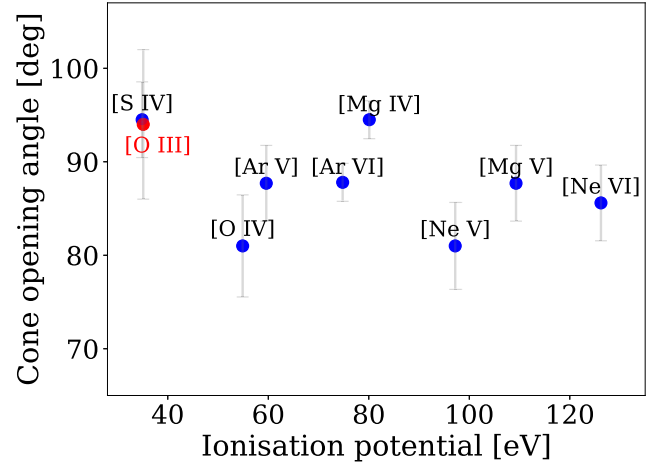


Figure 13. Cone opening angle (angle between positive and negative peaks) for each high-IP line plotted against their IP, including the [O III] 5007 line estimated from S. Juneau et al. (2022). We see no clear trend, indicating the cone opening angle is not affected by IP, i.e. that the AGN torus is thick to all emission lines and thus is not stratified in obscuring material density with IP.

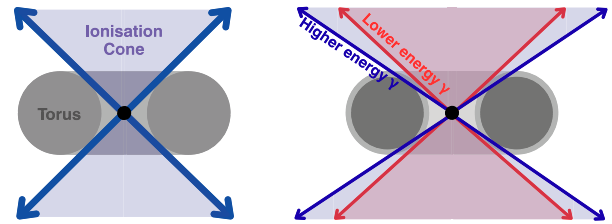


Figure 14. Schematic illustration of the two possible AGN torus morphologies considered. Each panel shows a cross-section of the torus with the central SMBH at the centre, and arrows indicating the boundaries of observable emission from the hollow ionization cones. Left: a torus with a polar axis that is largely unobscured, allowing ionizing photons of all energies to escape along similar angles, producing ionization cones with a roughly constant opening angle across IP. Right: a torus in which the polar axis contains more obscuring material, such as a two-phase medium with a low-density component, which would preferentially block low-IP photons and produce a narrower cone for lower energy emission lines. Note that a clumpy torus could produce the same effects, provided a low-density filling component is present in the second case.

low- and high-IP lines. Overall, the velocity dispersion maps reveal signatures of both the disc and outflow components, allowing a direct spatial comparison between the two. They further demonstrate that the intermediate-IP lines trace both kinematic structures simultaneously.

We also compute the median velocity dispersion of each emission line as a function of IP across all unmasked spaxels, shown in Fig. 6 (left), and Table 1. Both the median and mean (not plotted) dispersions exhibit a positive correlation with IP. This relation appears approximately linear for the low-IP, mostly disc-tracing lines (shown as blue diamond points). High-IP lines (shown as orange square points) show significantly higher velocity dispersions, and follow a similar or even flatter trend with IP but with significantly greater scatter. Similar velocity dispersion-IP trends have also been reported for other galaxies e.g. K. Dasyra et al. (2011) using archival Spitzer data, and L. Hermosa Muñoz et al. (2024) with MIRI/MRS.

The intermediate-IP lines (shown as magenta triangles) fall in between, with [S III] and [Ne III] both having relatively high median velocity dispersions consistent with the outflow-tracing lines, whereas [Ar III] shows a significantly lower median velocity dispersion, more along the trend of the disc-tracing lines. These three lines show hybrid behaviour with the positions of their points lying roughly in between the disc and outflow tracing values, once more suggesting that they are being efficiently excited in both circumnuclear structures.

The contrasting trends in the velocity dispersion-IP relation highlight the different kinematic environments of the disc and the outflow. In the disc, the roughly linear increase of velocity dispersion with IP may arise because lower IP (up to $\lesssim 20$ eV) emission preferentially traces star-forming clumps, which are present in the disc of this galaxy (O. Veenema et al. 2025), where enhanced turbulence and local velocity gradients broaden the line profiles, naturally producing the observed trend. In contrast, the nearly constant high velocity dispersion of the high-IP lines in the outflow is consistent with gas whose kinematics are dominated by bulk outflow motion with a roughly uniform level of turbulence across ionization phases, producing little variation with IP.

The mean average velocity dispersion for the low-IP, disc-tracing lines is 78 ± 11 km s⁻¹. For the three intermediate-IP lines that trace both disc and outflow components, it is 108 ± 12 km s⁻¹, and for the high-IP, purely outflow-tracing lines, 119 ± 13 km s⁻¹.

While fitting single-Gaussian profiles provides a useful first-order characterization of the gas kinematics, particularly for capturing large-scale rotational trends, it breaks down in inner regions where multiple, distinct kinematic components contribute to the line emission, resulting in non-Gaussian profiles. Examination of the line shapes reveals that, for the three intermediate-IP lines, the majority of spaxels exhibit double-peaked profiles, suggesting emission arising from two kinematically decoupled components.

4.2 Double-Gaussian fitting

To more thoroughly explore the kinematics traced by each emission line, we attempt to fit double-Gaussian profiles to each continuum-subtracted spectrum for every emission line in Table 1. Each fit comprises a sum of two Gaussians, each with the usual three free parameters (amplitude, centroid, and width), applied to the same spectral regions used in the single-Gaussian fitting. This approach yields two velocity maps per emission line, corresponding to the two fitted centroids.

4.2.1 Low- and high-IP lines: double Gaussian

We find that for most of the emission lines listed in Table 1, double-Gaussian fits typically assign one component to the dominant disc and outflow velocity field, while the second merely captures residual noise, producing unphysical velocity structures. This indicates that the majority of lines are well described by a single-Gaussian across most spaxels, particularly those tracing only the disc (Fig. A1). In contrast, several high-IP outflow-tracing lines, most notably [S IV], [O IV], and [Ne V], are significantly better reproduced by a double-Gaussian (Fig. A2), but with both components tracing the outflow kinematics, through having the same PA as that of the outflow (e.g. Fig. A4). While one

might alternatively attribute one of these components to emission from the AGN narrow-line region, we find this unlikely, as both Gaussian components exhibit highly similar and coherent velocity fields aligned with the outflow axis. Furthermore, both components show similar spatial distributions across the FoV. A similar behaviour was reported by S. Juneau et al. (2022), who, leveraging a larger FoV, showed that both components are associated with the extended ionization cone (due to limb brightening). Given the lack of discernible differences in the velocity structure and amplitude between the two components in our data, we also find no evidence that either component traces the narrow-line region.

A similar effect was reported by S. Juneau et al. (2022), who showed that in optical emission lines (H α , H β , [N II] 6548, 6584 Å, [S II] 6717, 6731 Å, and [O III] 4959, 5007 Å), spectra extracted from a purely disc region were adequately fit with single-Gaussians, whereas spectra from outflow regions were moderately improved by double-Gaussian fits. Our observations align with these findings for the high-IP lines mentioned above. For other high-IP transitions, double-Gaussian fits do not produce two physical velocity fields; we find that one Gaussian component traces the outflow, and the other residual noise, but this effect may simply reflect their lower S/N compared to [O IV], [S IV], and [Ne V]. Overall, these results align with the suggestion from S. Juneau et al. (2022) that the double-Gaussian profiles in these high-IP lines likely arise because the ionization cone is at least partially hollow. Hence, double-peaked Gaussian profiles would emerge naturally when both the near and far sides of the cone edge (one blueshifted, the other redshifted) are simultaneously visible along the line of sight of a given spaxel.

