

1 **Assessing the reproducibility of quartz OSL lifetime determinations derived using isothermal decay**

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8 **Highlights**

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10 Lifetime determinations made using isothermal decay measurements are tested.

11 Multiple-trap signal contributions can impact lifetime determinations.

12 Trap depth can be well constrained, but frequency factor is more variable.

13 Variations in the isothermal decay protocol can lead to variability in E and s.

14

15 **Abstract**

16

17 Signal stability is a key consideration when using luminescence dating techniques. The stability, or lifetime,
18 of a signal is one of the factors determining the upper age constraint for luminescence dating, and it has
19 been suggested that the signal being used for dating should have a lifetime at least ten times the age being
20 dated in order to limit age underestimation to an upper loss of 5%. Accurate derivations of signal stability
21 and associated kinetic parameters, such as trap depth and frequency factor, are also important parameters
22 for constraining rock cooling histories in thermochronometric techniques. This paper aims to assess the
23 reproducibility of lifetime determinations derived isothermal decay measurements. Variability arising from
24 changing the isothermal decay protocol used is tested. Simulating Arrhenius plot from fixed trap depth and
25 frequency values shows that whilst trap depth can be relatively well constrained, significant variability in
26 the frequency factor, hence signal lifetime, should be expected. This paper also uses luminescence signals
27 measured using different wavelengths to better understand the impact of signal from non-fast quartz OSL
28 components in lifetime calculations. The presence of contributions from non-fast OSL components in the
29 initial part of the OSL signal can result in the lifetime being calculated from charge contributions from
30 multiple traps, not solely the 325°C TL peak, as has been previously assumed. This effect can be reduced
31 however by stimulating luminescence signals with longer wavelengths to better isolate the signal from the
32 fast component.

33

34 **Keywords**

35

36 Isothermal decay, lifetime, quartz, OSL components, Arrhenius plot, E, s

37

38 **1. Introduction**

39

40 A key consideration when using luminescence dating for chronological applications is signal stability
41 (Spooner and Questiaux, 2000); it is important that trapped charge populations remain at the defects in the
42 crystal lattice of a mineral over sufficiently long periods of geological time relative to the timescale of
43 investigation (e.g. Preusser et al., 2009). It is this signal stability which is one of the factors controlling the
44 upper age constraint for luminescence dating, and Aitken (1985) recommends that the stability, or lifetime,
45 of the signal should be at least ten times the age being dated in order to limit age underestimation to an
46 upper loss of 5%. In recent years, efforts to extend the dateable age range using luminescence techniques
47 have seen the testing and application of a range of signals from multiple dosimeters. Methods have

48 included (but are not limited to) violet stimulated luminescence (e.g. Ankjægaard et al., 2013) and
49 thermally-transferred optically stimulated luminescence (e.g. Wang et al., 2006) signals in quartz, post
50 infra-red infra-red stimulated feldspar signals (e.g. Thomsen et al., 2008), and more novel approaches such
51 as the thermoluminescence dating of biogenic calcite (e.g. Duller et al., 2009). Key to the success of
52 techniques aiming to extend the age range is not only a high signal saturation point, but a sufficiently long
53 lifetime of the signal being used for dating. In addition, fundamental to optically stimulated luminescence
54 (OSL) thermochronometry techniques is the accurate derivation of thermal kinetic parameters such as
55 lifetime, including trap depth and frequency factor, because they are used to quantitatively constrain rock
56 cooling histories (e.g. King et al., 2016).

57
58 There are a number of approaches which can be used to empirically derive the lifetime of a luminescence
59 signal, including pulse annealing (e.g. Li and Tso, 1997) and isothermal decay (e.g. McKeever, 1985). Li and
60 Chen (2001) apply both methods to a quartz sample from Hong Kong, and find that both approaches
61 provide comparable results. In a recent study, Lowick and Valla (2018) obtained lifetime estimates from
62 isothermal decay measurements of Swiss glacial quartz samples and Risø calibration quartz, and reported
63 intra-sample variability of up to two orders of magnitude. They were unable to conclude whether this
64 variability was related to their measurement protocol or whether intra-sample variability is intrinsic. Similar
65 observations have been made by Buechi et al. (2017) and Trauerstein et al. (2017). Schmidt et al. (2018)
66 also investigated the reproducibility of lifetime measurements of the 110°C TL peak of a single quartz
67 sample using three different measurement techniques across eight different laboratories. They found that
68 variability between laboratories was greater than within-laboratory variability, but that results were least
69 dispersed when isothermal decay measurements were used. Averaged across laboratories, all three
70 methods produce trap depth results consistent within 1σ standard deviation (Schmidt et al., 2018).

71
72 Given the importance of establishing signal lifetime when dating older samples and/or when using OSL
73 thermochronometry techniques, this study aims to investigate the variability resulting from measurement
74 condition selections within isothermal decay protocols to better inform estimates of the associated
75 uncertainty. These parameters include the stimulation wavelength used for OSL measurement, and the
76 temperature and duration of elevated temperature holds used in the measurement protocol.

77

78 **2. Samples, Instrumentation, and Methodology**

79

80 *2.1 Samples*

81 A Malawian lake shoreline sedimentary quartz sample (MAL05-01-01) was selected for analyses, primarily
82 on the grounds of luminescence sensitivity and favourable luminescence properties (e.g. dose recovery
83 within $\pm 10\%$ uncertainties, limited IR response at room temperature, ability to recycle a laboratory dose).
84 The sample was prepared at Royal Holloway, University of London, according to routine sample preparation
85 procedures detailed by Thomas et al. (2009). In this study, medium aliquots (5mm in diameter) of sample
86 (grain size range 180-250 μm) were loaded onto 9.7 mm diameter aluminium discs and used for
87 measurement.

