

— Supporting Information —

# Exchange Bias in Magnetic Topological Insulator Superlattices

Jieyi Liu,<sup>\*,†,¶</sup> Angadjit Singh,<sup>‡,§</sup> Yu Yang Fredrik Liu,<sup>‡</sup> Adrian Ionescu,<sup>‡</sup> Balati  
Kuerbanjiang,<sup>†</sup> Crispin H. W. Barnes,<sup>‡</sup> and T. Hesjedal<sup>\*,†</sup>

<sup>†</sup>*Clarendon Laboratory, Department of Physics, University of Oxford, Parks Road, Oxford,  
OX1 3PU, United Kingdom*

<sup>‡</sup>*Cavendish Laboratory, University of Cambridge, J. J. Thomson Avenue, Cambridge,  
CB3 0HE, United Kingdom*

<sup>¶</sup>*Cavendish Laboratory, University of Cambridge, J. J. Thomson Avenue, Cambridge,  
CB3 0HE, United Kingdom*

<sup>§</sup>*Department of Physics, Royal Holloway, University of London, Egham Hill, Egham,  
TW20 0EX, United Kingdom*

E-mail: Jieyi.Liu@physics.ox.ac.uk; Thorsten.Hesjedal@physics.ox.ac.uk

# S1. Methods

## Sample growth

The [CST/DBT]<sub>10</sub> superlattices were grown on *c*-plane sapphire substrates using molecular beam epitaxy (MBE), with a base pressure of  $1 \times 10^{-10}$  mbar. Prior to the thin film growth, the substrates were cleaned in solvents and rinsed in DI water, degassed overnight in the preparation chamber, and heat-treated at 450 °C for up to 30 minutes in the growth chamber. Bi, Sb and Te (99.9999% purity) were evaporated from standard effusion cells, while Cr and Dy (99.99% purity) were evaporated from high temperature effusion cells. A beam flux ratio of  $\sim 1:10:100$  for (Cr+Dy):(Bi+Sb):Te was maintained throughout the growth. The Te overpressure helps to reduce Te vacancies and anti-site defects common for the (Bi,Sb)<sub>2</sub>Te<sub>3</sub> system.<sup>1</sup> As a seed layer, a 3-nm-thick Bi<sub>2</sub>Te<sub>3</sub> film was deposited onto the substrate at 250 °C, followed by an anneal at 300 °C under Te flux. The CST/DBT superlattice consists of 10 repeats of nominally 5-nm-thick DBT layers and 10-nm-thick CST layers, grown at 300 °C. Based on previous calibrations using a combination of Rutherford backscattering spectroscopy (RBS) and particle induced x-ray emission (PIXE), the nominal composition of the superlattice is [Cr<sub>0.41</sub>Sb<sub>1.59</sub>Te<sub>3</sub>/Dy<sub>0.62</sub>Bi<sub>1.38</sub>Te<sub>3</sub>]<sub>10</sub>. Finally, an amorphous Bi layer was deposited as a capping layer to protect the sample from oxidation. The overall thickness of the superlattice (excluding the seed layer and the capping layer) is  $\sim 150$  nm. Following this general recipe, we also grew 20-nm-thick CST and DBT films on *c*-plane sapphire substrates as references, with nominal doping concentrations identical to the ones of the respective layers of the TI superlattice sample.

## Device fabrication and transport measurements

The TI thin films were patterned into Hall bars measuring  $1000 \mu\text{m} \times 100 \mu\text{m}$  using standard optical lithography and Ar<sup>+</sup> ion milling, followed by the fabrication of Ohmic Ti/Au contacts (20 and 80-nm-thick) by thermal evaporation and lift-off. The devices were then mounted

onto 20-pin leadless chip carriers and the contacts were made by wire bonding. The chip carriers fit into the sample holder of a transport measurement probe, which was loaded into a  $^4\text{He}$  continuous-flow cryostat. The electrical transport properties of the devices were measured using a standard four-terminal AC lock-in scheme with an input current of  $1\text{ }\mu\text{A}$  at a frequency of  $77\text{ Hz}$ .

## MOKE measurements

Polar MOKE measurements were carried out in the same  $^4\text{He}$  cryostat using a stand-alone MOKE probe which has been described in detail in Ref.<sup>2</sup> A continuous-wave laser beam with a wavelength of  $850\text{ nm}$  was generated from a stabilized diode laser, and the light was polarized using a Glan-Taylor polarizer. The laser beam was then focused onto the TI film down to a  $2.2\text{ }\mu\text{m}$  diameter spot size. The Kerr rotation of the reflected beam was modulated at  $100\text{ kHz}$  by a photoelastic modulator before going through a second polarizer, and finally reaching the photodetector. The Kerr angle is proportional to the ratio between the AC (at  $100\text{ kHz}$ ) and DC photovoltages generated at the detector.

