

T-DUALITY TRIVIALIZES BULK-BOUNDARY CORRESPONDENCE: THE PARAMETRISED CASE

KEITH C. HANNABUSS, VARGHESE MATHAI, AND GUO CHUAN THIANG

ABSTRACT. We state a general conjecture that T-duality trivialises a model for the bulk-boundary correspondence in the parametrised context. We give evidence that it is valid by proving it in a special interesting case, which is relevant both to String Theory and to the study of topological insulators with defects in Condensed Matter Physics.

1. INTRODUCTION

Recently the last two authors introduced T-duality in the study of topological insulators [40]. As an application, it was shown in [41] that T-duality trivializes the *bulk-boundary correspondence* (in a K -theoretic sense as pioneered by [30, 31], see also [45, 16, 21, 26], and [6] for a KK -theory perspective) in two dimensions, even in the presence of disorder. A similar simplification was found in basic cases in higher dimensions, for both complex and real K -theory, in a follow-up paper [42]. The bulk-boundary correspondence of topological invariants is an important aspect of the analysis of the integer quantum Hall effect via noncommutative geometry [3, 13, 14], as well as its generalisations [10, 37]. For a recent interesting alternate approach in the general field of topological phases of matter, see [55].

The notion of a Brillouin torus of quasi-momenta familiar from Bloch theory is a central one in condensed matter physics. Besides being a primary source of topological invariants of physical interest, the torus structure permits one to perform T-duality transformations. This means that the topological invariants, especially the K -theoretic ones [18, 53, 36, 51, 32], can be analysed on the T-dual side, or in real space, where they may be more easily understood. In fact, the concept of a Brillouin torus admits a vast generalisation using the language of noncommutative geometry. T-duality, as a geometric version of a generalized Fourier transform, continues to make sense in the more general noncommutative and even parametrised setting.

The first two authors introduced *parametrised* noncommutative strict deformation quantization of C^* -algebras with a torus action, and related it to T-duality of principal torus bundles with H-flux [23, 24]. While these papers were initially motivated by dualities in string theory, they may also have applications in condensed matter physics. For instance, parametrised families of physical systems are interesting in that they can have observable geometric phases or holonomies associated to them. In principle, such phenomena can be combined in non-trivial ways with commonly used topological invariants such as those arising from Bloch band topology.

In this paper, we prove that the bulk-boundary correspondence is again trivialised in the *parametrised* context, for a special case with nontrivial H-flux. In this case, the T-dual is a parametrised deformation of a 3-torus, and may be identified with the integer Heisenberg group algebra. The latter has an interpretation as a deformed version of the Brillouin torus in 3D, and we apply this picture to a model of a topological insulator with a uniform distribution of lattice dislocations. In this condensed matter physics context, a second related trivialization result for the bulk-boundary correspondence is proven. Finally, we also state a general conjecture which subsumes our previous results and is of independent mathematical interest.

CONTENTS

1. Introduction	1
2. T-duality and bulk-boundary correspondence	2
2.1. Overview of T-duality and NCPT bundles	3
2.2. Boundary NCPT bundles and the bulk-boundary map	3
2.3. Ordinary Fourier transform of restriction is integration	5
3. T-duality trivialises bulk-boundary correspondence in parametrised context	5
3.1. Generalities on the integer Heisenberg group	6
3.2. Statement of main results	7
3.3. Factorization of T-duality	9
4. Geometry of the Heisenberg Nilmanifold, K -theory of integer Heisenberg group algebra and PV sequence	11
4.1. K -theory generators	11
5. Proof of main results	12
5.1. Explicit maps on generators	12
5.2. Assembling K -theory maps together	15
6. Topological insulators with screw dislocations and the Heisenberg group	16
6.1. $\text{Heis}^{\mathbb{Z}}$ -symmetric Hamiltonians and the noncommutative Brillouin zone	18
Appendix A. Cyclic cocycles on integer Heisenberg group algebra and pairing with K -theory	19
A.1. Line bundles over Nil_k	19
A.2. Cyclic cocycles from group cocycles	19
A.3. Pairing with K -theory	21
References	22

2. T-DUALITY AND BULK-BOUNDARY CORRESPONDENCE

We begin by briefly summarizing the results in [23, 24], which will provide the technical background for the general conjecture 2.1. This section can be skimmed over to get to

the main results of this paper in Section 3 more quickly. For an application to topological insulators with lattice dislocations, see Section 6.

2.1. Overview of T-duality and NCPT bundles. Suppose that $A(X)$ is a C^* -bundle over a locally compact space X with a fibrewise action of a torus \mathbb{T}^n , and that $A(X) \rtimes \mathbb{T}^n \cong CT(X, H_3)$, where $CT(X, H_3)$ is a continuous trace algebra with spectrum X and Dixmier-Douady class $H_3 \in H^3(X; \mathbb{Z})$. These C^* -bundles are called H_3 -twisted NCPT (noncommutative principal torus) bundles over X . Our first main result there is that any H_3 -twisted NCPT bundle $A(X)$ is equivariantly Morita equivalent to the parametrised deformation quantization of the continuous trace algebra

$$CT(Y_{H_2}, q^*(H_3))_\sigma,$$

where $q : Y_{H_2} \rightarrow X$ is a principal torus bundle with Chern class $H_2 \in H^2(X; H^1(\mathbb{T}^n; \mathbb{Z}))$, and $\sigma \in C_b(X, Z^2(\mathbb{Z}^n, U(1)))$ a defining parametrised deformation such that $[\sigma] = H_1 \in H^1(X; H^2(\mathbb{T}^n; \mathbb{Z}))$. Here \mathbb{Z}^n is the Pontryagin dual of \mathbb{T}^n . This enables us to prove that the continuous trace algebra

$$CT(X \times \mathbb{T}^n, H_1 + H_2 + H_3),$$

with Dixmier-Douady class $H_1 + H_2 + H_3 \in H^3(X \times \mathbb{T}^n; \mathbb{Z})$ where $H_j \in H^j(X; H^{3-j}(\mathbb{T}^n; \mathbb{Z}))$, has an action of the vector group \mathbb{R}^n that is the universal cover of the torus \mathbb{T}^n , and covering the \mathbb{R}^n -action on $X \times \mathbb{T}^n$. Moreover the crossed product can be identified up to \mathbb{T}^n -equivariant Morita equivalence,

$$CT(X \times \mathbb{T}^n, H_1 + H_2 + H_3) \rtimes \mathbb{R}^n \cong CT(Y_{H_2}, q^*(H_3))_\sigma. \quad (2.1)$$

That is, the T-dual of $(X \times \mathbb{T}^n, H_1 + H_2 + H_3)$ is the parametrised strict deformation quantization of $(Y_{H_2}, q^*(H_3))$ with deformation parameter σ , $[\sigma] = H_1$. From this we obtain the explicit dependence of the K -theory of $CT(Y_{H_2}, q^*(H_3))_\sigma$ in terms of the deformation parameter.

In particular, by equation (2.1) and the Connes–Thom isomorphism [12], we deduce the isomorphism

$$K^j(X \times \mathbb{T}^n, H_1 + H_2 + H_3) \cong K_{j+n}(CT(Y_{H_2}, q^*(H_3))_\sigma). \quad (2.2)$$

This is a more explicit version of a special case of noncommutative T-duality considered in [38, 39].

2.2. Boundary NCPT bundles and the bulk-boundary map. In this paper, we will regard $CT(Y_{H_2}, q^*(H_3))_\sigma$ as the *bulk NCPT bundle* over X . From the above discussion, we have

$$CT(Y_{H_2}, q^*(H_3))_\sigma \rtimes \mathbb{T}^n \sim CT(X, H_3),$$

so Takai duality yields a Morita equivalence

$$CT(Y_{H_2}, q^*(H_3))_\sigma \sim CT(X, H_3) \rtimes \mathbb{Z}^n. \quad (2.3)$$

We wish to “peel off” the action of a particular \mathbb{Z} subgroup which we denote by $\mathbb{Z}_b = \widehat{\mathbb{T}}_b$, where $\mathbb{T}_b \subset \mathbb{T}^n = \mathbb{T}^{n-1} \times \mathbb{T} \equiv \mathbb{T}_a^{n-1} \times \mathbb{T}_b = \widehat{\mathbb{Z}}_a^{n-1} \times \widehat{\mathbb{Z}}_b$. That is,

$$CT(X, H_3) \rtimes \mathbb{Z}^n \cong (CT(X, H_3) \rtimes \mathbb{Z}_a^{n-1}) \rtimes \mathbb{Z}_b.$$

Under the inclusion $\iota : X \times \mathbb{T}_a^{n-1} \rightarrow X \times \mathbb{T}^n$, we have the restriction

$$\iota^*(X \times \mathbb{T}^n, H_1 + H_2 + H_3) = (X \times \mathbb{T}_a^{n-1}, \iota^*H_1 + \iota^*H_2 + H_3).$$

As in (2.2), the restricted T-dual transformation (as noncommutative principal \mathbb{T}^{n-1} -bundles) is also an isomorphism in K -theory,

$$K^j(X \times \mathbb{T}_a^{n-1}, \iota^*H_1 + \iota^*H_2 + H_3) \xrightarrow{T_a} K_{j+n-1}(CT(Y_{\iota^*H_2}, q_a^*(H_3))_{\iota^*\sigma}). \quad (2.4)$$

On the right-hand-side of (2.4), $Y_{\iota^*H_2}$ is a principal \mathbb{T}_a^{n-1} -bundle over X whose Chern class is equal to $\iota^*H_2 \in H^2(X; H^1(\mathbb{T}_a^{n-1}; \mathbb{Z}))$, its bundle projection is $q_a : Y_{\iota^*H_2} \rightarrow X$, and $[\iota^*\sigma] = \iota^*H_1 \in H^1(X; H^2(\mathbb{T}^{n-1}; \mathbb{Z}))$ is the induced parametrised deformation. Note that the \mathbb{T}^n -bundle Y_{H_2} can also be regarded as a principal \mathbb{T}_b -bundle over $Y_{\iota^*H_2}$, with bundle projection q_b being the quotient under the action of the subgroup $\mathbb{T}_b \subset \mathbb{T}^n$; thus $q = q_a \circ q_b$. We will regard $CT(Y_{\iota^*H_2}, q_a^*(H_3))_{\iota^*\sigma}$ as the *boundary NCPT bundle* over X .

As in (2.3), we also have a Morita equivalence $CT(Y_{\iota^*H_2}, q_a^*(H_3))_{\iota^*\sigma} \sim CT(X, H_3) \rtimes \mathbb{Z}_a^{n-1}$. Taking crossed products with \mathbb{Z}_b , we obtain

$$CT(Y_{\iota^*H_2}, q_a^*(H_3))_{\iota^*\sigma} \rtimes \mathbb{Z}_b \sim CT(X, H_3) \rtimes \mathbb{Z}^n \sim CT(Y_{H_2}, q^*(H_3))_\sigma,$$

exhibiting the bulk NCPT bundle as a crossed product of the boundary NCPT bundle by \mathbb{Z}_b . Associated to this \mathbb{Z}_b action is the Pimsner–Voiculescu *bulk-boundary map* [44]

$$\partial : K_{j+n}(CT(Y_{H_2}, q^*(H_3))_\sigma) \rightarrow K_{j+n-1}(CT(Y_{\iota^*H_2}, q_a^*(H_3))_{\iota^*\sigma}).$$

In general, we expect the following to hold.