88

89 *2.2 Instrumentation*

90 Luminescence measurements were made at the Oxford Luminescence Dating laboratory, University of
91 Oxford using automated Risø TL/OSL DA-15 readers. The readers were fitted with $^{90}\text{Sr}/^{90}\text{Y}$ beta sources
92 providing dose rates of ~ 3.3 and ~ 4.0 Gy/min. Luminescence was stimulated using blue (wavelength 470
93 nm, stimulation power ~ 30 mW cm^{-2} at 90%), green, (525 nm, ~ 70 mW cm^{-2} at 90%) and IR (870 nm, ~ 90
94 mW cm^{-2} at 90%) LEDs, and ultraviolet luminescence signals were detected through a bialkali photo

95 multiplier tube fitted with 7.5 mm U340 filters. Calibration of the reader hotplates was undertaken in 2015
 96 and 2017, and measurements of the effective stimulation power of the LEDs were made in 2014 using an
 97 energy meter.

98
 99 *2.3 Signal lifetime determination*

100 Prior to measurement, aliquots were sensitised using at least ten cycles of beta irradiation and
 101 thermoluminescence measurement up to 450°C (heating rate 5°C/s) to minimise the impact of sensitivity
 102 change during measurement. An isothermal decay protocol following that of Wang et al. (2006) was used
 103 (table 1), using temperature and hold time combinations based on those of Bailey (1998a). Isothermal
 104 decay measurements were made using visible light (green and blue; table 1a) and infra-red light (table 1b).
 105 Stimulating quartz at an elevated temperature allows the signal from the fast component to be isolated and
 106 directly measured (e.g. Bailey, 1998b, Singarayer and Bailey, 2004, Jain et al., 2005, Fan et al., 2009). The
 107 same aliquots of sample were used for all measurements. To check the reproducibility of luminescence
 108 measurements, the blue/green-light protocol (table 1a) was repeated and when signals were compared,
 109 L_x/T_x values were found to be within uncertainties.

110
 111 Table 1: The isothermal decay measurement protocols used in this study. Where blue/green light stimulation was
 112 used, the protocol followed that of Wang et al., 2006 (a). Where elevated temperature IR stimulation was used, a
 113 modified protocol (following Bailey, 2010) was applied.

Isothermal Decay Protocol		
Step	a) Blue and green LEDs	b) IR LEDs
1	Beta dose (64 Gy)	Beta dose (128 Gy)
2	Pre-heat at 220°C for 10 s	Pre-heat at 220°C for 10 s
3	Pre-heat at T°C for t s (table 2)	Pre-heat at T°C for t s (table 2)
4		IR bleach at 20°C for 20 s
5		IRSL at 160°C for 200 s (L_x)
6	CW-OSL at 125°C for 100 s (L_x)	CW-OSL at 125°C for 100 s
7	Beta dose (32 Gy)	Beta dose (128 Gy)
8	Pre-heat at 200°C for 10 s	Pre-heat at 200°C for 10 s
9		IR bleach at 20°C for 20 s
10		IRSL at 160°C for 200 s (T_x)
11	CW-OSL at 125°C for 100 s (T_x)	CW-OSL at 125°C for 100 s

114
 115 Table 2: Duration (t, in seconds) of the raised temperature hold times at each temperature (T) used in step 3 of the
 116 isothermal decay measurement protocol (table 1). Hold durations are based on those applied by Bailey (1998a) who
 117 used t_1 , and t_{5-8} .

Duration (t, in seconds) of the raised temperature holds.								
T (°C)	t_1	t_2	t_3	t_4	t_5	t_6	t_7	t_8
240	0	10	100	400	1000	2000	5000	8000
245	0	7	72	290	725	1450	3450	6000
250	0	5	50	200	500	1000	2500	5000
255	0	3	32	130	325	750	1650	3000
260	0	3	30	120	300	600	1000	1500
265	0	2	21	85	212	425	775	1100
270	0	2	18	70	175	350	600	800
275	0	1	12	50	125	250	400	600
280	0	1	10	40	100	200	300	450
285	0	1	6	25	62	125	200	375
290	0	1	5	18	45	90	125	250

295	0	1	4	14	35	70	115	190
300	0	1	3	10	25	50	100	150

118
119
120
121

Assuming first order kinetics, the duration that an electron can be expected to remain trapped can be mathematically described by equation 1 (e.g. Aitken, 1985; McKeever, 1985. See also Schmidt et al., 2018).

$$\tau_T = s^{-1} \cdot e^{\frac{E}{k_B T}} \quad \text{Equation 1}$$

122
123
124 In equation 1, τ_T describes the lifetime (in s) of the trap. s is the frequency factor (s^{-1}), which Aitken (1985)
125 describes as the number of escape attempts of an electron from its trap per second. E is trap depth (in eV),
126 or the energy required to release the electron from the trap. k_B is Boltzmann's constant (equal to 8.62×10^{-5}
127 eV K^{-1}) and T is temperature (in K). In this study, values of E and s have been derived empirically. Individual
128 aliquots were stored at a range of temperatures for varying lengths of time (as described in table 2), and
129 measurement of the OSL signal remaining following each storage allowed estimates of the lifetime of the
130 OSL signal at each storage temperature to be made. E and s were derived from the Arrhenius plot:

$$\ln(\tau_T) = \ln(s^{-1}) + (E/k_B) \cdot (1/T) \quad \text{Equation 2}$$

131
132
133 Using a plot of $\ln \tau_T$ versus $1/T$, E is obtained from the gradient of the linear fit to this plot, and s from the y -
134 axis intercept (Bailey, 1998a; e.g. figure 1).
135
136