4.2.2 Intermediate-IP lines: double Gaussian

Beyond the purely disc- and cone-tracing lines, the intermediate-IP transitions ([Ne III], [Ar III], and [S III]) are particularly notable. These lines are not only well reproduced by double-Gaussian fits, but each Gaussian component traces a distinct kinematic structure. Specifically, the disc rotation is generally associated with the narrower Gaussian, while the broader component corresponds to the outflow (Fig. A3), consistent with the trends seen in Fig. 6 from the single-Gaussian analysis. On this basis, we defined the outflow component in each spaxel as the Gaussian with the larger velocity dispersion, while the disc is represented by the narrower component and reran the double-Gaussian fitting process. For clarity, we refer to these as component 1 (disc) and component 2 (cone), requiring $\sigma_{v\text{-disc}} < \sigma_{v\text{-cone}}$ for each spaxel.

This classification produces two velocity maps for each line (Fig. 7) that reproduce the expected kinematic signatures: the disc component aligns with the major axis rotation seen in the low-IP lines, while the outflow component follows the outflow kinematics traced by the high-IP lines (Fig. 2). We show this quantitatively in Fig. 5 (right), where we show the position angles from PAFIT for each Gaussian component of the intermediate-IP lines separately, with the disc and outflow component PAs aligning with what is seen from the solely disc/outflow tracing emission lines. Overall, these results provide strong evidence that the intermediate-IP lines have contributions from both the disc and outflow velocity fields.

We also examine the velocity dispersion distributions of each Gaussian component from the double-Gaussian fits, weighted

by amplitude, for all three intermediate-IP lines (Fig. 8). The histograms reveal two distinct populations (shown in purple and green for each component), with the median disc component dispersion consistently lower than that of the outflow component. Importantly, these two distributions remain distinct even without enforcing $\sigma_{v-\text{disc}} < \sigma_{v-\text{cone}}$, when components are instead classified by visual inspection of whether they trace the disc or the outflow, showing this separation is robust, and not simply due to construction. The disc component distributions are generally narrower with less scatter, indicative of ordered rotational kinematics. In contrast, the outflow components exhibit broader and more variable dispersions across the FoV, reflecting the turbulent nature of the ionized outflow.

We further refine our IP versus median velocity dispersion analysis for the intermediate-IP lines by plotting their median dispersions of each Gaussian component separately. As shown in Fig. 6 (right), each intermediate-IP line now includes its median disc component (diamond) and average outflow component (square) velocity dispersion. The low-IP lines, along with the disc components of the intermediate-IP lines, occupy the lower dispersion region of the plot, while the high-IP lines and outflow components lie at higher dispersions. Although separating the components introduces additional scatter, particularly among the low-IP group, this approach reinforces that gas associated with the circumnuclear disc exhibits systematically lower velocity dispersion than gas in the outflow. Consequently, the intermediate-IP lines continue to follow the same overall trends when their components are considered independently, strengthening the evidence that these transitions are excited by both the disc and outflow.

Further comparing the two Gaussian components across all spaxels for the three intermediate-IP lines, we find that the disc component amplitude is on average ~ 50 per cent larger than that of the outflow component. However, this does not imply that most of the flux originates in the disc, since the total flux depends on the integrated line profile and the outflow components are generally broader as discussed (Fig. 8). A similar result was reported by I. García-Bernete et al. (2021), who found that the outflow-tracing emission lines in NGC 5643 exhibit lower amplitudes and higher velocity dispersions than the emission lines associated with the disc. Overall, we measure comparable flux contributions (same order of magnitude) from both the disc and outflow components in the intermediate-IP lines. This balance is physically reasonable: if one component dominated the flux, the kinematic signature of the other would be much harder to detect.

Fig. 9 presents the ratio of the median cone-to-disc Gaussian component flux, defined as $(\langle A_{\text{cone}} \rangle \langle \sigma_{\text{cone}} \rangle) / (\langle A_{\text{disc}} \rangle \langle \sigma_{\text{disc}} \rangle)$ (where A and σ are the Gaussian amplitude and dispersions respectively), across all unmasked spaxels for the intermediate-IP lines. This demonstrates that the fluxes of the cone and disc components are generally comparable, while revealing a dichotomy in disc versus cone properties with IP: [Ne III] shows more cone emission, whereas [S III] and [Ar III] are predominantly excited in the disc. These trends support our interpretation that the intermediate-IP lines are excited by both disc and outflow components, with the relative contributions strongly influenced by IP, however neither component strongly outshining the other for these lines.

As illustrated in Figs 2 and 5, the disc exhibits a kinematic major axis with PA $\sim -12^\circ$, and previous studies (R. A. Riffel et al. 2009) have suggested that it is highly inclined and aligned close to perpendicular with the plane of the sky. This orientation

results in a high column density, contributing to the significant dust obscuration of the AGN. X-ray studies report line-of-sight column densities of $N_{\text{H}} \sim 10^{23} - 10^{24} \text{ cm}^{-2}$ for NGC 7582 (E. Rivers et al. 2015), consistent with a heavily obscured nucleus, while measurements also indicate substantial nuclear obscuration ($\tau_{9.8} \sim 2.67$; O. Veenema et al. 2025).

In contrast, investigations into the geometry of the ionization cone (S. Morris et al. 1985; S. Juneau et al. 2022) indicate that it is inclined at a substantial angle relative to our line of sight, meaning we are not observing gas motion directly along the axis of maximum outflow velocity. This orientation difference leads to a projection effect: while the disc's velocity field is almost fully projected along our line of sight (thus yielding radial velocities that closely reflect the true maximum rotational speed), the cone's projected radial velocities underestimate the true outflow speeds. Consequently, although the outflow components of the intermediate-IP lines already show larger speeds than the disc components (as shown in Fig. 2), these values are likely lower limits due to its inclination.

The double-Gaussian fitting has been highly informative, confirming that most emission lines primarily trace a single dynamical component, with only the three intermediate-IP lines showing strong signatures of both disc rotation and outflow motion. However, in some spaxels the method reaches its limits, making it difficult to fully separate and map the kinematics of both components. For instance, in the [Ne III] disc component velocity map shown in Fig. 7, several spaxels located to the north and south of the zero velocity axis deviate significantly from the expected disc rotation pattern and from the behaviour of surrounding spaxels. We note that these apparent deviations in the [Ne III] disc component are likely not indicative of true departures from ordered rotation, but rather reflect limitations of the moderately unconstrained double-Gaussian fitting in spaxels where the two components overlap. This highlights the need for more physically motivated modelling to robustly separate disc and outflow kinematics. Similar noisy regions are also seen in the [Ar III] and [S III] component velocity maps. Once more, these discrepancies likely also reflect spaxels where the fitting routine has struggled to accurately recover both components due to overlapping Gaussians that result in an overall line profile that resembles close to a single-Gaussian. This issue is exacerbated by the nature of automated fitting procedures, particularly when using models with a relatively high number of free parameters (six, in the case of the double-Gaussian fit), which can introduce degeneracies and lead to unreliable solutions in some spaxels. Thus, although the double-Gaussian fitting has performed surprisingly well overall in separating disc and outflow kinematics in the intermediate-IP lines, a more physically motivated and constrained modelling approach is warranted. To this end, we proceed by fitting an inclined thin disc rotation model to the data and subsequently fitting the then isolated cone component, and vice versa using a simplified, first-order cone model. This strategy reduces the number of free parameters in each case, and may allow for more robust, decoupled, and physically meaningful kinematic decomposition.