Conjecture 2.1. *The following diagram commutes,*

$$\begin{array}{ccc} K^j(X \times \mathbb{T}^n, H_1 + H_2 + H_3) & \xrightarrow[\sim]{T} & K_{j+n}(CT(Y_{H_2}, q^*(H_3))_\sigma) \\ \downarrow \iota^* & & \downarrow \partial \\ K^j(X \times \mathbb{T}^{n-1}, \iota^*H_1 + \iota^*H_2 + H_3) & \xrightarrow[\sim]{T_a} & K_{j+n-1}(CT(Y_{\iota^*H_2}, q_a^*(H_3))_{\iota^*\sigma}) \end{array} \quad (2.5)$$

which will show that the bulk-boundary correspondence is trivialised by T-duality in this parametrised context. Here, T_a is noncommutative T-duality with respect to $\mathbb{T}^{n-1} = \mathbb{T}_a^{n-1}$, and ι^* is the induced restriction map in (twisted) K -theory under the inclusion $\iota : X \times \mathbb{T}_a^{n-1} \rightarrow X \times \mathbb{T}^n$.

For a general reference on C^* -algebras and K -theory, see [5], for a general reference on C^* -crossed product algebras see [54], for a general reference on continuous trace C^* -algebras, see [46] and for their K -theory, see [50]. A general reference on Algebraic Topology is [25].

2.3. Ordinary Fourier transform of restriction is integration. T-duality can be thought of as a generalised Fourier transform which, instead of transforming ordinary functions, gives an isomorphism at the level of topological invariants. Consider the Fourier transform $\text{FT}_{\mathbb{T}^d} : f \mapsto \widehat{f}$ which takes $f : \mathbb{Z}^d \rightarrow \mathbb{C}$ to $\widehat{f} : \widehat{\mathbb{Z}^d} = \mathbb{T}^d \rightarrow \mathbb{C}$. This is implemented by the kernel $P(\mathbf{n}, \mathbf{k}) = e^{i\mathbf{n} \cdot \mathbf{k}}$, $\mathbf{n} \in \mathbb{Z}^d, \mathbf{k} \in \mathbb{T}^d$,

$$\widehat{f}(\mathbf{k}) = \sum_{\mathbf{n}} P(\mathbf{n}, \mathbf{k}) f(\mathbf{n}) = \sum_{\mathbf{n}} e^{i\mathbf{n} \cdot \mathbf{k}} f(\mathbf{n}),$$

while $P(\mathbf{n}, \mathbf{k})^{-1}$ implements the inverse transform with a similar formula (the Chern character for the *Poincaré line bundle* $\mathcal{P} \rightarrow \mathbb{T}^d \times \mathbb{T}^d$ is the analogous object in the Fourier–Mukai transform in T-duality).

Write $(\mathbf{m}, n_d) = \mathbf{n}$ and let ι be the inclusion of $\mathbb{Z}^{d-1} \rightarrow \mathbb{Z}^d$ taking $\mathbf{m} \mapsto (\mathbf{m}, 0)$. Let $\partial : \widehat{f} \mapsto \partial \widehat{f}$ be partial integration along the d -th circle in \mathbb{T}^d . Since this picks out only the part of \widehat{f} with Fourier coefficient $n_d = 0$, it follows that there is a commutative diagram

$$\begin{array}{ccc} f & \xrightarrow[\text{FT}_{\mathbb{T}^d}]{\sim} & \widehat{f} \\ \downarrow \iota^* & & \downarrow \partial \\ \iota^* f & \xrightarrow[\text{FT}_{\mathbb{T}^{d-1}}]{\sim} & \partial \widehat{f} \end{array} \quad (2.6)$$

If we view $C(\mathbb{T}^d)$ as a crossed product $C(\mathbb{T}^{d-1}) \rtimes_{\mathbb{Z}} \mathbb{Z}$ with trivial action of the d -th copy of \mathbb{Z} , and represent the torus K -theory classes by differential forms, the Pimsner–Voiculescu boundary map is implemented by integration (or push-forward) along $\mathbb{T}^d = \widehat{\mathbb{Z}^d}$. The general conjecture 2.1 is in this way suggested by the elementary diagram (2.6).

3. T-DUALITY TRIVIALISES BULK-BOUNDARY CORRESPONDENCE IN PARAMETRISED CONTEXT

In this paper, we will restrict ourselves to the special case where $X = \mathbb{T} = S^1$, $n = 2$, $H_3 = 0 = H_2$ and $[\sigma^1] = H_1 \in H^1(S^1, H^2(\mathbb{T}^2, \mathbb{Z})) \cong \mathbb{Z}$ the decomposable generator. When regarded as an element of $H^3(S^1 \times \mathbb{T}^2; \mathbb{Z})$, H_1 is the volume form of $S^1 \times \mathbb{T}^2 \cong \mathbb{T}^3$. Its corresponding deformation parameter $\sigma^1 \in C(S^1, Z^2(\mathbb{Z}^2, U(1))) \cong C(S^1, \mathbb{T})$ is the identity function; more explicitly, the $U(1)$ -valued multiplier on $\mathbb{Z}^2 \times \mathbb{Z}^2$ at the point $e^{2\pi i \theta} \in S^1$ is

$$\sigma^1(e^{2\pi i \theta}) : ((p, q), (r, s)) \mapsto e^{i\theta(ps - qr)}, \quad (p, q), (r, s) \in \mathbb{Z}^2. \quad (3.1)$$

For kH_1 , the deformation parameter is σ^k , i.e. $kH_1 = k[\sigma^1] = [\sigma^k]$.

In these cases, the top row of Conjecture 2.1 is

$$K^j(S^1 \times \mathbb{T}^2, kH_1) \xrightarrow{T} K_j(C(S^1 \times \mathbb{T}^2)_{\sigma^k}). \quad (3.2)$$

The general case is work in progress.

3.1. Generalities on the integer Heisenberg group. The right-hand-side of (3.2) (the bulk NCPT bundle) is the C^* -algebra $C(S^1 \times \mathbb{T}^2)_{\sigma^k}$, which is the parametrised deformation quantization of $C(S^1 \times \mathbb{T}^2)$ by σ^k . By section 5 [38] we see that for $k \neq 0$,

$$C(S^1 \times \mathbb{T}^2)_{\sigma^k} \cong C^*(\text{Heis}^{\mathbb{Z}}(k))$$

where the integer Heisenberg group $\text{Heis}^{\mathbb{Z}}(k)$ is defined as

$$\text{Heis}^{\mathbb{Z}}(k) = \left\{ \begin{pmatrix} 1 & a & \frac{c}{k} \\ 0 & 1 & b \\ 0 & 0 & 1 \end{pmatrix} \mid a, b, c \in \mathbb{Z} \right\}.$$

We also write $\text{Heis}^{\mathbb{Z}} := \text{Heis}^{\mathbb{Z}}(1)$.

As we will need to keep track of the various subgroups in $\text{Heis}^{\mathbb{Z}}$, we introduce the following notation. The normal subgroup of matrices with $a = 0$ (resp. $b = 0$) is denoted by \mathbb{Z}_{bc}^2 (resp. \mathbb{Z}_{ac}^2). The subset with $c = 0$ is denoted by \mathbb{Z}_{ab}^2 . The central subgroup with $a = 0 = b$ is denoted by \mathbb{Z}_c , while the subgroup with $b = 0 = c$ (resp. $a = 0 = c$) is denoted by \mathbb{Z}_a (resp. \mathbb{Z}_b). Thus $\text{Heis}^{\mathbb{Z}}(k)$ is a non-split central extension of \mathbb{Z}_{ab}^2 by \mathbb{Z}_c ,

$$0 \longrightarrow \mathbb{Z}_c \longrightarrow \text{Heis}^{\mathbb{Z}}(k) \longrightarrow \mathbb{Z}_{ab}^2 \longrightarrow 0, \quad (3.3)$$

where we have reused the symbol \mathbb{Z}_{ab}^2 for the quotient group. We will peel off the action of \mathbb{Z}_b as described in Section 2.2

If we label elements of $\text{Heis}^{\mathbb{Z}}(k)$ by the 3-tuple $(a, b, c) \in \mathbb{Z}^3$, the group multiplication is $(a_1, b_1, c_1) \cdot (a_2, b_2, c_2) = (a_1 + a_2, b_1 + b_2, c_1 + c_2 + ka_1b_2)$. Then the group 2-cocycle for the central extension (3.3) is $\sigma_k^{\text{group}}((a_1, b_1), (a_2, b_2)) = ka_1b_2$.

We can also write $\text{Heis}^{\mathbb{Z}}(k)$ as a semi-direct product in two ways,

$$\text{Heis}^{\mathbb{Z}}(k) \cong \mathbb{Z}_{bc}^2 \rtimes \mathbb{Z}_a \cong \mathbb{Z}_{ac}^2 \rtimes \mathbb{Z}_b. \quad (3.4)$$

For example, the action of $b \in \mathbb{Z}$ on \mathbb{Z}_{ac}^2 by conjugation takes $(a, c) \mapsto (a, c - kba)$. The action of $b \in \mathbb{Z}$ on \mathbb{Z}_{ac}^2 can also be expressed via an $\text{SL}(2, \mathbb{Z})$ matrix as follows,

$$(a, c) \begin{pmatrix} 1 & -kb \\ 0 & 1 \end{pmatrix} = (a, c - kba). \quad (3.5)$$

In particular, using (3.4), we get the (split) short exact sequence of groups,

$$0 \longrightarrow \mathbb{Z}_{ac}^2 \longrightarrow \text{Heis}^{\mathbb{Z}}(k) \longrightarrow \mathbb{Z}_b \longrightarrow 0. \quad (3.6)$$

The Pontryagin dual of \mathbb{Z}_{ac}^2 will be denoted by \mathbb{T}_{ac}^2 , and similarly for the other subgroups. For instance, we have $C^*(\mathbb{Z}_{ac}^2) \cong C(\mathbb{T}_{ac}^2)$.

The group C^* -algebra $C^*(\text{Heis}^{\mathbb{Z}}(k))$ is generated by three unitaries U, V, W subject to the relation $UV = W^k VU$ and W being central. We can view U, V, W as the respective images

in $C^*(\text{Heis}^{\mathbb{Z}}(k))$ of the generators of $\mathbb{Z}_a, \mathbb{Z}_b, \mathbb{Z}_c$. As $\text{Heis}^{\mathbb{Z}}(k)$ is a semi-direct product, we may write $C^*(\text{Heis}^{\mathbb{Z}}(k))$ as a crossed product

$$C^*(\text{Heis}^{\mathbb{Z}}) \cong C^*(\mathbb{Z}_{ac}^2) \rtimes \mathbb{Z}_b \cong C(\mathbb{T}_{ac}^2) \rtimes \mathbb{Z}_b, \quad (3.7)$$

with $C(\mathbb{T}_{ac}^2)$ the boundary NCPT bundle. It is convenient to think of $\mathbb{T}_a, \mathbb{T}_b, \mathbb{T}_c$ as unit circles in the complex plane, whose points are respectively labelled by the complex numbers u, v, w , then U, V, W are the identity functions on these unit circles. Then the generating automorphism α_k in the crossed product $C(\mathbb{T}_{ac}^2) \rtimes \mathbb{Z}_b \equiv C(\mathbb{T}_{ac}^2) \rtimes_{\alpha_k} \mathbb{Z}_b$ acts on $f \in C(\mathbb{T}_{ac}^2)$ by

$$(\alpha_k \cdot f)(u, w) = f(w^{-k}u, w). \quad (3.8)$$

Alternatively, it acts on the unitaries $U^a W^c \in C(\mathbb{T}_{ac}^2)$ by

$$U^a W^c \mapsto U^a W^{c-ka}. \quad (3.9)$$

From the parametrised viewpoint, $C^*(\text{Heis}^{\mathbb{Z}}(k))$ is also a twisted crossed product

$$C^*(\text{Heis}^{\mathbb{Z}}(k)) \cong C(\mathbb{T}_c) \rtimes_{\sigma^k} \mathbb{Z}_{ab}^2$$

with $U(C(\mathbb{T}_c))$ -valued cocycle σ^k (c.f. Eq. (3.1)). Thus, we can regard $C^*(\text{Heis}^{\mathbb{Z}}(k))$ as a continuous field of noncommutative tori parametrised by $S^1 = \mathbb{T}_c$, with the rotation angle of the noncommutative torus over a point $w = e^{2\pi i\theta} \in S^1$ being $2\pi k\theta$. To emphasize this parametric point of view, we will often write S^1 in place of \mathbb{T}_c .

