

Supporting Information for: 'SLAMS-2.0: a Stochastic, Lagrangian Aggregate Model of Sinking Particles for the Ocean's Biological Carbon Pump, version 2'

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Texts S1–S3 detail the datasets used in this study, including prescribed fields, related input variables, and observational constraints to assess model performance. **Texts S4–S5** detail the governing equations for particle state and mesozooplankton characteristics.

Figures S1–S6 illustrate the prescribed datasets and observational constraints, including a validation of prescribed NPP. **Figure S7** compares particle size distributions and mass concentrations produced by the sectional method and the SLAMS-2.0 coagulation scheme, while **Figures S8–S9** show the sensitivity of export POC flux and particle number concentration to one-at-a-time parameter perturbations.

Table S1 provides quantitative validation of prescribed NPP, and **Table S2** lists depth ranges for the key horizons used in model-data comparisons.

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Text S1. Observational Datasets for Model Forcing

For the simulations presented here, the following datasets are used (**Figure S1**).

- Surface ocean photosynthetically active radiation (PAR_0 , W m^{-2}) is derived from monthly climatologies of Aqua-MODIS Level 3 satellite data, available via NASA's Ocean Color portal (<https://oceancolor.gsfc.nasa.gov/13/order/>).

- Aeolian dust deposition flux ($\text{g clay m}^{-2} \text{s}^{-1}$) was sourced from the historical run (1850–2014) of the NCAR-CESM2 model, part of the Coupled Model Intercomparison Project Phase 6 (CMIP6). The dataset has monthly and 100 km spatial resolution and was obtained via the Earth System Grid Federation (<https://esgf-node.llnl.gov/projects/cmip6/>). Dust fluxes were converted to clay fluxes using a 0.5:1 clay-to-dust mass ratio, following Journet, Desboeufs, Caquineau, and Colin (2008, Table 1).

- Net primary production integrated over the euphotic zone (NPP, $\text{mg C m}^{-2} \text{d}^{-1}$) is obtained from the ESA BICEP v4.2 multi-sensor monthly climatology (Kulk et al., 2021). NPP is distributed over the daylight fraction of the modelled day and converted to $\text{mol C m}^{-2} \text{s}^{-1}$ for the time steps receiving PAR_0 . Among several satellite-based NPP products (**Figure S2**), BICEP was selected to force SLAMS-2.0 due to its closer agreement with available *in situ* ^{14}C uptake measurements at our study sites (**Figure S3** and **Table S1**).

- Chlorophyll *a* concentration integrated over the euphotic layer ($\text{Chl}a$, mg m^{-3}) is obtained from the multi-sensor monthly climatology of ESA's Ocean Colour Climate Change Initiative (OC-CCI) v6 and available at the OC-CCI website (<https://www.oceancolour.org>), after Sathyendranath et al. (2023).

- Mixed layer depth (MLD, m) is from the monthly climatology of De Boyer Montégut, Madec, Fischer, Lazar, and Iudicone (2004) that applies a variable density threshold equivalent to a 0.2°C temperature decrease. The dataset is available from the IFREMER/LOPS Mixed Layer Depth Climatology portal (<http://www.ifremer.fr/cerweb/deboyer/mld/>).

- Seawater temperature ($^{\circ}\text{C}$), salinity (PSU), dissolved O_2 concentration (mL L^{-1}), and macronutrient concentration (nitrate, phosphate, and silicate; mmol m^{-3}) are from World Ocean Atlas 2023 (WOA23) monthly climatologies (<https://www.ncei.noaa.gov/access/world-ocean-atlas-2023/>, Reagan et al., 2024). Since WOA23 monthly fields are depth-limited (e.g., temperature and O_2 to 1500 m; nutrients to 800 m), deeper layers were filled using corresponding annual climatologies to produce full-depth profiles.