4.3 Disc modelling

4.3.1 The thin inclined rotating disc

In this subsection, we aim to model the kinematics of the disc component independently, then subtract its contribution from the total intermediate-IP emission-line profiles. This approach

allows us to isolate and analyse the residual line profiles under the assumption that they primarily trace the kinematics of the outflowing gas. To proceed, we first define and fit a physically motivated model for the rotation of the circumnuclear disc. We adopt a thin, inclined rotating disc model, centred on the AGN position, with free parameters such as the inclination of the disc relative to our line of sight, i , and the PA of its major kinematic axis. These parameters must be either constrained from the data or fixed based on independent observational estimates to enable a meaningful fit. To model the disc, we analyse several low-IP emission lines (Fig. 2).

We previously verified the PA of the disc's major axis using PAFIT (D. Krajinović et al. 2006) on the single-Gaussian velocity maps, finding a $PA \sim -12 \pm 3^\circ$ being highly consistent across all low-IP lines (Fig. 5).

For the disc inclination relative to the line of sight, i , we fix this parameter at $i = 58^\circ$, consistent with values derived by previous studies (S. Morris et al. 1985; M. Wold et al. 2006; S. García-Burillo et al. 2021). While more sophisticated fitting techniques such as 3DBAROLO (E. Di Teodoro & F. Fraternali 2015) could be employed to independently constrain the inclination, we find that adopting this fixed value yields a disc rotation model that reproduces the observed velocity fields exceptionally well (see Fig. 10, top and its discussion), and we have seen no major discrepancies between any of our velocity analyses that suggest large differences with the literature.

With the PA and inclination fixed, we proceed to fit the projected 2D rotational velocity field for a thin disc, where the projected sky coordinates (x, y) are measured relative to the galaxy centre. We adopt a \tanh rotation curve to describe the line of sight velocity field, which captures the typical shape of a disc rotation profile rising steeply near the centre and flattening at larger radii, and gives a good empirical fit to our data. The functional form of the observed disc velocity field is therefore given by:

$$v_{\text{model}}(x, y) = v_{\text{sys}} + V_{\text{max}} \tanh\left(\frac{R}{R_{\text{turn}}}\right) \sin(i) \cos(\theta). \quad (1)$$

Here, v_{sys} is the systemic velocity of the galaxy due to the *Hubble* expansion. For nearby galaxies, this can be approximated as $v_{\text{sys}} \approx cz$, where c is the speed of light and z is the redshift. For NGC 7582, we adopt $z = 0.00525$ (V. Braito et al. 2017). However, we note that all kinematic maps presented in this work have already been corrected by subtracting this systemic velocity. In the model, R denotes the deprojected radius in the disc plane, and θ is the corresponding azimuthal angle in the disc plane, defined as:

$$R = \sqrt{x'^2 + \left(\frac{y'}{\cos i}\right)^2} \quad (2)$$

$$\theta = \arctan\left(\frac{y'}{\cos i \cdot x'}\right) \quad (3)$$

$$x' = x \cos \phi + y \sin \phi \quad (4)$$

$$y' = -x \sin \phi + y \cos \phi. \quad (5)$$

where ϕ is the PA of the kinematic major axis (fixed at -12°), and (x', y') are the rotated (x, y) coordinates aligned with the major axis.

Two additional free parameters in the model are the maximum (asymptotic) rotational velocity, V_{max} , and the turnover radius, R_{turn} , which sets the radial scale over the velocity profile transitions. These parameters are determined through an LM optimization procedure, applied to the low-IP single-Gaussian velocity maps under the assumption that they purely trace the

disc rotation, with the PA and i fixed as described previously. To check consistency, we also performed fits using this model allowing the PA and i to vary, and found that they converged to values consistent with our adopted $PA = -12^\circ$ and $i = 58^\circ$ within one standard error. We therefore fixed these parameters in all subsequent model fits to avoid potential degeneracies, given our confidence in their values. The resulting projected velocity field of the model disc from fitting to the [Ne II] velocity map is shown in Fig. 10 (top). We ran this procedure for all disc-tracing emission lines, with the best fit V_{max} and R_{turn} for each stated in Table 2. The values of R_{turn} and V_{max} are broadly consistent across all disc-tracing lines confirming that they are all tracing the same kinematics. We choose the [Ne II] parameters for our model moving forward as it has the highest S/N of all disc-tracing lines. Fitting this model to the [Ne II] single-Gaussian velocity map yields a reduced chi-squared of $\chi_r^2 = 2.22$.

4.3.2 Disc model application and results

With an analytical velocity field for the circumnuclear disc established, Fig. 10 (top), we can incorporate it into the analysis of the intermediate-IP line kinematics. In its most straightforward application, the model disc velocity field provides a prediction for the centroid wavelength of the disc component in an emission-line profile. If the double-Gaussian-line profile assumption is retained, this effectively fixes one of the model parameters – namely the disc centroid, λ_{disc} , reducing the number of free parameters to five: the centroid, dispersion, and amplitude of the outflow component, λ_{outflow} , σ_{outflow} , A_{outflow} , and the dispersion and amplitude of the disc component, σ_{disc} , A_{disc} .

While the approach of fixing only the centroid wavelength of the disc component provides a basic method for aiding in isolating the outflow contribution in intermediate-IP lines, it may not yield the most accurate results as there are still five free parameters. An alternative strategy involves also constraining one or both of the remaining disc Gaussian parameters too, informed by insights from our earlier double-Gaussian fits. As previously discussed, we found that the disc component typically exhibited an emission-line amplitude ~ 50 per cent higher than that of the outflow component on average. Furthermore, as shown in Fig. 8, the disc velocity dispersion is systematically lower than that of the outflow and approximately follows a Gaussian distribution across all spaxels.

These findings suggest it may be justifiable to fix either or both of σ_{disc} and A_{disc} to representative values. For instance, σ_{disc} could be set to the median dispersion from the disc component of the double-Gaussian fits (i.e. the dotted purple lines in Fig. 8), and A_{disc} to $1.5A_{\text{outflow}}$. Although fixing these parameters might marginally reduce the fidelity of the disc component model, it mitigates degeneracies commonly encountered in automated fitting routines. This, in turn, facilitates more reliable convergence on the outflow kinematics by minimizing contamination from the disc component in the intermediate-IP line profiles.

For completeness, and given the difficulty in intuitively determining which approach best decouples the kinematics, we test all methods and present the resulting velocity fields for the three intermediate-IP lines in Fig. B1, where each subfigure shows the resulting outflow velocity field found from fitting a single-Gaussian to it after initially subtracting the disc model applied in different ways. Each row shows the results for a specific intermediate-IP line, and each column shows a different model

fitting approach as given by its title. These disc models all adopt the AGN position as the kinematic centre and share the same parameters as the thin rotating disc model previously derived from the [Ne II] velocity field but of course projected on to the corresponding NIRSspec or MIRI/MRS FoV and spaxel scale for each specific line.

Overall, the derived outflow velocity fields for each intermediate-IP line (Fig. B1) exhibit the expected morphology, with PAs consistent within 1σ errors with the high-IP, outflow tracing velocity maps from the single-Gaussian fits ($PA \sim 57^\circ$). When compared to the outflow fields obtained from the full six parameter double-Gaussian fits (bottom row of Fig. 7), the results from most fixed disc parameter combinations are less noisy. This improvement is particularly evident in the case of [Ne III], which provides the clearest example of successfully decoupled disc and outflow kinematics.