3.2. Statement of main results. Associated to the \mathbb{Z}_b action on $C(\mathbb{T}_{ac}^2)$ is the Pimsner–Voiculescu boundary map

$$\partial : K_0(C^*(\text{Heis}^{\mathbb{Z}}(k))) \rightarrow K_1(C(\mathbb{T}_{ac}^2)) \cong K^1(S^1 \times \mathbb{T}_a).$$

The first main result in our paper is,

Theorem 3.1. *The following diagram commutes,*

$$\begin{array}{ccc} K^0(S^1 \times \mathbb{T}^2, kH_1) & \xrightarrow[\sim]{T} & K_0(C^*(\text{Heis}^{\mathbb{Z}}(k))) \\ \downarrow \iota^* & & \downarrow \partial \\ K^0(S^1 \times \mathbb{T}) & \xrightarrow[\sim]{T_a} & K^1(S^1 \times \mathbb{T}) \end{array} \quad (3.10)$$

which shows that the bulk-boundary correspondence is trivialised by T -duality in this parametrised context. Here, T_a is T -duality (Fourier–Mukai transform) with respect to $\mathbb{T} = \mathbb{T}_a$, and ι^* is the induced restriction map in (twisted) K -theory under the inclusion $\iota : S^1 \times \mathbb{T} \rightarrow S^1 \times \mathbb{T}^2$.

Theorem 3.1 is a special case of Conjecture 2.1. We will prove the commutation of the diagram (3.10) by first writing T as a composition of a T -duality isomorphism T_1 for circle bundles, followed by a Baum–Connes isomorphism [2] T_2 , which in this case we show is a consequence of an imprimitivity theorem [48] and the Connes–Thom isomorphism theorem

[12] (Proposition 3.3). Then we compute the effect of the various maps on explicit K -theory generators. The proof of Theorem 3.1 is assembled in Section 5.2. Along the way, we also prove a second related result:

Theorem 3.2. *The following diagram commutes,*

$$\begin{array}{ccc} K^1(\text{Nil}_k) & \xrightarrow[\sim]{T_2} & K_0(C^*(\text{Heis}^{\mathbb{Z}}(k))) \\ \downarrow \iota^* & & \downarrow \partial \\ K^1(S^1 \times \mathbb{T}) & \xrightarrow[\sim]{T_{ac}} & K^1(S^1 \times \mathbb{T}) \end{array} \quad (3.11)$$

where Nil_k is the Heisenberg nilmanifold, T_2 is the Baum–Connes isomorphism described in the previous paragraph, T_{ac} is the full T -duality isomorphism (Fourier–Mukai transform) with respect to $S^1 \times \mathbb{T} = \mathbb{T}^2 = \mathbb{T}_{ac}^2$, and ι is a fibre inclusion. Here Nil_k is regarded as a fibre bundle over the circle \mathbb{T}_b with typical fibre \mathbb{T}_{ac}^2 .

As we shall see, Theorems 3.1 and 3.2 can actually be summarised by a single diagram (3.24).

Let us elaborate on the description of Nil_k as a bundle over \mathbb{T}_b in Theorem 3.2. The real Heisenberg group $\text{Heis}^{\mathbb{R}}$ is defined as

$$\text{Heis}^{\mathbb{R}} = \left\{ \begin{pmatrix} 1 & a & c \\ 0 & 1 & b \\ 0 & 0 & 1 \end{pmatrix} \mid a, b, c \in \mathbb{R} \right\}$$

and $\text{Heis}^{\mathbb{Z}}(k)$ sits inside $\text{Heis}^{\mathbb{R}}$ as a discrete subgroup. The Heisenberg nilmanifold is

$$\text{Nil}_k = \text{Heis}^{\mathbb{R}} / \text{Heis}^{\mathbb{Z}}(k).$$

We can write $\text{Heis}^{\mathbb{R}}$ as a semi-direct product in two ways,

$$\text{Heis}^{\mathbb{R}} \cong \mathbb{R}_{bc}^2 \rtimes \mathbb{R}_a \cong \mathbb{R}_{ac}^2 \rtimes \mathbb{R}_b. \quad (3.12)$$

For the second semi-direct product, the action of $b \in \mathbb{R}$ on \mathbb{R}_{ac}^2 by conjugation takes $(a, c) \mapsto (a, c - ba)$. As in (3.13), the action of $b \in \mathbb{R}$ on \mathbb{R}_{ac}^2 can be expressed via an $\text{SL}(2, \mathbb{R})$ matrix as follows,

$$(a, c) \begin{pmatrix} 1 & -b \\ 0 & 1 \end{pmatrix} = (a, c - ba). \quad (3.13)$$

Note that the restriction of (3.13) to an action of $b \in \mathbb{Z}_b$ on the discrete subgroup \mathbb{Z}_{ac}^2 is $(a, \frac{c}{k}) \mapsto (a, \frac{c - kba}{k})$. Equivalently, this action takes $(a, c) \mapsto (a, c - kba)$ after relabelling the elements of \mathbb{Z}_{ac}^2 and $\text{Heis}^{\mathbb{Z}}(k)$ with integers a, b, c according to our convention in Section 3.1. Thus (3.13) subsumes (3.5).

From (3.12), we get the (split) short exact sequence of groups,

$$0 \longrightarrow \mathbb{R}_{ac}^2 \longrightarrow \text{Heis}^{\mathbb{R}} \longrightarrow \mathbb{R}_b \longrightarrow 0. \quad (3.14)$$

Using (3.14) and (3.6) and taking the quotient, we arrive at the description of Nil_k regarded as a (non-principal) torus fibre bundle over the circle $\mathbb{T} = \mathbb{R}/\mathbb{Z}$, or mapping torus, with fibre $F \sim \mathbb{T}_{ac}^2 = \mathbb{R}_{ac}^2/\mathbb{Z}_{ac}^2$,

$$\mathbb{T}_{ac}^2 \longrightarrow \text{Nil}_k \longrightarrow \mathbb{T}_b. \quad (3.15)$$

More explicitly, the $\text{SL}(2, \mathbb{Z})$ transformation $\begin{pmatrix} 1 & -k \\ 0 & 1 \end{pmatrix}$ gives the diffeomorphism of the fibre specifying the bundle, and the map ι is taken to be a fibre inclusion.

3.3. Factorization of T-duality. The Heisenberg nilmanifold $\text{Nil}_k = \text{Heis}^{\mathbb{R}}/\text{Heis}^{\mathbb{Z}}(k)$ has a left action of the real Heisenberg group $\text{Heis}^{\mathbb{R}}$. Since $\text{Heis}^{\mathbb{R}}$ is a central extension,

$$0 \rightarrow \mathbb{R} \rightarrow \text{Heis}^{\mathbb{R}} \rightarrow \mathbb{R}^2 \rightarrow 0,$$

we can write it as a twisted Cartesian product $\text{Heis}^{\mathbb{R}} = \mathbb{R} \times_{\omega} \mathbb{R}^2$, where ω is the standard symplectic form on the vector space \mathbb{R}^2 and twists the product on $\mathbb{R} \times \mathbb{R}^2$ in the usual way.

We have, on the one hand,

$$\begin{aligned} C(\text{Nil}_k) \rtimes \text{Heis}^{\mathbb{R}} &= (C(\text{Nil}_k) \rtimes \mathbb{R}) \rtimes_{\omega} \mathbb{R}^2 \sim CT(\mathbb{T}^3, kH_1) \rtimes_{\omega} \mathbb{R}^2 \\ &CT(\mathbb{T}^3, kH_1) \rtimes_{\omega} \mathbb{R}^2 \otimes \mathcal{K} \cong CT(\mathbb{T}^3, kH_1) \rtimes \mathbb{R}^2 \end{aligned} \quad (3.16)$$

where \sim refers to Morita equivalence and the isomorphism in the second line is a consequence of the Packer-Raeburn trick, [43]. The 3-torus $\mathbb{T}^3 = S^1 \times \mathbb{T}_{ab}^2$ can be regarded as a circle bundle over \mathbb{T}_{ab}^2 with H -flux kH_1 , and we observe that T-duality for this circle bundle [7] is

$$K_{\bullet}(C(\text{Nil}_k) \rtimes \mathbb{R}) \xrightarrow{\text{Connes-Thom}} K^{\bullet+1}(\text{Nil}_k) \xrightarrow{T_1^{-1}} K_{\bullet}(CT(\mathbb{T}^3, kH_1)). \quad (3.17)$$

Note that Nil_k is regarded here as a principal circle bundle over \mathbb{T}_{ab}^2 (see Section 4.1.4). Also by noncommutative T-duality for \mathbb{T}^3 as a torus bundle over S^1 , Section 5 [38],

$$K_{\bullet}(CT(\mathbb{T}^3, kH_1) \rtimes \mathbb{R}^2) \xrightarrow{\text{Connes-Thom}} K^{\bullet}(\mathbb{T}^3, kH_1) \xrightarrow{T} K_{\bullet}(C^*(\text{Heis}^{\mathbb{Z}}(k))). \quad (3.18)$$

On the other hand, by the symmetric imprimitivity Theorem [48], the crossed product C^* -algebra $C(\text{Nil}_k) \rtimes \text{Heis}^{\mathbb{R}}$ is strongly Morita equivalent to $C(\text{Heis}^{\mathbb{R}} \backslash \text{Heis}^{\mathbb{R}}) \rtimes \text{Heis}^{\mathbb{Z}}(k) = C^*(\text{Heis}^{\mathbb{Z}}(k))$, and in particular,

$$K_{\bullet}(C(\text{Nil}_k) \rtimes \text{Heis}^{\mathbb{R}}) \cong K_{\bullet}(C^*(\text{Heis}^{\mathbb{Z}}(k))). \quad (3.19)$$

By the Connes-Thom isomorphism Theorem (c.f. Corollary 7 of [12], Corollary 2 of [17]), there is a natural isomorphism

$$K_{\bullet}(C(\text{Nil}_k) \rtimes \text{Heis}^{\mathbb{R}}) \cong K_{\bullet+1}(C(\text{Nil}_k)). \quad (3.20)$$

From (3.19) and (3.20), we conclude that

$$K^{\bullet+1}(\text{Nil}_k) = K_{\bullet+1}(C(\text{Nil}_k)) \xrightarrow{T_2} K_{\bullet}(C^*(\text{Heis}^{\mathbb{Z}}(k))), \quad (3.21)$$

which can also be understood as the the assembly map [2] composed with Poincaré duality in this context. Equations (3.16), (3.17), (3.18), (3.21) imply the claimed factorization of T-duality,