137 Published isothermal decay measurement protocols have typically been less extensive than those tested in
138 this study (table 1) and are predominately made using blue light stimulation of ultra-violet OSL signals. For
139 example, Bailey (1998a) derives empirical estimates of τ_T at 240, 250, 260, 270, 280, and 300°C, using
140 elevated temperature holds of t_1, t_6, t_7, t_8 in table 2. Lowick and Valla (2018) use τ_T s of 220 240, 260 and
141 280°C and holds similar in magnitude to t_5, t_6, t_7, t_8 in table 2. The extended combination of temperature
142 and hold times applied in this studied, measured using blue, green and infra-red light stimulation allows an
143 exploration of the impact of varying the isothermal decay protocol on the calculated kinetic parameters. To
144 investigate this variability, two key elements of the protocol are varied:

- 145
- 146 i) Stimulation wavelength. Stimulating quartz with infra-red light at an elevated temperature allows the
147 signal from the fast component to be isolated (e.g. Bailey, 1998b; Singarayer and Bailey, 2004; Jain et
148 al., 2005). Integrating the elevated temperature IR protocol of Bailey (2010) into the isothermal decay
149 protocol used in this study (e.g. table 1b) allows the lifetime from the fast-only signal (i.e. the 325°C
150 thermoluminescence peak) to be calculated. In contrast, continuous-wave OSL signals stimulated with
151 blue and green light will contain contributions from various OSL components in the initial part of the
152 signal. However, these contributions will differ because of the wavelength dependence of photo-
153 ionisation cross-sections with stimulation wavelength (Singarayer and Bailey, 2004). Therefore,
154 comparison of lifetimes measured from the same aliquot using different stimulation wavelengths will
155 allow the impact of contributions from other, non-fast OSL component contributions in lifetime
156 calculations to be explored.
 - 157 ii) Elevated temperature and hold time combinations. The effect of varying these parameters in the
158 measurement protocol will be explored changing the elevated temperature hold (step 3 in table 1), by
159 i) varying the combinations of hold times for τ_T derivation and ii) using different τ_T intervals for E and s
160 derivation (e.g. temperatures every 5°C, 10°C etc. over the temperature range 240-300°C).
161

162 *2.4 Component fitting of the quartz OSL signals*

163 Curve deconvolution of visible light stimulated composite-wave OSL signals was undertaken to better
164 understand the contribution of signal from the fast and medium quartz components to the bulk
165 luminescence signal. Assuming first order kinetics, the luminescence signal from each component can be
166 described mathematically (Bulur, 2000) by:

167
168
$$L_{(t)} = n_0 b \exp(-bt)$$
 Equation 3

169
170 where, $L_{(t)}$ is the luminescence intensity as a function of stimulation time t (s), n_0 is the initial trapped
171 charge population, b is the de-trapping probability, where $b = \sigma I_0$, σ is the photo ionisation cross-section
172 (cm^2) and I_0 is the light intensity (in photons $\text{s}^{-1}.\text{cm}^{-2}$). The sum of four exponentials was fitted to each
173 blue/green OSL signal using SigmaPlot (following the method outlined in Durcan and Duller, 2011), and the
174 goodness of fit was assessed through the calculation of the residual between measured and modelled
175 signals and the R^2 value. Calculation of the % contribution of each component to the total OSL signal was
176 made using the derived trapped charge populations (e.g. Durcan and Duller, 2011).

177
178 **3. Results and Discussion**

179
180 *3.1 The impact of varying quartz OSL component contributions on lifetime determination*

181 The isothermal decay protocols in table 1 were applied to the same aliquots of sample using different
182 wavelengths of stimulation light to determine whether contributions from different quartz OSL
183 components impact lifetime determinations. Elevated temperature IR stimulation was used to isolate signal
184 from the fast component, with the aim of calculating a lifetime from a signal originating from the fast
185 component trap only. Due to the wavelength dependency of quartz photo-ionisation cross sections, green
186 and blue light was used to investigate the effects of differing contributions of components in the initial part
187 of the OSL signals on lifetime determination. Table 3 summarises E , s and τ_{20} (lifetime at ambient
188 temperature) calculated from the same aliquots using blue, green and IR LEDs. There is a relationship
189 between E and wavelength, with longer wavelengths providing higher trap depths and longer lifetimes. Of
190 particular interest is i) the relatively high values of E , s , and lifetime using the elevated temperature IR
191 protocol and ii) the discrepancy between the green and blue light stimulated E and lifetimes values. The IR-
192 E values of 1.88 ± 0.16 and 1.99 ± 0.18 are higher than published estimates of E for the 325°C quartz TL
193 peak, which are typically in the region of 1.65 eV (e.g. Wintle and Murray, 1998; Murray and Wintle 1999;
194 Preusser et al., 2009; Lowick and Valla, 2018). These are associated with excessively high s values, hence
195 high lifetimes. Investigation of the relationship between elevated temperature IR luminescence and hold
196 time (figure 2) shows that this relationship is not well described by an exponential function. The measured
197 data after longer hold times (e.g. t_4 - t_8 in table 2) in particular, deviate from the expected exponential fit.
198 This may be a function of counting statistics and/or the effects of thermal transfer, however this
199 phenomena requires further investigation, currently beyond the scope of this paper. The Arrhenius plot
200 constructed from elevated-temperature IR luminescence signals (figure 2) is associated with large
201 uncertainties (typically 20-50%) and the relationship between $\ln(\tau)$ and temperature is not well
202 approximated with a linear fit. Therefore, the E , s and τ_T values presented in table 3 based on the elevated
203 temperature IR signals should be considered unreliable at present.