The most reliable outflow fits across each intermediate-IP line (Fig. B1) are obtained either by fixing all disc-Gaussian parameters (first column) or by allowing only the disc amplitude to vary (third column). These approaches yield velocity maps with minimal noise and far fewer spaxels that diverge drastically from the kinematics of their surroundings. This demonstrates that fixing the disc velocity dispersion to the median value for each line, as determined from the initial double-Gaussian fits (Fig. 8), is crucial for robustly recovering the outflow velocities. In contrast, when $\sigma_{v-\text{disc}}$ is left free (second and fourth columns in Fig. B1), we observe a significant increase in noisy or poorly constrained spaxels, especially in [Ar III] and [S III]. We therefore conclude that $\sigma_{v-\text{disc}}$ introduces substantial degeneracies, and fixing it to a physically motivated value is both appropriate and necessary to cleanly separate disc rotation from outflow kinematics when applying a thin inclined disc model.

This result is consistent with E. Bellocchi et al. (2019), who showed that the velocity dispersion is the parameter most sensitive to fitting methodology, while the velocity centroid remains comparatively robust. Allowing σ to vary introduces larger systematic deviations and increased scatter, particularly at low S/N. Their results support our choice to fix $\sigma_{v-\text{disc}}$ to a physically motivated value in order to minimize degeneracies and obtain more reliable gas kinematics.

4.4 Biconical outflow modelling

We have shown that incorporating a fixed disc model to constrain two or three parameters provides a more coherent picture of both the disc and outflow velocity fields in the intermediate-IP lines than a fully unconstrained double-Gaussian fit. An alternative strategy is to reverse the fitting order by first modelling the outflow kinematics independently, then fitting the disc Gaussian profile to the residuals. We explore this approach in this subsection.

4.4.1 First-order outflow model

Unlike ordered disc rotation, modelling outflow kinematics is less straightforward, particularly on the sub-kiloparsec scales probed closest to the AGN in this study. Ionized outflows, including the one in NGC 7582, are thought to be driven by AGN accretion powered radiation, which becomes collimated by the dusty torus, producing axially symmetric outflows perpendicular to the torus plane (S. García-Burillo et al. 2021; S. Juneau et al. 2022). Gas

in the ISM that lies along the path of these outflows becomes strongly ionized as it absorbs this radiative energy from the AGN. Observationally, it is unclear whether more powerful outflows are more frequently detected in heavily dust-obscured Type 2 AGN (G. Tozzi et al. 2024; G. R. Hauschild-Roier et al. 2025), none the less suggesting that interactions between accreted dust and the torus play a key role in launching them, which has been shown also through simulations (N. H. Soliman & P. F. Hopkins 2023). This makes modelling outflows challenging: understanding of their physical origin is incomplete, and their kinematic signatures are less intuitive than those of rotating discs.

We begin by analysing the kinematics of the outflow-tracing high-IP lines. Many such lines are present in the NIRSspec and MIRI/MRS wavelength ranges, but the cleanest tracer of the outflow with the best S/N is [O IV], as shown in our single-Gaussian velocity maps in Fig. 2. On these scales, the kinematics in the [O IV] map might, at a glance, resemble disc-like rotation, with a major kinematic PA of roughly 42° . The maps in Fig. 2 are flux-masked, showing only spaxels above the 60th percentile in flux for each line, thus the morphology of each subplot reflects the spatial extent of the emission. For [O IV] and other high-IP transitions, the structure appears to originate at the AGN and broadens with distance, as seen from the increasing number of unmasked spaxels further along the kinematic major axis. The emission aligns with the known ionization cone rather than the disc major axis, and the velocity dispersions are systematically larger than in disc-tracing lines, consistent with turbulent or bulk motions in an outflow. Moreover, the high-IP ions require hard radiation fields and so should be preferentially produced in the AGN photoionized cone, not from stellar feedback in the circumnuclear disc. Finally, the rotation-like appearance arises from the limited FoV (~ 200 pc), whereas the cones extend to $\gtrsim 3$ kpc (S. Juneau et al. 2022; C. Marconcini et al. 2023). Together, these arguments show that the high-IP lines do not trace a disc rotation, but instead the inner regions of the AGN-driven ionized outflow.

We present in Fig. 11 the velocity of each spaxel, derived from single-Gaussian fits, as a function of distance from the AGN along the kinematic major axis for [O IV]. This profile reveals a strong dependence of gas velocity on position along the major axis of the ionization cone. The gas reaches a peak velocity on either side of the nucleus. This peak is particularly pronounced and occurs between approximately 1–2 arcsec from the kinematic centre, beyond which the speed gradually declines. We also see asymmetry between the redshifted and blueshifted sides of the cone, with the redshifted side exhibiting maximum speeds that are $\sim 50 \text{ km s}^{-1}$ greater. This might indicate stronger acceleration on the far side of the cone or interaction with the galactic disc. We also examined the velocity structure as a function of distance perpendicular to the cone axis, i.e. toward the cone edges, and found no statistically significant variation across the FoV. This suggests a relatively uniform acceleration mechanism throughout the ionization cone, with no enhanced mechanical energy injection at larger angular distances from the cone axis. For consistency, we also checked this for [S IV] and [Ne V] (not shown), and confirm we see the same profile shape and trends as discussed for [O IV].

Several analytical formulae for AGN outflow kinematics have been adopted (e.g. V. Das et al. 2005; V. Das, D. Crenshaw & S. Kraemer 2007; F. Müller-Sánchez et al. 2011; T. Fischer et al. 2013; M. Durré & J. Mould 2019), typically featuring a monotonic rise in velocity over a few hundred parsecs to a maximum speed, followed by a decline toward the galaxy’s systemic velocity by ~ 1 kpc. The high-IP lines in NGC 7582 show a similar pattern:

the outflow accelerates to a peak and then gradually decelerates. However, the maximum speeds on the redshifted and blueshifted sides differ, necessitating an asymmetrical profile to more accurately capture the outflow velocity field heuristically.

Overall, as a simple approach, we model the outflow velocity field as a simple 1D velocity field, where velocities vary as a function of distance along the major axis of the cone from the AGN. This can be determined by fitting an asymmetrical decaying profile to the velocity curve shown in Fig. 11. Mathematically, this is expressed as:

$$v(r) = \begin{cases} v_{\text{red}} \cdot \left(\frac{r}{r_d}\right) \exp\left(-\frac{r}{2r_d}\right), & r \geq 0 \\ v_{\text{blue}} \cdot \left(\frac{r}{r_d}\right) \exp\left(\frac{r}{2r_d}\right), & r < 0 \end{cases} \quad (6)$$

where r denotes the distance along the major kinematic axis of the cone, and v_{red} , v_{blue} , and r_d represent the maximum asymptotic speeds on the redshifted and blueshifted sides, and the scale over which the velocity profile rises (the same for both the blue- and redshifted sides), respectively. We find that allowing for asymmetric maximum speeds on either side of the cone significantly improves the fit, with equation (6) outperforming asymmetric \tanh and \arctan profiles. Since our goal is to empirically characterize the velocity field of the outflow, this fitting formula is both sufficient and appropriate for our purposes.

For our model based on the [O IV] kinematics, we obtain a best fit of $v_{\text{red}} = 260 \pm 30 \text{ km s}^{-1}$, $v_{\text{blue}} = -160 \pm 30 \text{ km s}^{-1}$, and $r_d = 130 \pm 20 \text{ pc}$. Fitting this model to the [O IV] single-Gaussian velocity map yields a reduced chi-squared of $\chi_r^2 = 3.39$ (which is worse than the disc model to the low-IP lines, and likely due to our outflow model being more simplistic, with less free parameters). When applied to the [Ne V] and [S IV] velocity profiles, we find consistency within the 2σ uncertainties. Therefore, we adopt these parameters for all subsequent outflow modelling.