Proposition 3.3 (Factorization of T-duality). *In the notation above, the following diagram commutes,*

$$\begin{array}{ccc}
 K^\bullet(\mathbb{T}^3, kH_1) & \xrightarrow[\sim]{T} & K_\bullet(C^*(\text{Heis}^{\mathbb{Z}}(k))) \\
 \searrow \scriptstyle T_1 \sim & & \nearrow \scriptstyle \sim T_2 \\
 & K^{\bullet+1}(\text{Nil}_k) &
 \end{array} \tag{3.22}$$

For the 2-torus $S^1 \times \mathbb{T} = \mathbb{T} \times_{\substack{c \\ a}} \mathbb{T}$, T-duality T_{ac} is just the ordinary Fourier–Mukai transform. It factorises as (the inverse¹ of) T-duality T_c with respect to one circle factor $S^1 = \mathbb{T}_{\substack{c \\ a}}$, followed by T-duality T_a with respect to the other circle factor \mathbb{T}_a ,

$$\begin{array}{ccc}
 & K^\bullet(S^1 \times \mathbb{T}) & \\
 \nearrow \scriptstyle \sim T_c & & \searrow \scriptstyle \sim T_{ac} \\
 K^{\bullet+1}(S^1 \times \mathbb{T}) & \xrightarrow[\sim]{T_a} & K^\bullet(S^1 \times \mathbb{T})
 \end{array} \tag{3.23}$$

We use the factorisations in (3.22) and (3.23) to rewrite and combine the commutative diagrams (3.10) and (3.11) as

$$\begin{array}{ccccc}
 K^0(S^1 \times \mathbb{T}^2, kH_1) & \xrightarrow[\sim]{T} & K_0(C^*(\text{Heis}^{\mathbb{Z}}(k))) & & \\
 \downarrow \scriptstyle \iota^* & \searrow \scriptstyle \sim T_1 & \nearrow \scriptstyle \sim T_2 & & \downarrow \scriptstyle \partial \\
 & K^1(\text{Nil}_k) & & & \\
 & \downarrow \scriptstyle \iota^* & & & \\
 & K^1(S^1 \times \mathbb{T}) & & & \\
 \nearrow \scriptstyle \sim T_c & & \searrow \scriptstyle \sim T_{ac} & & \\
 K^0(S^1 \times \mathbb{T}) & \xrightarrow[\sim]{T_a} & K^1(S^1 \times \mathbb{T}) & &
 \end{array} \tag{3.24}$$

¹The Fourier transform and its inverse are the same up to a sign choice in the exponent of the kernel $P(\mathbf{n}, \mathbf{k}) = e^{\pm i \mathbf{n} \cdot \mathbf{k}}$. Similarly the Fourier–Mukai transform and the inverse are the same up to a sign choice for first Chern class of the Poincaré line bundle.

4. GEOMETRY OF THE HEISENBERG NILMANIFOLD, K -THEORY OF INTEGER HEISENBERG GROUP ALGEBRA AND PV SEQUENCE

4.1. K -theory generators.

4.1.1. *Generators for $K^0(S^1 \times \mathbb{T}^2, kH_1)$.* Let us coordinatise $S^1 =: \mathbb{T}$ by x_c and $\mathbb{T}^2 \equiv \mathbb{T}_{ab}^2 = \mathbb{T}_a \times \mathbb{T}_b$ by (x_a, x_b) . By equation (4.15) [7, 8] (see also section 4.1.4 [9]), we see that there is a natural isomorphism when $k \neq 0$,

$$\begin{aligned} K^0(S^1 \times \mathbb{T}^2, kH_1) &\cong \tilde{K}^0(S^1 \times \mathbb{T}^2) \cong \mathbb{Z}[\mathcal{L}_{ca}] \oplus \mathbb{Z}[\mathcal{L}_{cb}] \oplus \mathbb{Z}[\mathcal{L}_{ab}] \\ &\cong H^2(S^1 \times \mathbb{T}^2; \mathbb{Z}) \cong \mathbb{Z}[dx_c \wedge dx_a] \oplus \mathbb{Z}[dx_c \wedge dx_b] \oplus \mathbb{Z}[dx_a \wedge dx_b] \end{aligned}$$

where $\tilde{K}^0(S^1 \times \mathbb{T}^2)$ denotes the reduced K -theory, while for $i \neq j \in \{a, b, c\}$, \mathcal{L}_{ij} denotes the line bundle supported on the subtorus $\mathbb{T}_i \times \mathbb{T}_j$ of $S^1 \times \mathbb{T}^2$ with first Chern class equal to $[dx_i \wedge dx_j]$.

Note that unlike in the untwisted ($k = 0$) case, there is no generator for the “rank” in twisted K -theory. Strictly speaking, \mathcal{L}_{ij} should refer to the rank-0 virtual bundle $\tilde{\mathcal{L}}_{ij} := \mathcal{L}_{ij} - \mathbf{1}$. We will write the latter when this clarification is required.

4.1.2. *Generators for $K^\bullet(S^1 \times \mathbb{T}^2)$.* The K^0 group of the 2-torus $S^1 \times \mathbb{T} = \mathbb{T}_c \times \mathbb{T}_a$ is simply

$$K_0(C(S^1 \times \mathbb{T})) = \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[P_{\text{Bott}} - \mathbf{1}] = \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[dx_c \wedge dx_a],$$

where we have used the Chern character isomorphism in the second equality to represent the second generator in terms of a differential form. Similarly, the K^1 group is

$$K_1(C(S^1 \times \mathbb{T})) \cong H^1(S^1 \times \mathbb{T}) = \mathbb{Z}[U_c] \oplus \mathbb{Z}[U_a] = \mathbb{Z}[dx_c] \oplus \mathbb{Z}[dx_a],$$

where $U_c, U_a : S^1 \times \mathbb{T} \rightarrow U(1)$ are continuous maps that are also generators of $H^1(S^1 \times \mathbb{T}; \mathbb{Z})$ with odd degree 1 Chern class $[dx_c]$ and $[dx_a]$ respectively.

4.1.3. *Generators for $K_0(C^*(\text{Heis}^{\mathbb{Z}}(k)))$.* The Bott projections for the commutative subalgebras $C^*(U, W) \cong C(\mathbb{T}_{ac}^2)$ and $C^*(V, W) = C^*(\mathbb{T}_{bc}^2)$ are denoted by $[P_{ac}]$ and $[P_{bc}]$ respectively (see e.g. [1] for explicit formulae).

By results of [1, 34] (see also the derivation of (5.1) in Section 5.1.1), $K_0(C^*(\text{Heis}^{\mathbb{Z}}(k)))$ is generated by the free module of rank one, $[\mathbf{1}]$, and also by the projections $[P_{ac}]$ and $[P_{bc}]$; that is,

$$K_0(C^*(\text{Heis}^{\mathbb{Z}}(k))) \cong \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[P_{ac}] \oplus \mathbb{Z}[P_{bc}].$$

Note that there is no “Bott projection” P_{ab} in $C^*(\text{Heis}^{\mathbb{Z}}(k))$ since $C^*(U, V)$ is not a subalgebra of $C^*(\text{Heis}^{\mathbb{Z}}(k))$.

4.1.4. K -theory of Heisenberg nilmanifold. The Heisenberg nilmanifold

$$\text{Nil}_k = \text{Heis}^{\mathbb{R}} / \text{Heis}^{\mathbb{Z}}(k)$$

is a principal circle bundle over \mathbb{T}_{ab}^2 with Chern class equal to $k[dx_a \wedge dx_b]$, and is also a classifying space for $\text{Heis}^{\mathbb{Z}}(k)$.

The generators for $K^1(\text{Nil}_k)$ are described as follows:

$$\begin{aligned} K^1(\text{Nil}_k) &\cong \mathbb{Z}[Y] \oplus \mathbb{Z}[U_a] \oplus \mathbb{Z}[U_b] \\ &\cong H^{\text{odd}}(\text{Nil}_k) \cong \mathbb{Z}[dx_a \wedge dx_b \wedge \widehat{A}] \oplus \mathbb{Z}[dx_a] \oplus \mathbb{Z}[dx_b] \end{aligned}$$

where \widehat{A} is a connection on the principal circle bundle Nil_k over \mathbb{T}_{ab}^2 with curvature $d\widehat{A} = kdx_a \wedge dx_b$. Here $Y : \text{Nil}_k \rightarrow SU(2)$ is a degree 1 continuous map with odd degree-3 Chern class $[dx_a \wedge dx_b \wedge \widehat{A}]$ which is the generator of $H^3(\text{Nil}_k; \mathbb{Z})$, while $U_a, U_b : \text{Nil}_k \rightarrow U(1)$ are continuous maps that are also generators of $H^1(\text{Nil}_k; \mathbb{Z})$ with odd degree-1 Chern classes $[dx_a]$ and $[dx_b]$ respectively.

5. PROOF OF MAIN RESULTS

5.1. Explicit maps on generators.

5.1.1. *Bulk-boundary map ∂ .* Recall from (3.8) that $C^*(\text{Heis}^{\mathbb{Z}}(k)) \cong C(S^1 \times \mathbb{T}_a) \rtimes_{\alpha_k} \mathbb{Z}_b$. Its K -theory groups can be calculated along the lines of [1] (which handled the $k = 1$ case). The PV sequence ([44]) associated to the action α_k of \mathbb{Z} on $C(S^1 \times \mathbb{T}_a)$ is

$$0 \rightarrow \text{coker}_{1-\alpha_{k*}}(K_{\bullet}(C(S^1 \times \mathbb{T}_a)) \rightarrow K_{\bullet}(C^*(\text{Heis}^{\mathbb{Z}}(k))) \xrightarrow{\partial} \ker_{1-\alpha_{k*}}(K_{\bullet-1}(C(S^1 \times \mathbb{T}_a)) \rightarrow 0$$

Using the explicit expression (3.8), we can compute the induced map α_{k*} on K -theory as follows. It can be shown that α_{k*} acts trivially on $K_0(C(S^1 \times \mathbb{T}_a)) = \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[P_{ac}]$. Thus

$$\begin{aligned} \ker_{1-\alpha_{k*}}(K_0(C(S^1 \times \mathbb{T}_a))) &= \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[P_{ac}], \\ \text{Im}_{1-\alpha_{k*}}(K_0(C(S^1 \times \mathbb{T}_a))) &= [\mathbf{0}], \\ \text{coker}_{1-\alpha_{k*}}(K_0(C(S^1 \times \mathbb{T}_a))) &= \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[P_{ac}]. \end{aligned}$$

Next, $K^1(S^1 \times \mathbb{T}_a) \cong K_1(C(S^1 \times \mathbb{T}_a)) = \mathbb{Z}[W] \oplus \mathbb{Z}[U]$, and we recall from (3.9) that $\alpha_k(U^a W^c) = U^a W^{c-ka}$. Therefore,

$$\begin{aligned} \ker_{1-\alpha_{k*}}(K_1(C(S^1 \times \mathbb{T}_a))) &= \mathbb{Z}[W], \\ \text{Im}_{1-\alpha_{k*}}(K_1(C(S^1 \times \mathbb{T}_a))) &= \mathbb{Z}[W^k], \\ \text{coker}_{1-\alpha_{k*}}(K_1(C(S^1 \times \mathbb{T}_a))) &= \mathbb{Z}_k[W] \oplus \mathbb{Z}[U]. \end{aligned}$$

Thus, for $\bullet = 0$, the PV sequence is

$$0 \rightarrow \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[P_{ac}] \rightarrow \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[P_{ac}] \oplus \mathbb{Z}[P_{bc}] \xrightarrow{\partial} \mathbb{Z}[W] \rightarrow 0, \quad (5.1)$$

with ∂ taking $[P_{bc}] \mapsto -[W]$ and annihilating $[\mathbf{1}]$ and $[P_{ac}]$. Equivalently, with

$$[\tilde{P}_{ic}] := [P_{ic}] - [\mathbf{1}], \quad i = a, b,$$

we can write

$$\partial([\tilde{P}_{bc}]) = -[W], \quad \partial([\tilde{P}_{ac}]) = 0 = \partial([\mathbf{1}]).$$

For $\bullet = 1$, the PV sequence ([44]) is

$$0 \rightarrow \mathbb{Z}_k[W] \oplus \mathbb{Z}[U] \rightarrow \mathbb{Z}_k[W] \oplus \mathbb{Z}[U] \oplus \mathbb{Z}[V] \oplus \mathbb{Z}[V_a] \xrightarrow{\partial} \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[P_{ac}] \rightarrow 0,$$

where V_a is a unitary constructed as in pp. 7 of Ref. [1] (it is roughly the analogue of the top class in $H^3(\mathbb{T}^3)$). The boundary map takes $[V] \mapsto -[\mathbf{1}]$, $[V_a] \mapsto [P_{ac}]$, and annihilates $[W]$ and $[U]$.