- Seawater density (ρ_w , g cm^{-3}) and dynamic viscosity (μ , $\text{g cm}^{-1} \text{s}^{-1}$) were calculated from WOA23 temperature and salinity fields. Density was computed using the Gibbs-SeaWater (GSW) Oceanographic Toolbox (MATLAB, v3.06, <https://www.teos-10.org/>) based on the Thermodynamic Equations of Seawater–2010 (TEOS-10) McDougall and Barker (2011). Dynamic viscosity was derived using routines from the MIT seawater properties library (MATLAB, v3.1.5, <https://web.mit.edu/seawater/>) following the formulations of Sharqawy, Lienhard, and Zubair (2010) and Nayar, Sharqawy, Banchik, and Lienhard (2016).

- Calcite saturation state is calculated using CO2SYS (MATLAB, v.1.1, Lewis & Wallace, 1998; Van Heuven et al., 2011), with default dissociation constants and incorporating annual climatological data of alkalinity, CO_2 , temperature, salinity, silicate, and phosphate from the Global Ocean Data Analysis Project version 2 (GLODAPv2.2016b, <https://>

www.nodc.noaa.gov/archive/arc0107/0162565/2.2/data/0-data/mapped/, Olsen et al., 2016).

- Lastly, mesozooplankton biomass (mg C m^{-3}) is obtained from the CMIP6 Earth System Model IPSL-CM6A-LR, developed by the Institut Pierre Simon Laplace (IPSL), France. Data were downloaded from the Earth System Grid Federation (ESGF) Lawrence Livermore National Laboratory node (http://esgf-node.llnl.gov/search/cmip6/?mip_era=CMIP6&activity_id=CMIP&institution_id=IPSL&source_id=IPSL-CM6A-LR/). IPSL-CM6A-LR uses the biogeochemical ecosystem model PISCES-v2.0 (Aumont et al., 2015), which resolves the full water column. We use the monthly climatology from the historical IPSL-CM6A-LR simulation (1997–2014; r1i1p1f1). PISCES-v2.0 was selected because it shows strong agreement with observational mesozooplankton biomass when evaluated against MAREDAT using a Generalized Linear Mixed Model approach (Petrik et al., 2022; Moriarty & O’Brien, 2013). We note that CMIP6 models define mesozooplankton as organisms in the 20–200 μm size range (definitions here: <https://clipc-services.ceda.ac.uk/dreq/mipVars.html>), whereas the MAREDAT classification uses a 200–2000 μm definition (Sieburth et al., 1978).

Text S2. Derived Physical and Optical Fields

The prescribed forcing fields described in **Text S1** are used to calculate the physical and light-related quantities described below.

Water Physics

In SLAMS-2.0, water physics is parameterised through two foundational variables —dynamic viscosity and turbulent kinetic energy (TKE) dissipation rate— from which additional hydrodynamic quantities are derived.

Prescribed dynamic viscosity, μ ($\text{g cm}^{-1} \text{s}^{-1}$), is used to derive the kinematic viscosity, ν ($\text{cm}^2 \text{s}^{-1}$), as

$$\nu = \frac{\mu}{\rho_w}. \quad (1)$$

The TKE dissipation rate, ϵ ($\text{m}^2 \text{s}^{-3}$), characterises the rate at which energy from wind- and tide-driven stirring cascades to small scales and is dissipated through molecular viscosity, generating turbulent mixing. In SLAMS-2.0, ϵ is sampled via Monte Carlo from a log-normal distribution, as form commonly assumed in global ocean models (Pearson & Fox-Kemper, 2018):

$$f(\epsilon) = \frac{1}{\epsilon \cdot \sigma \cdot \sqrt{2\pi}} \cdot e^{-\left(\frac{(\ln \epsilon - \mu)^2}{2\sigma^2}\right)}, \quad (2)$$

with typical values of $\log_{10}(\epsilon) \in [-7, -4]$ $\text{m}^2 \text{s}^{-3}$ for shelf seas and coastal waters (depth < 200 m), and $[-9, -6]$ $\text{m}^2 \text{s}^{-3}$ for the open ocean (Saiz & Kiørboe, 1995). Here, μ is the logarithmic mean, $\mu = \ln[\sum(\epsilon)/n]$, and σ the variance of the distribution, set to 1.2. Although ϵ varies spatially and with depth (Pearson & Fox-Kemper, 2018), SLAMS-2.0 assumes a vertically uniform distribution since variations with depth are small.