4.4.2 Outflow model application and results

Our 1D outflow model can then be directly compared to the observed [O IV] velocity map, as shown in Fig. 10 (bottom), and then reprojected on to the IFS FoV of other line observations to account for the outflow velocity component as desired. We find that the residuals between this simple 1D model and the [O IV] single-Gaussian velocity map used to generate it are small, typically comparable to or below the level of instrumental velocity dispersion, and that the model reproduces the observed velocities well as shown by Fig. 10 (bottom).

We can create model outflow velocity maps for the intermediate-IP lines and, analogously to the previous section, use this to keep various outflow Gaussian parameters fixed. We use the median velocity dispersion from the outflow Gaussian distribution (green dashed lines in Fig. 8), in fits where the outflow dispersion is fixed, and for fixed amplitude, we again set $A_{\text{disc}} = 1.5A_{\text{outflow}}$, finding this typically gives the best fits. The resulting disc velocity fits for the different parameter combinations and for the different intermediate-IP lines are presented in Fig. B2.

These maps demonstrate that subtracting the outflow component in different ways prior to fitting the disc Gaussian yields velocity maps that closely trace the expected disc kinematics. For all three intermediate-IP lines, this procedure produces velocity maps with the same overall morphology as those derived from the disc Gaussian component of the full double-Gaussian fits

(Fig. 7), but with much reduced noise. This highlights that this approach is also more robust than a full six-parameter double-Gaussian fit. We ran PAFIT on the recovered disc velocity field and find a position angle consistent with $\text{PA} \sim -12^\circ$ within two standard errors for all three lines in all four methods presented in Fig. B2, confirming that we are accurately resolving the disc. We have therefore demonstrated that, by first modelling the outflow in a simple 1D prescription, the intrinsic disc velocity field can be reliably recovered.

We observe the same trend as in the disc modelling used to fit the outflow (Fig. B1). Specifically, the most successful fits in Fig. B2 occur in the first and third columns, where either all outflow model parameters are fixed or only the outflow amplitude is allowed to vary. In contrast, permitting the outflow velocity dispersion to vary consistently produces poorer overall fits to the recovered disc velocity maps for each intermediate-IP line. Thus, in both approaches – whether subtracting the disc to fit the outflow, or subtracting the outflow to fit the disc – our results indicate that allowing both velocity dispersions to vary simultaneously significantly degrades the fits. It is therefore advantageous to fix at least one velocity dispersion parameter to a representative value in order to achieve more reliable kinematic decompositions. This should be an appropriate physically representative value, such as the median of the velocity dispersion distributions found from the double-Gaussian fits (Fig. 8).

4.5 Simultaneous disc and outflow constraints

Our disc-first and outflow-first modelling approaches show that the parameters exerting the strongest influence on the derived velocity fields are the velocity dispersions of the disc and outflow, $\sigma_{v\text{-disc}}$ and $\sigma_{v\text{-outflow}}$, respectively. To reduce the degeneracy and constrain the modelling, we now fix these dispersions and fit only the remaining four parameters of the double-Gaussian model: the amplitudes and centroid wavelengths of the outflow and disc components. We adopt, for each intermediate-IP line, the median dispersion values obtained from the fully unconstrained double-Gaussian fits (Fig. 8). With these dispersions fixed, we refit the data allowing the other four parameters to vary freely to derive the corresponding disc and outflow velocity maps, shown in Fig. B3.

These revised velocity maps from double-Gaussian fitting represent a notable improvement over the fully unconstrained fits shown in Fig. 7 for all intermediate-IP lines. In particular, the derived disc velocity fields exhibit substantially reduced noise and more easily identifiable position angles. The outflow kinematics also show higher velocities relative to the disc.

The outflow spaxels exhibit higher noise than those in the disc due to fixing the velocity dispersion. As shown in Fig. 8, the outflow has a broader velocity dispersion distribution, with fewer spaxels near the median compared to the disc. Consequently, assuming a constant dispersion is less accurate for the outflow and introduces noise across more spaxels than it does for the disc.

4.6 Escaping gas within the outflow?

In this section, we further examine the properties of the gas in the ionized outflow. M. Wold et al. (2006) modelled the gravitational potential of the circumnuclear stellar bulge (i.e. the star-forming disc) of NGC 7582 using a multi-Gaussian expansion (M. Cappellari 2002) fit to the Near Infrared Camera and Multi-Object Spectrometer (NICMOS) $F160W$ image. From their table

1, the summed luminosity of the four bulge components is $2.06 \times 10^9 L_{\odot}$. Adopting their stellar mass-to-light ratio of 3.8 gives a total stellar bulge mass of $7.83 \times 10^9 M_{\odot}$. They also estimated the central Supermassive Black Hole (SMBH) mass to be $\sim 5.5 \times 10^7 M_{\odot}$, yielding a total enclosed mass within the inner ~ 200 pc of $7.89 \times 10^9 M_{\odot}$. This estimate neglects the contributions from gas and dark matter, which are assumed to be negligible compared to the stellar component (disc + bulge) and SMBH masses on these spatial scales. Using this mass, we estimate the escape velocity at a distance of 200 pc from the AGN to be $v_{\text{esc}} \sim 580 \text{ km s}^{-1}$.

From our outflow velocity modelling (Fig. 11), the ionized gas within the cone reaches projected velocities of up to $\sim 200 \text{ km s}^{-1}$ at this radius. Because the outflow is a hollow bicone with a finite opening angle, individual gas parcels follow different trajectories and therefore intersect our line of sight at slightly different inclinations. As a result, the observed velocity field represents a superposition of projected velocities spanning a range of true outflow directions, rather than a single inclination. This naturally produces the observed spread in projected velocities, yielding average values of $\sim 200 \text{ km s}^{-1}$ at 200 pc for the mid-IR outflow-tracing lines. The corresponding deprojected velocities along the cone axis are higher, but remain highly sensitive to the inclination of the specific outflowing gas relative to our line of sight.

Several studies have estimated the outflow inclination in NGC 7582. Notably, C. Marconcini et al. (2025a) performed a detailed kinematic analysis of ionized AGN winds using optical emission lines, finding a best-fitting inclination of $i \sim 70^{\circ}$ for the blueshifted outflow. Similarly, C. López-Cobá et al. (2020), modelling nearby galaxies with VLT/MUSE IFS data, found $i \sim 72^{\circ}$. These estimates are in good agreement and broadly consistent with the ionization cone models of S. Morris et al. (1985) and T. Ricci et al. (2018), which, although not quantitatively constraining the inclination, support a similarly inclined geometry. Adopting $i \sim 70^{\circ} - 72^{\circ}$ implies that the deprojected outflow velocities are roughly a factor of 3 higher than the observed values. The mid-IR outflow-tracing lines analysed here exhibit typical velocities of $\sim 200 \text{ km s}^{-1}$ at 200 pc, corresponding to true deprojected velocities of $\sim 600 \text{ km s}^{-1}$ along the cone axis.

This indicates that a significant fraction of gas may escape the gravitational potential of the galaxy and be deposited into the circumgalactic (CGM) or intergalactic medium (IGM) rather than being re-accreted. Indeed, C. Marconcini et al. (2025a) examined the ratio of outflow velocity to escape velocity as a function of distance from the AGN using tracers at visible wavelengths (their fig. 3), finding values of $\sim 1-1.5$, in good agreement with the deprojected mid-IR outflow velocity to escape velocity ratio inferred here.

4.7 The effect of IP on the opening angle of the AGN ionization cone in context of the dusty torus

We further investigate the morphology of the AGN ionization cones and their connection to the properties of the unresolved dusty torus by analysing how the cone opening angle may be affected by IP for each high-IP line. The study by S. Juneau et al. (2022) reported a hollow ionization cone structure, for which we found supporting evidence in the form of double-peaked emission-line profiles in several high-IP lines, with both Gaussian components tracing the outflow velocity field.