5.1.2. Circle bundle T -duality T_1 . Viewing $S^1 \times \mathbb{T}^2$ as a circle bundle over \mathbb{T}_{ab}^2 , we use the results of [7, 8] to deduce that T_1 , the T -duality isomorphism for circle bundles, maps $K^0(S^1 \times \mathbb{T}^2, kH_1)$ to $K^1(\text{Nil}_k)$ as follows:

$$\begin{aligned} T_1([dx_c \wedge dx_a]) &= [dx_a] \\ T_1([dx_c \wedge dx_b]) &= [dx_b] \\ T_1([dx_a \wedge dx_b]) &= [dx_a \wedge dx_b \wedge \hat{A}]. \end{aligned}$$

From this, we deduce that on the (reduced) K -theory classes,

$$\begin{aligned} T_1([\tilde{\mathcal{L}}_{ca}]) &= [U_a] \\ T_1([\tilde{\mathcal{L}}_{cb}]) &= [U_b] \\ T_1([\tilde{\mathcal{L}}_{ab}]) &= [Y], \end{aligned}$$

giving a complete description of the isomorphism T_1 on the level of generators.

5.1.3. Baum–Connes map T_2 . The map T_2 can be identified with the Baum–Connes assembly map [2] (together with Poincaré duality) for the discrete group $\text{Heis}^{\mathbb{Z}}(k)$, which is known to be an isomorphism, see Proposition 2.5 in [49],

$$\mu : K^1(\text{Nil}_k) \rightarrow K_0(C^*(\text{Heis}^{\mathbb{Z}}(k))).$$

In fact, Rosenberg shows that this Baum–Connes map can be deduced from the Baum–Connes map for the torus \mathbb{T}^2 (which is essentially the Fourier transform) via a PV sequence ([44], often used in our paper) and the 5-Lemma.

Upon applying the von Neumann trace $\tau : C^*(\text{Heis}^{\mathbb{Z}}(k)) \rightarrow \mathbb{C}$, we see that

$$\tau(\mu(Y)) = \int_{\text{Nil}_k} Ch^{odd}(Y) = 1.$$

Note also that

$$\tau(\mu(U_a)) = 0 = \tau(\mu(U_b)).$$

Now there are two nontrivial line bundles \mathcal{P}_{ac} and \mathcal{P}_{bc} over Nil_k , which are explicitly constructed in Appendix A.1. Their respective Chern classes are $[dx_a \wedge \hat{A}]$ and $[dx_b \wedge \hat{A}]$,

which themselves correspond to group 2-cocycles σ_{ac} and σ_{bc} on the discrete group $\text{Heis}^{\mathbb{Z}}(k)$. These give rise to cyclic 2-cocycles τ_{ac} and τ_{bc} on the smooth subalgebra of $C^*(\text{Heis}^{\mathbb{Z}}(k))$ [15] (see Appendix A.2), and hence define maps² $K_0(C^*(\text{Heis}^{\mathbb{Z}}(k))) \rightarrow \mathbb{Z}$.

We then compute, using the Connes–Moscovici higher index theorem [15],

$$\tau_{ac}(\mu(U_b)) = \int_{\text{Nil}_k} dx_a \wedge \widehat{A} \wedge Ch^{odd}(U_b) = \int_{\text{Nil}_k} dx_a \wedge \widehat{A} \wedge dx_b = -1$$

and

$$\tau_{bc}(\mu(U_a)) = \int_{\text{Nil}_k} dx_b \wedge \widehat{A} \wedge Ch^{odd}(U_a) = \int_{\text{Nil}_k} dx_b \wedge \widehat{A} \wedge dx_a = 1.$$

On the other hand, Proposition A.1 says that

$$\begin{aligned} \tau_{ac}(P_{ac}) &= 1, \quad \tau_{ac}(P_{bc}) = 0, \quad \tau_{ac}(\mathbf{1}) = 0 \\ \tau_{bc}(P_{bc}) &= 1, \quad \tau_{bc}(P_{ac}) = 0, \quad \tau_{bc}(\mathbf{1}) = 0 \\ \tau(P_{ac}) &= \tau(P_{bc}) = \tau(\mathbf{1}) = 1. \end{aligned}$$

This implies that

$$\begin{aligned} [\mu(Y)] &= [\mathbf{1}] \\ [\mu(U_a)] &= [\widetilde{P}_{bc}] \\ [\mu(U_b)] &= -[\widetilde{P}_{ac}]. \end{aligned}$$

That is,

$$\begin{aligned} T_2([Y]) &= [\mathbf{1}] \\ T_2([U_a]) &= [\widetilde{P}_{bc}] \\ T_2([U_b]) &= -[\widetilde{P}_{ac}]. \end{aligned}$$

5.1.4. *Torus bundle T-duality* $T = T_2 \circ T_1$. We deduce that

$$\begin{aligned} T([\widetilde{\mathcal{L}}_{ca}]) &= [\widetilde{P}_{bc}] \\ T([\widetilde{\mathcal{L}}_{cb}]) &= -[\widetilde{P}_{ac}] \\ T([\widetilde{\mathcal{L}}_{ab}]) &= [\mathbf{1}], \end{aligned}$$

giving a complete description of T on the level of generators.

²When $k \neq 1$, there is also a nontrivial line bundle \mathcal{P}_{ab} over Nil_k whose (integral) Chern class is k -torsion. It is obtained from pulling back $\text{Nil}_1 \rightarrow \mathbb{T}^2$ along the bundle projection $\text{Nil}_k \rightarrow \mathbb{T}^2$. We do not get another cyclic 2-cocycle from this.

5.1.5. *Restriction maps ι^* .* We have $K^0(S^1 \times \mathbb{T}) \cong \mathbb{Z}[\mathbf{1}] \oplus \mathbb{Z}[\tilde{\mathcal{L}}_{ca}]$, and the restriction map on K -theory induced by $\iota : S^1 \times \mathbb{T} \rightarrow S^1 \times \mathbb{T}^2 \rightarrow$ is easy to compute:

$$\begin{aligned}\iota^*([\tilde{\mathcal{L}}_{ca}]) &= [\tilde{\mathcal{L}}_{ca}] \\ \iota^*([\tilde{\mathcal{L}}_{cb}]) &= [\mathbf{0}] \\ \iota^*([\tilde{\mathcal{L}}_{ab}]) &= [\mathbf{0}].\end{aligned}$$

For the other inclusion map $\iota : S^1 \times \mathbb{T} = \mathbb{T}_{ac}^2 \rightarrow \text{Nil}_k$ (see (3.15)), the restriction map ι^* in K -theory is simply

$$\begin{aligned}\iota^*[Y] &= [\mathbf{0}] \\ \iota^*[U_a] &= [U_a] \\ \iota^*[U_b] &= [\mathbf{0}].\end{aligned}\tag{5.2}$$

5.1.6. *Commutative T-duality maps T_a, T_c, T_{ac} .* The map

$$T_a : K^0(S^1 \times \mathbb{T}) \longrightarrow K^1(S^1 \times \mathbb{T}) \cong K_1(C(S^1 \times \mathbb{T}))$$

in (3.10) refers to the 1D commutative T-duality map with respect to $\mathbb{T} = \mathbb{T}_a$. It takes

$$T_a : [\mathbf{1}] \mapsto -[U], \quad [\tilde{\mathcal{L}}_{ca}] \sim [dx_c \wedge dx_a] \mapsto -[dx_c] \sim -[W]$$

(the parametrising circle $S^1 = \mathbb{T}_c$ goes along for the ride).

Similarly, the map

$$T_c : K^0(S^1 \times \mathbb{T}) \longrightarrow K^1(S^1 \times \mathbb{T})$$

in (3.23) refers to the 1D commutative T-duality map with respect to $S^1 = \mathbb{T}_c$. It takes

$$T_c : [\tilde{\mathcal{L}}_{ca}] \mapsto [U_a], \quad [\mathbf{1}] \mapsto -[U_c].$$

Finally, the full T-duality map

$$T_{ac} : K^1(S^1 \times \mathbb{T}) \longrightarrow K^1(S^1 \times \mathbb{T})$$

in (3.23) is the composition $T_a \circ T_c^{-1}$, which takes

$$T_{ac} : [U_a] \mapsto -[W], \quad [U_c] \mapsto -[U].$$

5.2. Assembling K -theory maps together.

Proof of Theorem 3.1: T-duality trivializes the bulk-boundary correspondence, first version.
We can now verify the commutativity of the diagram (3.10).

$$\begin{aligned}T_a \circ \iota^*([\tilde{\mathcal{L}}_{ca}]) &= T_a([\tilde{\mathcal{L}}_{ca}]) \\ &= -[W] \\ &= \partial([\tilde{P}_{bc}]) \\ &= \partial \circ T([\tilde{\mathcal{L}}_{ca}]),\end{aligned}$$

and

$$\begin{aligned} T_a \circ \iota^*([\tilde{\mathcal{L}}_{cb}]) &= [\mathbf{0}] = \partial[\tilde{P}_{ac}] = \partial \circ T([\tilde{\mathcal{L}}_{cb}]) \\ T_a \circ \iota^*([\tilde{\mathcal{L}}_{ab}]) &= [\mathbf{0}] = \partial[\mathbf{1}] = \partial \circ T([\tilde{\mathcal{L}}_{ab}]), \end{aligned}$$

so $T_a \circ \iota^* = \partial \circ T$. □

Proof of Theorem 3.2: T-duality trivializes the bulk-boundary correspondence, second version.
We can also verify the commutativity of the diagram (3.11).

$$\begin{aligned} T_{ac} \circ \iota^*([U_a]) &= T_{ac}([U_a]) \\ &= -[W] \\ &= \partial([\tilde{P}_{bc}]) \\ &= \partial \circ T_2([U_a]), \end{aligned}$$

and

$$\begin{aligned} T_{ac} \circ \iota^*([U_b]) &= [\mathbf{0}] = \partial(-[\tilde{P}_{ac}]) = \partial \circ T_2([U_b]) \\ T_{ac} \circ \iota^*([Y]) &= [\mathbf{0}] = \partial[\mathbf{1}] = \partial \circ T_2([Y]), \end{aligned}$$

so $T_{ac} \circ \iota^* = \partial \circ T_2$. □

Alternatively, Theorem 3.2 follows directly from Theorem 3.1, the easily verified commutativity of

$$\begin{array}{ccc} K^0(S^1 \times \mathbb{T}^2, kH_1) & \xrightarrow[\sim]{T_1} & K^1(\text{Nil}_k) \\ \downarrow \iota^* & & \downarrow \iota^* \\ K^0(S^1 \times \mathbb{T}) & \xrightarrow[\sim]{T_c} & K^1(S^1 \times \mathbb{T}) \end{array}$$

and the factorisations $T_2 = T \circ T_1^{-1}$ and $T_{ac} = T_a \circ T_c^{-1}$; in other words, the combined diagram (3.24) commutes.