From ν and ϵ , two additional hydrodynamic quantities are derived: the water shear rate, γ (s^{-1}), which quantifies velocity gradients caused by turbulent mixing:

$$\gamma = \left(\frac{\epsilon}{\nu}\right)^{1/2}, \quad (3)$$

and the Kolmogorov length scale, η (m), which represents the smallest turbulent eddy size below which fluid viscosity dominates and all turbulent kinetic energy is dissipated (as heat) instead of creating turbulent motion,

$$\eta = \left(\frac{\nu^3}{\epsilon}\right)^{1/4}. \quad (4)$$

Water Light-Related Fields

SLAMS-2.0 requires the calculation of three light-related fields: (i) photoperiod, (ii) vertical attenuation of PAR, and (iii) euphotic zone depth. Together, these fields control the timing and depth of phytoplankton growth and zooplankton vertical migration.

The photoperiod, or number of daylight hours at the ocean surface, $n_h(\phi, t)$, depends on latitude, ϕ (radians) and Julian day, t , following the solar geometry model of Brock (1981):

$$n_h(\phi, t) = \frac{2}{15} \cdot \arccos[-\tan(\phi) \cdot \tan(\delta(t))] \cdot \frac{360}{2\pi}, \quad (5)$$

where $\delta(t)$ is the solar declination angle (in radians) and is given by:

$$\delta(t) = 23.45 \cdot \sin\left(\frac{2\pi \cdot (284 + t)}{365}\right) \cdot \frac{2\pi}{360}. \quad (6)$$

Prescribed PAR_0 is attenuated with depth following the Beer–Lambert law:

$$\text{PAR}(z) = \text{PAR}_0 \cdot e^{-(k_{\text{PAR}} \cdot z)}, \quad (7)$$

where k_{PAR} (m^{-1}) is the PAR attenuation coefficient and is calculated following Morel et al. (2007) as a single value for the euphotic zone as a function of prescribed MLD and the

diffuse attenuation coefficient at 490 nm, $k_d(490)$:

$$k_{\text{PAR}} = 0.0864 \cdot 0.884 \cdot k_d(490) - 0.00137 \cdot (k_d(490))^{-1}, \quad \text{if } \text{MLD} \leq (k_d(490))^{-1}, \quad (8a)$$

$$k_{\text{PAR}} = 0.0665 \cdot 0.874 \cdot k_d(490) - 0.00121 \cdot (k_d(490))^{-1}, \quad \text{if } \text{MLD} > (k_d(490))^{-1}, \quad (8b)$$

where $k_d(490)$ is derived from prescribed $\text{Chl}a$ using the Case 1 water regression of Morel (1988):

$$k_d(490) = 0.0166 + 0.07242 \cdot \text{Chl}a^{0.68955}. \quad (9)$$

The euphotic zone depth z_{eu} is defined as the depth at which PAR decreases to 0.1% of its surface value (as recommended by Buesseler et al., 2020).