204
205 Table 3: Calculated E , s and τ_{20} for the same two aliquots of MAL05-01-01 (discs 1 and 2) measured using different LED
206 wavelengths (λ), where 870 nm = IR, 525 nm = green and 470 nm = blue LEDs.

Disc	λ (nm)	E (eV)	s (s)	τ_{20} (a)
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1	870	1.88 ± 0.16	9.65*10 ¹⁵	7.68*10 ⁸
	525	1.67 ± 0.08	8.51*10 ¹³	1.87*10 ⁷
	470	1.42 ± 0.13	2.62*10 ¹²	1.54*10 ⁶
2	870	1.99 ± 0.18	1.38*10 ¹⁷	3.72*10 ⁹
	525	1.52 ± 0.08	6.37*10 ¹¹	1.42*10 ⁵
	470	1.41 ± 0.13	4.25*10 ¹¹	1.17*10 ⁵

207 To test whether the discrepancies in E, s and τ_{20} by green/blue wavelength (table 3) are influenced by
208 differing contributions of signal from the different OSL components, $E_{(t)}$, $s_{(t)}$ and $\tau_{20(t)}$ plots are presented in
209 figure 3, alongside the % contribution of the fast, medium and slow components to the bulk signal at 0.1 s
210 integrals. These data demonstrate that for both blue and green light stimulated signals, E, s and τ_{20} vary
211 with the time integral used for calculation; falling as time is extended. Estimates for green-E values are
212 consistent between 0.1 and 1 s (figure 3b), although there is a tendency for s to decrease during this
213 period. In contrast, blue-E is lower than green-E, and also decreases with time (figure 3). s remains
214 relatively stable within the first 0.2 s. With blue-light stimulation, the fast component contributes less than
215 50% of the signal in the first 0.1 s of stimulation, and it is suggested that the decrease in E and hence
216 lifetime is due to the incorporation of signal from different traps (presumably the medium component, but
217 perhaps the slow1 component identified by Jain et al. (2003) and Durcan and Duller (2011)). It is proposed
218 that for this sample at least, signal contributions from non-fast component traps ‘muddy’ the calculated
219 lifetime, particularly in the presence of a large contribution from the medium component (as is the case
220 with this sample, e.g. figure 3a). This is observed particularly in the blue stimulated signals due to the
221 wavelength dependencies of quartz photo-ionisation cross-sections (Singarayer and Bailey, 2004), where
222 cross-section values for the fast and medium components converge with lower stimulation wavelengths.
223 For samples where the fast component is dominant, lifetimes calculated from the initial part of the signal
224 perhaps can be assumed to originate mostly from a single trap (e.g. the 325°C peak), which has been
225 concluded in studies by Rhodes (1990) and Wintle and Murray (1998). However, for samples which have a
226 significant contribution(s) from different component(s) in the initial part of the OSL signal, calculated
227 lifetimes will reflect multiple-trap contributions of charge, as seen in these data (figure 3). If this is the case,
228 the use of stimulation light with longer wavelengths is an option for better isolating signal from the fast
229 component and the relatively lower de-trapping rates of the other components (in comparison to the fast)
230 with longer stimulation wavelengths (c.f. Bailey et al., 2011).

232 3.2 Apparent lifetime variability induced during the isothermal decay protocol

233 As discussed, in the isothermal decay protocol τ_T is derived from the luminescence signal measured after a
234 series of elevated temperature holds. To investigate the variability in measured E, s and τ_{20} that may be
235 introduced by variations in the isothermal decay protocol, a jackknife approach was taken using the
236 measured datasets. E, s and τ_{20} were re-calculated using i) different combinations of raised temperature
237 hold times for τ_T calculation and ii) different temperature intervals for construction of the Arrhenius plot
238 from which τ_T was calculated. This approach was applied to luminescence signals measured with blue and
239 green light for both aliquots of sample, and an example of the induced variability for disc 1, measured with
240 green light is shown in figure 2. This approach was not applied to elevated temperature IR signals, given the
241 difficulties in deriving a lifetime, discussed in section 3.1.

242
243 The different symbols in figure 4 represent E, s and τ_{20} values derived from the same luminescence signals,
244 but where τ_T has been calculated from different hold time combinations and temperature intervals.
245 Individual E, s and τ_{20} values are normalised to values calculated where temperature interval is 5°C and
246 holds t_{1-8} are used. For the most part, E is consistently calculated to within ±5%, although E values are most