Remark 5.1. For the degree-shifted versions of (3.10) and (3.11), a torsion subgroup \mathbb{Z}_k appears in the two K -theory groups on the top rows (when $k \neq 1$). We can ignore these for the commutativity of (3.10) since the K -theory groups in the bottom rows are torsion-free; the use of cyclic theory continues to work in this case, but we have left out the explicit computations.

6. TOPOLOGICAL INSULATORS WITH SCREW DISLOCATIONS AND THE HEISENBERG GROUP

In [47], it was proposed that screw dislocations (Fig. 1) in a 3D time-reversal invariant topological insulator can host topologically protected modes traversing the bulk bandgap. This phenomenon was demonstrated numerically in a tight-binding model, where a pair of screw dislocations was introduced into a unit cell subjected to periodic boundary conditions. A subsequent investigation into such dislocation-bound 1D modes can be found in [27].

Prior to these papers, the effect of a single screw dislocation on Landau levels was studied in [20], using a differential-geometric theory of defects as described in [4, 29, 33]. The latter framework is very general, and allowed the authors of [52] to consider the quantum dynamics of a free particle in the presence of a cylindrically symmetric *distribution* of parallel screw dislocations. It was found that the energy levels were quantized to *elastic Landau levels* reminiscent of those of a charged particle in the presence of a uniform magnetic field.

In general, defects such as screw dislocations break the \mathbb{Z}^3 translation symmetry of the original 3D Euclidean lattice, which is needed to define the Fu–Kane–Mele \mathbb{Z}_2 -invariants [19] characterizing a 3D time-reversal invariant topological insulator. This raises the question as to whether the standard “undeformed” Fu–Kane–Mele invariants continue to be appropriate. Furthermore, the notion of a unit cell in \mathbb{R}^3 , used in [47], also requires this symmetry to be well-defined. We can circumvent this general difficulty by requiring the defects themselves to be distributed in a sufficiently regular manner (this is implicit in the imposition of periodic boundary conditions in [47, 27]). This regularity allows us to define topological invariants which are, in a precise sense, deformed versions of the usual ones arising from (commutative) Bloch theory. We explain this deformation in the simpler case without time-reversal symmetry, leaving the time-reversal invariant case for a subsequent work.

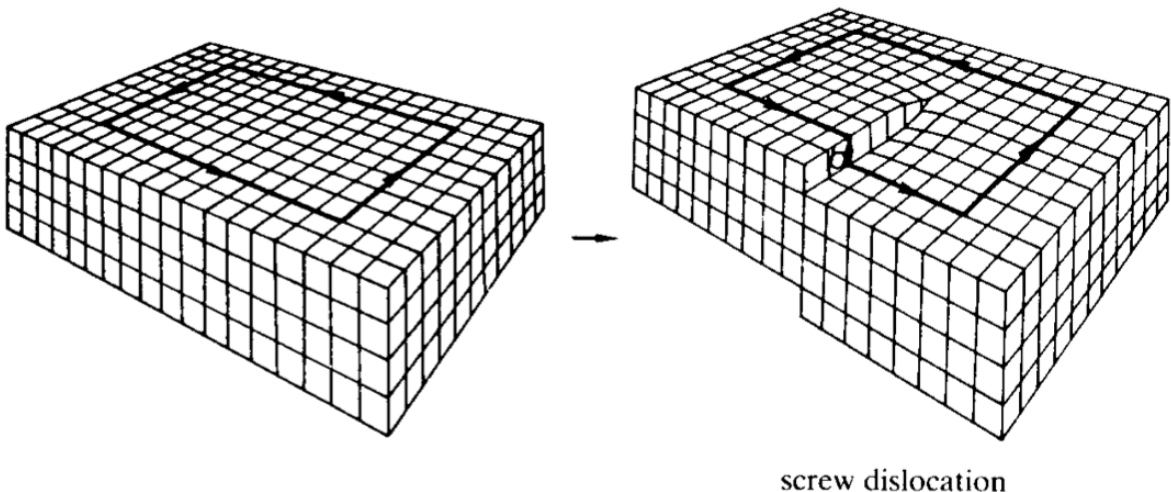


FIGURE 1. An elementary screw dislocation (Source: pp. 786 in [33]) with Burgers vector in the vertical direction. A circuit of translations in the horizontal direction enclosing the dislocation ends at a lattice site which differs from the starting point by a vertical translation.

Consider at first a 3D lattice of atomic sites in \mathbb{R}^3 with lattice translations along axes labelled by $\mathbb{R}_a, \mathbb{R}_b, \mathbb{R}_c$. Then introduce parallel elementary screw dislocations along the \mathbb{R}_c direction, which are distributed uniformly so that the dislocations are located on a 2D lattice when projected onto the $\mathbb{R}_a\text{--}\mathbb{R}_b$ plane. Let U and V be generating translations for this latter lattice in the a and b directions respectively, taken along the atomic bonds in the *distorted* 3D lattice. Then U and V no longer commute, but instead obey $UV = WVU$ with

W being a translation in the c -direction by one atom — these are the defining commutation relations for the integer Heisenberg group. Thus one may consider for instance, a tight-binding Hamiltonian on the above distorted 3D lattice symmetric under “translations” by $\text{Heis}^{\mathbb{Z}}$ as above.

6.1. $\text{Heis}^{\mathbb{Z}}$ -symmetric Hamiltonians and the noncommutative Brillouin zone. Ordinary lattices in \mathbb{R}^3 have an abelian group \mathbb{Z}^3 of translation symmetries which allows the use of Bloch theory for defining topological invariants of gapped Hamiltonians with these symmetries. In this commutative case, the Brillouin zone is a 3-torus which is the Pontryagin dual of \mathbb{Z}^3 , and the topological invariants live in various K -theory groups of this torus [18].

A noncommutative generalization is needed for the integer quantum Hall effect, where the *magnetic* translation symmetries do not commute but instead generate a noncommutative torus [3], which is a deformation of the ordinary 2-torus. The analysis of Hamiltonians symmetric under $\text{Heis}^{\mathbb{Z}}$ follows along much the same lines. Namely, the “noncommutative Brillouin zone” is the integer Heisenberg group C^* -algebra $C^*(\text{Heis}^{\mathbb{Z}})$, which as we have already seen in Section 3.1, is a parametrised deformation of the ordinary 3-torus.

In the ordinary commutative case, the 3-torus has three independent first Chern classes, represented by the Bott projections P_{ac}, P_{bc}, P_{ab} which are each non-trivial on one of the three independent choices of 2-subtori. Each of these Chern classes corresponds to a 3D Chern insulator which can be thought of as layers of a standard 2D Chern insulator. The quantum Hall state in 3D is again characterised by three integers [35], which come from the noncommutative Chern classes for the 3D noncommutative torus.

However, when we consider $\text{Heis}^{\mathbb{Z}}$ symmetry, the spatial direction along \mathbb{R}_c is singled out, and we see its signature in the “disappearance” of the transverse Bott projection P_{ab} from the generators of the K -theory of $C^*(\text{Heis}^{\mathbb{Z}})$ calculated in Section 4.1.3. Furthermore, the bulk-boundary homomorphism ∂ maps only onto $\mathbb{Z}[W] \in K_1(C(S^1 \times \mathbb{T})) = K_1(C^*(\mathbb{Z}_{ac}^2))$ but not in $\mathbb{Z}[U]$ (if we took the boundary to be transverse to the a -direction, then ∂ will again land in $\mathbb{Z}[W]$ but not in $\mathbb{Z}[V]$). Recall that W is the generator of translations in the c -direction. This suggests that edge modes propagate only in the c -direction, which is consistent with the intuition in [47, 27] that topologically protected 1D propagating modes develop along screw dislocations.

Remark 6.1. The continuum limit of a uniform distribution of parallel screw dislocations was studied in [11], leading to the real Heisenberg group manifold $\text{Heis}^{\mathbb{R}}$ (which is topologically still \mathbb{R}^3) as the so-called *material manifold*. The discrete subgroup $\text{Heis}^{\mathbb{Z}}$ is a lattice in $\text{Heis}^{\mathbb{R}}$, and the quotient nilmanifold $\text{Heis}^{\mathbb{R}}/\text{Heis}^{\mathbb{Z}}$ is the appropriate fundamental domain, or “Wigner–Seitz cell”, to use here.

The restriction map in (3.11), which is the T-dualized bulk-boundary homomorphism, now has a direct interpretation: it is a simple restriction of K -theory classes from the deformed bulk fundamental domain (a nilmanifold) to the boundary fundamental domain (a torus).

APPENDIX A. CYCLIC COCYCLES ON INTEGER HEISENBERG GROUP ALGEBRA AND PAIRING WITH K -THEORY

A.1. **Line bundles over Nil_k .** The group cohomology $H_{\text{group}}^2(\text{Heis}^{\mathbb{Z}}(k); \mathbb{Z}) \cong H^2(\text{Nil}_k; \mathbb{Z}) = \mathbb{Z}^2 \oplus \mathbb{Z}_k$ is generated by the following central extensions of $\text{Heis}^{\mathbb{Z}}(k)$ by \mathbb{Z} . First, there is

$$\widetilde{\text{Heis}^{\mathbb{Z}}(k)}^{bc} := \left\{ \begin{pmatrix} 1 & a & \frac{c}{k} & \frac{d}{2k} \\ 0 & 1 & b & \frac{1}{2}b^2 \\ 0 & 0 & 1 & b \\ 0 & 0 & 0 & 1 \end{pmatrix} \mid a, b, c, d \in \mathbb{Z} \right\}$$

with quotient map onto $\text{Heis}^{\mathbb{Z}}$ being restriction to the upper left 3×3 submatrix. Second, there is

$$\widetilde{\text{Heis}^{\mathbb{Z}}(k)}^{ac} := \left\{ \begin{pmatrix} 1 & a & \frac{1}{2}a^2 & \frac{d}{2k} \\ 0 & 1 & a & \frac{c}{k} \\ 0 & 0 & 1 & b \\ 0 & 0 & 0 & 1 \end{pmatrix} \mid a, b, c, d \in \mathbb{Z} \right\}$$

with quotient map being restriction to the lower right 3×3 submatrix. Third, there is

$$\widetilde{\text{Heis}^{\mathbb{Z}}(k)}^{ab} := \left\{ \begin{pmatrix} 1 & a & \frac{c}{k} & d \\ 0 & 1 & b & b \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \mid a, b, c, d \in \mathbb{Z} \right\}.$$

with quotient map being restriction to the upper left 3×3 submatrix.