We now examine the hollow cone in greater detail by analysing the high-IP line flux maps to identify signatures of the cone edges

as a function of angle at a fixed distance from the AGN. The edges of the hollow cone appear brighter than its interior, consistent with a structure in which the bicone is largely evacuated because gas inside is being rapidly driven outwards or compressed toward the cone walls by the outflow. This morphology is visible in many high-IP flux maps shown in Fig. 3.

To investigate this further, we analyse the flux within a hypothetical arc of fixed width at a constant distance from the AGN (i.e. an annulus). We adopt an inner radius of 0.8 arcsec (~ 80 pc) and an outer radius of 1.2 arcsec (~ 120 pc). This choice provides a suitable compromise to be sufficiently distant from the AGN to minimize PSF (point spread function) contamination, even in the longest MIRI/MRS channel, while also remaining fully within the FoV of all emission lines at all position angles. The annulus is illustrated in Fig. 12 (top) using the [Mg IV] integrated flux map as an example, with its boundaries shown with white solid lines, and the $\pm 90^{\circ}$ directions marked by cyan and purple lines, respectively. For convenience, we define west as the zero-angle reference, though this choice does not affect the results that follow.

The edges of the ionization cones were identified as bright and kinematically distinct features in the optical [O III] line by S. Juneau et al. (2022) (their figs 4, 6, and 7). Based on their results, we estimate the angular positions of the upper (northern) and lower (southern) edges of the westward, front-facing ionization cone. To account for uncertainty, we include a $\pm 5^{\circ}$ margin around each edge. The resulting regions correspond to $+12^{\circ}$ to $+48^{\circ}$ and -44° to -80° from west for the northern and southern edges, respectively. These are shown graphically in Fig. 12 (top) as black lines, where solid lines denote the outer boundaries and dashed lines the inner boundaries of each edge.

We then examine the average flux as a function of angle within the annulus. As an example, Fig. 12 (bottom) shows this distribution for [Mg IV], where the flux is normalized to the peak average flux within the annulus at each angle. The same edge positions are indicated by the black lines, and the shaded grey regions between each pair of lines mark the angular ranges where the cone-edge emission is expected to peak. The angles corresponding to the maximum mean flux within each region are identified as the cone-edge positions, marked by the red circle for the upper edge and the orange circle for the lower edge.

This analysis is repeated for all high-IP lines using the same annulus (fixed dimensions for inner and outer radius from the AGN), identifying the upper and lower cone edges as the flux peaks within the grey regions. For brevity, we show only the [Mg IV] result in Fig. 12, as all high-IP lines exhibit similar flux-angle profiles for this annulus. In all cases, the upper cone edge displays the strongest peak, while the lower edge is consistently weaker.

We calculate the cone opening angle as the angular separation between the two edge peaks, equal to the difference in angles of the red and orange circles for each high-IP line. The resulting opening angles are plotted as a function of IP in Fig. 13, with the uncertainties derived from the finite angular resolution set by the spaxel size of each IFU. We also include the [O III] 5007 line (IP = 35.1 eV) estimated from its flux map presented by S. Juneau et al. (2022), their fig. 4) as an additional point in red.

Fig. 13 shows no clear trend between IP and the measured cone opening angle, which remains consistently between $\sim 80^{\circ} - 95^{\circ}$, with average opening angle across all nine emission lines being $88 \pm 3^{\circ}$. Although a linear fit to the data points yields a slight negative gradient, this trend is not statistically significant

within the associated uncertainties. We therefore conclude that the biconical outflow maintains a constant opening angle across all high-IP lines. This implies that the AGN torus is relatively geometrically well defined, constraining the escape of ionizing photons similarly across this energy range, consistent with a torus in which the polar regions are largely unobscured.

This result does not require the torus to have a smooth, steeply declining density profile, as a clumpy torus with optically thick clouds would produce a similar constant opening angle (M. Nenkova et al. 2008a, b). In such a scenario, photons escaping along the polar axis always encounter the same effective optical depth, regardless of the detailed cloud distribution.

However, our results do constrain the torus central region: the polar axis must contain relatively little obscuring material, allowing photons of all energies to escape along similar angles. This corresponds to the torus configuration on the left in Fig. 14. A torus filled with a two-phase medium including a low-density component along the axis, as in the schematic on the right, would produce a narrower low-IP bicone, which is not observed.

5 CONCLUSIONS

In this paper, we analysed the inner regions of NGC 7582 using *JWST* NIRSpec and MIRI/MRS integral-field spectroscopic data, focusing on the kinematics of ionic emission lines based on their IP. Our goal was to disentangle and characterize the circumnuclear disc/ring rotation from the AGN-driven outflow, in the context of an active galaxy with complex nuclear morphology. Our key findings are as follows:

(i) The circumnuclear gas in NGC 7582 is kinematically stratified: low-IP emission lines trace the ordered circumnuclear disc/ring rotation with a $PA \sim -12 \pm 3^\circ$, while high-IP lines predominantly follow the AGN photoionized biconical outflow with a $PA \sim 54 \pm 10^\circ$ (Fig. 5).

(ii) The outflow-tracing gas exhibits systematically higher velocity dispersions than the disc-tracing gas on average and even on a per spaxel basis, consistent with turbulent or bulk motions in an ionized wind compared to ordered rotation (Fig. 8).

(iii) For lines tracing the disc component, velocity dispersion increases with IP, whereas the outflow shows a less well-defined trend, with a weak positive correlation between dispersion and IP (Fig. 6).

(iv) The low-IP emission lines are well described by single-Gaussian profiles, whereas the high-IP lines with the highest S/N ([S IV], [O IV], and [Ne V]) show moderately improved fits when modelled with two Gaussian components. Both components trace the outflow kinematics, consistent with emission arising from a hollow ionization cone.

(v) Three intermediate-IP species ([Ne III], [Ar III], and [S III]) are better fit by a double-Gaussian profile, but exhibit contributions from both kinematic components (Fig. 7). The outflow component is characterized by ~ 50 per cent higher velocity dispersion, ~ 50 per cent lower amplitude, and higher maximum speeds on average compared to the disc component.

(vi) Kinematic models; either a thin inclined rotating disc or a simple 1D outflow model are remarkably effective at isolating the complementary kinematic component in the intermediate-IP lines, allowing us to decouple the disc from the outflow kinematics (Fig. 10). We find that fixing the velocity dispersion of at least one of the components moderately improves the fit, leading to clearer derived velocity fields.

(vii) We compare the observed outflow velocities with the gravitational potential of the SMBH and stellar bulge at 200 pc (as calculated by M. Wold et al. 2006), finding that the gas reaches deprojected speeds comparable to the local escape velocity (~ 600 versus ~ 580 km s $^{-1}$). This suggests that the outflow is capable of ejecting material into the CGM or even the IGM, where it is unlikely to be re-accreted.

(viii) By examining the angular flux distribution at a fixed distance from the AGN for nine high-IP lines, we measure the opening angle of the front-facing ionization cone and find no dependence on IP, with average opening angle of $88 \pm 3^\circ$. This constant opening angle with IP indicates that the polar regions of the AGN torus contain relatively little obscuring material, allowing ionizing photons of all energies to escape along similar ranges in angle.

Finally, this work has demonstrated the exceptional capabilities of *JWST* IFS for studying galactic kinematics in nearby galaxies in the mid-IR. The combination of reduced extinction at these wavelengths and *JWST*'s sensitivity allows for the most accurate and precise measurements of gas kinematics to date.

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DATA AVAILABILITY

The data cubes used here are publicly available from the Mikulski Archive for Space Telescopes (MAST),¹ with doi:10.17909/9sc5-q436.