247 consistent when a shorter hold time (t_2 in table 2) is included immediately after the first elevated
248 temperature hold (for $0s$; t_1). Of the hold time combinations tested, the most consistent combinations for E
249 calculation are t_{1-8} , t_{1-4} and t_{1-3} (figure 4). This suggests that the shorter hold times are relatively more
250 important than the longer duration holds, presumably because they better constrain the exponential fit of
251 the luminescence data for τ_T derivation in the area of the curve where the gradient is steepest (in absolute
252 value). If the data could be fitted precisely (e.g. $r^2=1$) with an exponential function, the combination of hold
253 times should not matter. However, uncertainties in the sensitised luminescence signals (L_x/T_x) lead to
254 uncertainties in the mathematical fit, and focusing efforts on measuring luminescence signals after shorter,
255 rather than longer, hold times should allow a better fitting of the relationship between luminescence signal
256 and hold time, due to the lower uncertainties associated with counting statistics. For clarity, uncertainties
257 associated with E are not plotted in figure 4, however are in the region of $\sim 5\%$. Whilst the majority of E
258 estimations are within uncertainties of each other, there is a tendency for lower E values as temperature
259 interval is increased. The reason for the tendency for E to increase with greater number of empirically
260 derived τ_T values is unknown. Presumably, more data leads to better fitting of the linear function in the
261 Arrhenius plot, however that the gradient of this fit should consistently vary in one direction requires
262 further exploration. That said, these data demonstrate that for the majority of elevated temperature hold
263 combinations, E can be consistently determined (figure 4a). Greater confidence can be placed in
264 determinations where emphasis is placed on shorter elevated temperature hold times for τ_T derivation and
265 higher temperature intervals (ideally 5°C or 10°C).

266
267 When hold time combinations and temperature intervals are varied, greater variability is observed for
268 estimates of s and therefore τ_T (figure 4). Again, the combination of focus on shorter elevated temperature
269 hold times and narrower temperature intervals results in greater consistency between estimates. That s
270 estimates display more variability is not surprising. s is derived through extrapolation from the empirical
271 data to the ordinate on the Arrhenius plot. Therefore even small uncertainties in empirically derived τ_T
272 results in the potential for greater variability approaching the ordinate. To minimise variability in s and τ_{20} ,
273 it is important that τ_T is calculated as precisely as possible. E and s are used to calculate τ_{20} , hence the
274 trends in E and s variability are also seen in figure 4c. Focusing on elevated temperature hold time
275 combinations which focus on the shorter times (e.g. t_{1-8} , t_{1-4} and t_{1-3} ; table 2), lifetimes calculated from τ_{20}
276 determinations from intervals of 20°C and 30°C appear to underestimate τ_{20} by up to 50% in comparison to
277 the 5°C interval. For this sample at least, calculated signal lifetime estimates based on Arrhenius plots
278 constructed using wider temperature intervals ($>20^\circ\text{C}$) should be treated as minima.

279 280 3.3 τ_T uncertainty and the impact on lifetime calculation

281 To investigate the impact of uncertainty of τ_{20} estimates when temperature interval and τ_T uncertainty are
282 varied, values for E and s were fixed and Arrhenius plots were simulated, producing 5,000 plots where the
283 relationship between $\ln(\tau_T)$ and temperature was fit perfectly with a linear function. Uncertainties of ± 1 , ± 5
284 and $\pm 10\%$ were applied to individual τ_T values, and τ_T temperature interval was varied between 2°C and
285 30°C (table 4). The trends observed in table 4 confirm those observed in the empirical data (figure 4). Trap
286 depth is relatively insensitive to simulated temperature interval, and can be calculated consistently to
287 within $\pm 6\%$ of the known value in the presence of τ_T uncertainties of up to $\pm 10\%$ (table 4). Where individual
288 uncertainties on τ_T do not exceed $\pm 5\%$, in the majority of cases, E can be calculated to within $\pm 3\%$ of the
289 fixed value. Simulated estimates of s , and therefore τ_{20} , are more variable, as is the case in the empirical
290 evidence, with simulations producing positively skewed distributions for both parameters (e.g. table 4;
291 figure 5). Relatively large, asymmetric uncertainties in s should be expected given the long extrapolation
292 required in the Arrhenius plot to derive s . For example, simulations with τ_T uncertainties of $\pm 5\%$ and a

293 temperature interval of 20°C results in deviations from the fixed value of s and the calculated τ_{20} of ~-40%
 294 at the 5th percentile and >240% at the 95th percentile. To achieve an estimate of τ_{20} to within 10% of the
 295 target τ_{20} value (from fixed E and s), a temperature interval of 2°C and τ_T uncertainties of $\pm 1\%$ are required
 296 (table 4), which are both unlikely and impracticable in terms of measurement.

298 Table 4: Summary of data from simulated Arrhenius plots ($n = 5000$). Estimates of E , s and τ_{20} are normalised to the
 299 fixed values of $E = 1.65$ eV, $s = 1.25 \cdot 10^{12} \text{ s}^{-1}$ and therefore $\tau_{20} = 7.64 \cdot 10^8 \text{ a}$, and the 5th, 50th and 95th percentiles are
 300 reported. Fixed E and s values were chosen to be representative of frequently reported averages (c.f. Preusser et al.,
 301 2009). Interval refers to the temperature interval used over the range of 240–300 °C, and uncertainty is the
 302 uncertainty attributed to individual τ_T s. The ratio of estimated (est) to fixed values are shown.