Their group cocycles are, respectively,

$$\begin{aligned} -\sigma_{bc}((a_1, b_1, c_1), (a_2, b_2, c_2)) &= 2c_1b_2 + ka_1b_2^2 \sim -b_1c_2 + c_1b_2 - ka_1b_1b_2 \\ \sigma_{ac}((a_1, b_1, c_1), (a_2, b_2, c_2)) &= 2a_1c_2 + ka_1^2b_2 \sim a_1c_2 - c_1a_2 - ka_1a_2b_2 \\ \sigma_{ab}((a_1, b_1, c_1), (a_2, b_2, c_2)) &= a_1b_2. \end{aligned}$$

where \sim means “cohomologous to”. The cocycles σ_{bc} and σ_{ac} are the free generators, whereas one can check that $k \cdot \sigma_{ab}$ is a coboundary so that σ_{ab} generates \mathbb{Z}_k . Furthermore, we observe that σ_{bc} vanishes on \mathbb{Z}_{ac}^2 while σ_{ac} vanishes on \mathbb{Z}_{bc}^2 .

These constructions also give rise to central extensions of $\text{Heis}^{\mathbb{R}}$ by \mathbb{R} by allowing $a, b, c, d \in \mathbb{R}$. The non-trivial line bundles over $\text{Nil}_k = B\text{Heis}^{\mathbb{Z}}(k)$ are obtained by quotienting these real extensions by the discrete ones.

A.2. **Cyclic cocycles from group cocycles.** According to Connes–Moscovici [15], we can construct cyclic 2-cocycles τ_{bc}, τ_{ac} on (a smooth subalgebra of) $C^*(\text{Heis}^{\mathbb{Z}}(k))$ from the group 2-cocycles σ_{bc}, σ_{ac} on $\text{Heis}^{\mathbb{Z}}(k)$. Explicitly, τ_{bc} is

$$\tau_{bc}(f_0, f_1, f_2) = \sum_{\substack{\gamma_0\gamma_1\gamma_2=\text{id} \\ \gamma_0, \gamma_1, \gamma_2 \in \text{Heis}^{\mathbb{Z}}(k)}} \sigma_{bc}(\gamma_1, \gamma_2) f_0(\gamma_0) f_1(\gamma_1) f_2(\gamma_2), \quad f_i \in \mathbb{C}(\text{Heis}^{\mathbb{Z}}(k)), \quad (\text{A.1})$$

extended to the smooth subalgebra of $C^*(\text{Heis}^{\mathbb{Z}}(k))$, and similarly for τ_{ac} . These cyclic cocycles then pair with $K_0(C^*(\text{Heis}^{\mathbb{Z}}(k)))$ in the usual way, by extending the formula (A.1) to smooth matrix algebras over $\mathbb{C}(\text{Heis}^{\mathbb{Z}}(k))$,

$$\tau_{bc}(f_0 \otimes A_0, f_1 \otimes A_1, f_2 \otimes A_2) = \text{tr}(A_0 A_1 A_2) \sum_{\substack{\gamma_0 \gamma_1 \gamma_2 = \text{id} \\ \gamma_0, \gamma_1, \gamma_2 \in \text{Heis}^{\mathbb{Z}}(k)}} \sigma_{bc}(\gamma_1, \gamma_2) f_0(\gamma_0) f_1(\gamma_1) f_2(\gamma_2), \quad (\text{A.2})$$

and taking

$$\langle \tau_{bc}, [P] \rangle = \tau_{bc}(P, P, P), \quad [P] \in K_0(C^*(\text{Heis}^{\mathbb{Z}}(k))).$$

The aim is to show that τ_{bc}, τ_{ac} are linearly independent and non-zero. Note that σ_{ab} is torsion, so we do not get another cyclic cocycle τ_{ab} from it.

Corresponding to the two commutative subalgebras $C^*(\mathbb{Z}_{bc}^2), C^*(\mathbb{Z}_{ac}^2)$, there are two projections P_{bc}, P_{ac} (in their matrix algebras) which are the Bott projections P_{Bott} when $C^*(\mathbb{Z}_{bc}^2)$ and $C^*(\mathbb{Z}_{ac}^2)$ are identified with $C(\mathbb{T}^2)$. As an element of $C^*(\text{Heis}^{\mathbb{Z}}(k))$, P_{bc} can be thought of as the function $\mathbb{Z} \rightarrow C^*(\mathbb{Z}_{bc}^2)$ which is P_{Bott} when $a = 0 \in \mathbb{Z}$ and zero otherwise; similarly for P_{ac} .

Recall that there is also a standard cyclic 2-cocycle ψ on $C^\infty(\mathbb{T}^2)$, defined by

$$\psi(f, g, h) = 2\pi i \tau(f[\partial_1 g, \partial_2 h]), \quad f, g, h \in C^\infty(\mathbb{T}^2),$$

where $\tau(\sum_{m,n \in \mathbb{Z}} a_{m,n} U^m W^n) = a_{0,0}$ with U, W the two commuting unitaries generating $C^\infty(\mathbb{T}^2)$. The derivations are such that $\partial_1(U) = U, \partial_2(W) = W$, and $\partial_1(W) = 0 = \partial_2(U)$. It is known that the pairing of ψ with $K_0(C(\mathbb{T}^2))$ is given by

$$\langle \psi, [P_{\text{Bott}}] \rangle = 1, \quad \langle \psi, [\mathbf{1}] \rangle = 0,$$

and that $\tau(P_{\text{Bott}}) = 1 = \tau(\mathbf{1})$.

We can rewrite the formula for ψ in a form which resembles (A.1). Let

$$\begin{aligned} f &= \sum_{m,n \in \mathbb{Z}} f_{m,n} U^m W^n, \\ g &= \sum_{p,q \in \mathbb{Z}} g_{p,q} U^p W^q, \\ h &= \sum_{r,s \in \mathbb{Z}} h_{r,s} U^r W^s, \end{aligned}$$

then

$$\psi(f, g, h) = \sum_{p,q,r,s \in \mathbb{Z}} (ps - rq) \cdot f_{-p-r, -q-s} g_{p,q} h_{r,s}, \quad (\text{A.3})$$

which is then extended as in (A.2) to a pairing with $K_0(C(\mathbb{T}^2))$. Let $\gamma_0 \equiv (m, n), \gamma_1 \equiv (p, q), \gamma_2 \equiv (r, s)$ be elements of \mathbb{Z}^2 , so $c(\gamma_1, \gamma_2) = ps - rq$ defines a 2-cocycle on $\mathbb{Z}^2 \times \mathbb{Z}^2$. Equation (A.3) can be rewritten as

$$\psi(f, g, h) = \sum_{\substack{\gamma_0 \gamma_1 \gamma_2 = \text{id} \\ \gamma_0, \gamma_1, \gamma_2 \in \mathbb{Z}^2}} c(\gamma_1, \gamma_2) f(\gamma_0) g(\gamma_1) h(\gamma_2),$$

and the statement that $\langle \psi, [P_{\text{Bott}}] \rangle = 1$ becomes

$$\langle \psi, [P_{\text{Bott}}] \rangle = \sum_{\substack{\gamma_0 \gamma_1 \gamma_2 = \text{id} \\ \gamma_0, \gamma_1, \gamma_2 \in \mathbb{Z}^2}} c(\gamma_1, \gamma_2) \text{tr} (P_{\text{Bott}}(\gamma_0) P_{\text{Bott}}(\gamma_1) P_{\text{Bott}}(\gamma_2)) = 1. \quad (\text{A.4})$$

A.3. Pairing with K -theory. We can now compute the pairing of τ_{bc} with $[P_{ac}]$ and $[P_{bc}]$:

$$\begin{aligned} \tau_{bc}(P_{ac}, P_{ac}, P_{ac}) &= \sum_{\substack{\gamma_0 \gamma_1 \gamma_2 = \text{id} \\ \gamma_0, \gamma_1, \gamma_2 \in \text{Heis}^{\mathbb{Z}}(k)}} \sigma_{bc}(\gamma_1, \gamma_2) \text{tr} (P_{ac}(\gamma_0) P_{ac}(\gamma_1) P_{ac}(\gamma_2)) \\ &= \sum_{\substack{\gamma_0 \gamma_1 \gamma_2 = \text{id} \\ \gamma_0, \gamma_1, \gamma_2 \in \mathbb{Z}_{ac}^2}} \sigma_{bc}(\gamma_1, \gamma_2) \text{tr} (P_{ac}(\gamma_0) P_{ac}(\gamma_1) P_{ac}(\gamma_2)), \end{aligned}$$

since P_{ac} is supported on \mathbb{Z}_{ac}^2 . But σ_{bc} restricted to $(\gamma_1, \gamma_2) \in \mathbb{Z}_{ac}^2 \times \mathbb{Z}_{ac}^2$ vanishes, as we had found at the end of the previous subsection, so we conclude that

$$\langle \tau_{bc}, [P_{ac}] \rangle \equiv \tau_{bc}(P_{ac}, P_{ac}, P_{ac}) = 0.$$

Next,

$$\begin{aligned} \langle \tau_{bc}, [P_{bc}] \rangle \equiv \tau_{bc}(P_{bc}, P_{bc}, P_{bc}) &= \sum_{\substack{\gamma_0 \gamma_1 \gamma_2 = \text{id} \\ \gamma_0, \gamma_1, \gamma_2 \in \text{Heis}^{\mathbb{Z}}(k)}} \sigma_{bc}(\gamma_1, \gamma_2) \text{tr} (P_{bc}(\gamma_0) P_{bc}(\gamma_1) P_{bc}(\gamma_2)) \\ &= \sum_{\substack{\gamma_0 \gamma_1 \gamma_2 = \text{id} \\ \gamma_0, \gamma_1, \gamma_2 \in \mathbb{Z}_{bc}^2}} \sigma_{bc}(\gamma_1, \gamma_2) \text{tr} (P_{bc}(\gamma_0) P_{bc}(\gamma_1) P_{bc}(\gamma_2)) \\ &= \sum_{\substack{\gamma_0 \gamma_1 \gamma_2 = \text{id} \\ \gamma_0, \gamma_1, \gamma_2 \in \mathbb{Z}^2}} c(\gamma_1, \gamma_2) \text{tr} (P_{\text{Bott}}(\gamma_0) P_{\text{Bott}}(\gamma_1) P_{\text{Bott}}(\gamma_2)) \\ &= 1, \end{aligned}$$

where the second equality follows from the fact that P_{bc} is supported on \mathbb{Z}_{bc}^2 (where it is P_{Bott}), the third equality follows from $\sigma_{bc}|_{\mathbb{Z}_{bc}^2 \times \mathbb{Z}_{bc}^2} = c$, and the fourth equality is (A.4).

To summarize,

Proposition A.1.

$$\langle \tau_{bc}, [P_{bc}] \rangle = 1, \quad \langle \tau_{bc}, [P_{ac}] \rangle = 0, \quad \langle \tau_{bc}, [\mathbf{1}] \rangle = 0.$$

In a similar vein, we also obtain

$$\langle \tau_{ac}, [P_{ac}] \rangle = 1, \quad \langle \tau_{ac}, [P_{bc}] \rangle = 0, \quad \langle \tau_{ac}, [\mathbf{1}] \rangle = 0.$$

It is known, from [1, 34] for instance, that $K_0(C^(\text{Heis}^{\mathbb{Z}}(k))) \cong \mathbb{Z}^3$ is generated by $[P_{ac}]$, $[P_{bc}]$, and the trivial projection $[\mathbf{1}]$ (the identity). We can define the 0-cocycle τ on the smooth subalgebra of $C^*(\text{Heis}^{\mathbb{Z}}(k))$ by $\tau(\sum_{r,s,t \in \mathbb{Z}} a_{r,s,t} U^r V^s W^t) = a_{0,0,0}$, whence we see that*

$$\langle \tau, [P_{ac}] \rangle = \langle \tau, [P_{bc}] \rangle = \langle \tau, [\mathbf{1}] \rangle = 1.$$

Thus τ_{bc}, τ_{ac} are linearly independent, and together with τ , can be used to distinguish elements of $K_0(C^*(\text{Heis}^{\mathbb{Z}}(k)))$ uniquely.