## First-principles calculations

The DFT calculations were performed using the Vienna ab initio Simulation Package (VASP). The Perdew-Burke-Ernzerhof generalized gradient approximation (PBE-GGA) was applied for the exchange correlation functional. The projector augmented wave (PAW) method was employed, together with the plane wave basis set which had a kinetic energy cut-off of  $520\text{ eV}$ . The structure relaxation was performed under the conjugated gradient optimization method, using the following convergence criteria: the residual forces must be less than  $5 \times 10^{-4}\text{ eV}/\text{\AA}$ , and the total energy must have converged to within  $1 \times 10^{-5}\text{ eV}$ . A dense  $k$ -point grid was applied to compute the Brillouin-zone integrals, e.g., a  $\Gamma$ -centered  $6 \times 6 \times 1$  grid was used for the calculations of an 30-atom cell.

## Supplemental Figures

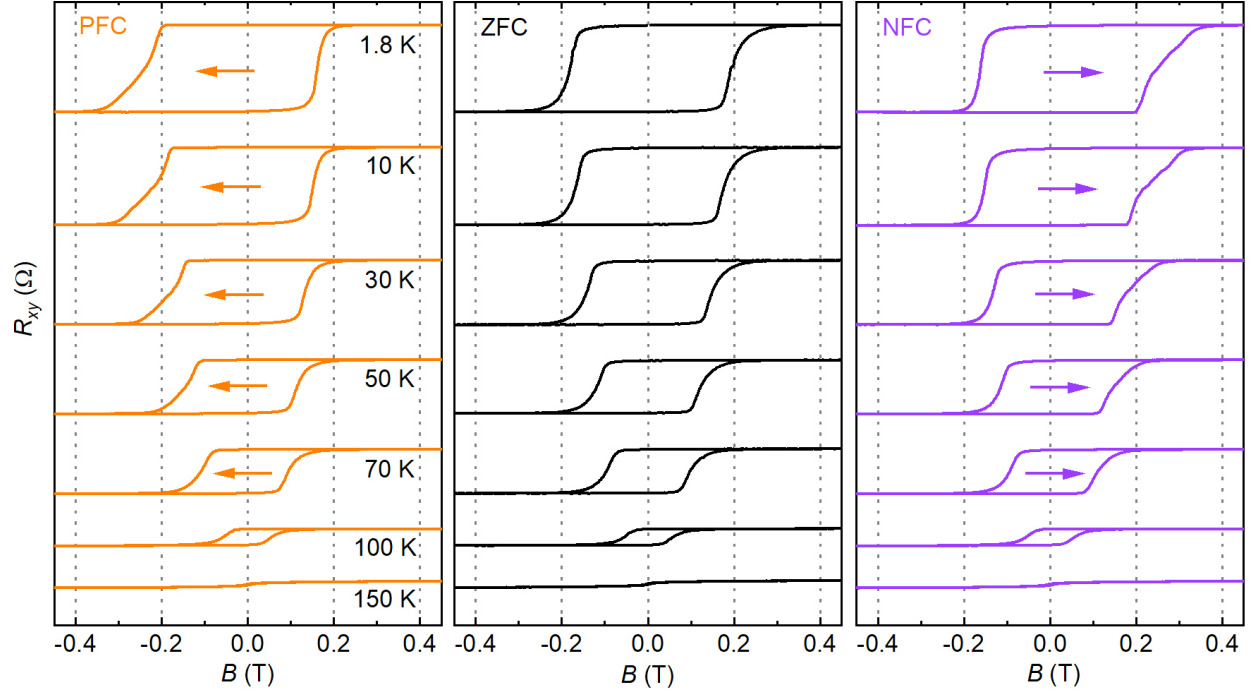


Figure S1: Exchange-biased hysteresis loops of the TI superlattice at different temperatures probed by electrical transport measurements. Positive-field-cooling (PFC) at +6 T results in a shift of the unbiased loop to the left (orange), and negative-field-cooling (NFC) at -6 T results in a shift to the right (purple). The unbiased loop (black), resulting from zero-field-cooling (ZFC), is shown as a reference.