	Interval	2°C			5°C			10°C			20°C			30°C		
		±1%	±5%	±10%	±1%	±5%	±10%	±1%	±5%	±10%	±1%	±5%	±10%	±1%	±5%	±10%
$E_{\text{est}}/E_{\text{fixed}}$	Uncertainty															
	Percentile															
	5 th	1.00	0.99	0.97	1.00	0.98	0.96	1.00	0.98	0.95	0.99	0.97	0.94	0.99	0.97	0.94
$S_{\text{est}}/S_{\text{fixed}}$	5 th	0.92	0.64	0.41	0.87	0.52	0.27	0.84	0.43	0.17	0.82	0.37	0.13	0.81	0.35	0.12
	50 th	1.00	1.01	1.01	1.00	1.00	1.00	1.00	1.01	1.01	1.00	1.01	1.00	1.00	1.00	1.02
	95 th	1.09	1.58	2.54	1.14	1.95	3.95	1.19	2.34	5.22	1.22	2.78	7.34	1.23	2.96	8.69
$\tau_{20_est}/\tau_{20_fixed}$	5 th	0.95	0.70	0.48	0.91	0.59	0.33	0.88	0.50	0.23	0.86	0.44	0.18	0.85	0.42	0.16
	50 th	1.02	1.03	1.03	1.02	1.02	1.01	1.02	1.03	1.03	1.02	1.03	1.01	1.02	1.03	1.05
	95 th	1.10	1.50	2.24	1.15	1.80	3.26	1.18	2.11	4.12	1.22	2.43	5.46	1.22	2.56	6.19

303
 304 **4. Conclusion and Recommendations**
 305

306 This paper has sought to explore potential sources of variability which may arise through the use of the
 307 isothermal decay protocol and Arrhenius plots for lifetime determination of quartz OSL signals typically
 308 used in dating applications. Focus has been upon variations in the combination of elevated temperature
 309 hold times used for τ_T measurement, the number of τ_T points used to construct the Arrhenius plot, and the
 310 OSL stimulation wavelength. Using both empirical and modelled datasets, this paper has demonstrated that
 311 E can be calculated with relative consistency, in contrast to s , which is subject to more variability. This
 312 variability arises from the long extrapolation distance on the Arrhenius plot. Optimising combinations of
 313 hold times and temperature intervals used within the isothermal decay protocol can reduce variability, and
 314 this has been demonstrated with empirical and simulated data. The measurement of luminescence signals
 315 using various different wavelengths has shown that contributions from at least one non-fast OSL
 316 component can impact lifetime calculation. Stimulating luminescence with light of longer wavelengths (e.g.
 317 green light), and utilising the different wavelength dependencies of the fast and non-fast OSL components
 318 (Singarayer and Bailey, 2004), the signal from the fast OSL component can be better isolated from slower
 319 components. Reducing the contribution of non-fast component signal will allow more reliable estimates of
 320 lifetime.

322 On the basis of the presented datasets, the following recommendations are proposed for quartz OSL signal
 323 lifetime determination using isothermal decay:

- 324
 325 i) Before estimates of the kinetic parameters of an OSL signal are made, the composition of the OSL
 326 under investigation should be assessed. Contributions of signal from multiple traps and/or poor
 327 mathematical fits to empirical data may render lifetime estimates inaccurate.

- 328 ii) In the presence of high signal contributions from non-fast OSL components, lifetimes may be affected
329 by signal from more than one trap. These signal contributions can be minimised by stimulating
330 luminescence with light of longer wavelengths to better isolate signal from the fast component from
331 other quartz OSL signals.
- 332 iii) Lifetime calculations should be made for multiple aliquots of sample. Due to the nature of the
333 mathematical fitting undertaken in the calculation of the lifetime, variability particularly in s should be
334 expected. E can be calculated more reproducibly and can therefore be used to inform assessments
335 of intra-sample variability.
- 336 iv) Improved reproducibility of lifetime estimates can be achieved by altering the isothermal decay
337 protocol. Elevated temperature hold times should incorporate shorter times, where the gradient of
338 the exponential fit is highest in absolute value. Temperature intervals of less than 10°C also provided
339 improvements in reproducibility in this study.
- 340 v) Relatively large, asymmetric uncertainties are associated with estimates of s . Luminescence
341 practitioners should bear this in mind, and report E , s and τ values with uncertainties.
- 342 vi) Simple modelling from fixed values of E and s can indicate the distribution of expected E and s values
343 based on the protocol adopted for empirical data collection, and can be used to inform the precision
344 of E , s and lifetime determinations.

345

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347

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353

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419

420 **Figure Captions**

421

422 **Figure 1:** Left, sensitivity corrected green stimulated luminescence recorded after various storage times at
 423 temperatures 240, 250, 260, 270, 280, 290, and 300°C (for disc 1), and the Arrhenius plot constructed from
 424 these data (middle). Right, the relationship between hold time and luminescence at a hold time of 240°C,
 425 when disc 1 was stimulated with blue, green and IR light. Inset, the Arrhenius plot constructed from the
 426 luminescence data stimulated with different wavelengths.

427

428 **Figure 2:** Sensitivity corrected luminescence recorded after various storage times at temperatures 240°C,
 429 270°C and 300°C, used to calculate τT (disc 1). Luminescence was stimulated using infra-red light at 160°C
 430 (e.g. table 1b). The exponential fits to the data are shown. Note, the first elevated temperature holds were
 431 for 0 s, but are plotted at 0.01 s for display purposes. Inset, the Arrhenius plot constructed from elevated
 432 temperature IR stimulated luminescence signals (disc 1).

433

434 **Figure 3:** E , s and τ_{20} by time integrals of 0.1 s using blue-stimulated (a) and green stimulated (b)
 435 luminescence signals (disc 1). Also shown, the % contribution of fast, medium and slow component signal
 436 per time integral, obtained by curve fitting of continuous-wave OSL signals.