Remark A.2. Hadfield [22] also studies these pairings in the $k = 1$ case, however his calculations there are not complete, and do not exploit the simplifications that we do.

Acknowledgements. This work was supported by the Australian Research Council via ARC Discovery Project grants DP150100008 and DP130103924.

REFERENCES

- [1] Anderson, J., Paschke, W. The rotation algebra. *Houston J. Math.* **15**(1) 1–26 (1989)
- [2] Baum, P., Connes, A., Higson, N.: Classifying space for proper actions and K -theory of group C^* -algebras. *Contemp. Math.* **167** 240–291 (1994)
- [3] Bellissard, J., van Elst, A., Schulz-Baldes, H.: The noncommutative geometry of the quantum Hall effect. *J. Math. Phys.* **35**(10) 5373–5451 (1994)
- [4] Bilby, B.A., Bullough, R., Smith, E.: Continuous distributions of dislocations: a new application of the methods of non-Riemannian geometry. *Proc. Roy. Soc. London A* **231**(1185) 263–273 (1955)
- [5] Blackadar, B.: K -theory for operator algebras. *Math. Sci. Res. Inst. Publ.* **5** Cambridge Univ. Press, Cambridge (1998)
- [6] Bourne, C., Carey, A.L., Rennie, A.: The Bulk-Edge Correspondence for the Quantum Hall Effect in Kasparov theory. *Lett. Math. Phys.* **105**(9) 1253–1273 (2015) [[arXiv:1411.7527](#)]
- [7] Bouwknegt, P., Evslin, J., Mathai, V.: T-duality: Topology Change from H-flux. *Commun. Math. Phys.* **249**(2) 383–415 (2004) [[arXiv:hep-th/0306062](#)].
- [8] Bouwknegt, P., Evslin, J., Mathai, V.: On the Topology and Flux of T-Dual Manifolds. *Phys. Rev. Lett.* **92** 181601 (2004) [[arXiv:hep-th/0312052](#)].
- [9] Bunke, U., Schick, T.: On the topology of T-duality. *Rev. Math. Phys.* **17**(1) 77–112 (2005)
- [10] Carey, A., Hannabuss, K., Mathai, V., McCann, P.: Quantum Hall Effect on the hyperbolic plane. *Commun. Math. Phys.* **190**(3) 629–673 (1998) [[arXiv:dg-ga/9704006](#)].
- [11] Christodoulou, D., Kaelin, I.: On the mechanics of crystalline solids with a continuous distribution of dislocations. *Adv. Theor. Math. Phys.* **17**(2) 399–477 (2014)
- [12] Connes, A.: An analogue of the Thom isomorphism for crossed products of a C^* -algebra by an action of \mathbb{R} . *Adv. Math.* **39**(1) 31–55 (1981)
- [13] Connes, A.: Non-commutative differential geometry. *Publ. Math. Inst. Hautes Étude Sci.* **62**(1) 41–144 (1985)
- [14] Connes, A.: *Noncommutative Geometry*. Acad. Press, San Diego (1994)
- [15] Connes, A. and Moscovici, H.: Cyclic cohomology, the Novikov conjecture and hyperbolic groups. *Topology* **29**(3) 345–388 (1990)
- [16] Elbau, P., Graf, G.M.: Equality of bulk and edge Hall conductance revisited. *Commun. Math. Phys.* **229**(3) 415–432 (2002)
- [17] Fack, T., Skandalis, G.: Connes’ analogue of the Thom isomorphism for the Kasparov groups. *Invent. Math.* **64**(1) 7–14 (1981)
- [18] Freed, D.S., Moore, G.W.: Twisted equivariant matter. *Ann. Henri Poincaré* **14**(8) 1927–2023 (2013)
- [19] Fu, L., Kane, C.L., Mele, E.J.: Topological insulators in three dimensions. *Phys. Rev. Lett.* **98**(10) 106803 (2007)
- [20] Furtado, C., Moraes, F.: Landau levels in the presence of a screw dislocation. *Europhys. Lett.* **45**(3) 279–282 (1999)

- [21] Graf, G.M., Porta, M.: Bulk-edge correspondence for two-dimensional topological insulators. *Commun. Math. Phys.* **324**(3) 851–895 (2013)
- [22] Hadfield, T.D.: Fredholm modules over certain group C^* -algebras. Thesis (Ph.D.)-University of California, Berkeley (2001) 161 pp. ISBN: 978-0493-58344-0
- [23] Hannabuss, K.C., Mathai, V.: Noncommutative principal torus bundles via parametrised strict deformation quantization. *AMS Proc. Sympos. Pure Math.* **81** 133–148 (2010) [[arXiv:0911.1886](#)]
- [24] Hannabuss, K.C., Mathai, V.: Parametrised strict deformation quantization of C^* -bundles and Hilbert C^* -modules, *J. Aust. Math. Soc.* **90**(1) 25–38 (2011) [[arXiv:1007.4696](#)]
- [25] Hatcher, A.: Algebraic topology. Cambridge Univ. Press, Cambridge (2002)
- [26] Hatsugai, Y.: Chern number and edge states in the integer quantum Hall effect. *Phys. Rev. Lett.* **71**(22) 3697 (1993)
- [27] Imura, K.-I., Takane, Y., Tanaka, A.: Weak topological insulator with protected gapless helical states. *Phys. Rev. B* **84** 035443 (2011)
- [28] Kane, C.L., Mele, E.J.: \mathbb{Z}_2 Topological Order and the Quantum Spin Hall Effect. *Phys. Rev. Lett.* **95**(14) 146802 (2005)
- [29] Katanaev, M.O., Volovich, I.V.: Theory of defects in solids and three-dimensional gravity. *Ann. Phys.* **216**(1) 1–28 (1992)
- [30] Kellendonk, J., Richter, T., Schulz-Baldes, H.: Edge current channels and Chern numbers in the integer quantum Hall effect. *Rev. Math. Phys.* **14**(1) 87–119 (2002)
- [31] Kellendonk, J., Schulz-Baldes, H.: Boundary Maps for C^* -Crossed Products with an Application to the Quantum Hall Effect. *Commun. Math. Phys.* **249**(3) 611–637 (2004)
- [32] Kellendonk, J.: On the C^* -algebraic approach to topological phases for insulators. [arXiv:1509.06271](#).
- [33] Kleinert, H.: Gauge fields in condensed matter, vol. II, World Scientific (1989)
- [34] Kodaka, K.: K -theory for the C^* -algebras of the discrete Heisenberg group. *Tokyo J. Math.* **9**(2) 365–372 (1986)
- [35] Kohmoto, M., Halperin, B.I., Wu, Y.-S.: Diophantine equation for the three-dimensional quantum Hall effect. *Phys. Rev. B* **45** 13488–13493 (1992)
- [36] Loring, T.A.: K -theory and pseudospectra for topological insulators. *Ann. Physics* **356** 383–416 (2015)
- [37] Marcolli, M., Mathai, V.: Twisted index theory on good orbifolds. II. Fractional quantum numbers. *Commun. Math. Phys.* **217**(1) 55–87 (2001) [[arXiv:math/9911103](#)]
- [38] Mathai, V., Rosenberg, J.: T-duality for torus bundles with H-fluxes via noncommutative topology. *Commun. Math. Phys.* **253**(3) 705–721 (2005) [[arXiv:hep-th/0401168](#)]
- [39] Mathai, V., Rosenberg, J.: T-duality for torus bundles with H-fluxes via noncommutative topology, II; the high-dimensional case and the T-duality group. *Adv. Theor. Math. Phys.* **10**(1) 123–158 (2006) [[arXiv:hep-th/0508084](#)].
- [40] Mathai, V., Thiang, G.C.: T-duality and topological insulators. *J. Phys. A: Math. Theor. (Fast Track Communications)* **48**(42) 42FT02 (2015) [[arXiv:1503.01206](#)]
- [41] Mathai, V., Thiang, G.C.: T-duality trivializes bulk-boundary correspondence. 28pp. [[arXiv:1505.05250](#)]
- [42] Mathai, V., Thiang, G.C.: T-duality trivializes bulk-boundary correspondence: some higher dimensional cases. 15pp. [[arXiv:1506.04492](#)]
- [43] Packer, J., Raeburn, I.: Twisted crossed products of C^* -algebras. *Math. Proc. Cambridge Philos. Soc.* **106**(2) 293–311 (1989)
- [44] Pimsner, M., Voiculescu, D.: Exact sequences for K -groups and EXT -groups of certain cross-product C^* -algebras. *J. Operator Theory* **4** 93–118 (1980)
- [45] Prodan, E.: Virtual Topological Insulators with Real Quantized Physics. *Phys. Rev. B* **91** 245104 (2015)
- [46] Raeburn, I., Williams, D.P.: Morita equivalence and continuous-trace C^* -algebras. *Math. Surveys Monogr.* **60** Amer. Math. Soc. Providence, RI (1998)

- [47] Ran, Y., Zhang, Y., Vishwanath, A.: One-dimensional topologically protected modes in topological insulators with lattice dislocations. *Nature Physics* **5** 298–303 (2009)
- [48] Rieffel, M.A.: Strong Morita equivalence of certain transformation group C^* -algebras. *Math. Annalen.* **222**(1) 7–22 (1976)
- [49] Rosenberg, J.: C^* -algebras, positive scalar curvature, and the Novikov Conjecture. *Publ. Math. Inst. Hautes Étude Sci.* **58**(1) 197–212 (1983)
- [50] Rosenberg, J.: Continuous-trace algebras from the bundle theoretic point of view. *J. Aust. Math. Soc. (Ser. A)* **47**(3) 368–381 (1989)
- [51] Sheinbaum, D., Adem, A., Semenoff, G.: Topology of Fermi Surfaces and Anomalies. [arXiv:1509.01635](#).
- [52] Silva Netto, A.L., Furtado, C.: Elastic Landau levels. *J. Phys. Condens. Matter* **20**(12) 125209 (2008)
- [53] Thiang, G.C.: On the K -theoretic classification of topological phases of matter. *Ann. Henri Poincaré* (Online First) [[arXiv:1406.7366](#)]
- [54] Williams, D.P.: Crossed products of C^* -algebras. *Math. Surveys Monogr.* **134** Amer. Math. Soc., Providence (2007)
- [55] Witten, E.: Fermion Path Integrals And Topological Phases. [arXiv:1508.04715](#).

(Keith Hannabuss) MATHEMATICAL INSTITUTE, 24-29 ST. GILES', OXFORD, OX1 3LB, AND BALLIOL COLLEGE, OXFORD, OX1 3BJ, U.K.

E-mail address: kch@balliol.ox.ac.uk

(Varghese Mathai) DEPARTMENT OF PURE MATHEMATICS, SCHOOL OF MATHEMATICAL SCIENCES, UNIVERSITY OF ADELAIDE, ADELAIDE, SA 5005, AUSTRALIA

E-mail address: mathai.varghese@adelaide.edu.au

(Guo Chuan Thiang) DEPARTMENT OF PURE MATHEMATICS, SCHOOL OF MATHEMATICAL SCIENCES, UNIVERSITY OF ADELAIDE, ADELAIDE, SA 5005, AUSTRALIA

E-mail address: guo.thiang@adelaide.edu.au