Text S3. Observational Datasets for Model Optimisation

We use observational data from six contrasting oceanographic sites that host, or have hosted, long-term time-series programs on particulate fluxes, following Rufas, Khatiwala, Bisson, Martin, and Bouman (2025): HOT/ALOHA (Hawaii Ocean Time-series station ALOHA); BATS/OFP (Bermuda Atlantic Time-Series/Oceanic Flux Program site); EqPac (U.S. Joint Global Ocean Flux Study Equatorial Pacific process study site); PAP-SO (Porcupine Abyssal Plain time-Series Observatory); OSP (Ocean Station Papa); and HAUSGARTEN (the Long-Term Ecological Research observatory HAUSGARTEN). Together, these sites span a broad range of productivity regimes across distinct biogeographical provinces: from oligotrophic subtropical gyres with persistently low productivity (HOT/ALOHA, BATS/OFP), to productive subpolar systems with strong seasonality and recurrent spring and autumn phytoplankton blooms (PAP-SO, OSP), to high-productivity equatorial upwelling zones with sustained year-round productivity (EqPac), to high-latitude polar environments where productivity is lower overall and peaks during the summer months (HAUSGARTEN).

For these six sites, we obtained three types of observational constraints: (i) *in situ* flux measurements of POC, PIC, and bSi derived from sediment traps and radionuclide disequilibria (mainly ^{234}Th , but also ^{210}Po , ^{228}Th , and ^{230}Th); (ii) *in situ* particle imagery from the Underwater Vision Profiler 5 (UVP5), providing particle number concentrations and particle size distributions; (iii) satellite-based estimates of POC export flux. Each observational approach has distinct advantages. Sediment traps, which physically intercept sinking particles, provide the most direct estimates of fluxes and are widely considered the gold standard for flux collection (e.g., Buesseler et al., 2007; Le Moigne et al., 2013).

Radionuclide disequilibria, by contrast, allow fluxes to be inferred from standard water sampling methods (e.g., Conductivity, Temperature, and Depth (CTD) deployments) during a single sampling event, enabling higher spatial resolution and more logistically efficient broad-scale surveys, as opposed to time-series measurements at single points. UVP5 imagery offers the most direct observations of particle size structure in the upper 2000 m of the water column. In contrast, satellite-based export estimates are indirect and surface-limited, and are used primarily as a consistency check in this study.

In Situ Particulate Flux Data

SLAMS-2.0 has been optimised using *in situ* observational data of POC flux from sediment traps and radionuclide disequilibria compiled by Rufas et al. (2025), along with newly compiled PIC and bSi flux data from the same set of 72 sources (**Figure S4**, and Table S8 in Rufas et al., 2025). Our six study locations were selected for their availability of concurrent measurements of POC, PIC, and bSi flux in a time-series fashion. The data span 1978–2022 and include long-term time-series, annual process studies, and short-term sampling campaigns.

Because flux collection is generally limited to a few discrete depths, we selected four depth horizons for model–data comparison that capture the main remineralisation zone of POC flux: the base of the euphotic zone (export depth), the upper and lower mesopelagic, and the base of the mesopelagic zone (sequestration depth) (**Table S2** for site-specific depths). This yields a potential of up to 864 observational constraints across tracers, depths, months, and sites.

From the calculated euphotic zone depth (z_{eu} ; **Text S2**), a $\pm 10\%$ depth margin was applied to account for uncertainty (after McKinna et al., 2019). Observed fluxes were

matched to depths within this range using a flexible window to accommodate measurement variability, according to three scenarios: (i) a single matching depth was expanded by an extra $\pm 10\%$ to reduce reliance on a single measurement/research project/collection method; (ii) if multiple depths matched, the window was adjusted to span from the shallowest to the deepest matching depth; (iii) if no depths fell within the range, the nearest depth shallower than 200 m was used, initially searching within $\pm 20\%$ of z_{eu} , then applying (i) or (ii) to define the final bounds. These adjusted ranges (**Table S2**) were also used to extract modelled fluxes.

Mesopelagic depth zones (upper, lower, base) were defined according to the distribution of flux measurements at each site, generally targeting 300, 600, and 1000 m, and defining windows of ~ 20 m around the nearest sediment trap depth. Notably, at HOT/ALOHA and BATS/OFP, the base of mesopelagic zone extends beyond the conventional 1000 m depth, reaching 1500 m, due to the absence of sediment traps between 500–1500 (BATS/OFP) and 800–1500 m (HOT/ALOHA). These mesopelagic depth zones were likewise applied to extract fluxes from the model outputs.