437

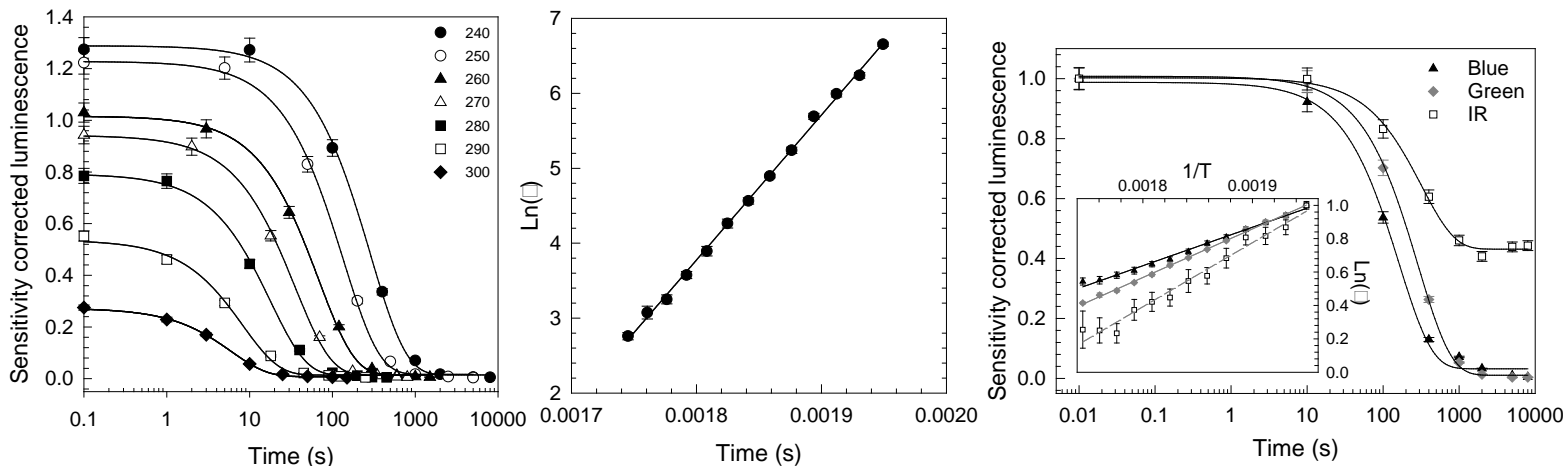
438 **Figure 4:** Variability in E , s , and τ_{20} as elevated temperature hold time and temperature interval are varied
 439 (green stimulated luminescence signals from disc 1). In all iterations, the temperature range is 240-300°C,
 440 where an interval of 5°C means τT was calculated at 240, 245, 250 ... 300°C etc.. Different hold time
 441 combinations are indicated by the different symbols. Values are normalised to the value calculated from a
 442 temperature interval of 5°C and hold time combinations t1-8.

443

444 **Figure 5:** Distributions of E , s and τ_{20} calculated from 5,000 simulated Arrhenius plots, using fixed values of
 445 E and s ($E = 1.65$ eV, $s = 1.25 \cdot 10^{12}$ s, and therefore $\tau_{20} = 7.64 \cdot 10^8$ a, shown by the vertical reference lines.
 446 An uncertainty of $\pm 5\%$ is ascribed to τT values, and the temperature interval used and 20°C.

447

448 **Figure 1**



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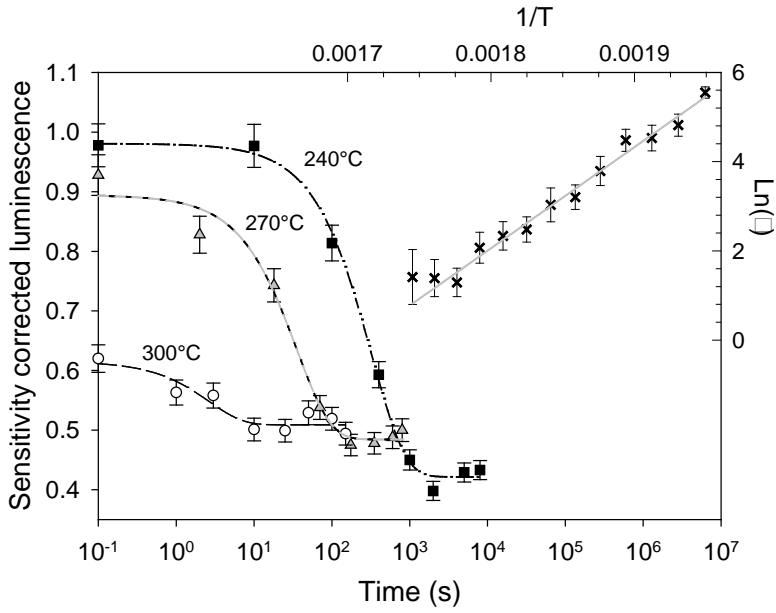
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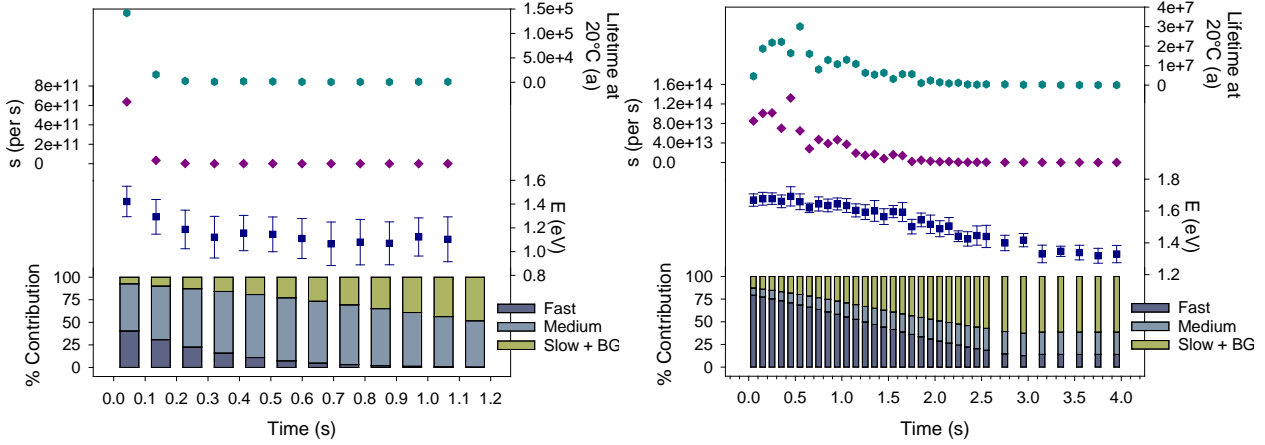
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458 **Figure 2**