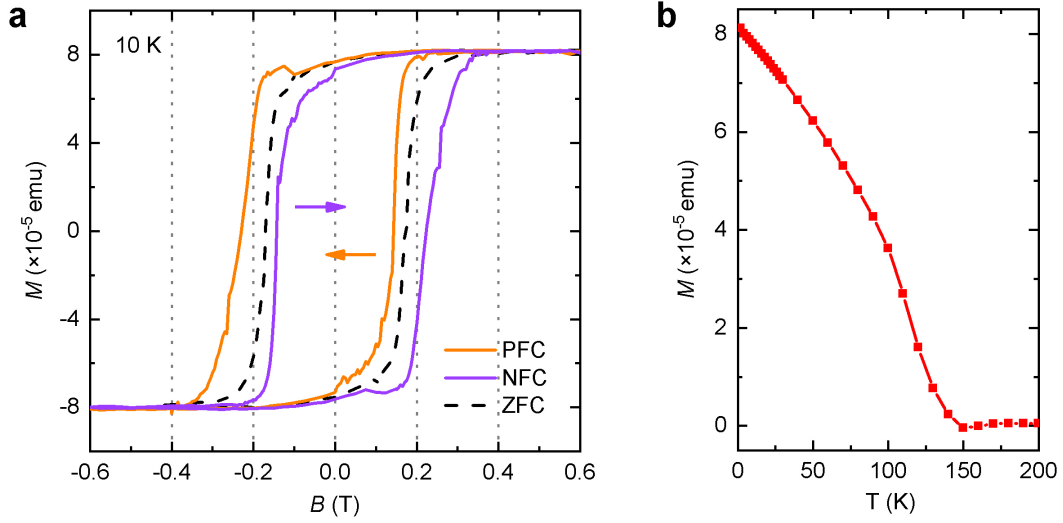


Figure S2: Exchange-biased magnetic response of the TI superlattice probed by SQUID magnetometry. (a) Hysteresis loops of the superlattice at 10 K. The measurements were conducted after positive-field-cooling (PFC) at +6 T (orange), after negative-field-cooling (NFC) at -6 T (purple), and after zero-field-cooling (ZFC) (black) from room temperature. (b) Temperature-dependent magnetization of the superlattice (ZFC case), measured at remanence.

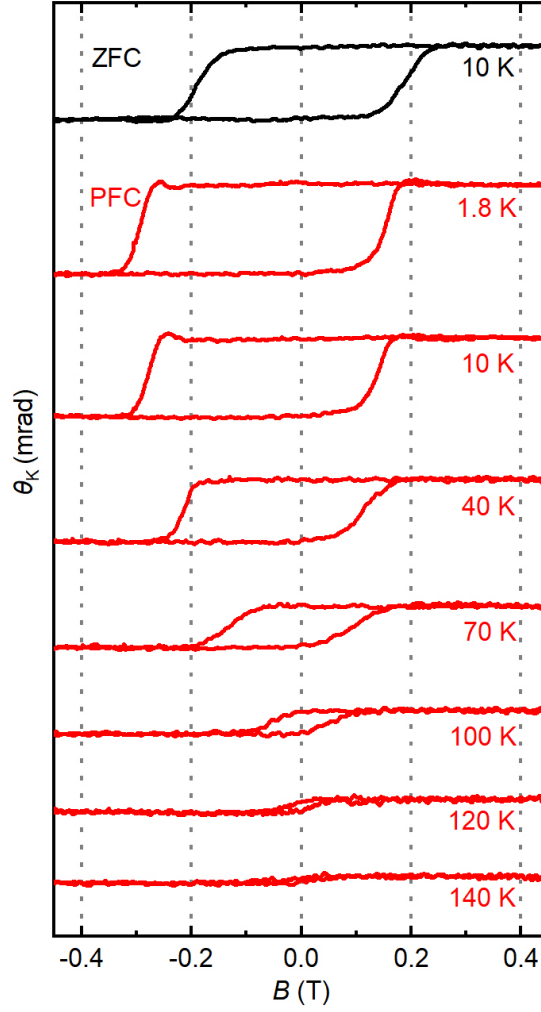


Figure S3: Exchange-biased hysteresis loops of the TI superlattice at different temperatures probed by MOKE measurements. Positive-field-cooling (PFC) at +6 T results in a shift of the unbiased loop to the left (red loops). The exchange bias effect is observed below 100 K, and hysteresis loops are present below 140 K. The zero-field-cooled (ZFC) hysteresis loop at 10 K is shown as a reference (black loop on top).

## References

- (1) Das, D.; Malik, K.; Deb, A. K.; Dhara, S.; Bandyopadhyay, S.; Banerjee, A. Defect induced structural and thermoelectric properties of  $\text{Sb}_2\text{Te}_3$  alloy. *J. Appl. Phys.* **2015**, *118*, 045102.
- (2) Liu, J.; Singh, A.; Llandro, J.; Duffy, L. B.; Stanton, M. R.; Holmes, S. N.; Applegate, M. J.; Phillips, R. T.; Hesjedal, T.; Barnes, C. H. W. A low-temperature Kerr effect microscope for the simultaneous magneto-optic and magneto-transport study of magnetic topological insulators. *Meas. Sci. Technol.* **2019**, *30*, 125201.