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APPENDIX A: EXAMPLE OF SPECTRA AND FITS FOR INDIVIDUAL SPAXELS

We present representative single-spaxel spectra and Gaussian fits illustrating the characteristic profiles of disc-, outflow-, and mixed disc/outflow-tracing emission lines. Specifically, we show examples for [Ar II], a primary disc tracer, [S IV], a strong outflow tracer, and [Ne III], an intermediate-IP line tracing both components (Figs A1–A3, respectively).

Fig. A1 demonstrates that typical [Ar II] spaxel spectra (and other disc-tracing lines) are well described by a single-Gaussian profile plus continuum. Introducing a second Gaussian component does not significantly improve the fit, as the overall χ^2 decreases by less than unity, indicating no statistically meaningful gain in goodness of fit.

In contrast, Fig. A2 shows a clear double-peaked profile for [S IV], a feature present in most spaxels when both edges of the cone are in the line of sight, and also observed in the other high-S/N outflow tracers, [O IV] and [Ne v], and is likely due to the bicone being hollow as discussed in Section 4.2. These profiles are significantly better described by double-Gaussian fits than by single-Gaussian models. The two fitted components have comparable amplitudes and widths (more so than a disc component), and their corresponding velocity maps exhibit nearly identical kinematic structures, including similar major-axis PAs and velocity extrema, confirming that both components trace the outflow kinematics (see Fig. A4).

Finally, Fig. A3 shows a representative [Ne III] spaxel spectrum (in the line of sight of both the cone and the disc) exhibiting an asymmetric profile with a pronounced wing, which is significantly better fit by a double- rather than single-Gaussian model, as indicated by a substantial reduction in the χ^2 . The two components typically have distinct parameters: a broader, lower amplitude component tracing the outflow, and a narrower, higher amplitude component tracing rotation in the circumnuclear disc. This behaviour is consistently observed across the IFU for [Ne III], [Ar III], and [S III], leading us to conclude that these lines robustly trace both disc and outflow kinematics.

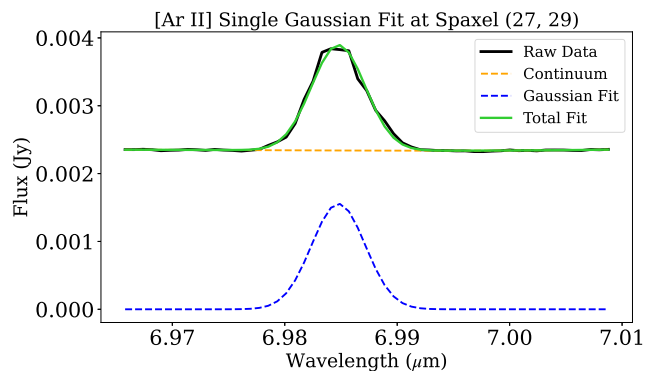


Figure A1. Representative single spaxel spectra around the [Ar II] line. This line strongly traces the circumnuclear star-forming disc and is well fitted by a single-Gaussian plus continuum fit.

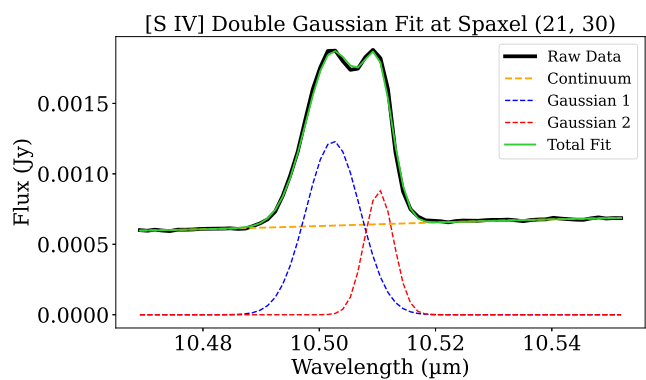


Figure A2. Representative single spaxel spectra around the [S IV] line. This line strongly traces the AGN-driven outflow, and a double-Gaussian plus continuum profile provides a significantly better fit than a single-Gaussian plus continuum. Both best-fitting Gaussian functions trace the outflow kinematics (see Fig. A4).

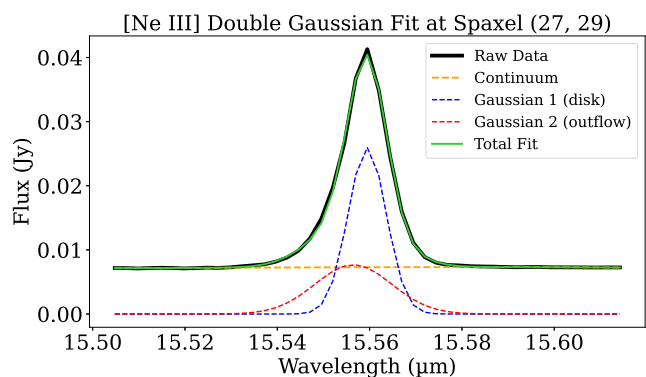


Figure A3. Representative single spaxel spectra around the [Ne III] line. This line strongly traces both the circumnuclear star-forming disc and AGN-driven outflow, and a double-Gaussian plus continuum profile provides a significantly better fit than a single-Gaussian plus continuum. However, these two Gaussian fits differ in that one has a significantly higher amplitude, lower width (dispersion), and strongly follows the disc kinematics, whereas the other has significantly lower amplitude, larger width (dispersion), and follows the outflow kinematics.

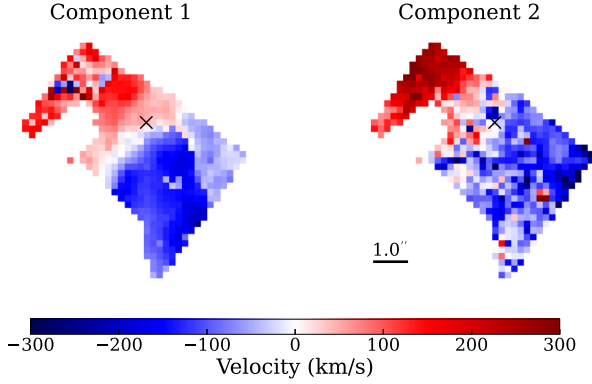


Figure A4. [S IV] velocity maps from double-Gaussian fitting. This line has an improved fit when using a double over a single-Gaussian-line profile, however both velocity maps strongly trace the outflow likely due to the bicone being hollow. The same effect is seen with other strongly outflow tracing lines with high S/N, [O IV] and [Ne V].

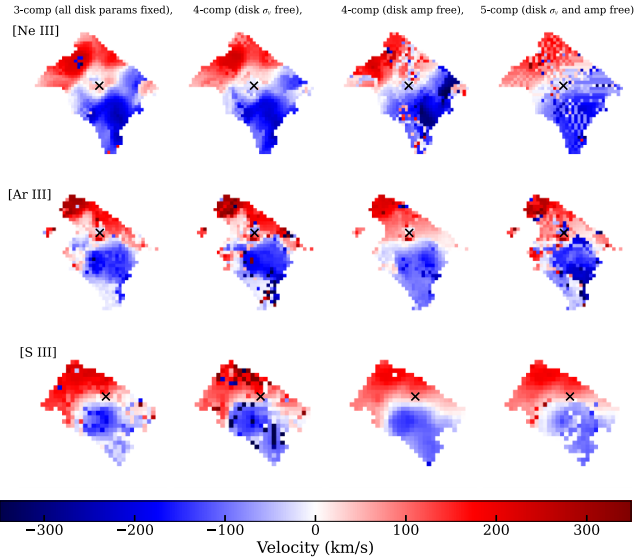


Figure B1. Resultant intermediate-IP line outflow velocity maps after incorporating the thin disc velocity model (Fig. 10, top), in different ways. Left column shows the resulting single-Gaussian fit velocity map after fixing all of the disc Gaussian parameters from the thin inclined disc model. Second column shows the velocity map when the disc velocity dispersion is an extra free parameter. Third column shows the velocity map when the disc amplitude is an extra free parameter. Fourth column shows the velocity map when both the disc dispersion and amplitude are both extra free parameters. Each row shows a different intermediate-IP line as given by the label on the left of each row.