Data were quality-controlled prior to analysis. Fluxes were log-transformed to account for their known log-normal distribution (e.g., Cael et al., 2021) and then filtered using the interquartile range (IQR) method: the first and third quartiles were calculated, and any value falling outside a conservative threshold of $3 \times \text{IQR}$ was flagged as an outlier. Only one outlier was identified: a PIC flux measurement at OSP at the base of the euphotic zone, derived from radionuclide estimates from the EXPORTS campaign (in Roca-Martí et al., 2021), which was over an order of magnitude lower than the long-term sediment trap record at the same site from Wong, Waser, Whitney, Johnson, and Page

(2002). Additionally, we excluded monthly flux profiles with fewer than two observations. After filtering, the dataset comprised 565 climatological mean data points (**Figure S5**). Inspection of **Figure S5A** shows that peak POC fluxes increase from subtropical to equatorial and subpolar sites, before declining again at the polar site. Each site displays a distinct particulate flux regime, explained below.

In the subtropical gyres (HOT/ALOHA and BATS/OFP), export fluxes are modest year-round ($\sim 20\text{--}30 \text{ mg C m}^{-2} \text{ d}^{-1}$), with HOT/ALOHA more stable and BATS/OFP showing occasional peaks ($\sim 60 \text{ mg C m}^{-2} \text{ d}^{-1}$ in March) (**Figure S5A**). POC flux attenuates sharply within the upper mesopelagic, over an order of magnitude, consistent with warm, low-productivity waters that have high remineralisation rates (Martin et al., 1987; Conte, Maureen et al., 2001; Marsay et al., 2015). Corresponding PIC:POC flux ratios at the base of mesopelagic (**Figure S6A**) are relatively high at BATS/OFP (>1.0) but low at HOT/ALOHA (<0.5), while Si:POC flux ratios (**Figure S6B**) remain very low at HOT/ALOHA (~ 0.04) and moderately low at BATS/OFP (~ 0.30). This indicates carbonate-rich coccolithophore export with little diatom contribution, supported by local observations during the VERTIGO program (Lamborg et al., 2008) and the OFP time-series (Conte et al., 2019).

At the equatorial upwelling site (EqPac), episodic export pulses exceed $250 \text{ mg C m}^{-2} \text{ d}^{-1}$ (**Figure S5A**), interspersed with long periods of low or missing data, likely reflecting variability in upwelling and El Niño/Southern Oscillation cycles (Kawahata & Gupta, 2003). Although POC flux attenuates strongly with depth, sequestration fluxes remain relatively high (annual mean $\sim 5 \text{ mg C m}^{-2} \text{ d}^{-1}$). High sequestration depth PIC:POC (>1.0) (**Figure S6A**) and moderate-to-high Si:POC (~ 0.6) (**Figure S6B**) suggest mixed coc-

cololithophore–diatom export (Arteaga & Rousseaux, 2023), with carbonates relatively well-preserved at depth.

In the subpolar North Atlantic and Pacific, PAP-SO and OSP both exhibit strong bloom-driven seasonality, with spring export peaks exceeding $150 \text{ mg C m}^{-2} \text{ d}^{-1}$ (**Figure S5A**). Export at PAP-SO penetrates more effectively to depth, but annual sequestration flux is higher at OSP (~ 5 vs. $\sim 3 \text{ mg C m}^{-2} \text{ d}^{-1}$) (**Figure S5A**). Both sites are characterised by low mid-water PIC:POC ratios (~ 0.10 at PAP-SO; ~ 0.40 at OSP), reflecting weak calcification (**Figure S6A**). However, at sequestration depth Si:POC diverges: PAP-SO remains relatively low (~ 0.30), whereas OSP is high (~ 0.90), indicating stronger diatom control at OSP (**Figure S6B**). Both sites have a mixed community dominated by diatoms in the spring bloom that transitions to coccolithophores later in spring (Boyd & Harrison, 1999; Lampitt et al., 2023).