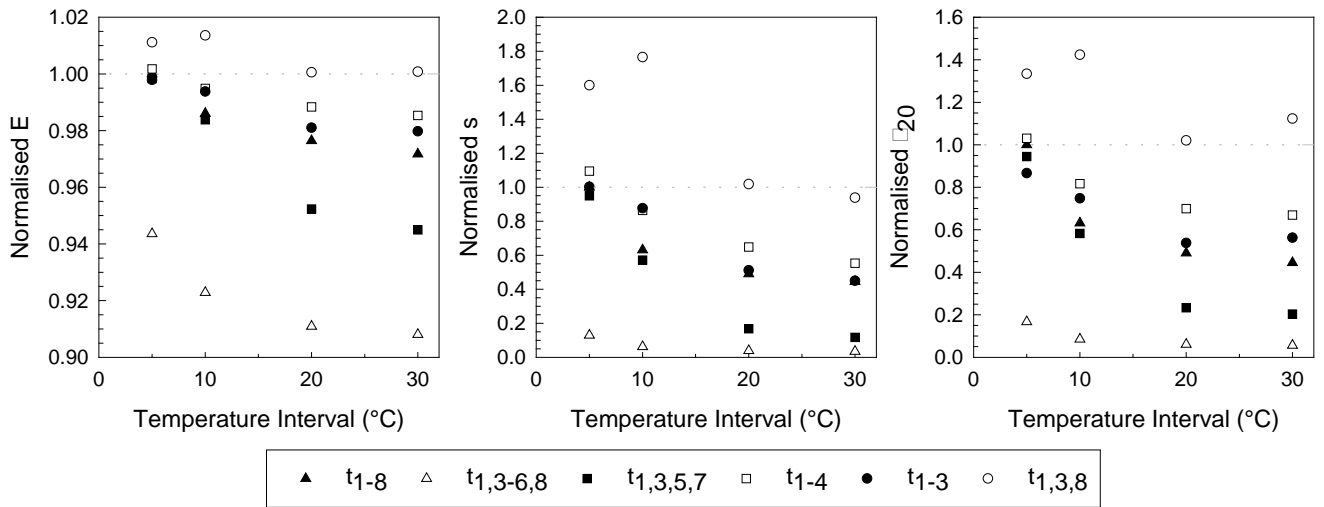
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Figure 3



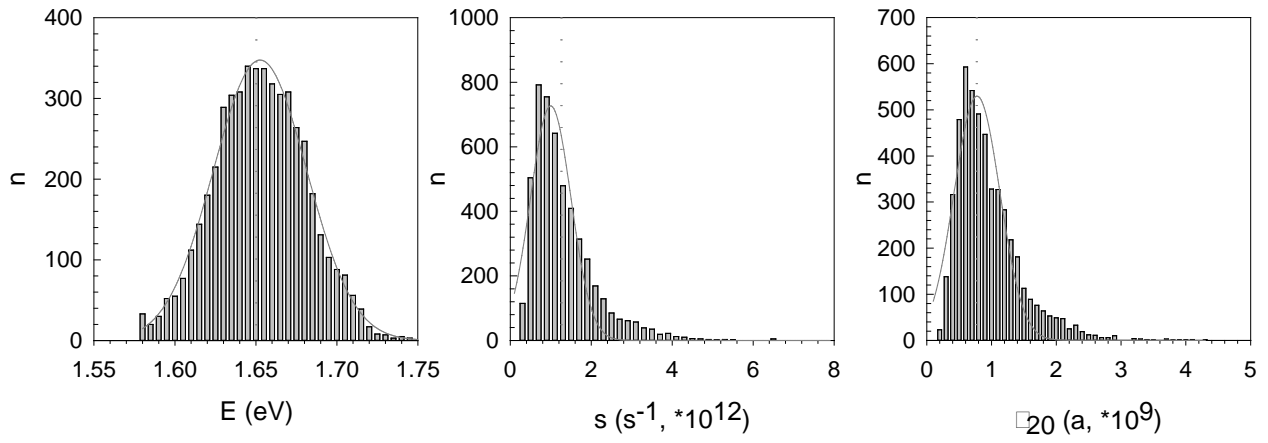
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Figure 4



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471 **Figure 5**



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