APPENDIX B: RESULTANT FITS FROM DISC AND OUTFLOW MODELLING

We include the resultant intermediate-IP line velocity maps after modelling a single component (disc or outflow) in different ways in Figs B1 and B2, and from fitting a double-Gaussian with fixed velocity dispersion in Fig. B3.

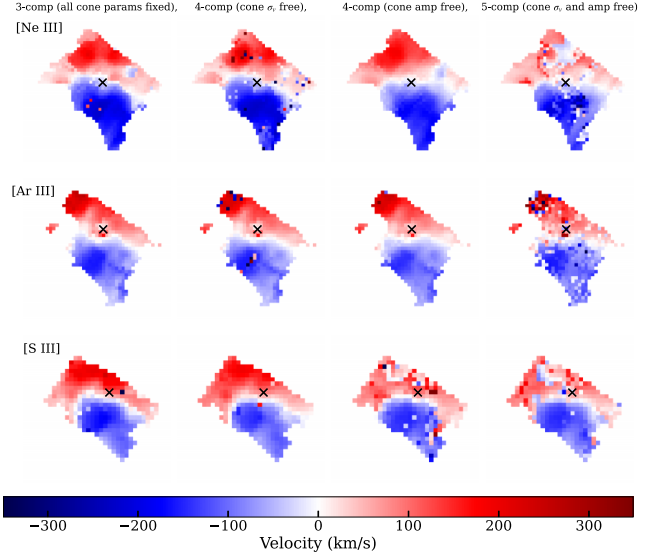


Figure B2. Resultant intermediate-IP line disc velocity maps after incorporating the simple outflow velocity model (Fig. 10, bottom), in different ways. Left column shows the resulting single-Gaussian fit velocity map after fixing all of the outflow Gaussian parameters from the simple outflow model. Second column shows the velocity map when the outflow velocity dispersion is an extra free parameter. Third column shows the velocity map when the outflow amplitude is an extra free parameter. Fourth column shows the velocity map when both the outflow dispersion and amplitude are extra free parameters. Each row shows a different intermediate-IP line as given by the label on the left of each row.

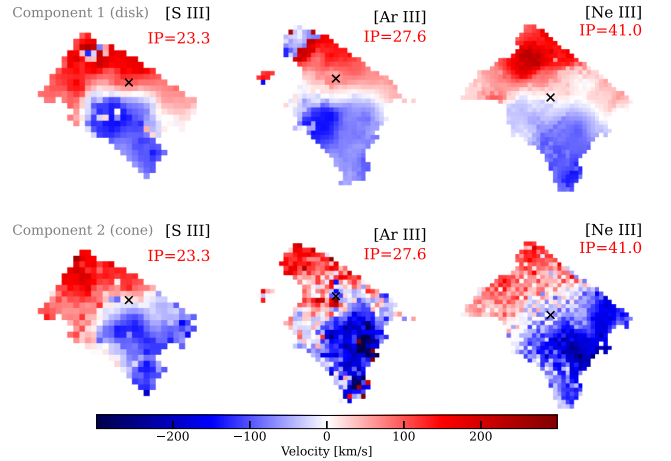


Figure B3. Velocity maps for the intermediate-IP emission-line fit using double-Gaussians but with both the disc and outflow velocity dispersions fixed to the median values from Fig. 8, showing moderate improvement on the naive double-Gaussian fits shown in Fig. 7. Top: disc Gaussian component, and bottom: outflow Gaussian component. The kinematics trace (disc or outflow) was set based on the criteria $\sigma_{v\text{-disc}} < \sigma_{v\text{-cone}}$ for each spaxel. Crosses mark the AGN position. North is up, and east is to the left.

¹Department of Physics, University of Oxford, Keble Road, Oxford, OX1 3RH, UK

²School of Sciences, European University Cyprus, Diogenes street, Engomi, 1516 Nicosia, Cyprus

³Centro de Astrobiología (CAB), CSIC-INTA, Camino Bajo del 497 Castillo s/n, E-28692 Villanueva de la Cañada, Madrid, Spain

⁴Instituto de Física Fundamental, CSIC, Calle Serrano 123, E-28006 Madrid, Spain

⁵Instituto de Astrofísica de Canarias, Calle Vía Láctea, s/n, E-38205, La Laguna, Tenerife, Spain

⁶Departamento de Astrofísica, Universidad de La Laguna, E-28206 La Laguna, Tenerife, Spain

⁷Departamento de Física de la Tierra y Astrofísica, Fac. de CC Físicas, Universidad Complutense de Madrid, E-28040 Madrid, Spain

⁸Instituto de Física de Partículas y del Cosmos IPARCOS, Fac. CC. Físicas, Universidad Complutense de Madrid, E-28040 Madrid, Spain

⁹School of Mathematics, Statistics and Physics, Newcastle University, Newcastle upon Tyne NE1 7RU, UK

¹⁰Observatoire de Paris, LUX, Collège de France, CNRS, PSL University, Sorbonne University, F-75014 Paris, France

¹¹Max Planck Institute for Extraterrestrial Physics, Giessenbachstrasse 1, D-85748 Garching, Germany

¹²Department of Astronomy & Astrophysics, University of California, San Diego, La Jolla, CA 92093, USA

¹³Observatorio Astronómico Nacional (OAN-IGN) – Observatorio de Madrid, Alfonso XII, 3, E-28014 Madrid, Spain

¹⁴Instituto de Radioastronomía y Astrofísica (IRyA), Universidad Nacional Autónoma de México, Antigua Carretera a Patzcuaro #8701, Ex-Hda. San José de la Huerta, C.P. 58089 Morelia, Michoacán, México

¹⁵Department of Physics & Astronomy, University of Alaska Anchorage, Anchorage, AK 99508-4664, USA

¹⁶Department of Physics, University of Alaska Fairbanks, Fairbanks, AK 99775-5920, USA

¹⁷Department of Physics and Astronomy, The University of Texas at San Antonio, 1 UTSA Circle, San Antonio, TX 78249, USA

¹⁸School of Physics and Astronomy, University of Southampton, Southampton SO17 1BJ, UK

¹⁹Telespazio UK for the European Space Agency (ESA), ESAC, Camino Bajo del Castillo s/n, E-28692 Villanueva de la Cañada, Madrid, Spain

²⁰Space Telescope Science Institute, San Martin Drive, Baltimore, MD 21218, USA

²¹National Astronomical Observatory of Japan, National Institutes of Natural Sciences (NINS), 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan

²²Department of Astronomy, University of Geneva, ch. d'Ecogia 16, CH-1290 Versoix, Switzerland

²³Instituto de Estudios Astrofísicos, Facultad de Ingeniería y Ciencias, Universidad Diego Portales, Av. Ejército Libertador 441, Santiago 8370191, Chile

²⁴Departamento de Física, CCNE, Universidade Federal de Santa Maria, 97105-900 Santa Maria, RS, Brazil

²⁵Centro de Astrobiología (CAB), CSIC-INTA, Ctra. de Ajalvir km 4, Torrejón de Ardoz, E-28850, Madrid, Spain

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