Finally, at the polar site HAUSGARTEN, export fluxes are low for most of the year ($< 15 \text{ mg C m}^{-2} \text{ d}^{-1}$), punctuated by short, intense summer pulses ($\sim 35 \text{ mg C m}^{-2} \text{ d}^{-1}$) that sink efficiently to depth (**Figure S5A**). PIC:POC ratios at the base of mesopelagic remain low (annual mean ~ 0.30) (**Figure S6A**) and Si:POC is quite modest (~ 0.23) (**Figure S6B**), pointing to a short productive season dominated by diatom-rich communities associated with sea-ice retreat and light availability (Hop et al., 2006; Nöthig et al., 2015).

In summary, the combined mineral-to-POC ratios support the idea that carbonate-rich export (coccolithophores) dominates in oligotrophic subtropical gyres (BATS/OFP, HOT/ALOHA), while biogenic silica-rich export (diatoms) is more important in the high-nutrient regions of high-latitude (PAP-SO, OSP, HAUSGARTEN) and upwelling (EqPac)

systems. Deep-water (sequestration depth) combined mineral-to-POC ratios further suggest that diatom-derived silica is relatively well-preserved in the Pacific Ocean sites (Eq-Pac, OSP) compared to subtropical gyres and the Atlantic Ocean sites, while carbonate preservation appears stronger at the Atlantic sites (BATS/OFP, PAP-SO) and EqPac but weaker in most high-latitude systems. These patterns are consistent with the ‘Silicate Ocean’ versus ‘Carbonate Ocean’ framework at the mesopelagic/bathypelagic boundary proposed by Honjo, Manganini, Krishfield, and Francois (2008).

In Situ Particle Imaging Data

UVP5 data were compiled by Kiko et al. (2022) and we extracted them from the EcoPart repository (<https://ecopart.obs-vlfr.fr>) following the procedure in Figure S4 in Rufas et al. (2025). UVP5-based observations of particle number concentration span the period 2008–2021 and are binned into 16 particle size classes, consistent with SLAMS-2.0. While valuable for characterising particle size structure, we note some caveats with these data: they represent instantaneous snapshots rather than time-integrated measurements, as sediment traps do; uncertainty estimates are not provided by the EcoTaxa platform, and may be substantial; and the UVP5 cannot detect particles smaller than 100 μm , which constitute the majority of particles.

Satellite-Based Estimates of POC Export Flux

Satellite-based estimates of POC export flux (i.e., at the base of the euphotic zone) were obtained from ESA’s BICEP project (Jönsson et al., 2022). These have been derived globally using three established export production algorithms (Dunne et al., 2005; Henson et al., 2011; Z. Li & Cassar, 2016) driven by satellite observations of sea surface temperature (SST) (GHRSSST/OSTIA, UK Met Office, 2005), Chl*a* and the diffuse atten-

uation coefficient at 490 nm (OC-CCI, Sathyendranath et al., 2020), and NPP (BICEP, Kulk et al., 2021).

These algorithms differ in their structure and data sources: Dunne et al. (2005)'s is based on a statistical fit to SST and a phytoplankton productivity metric (either NPP or Chl a normalised by euphotic depth) after a global compilation of POC flux export measurements from a combination of shallow sediment traps and ^{234}Th -derived POC fluxes, with outputs constrained within fixed bounds; Henson et al. (2011) uses a simpler exponential fit to SST and NPP using observations of ^{234}Th -derived fluxes from various authors, and is parameterised for export at ~ 100 m depth, without explicitly incorporating euphotic depth; and Z. Li and Cassar (2016) applies a machine-learning approach (genetic programming) to optimise the empirical relationship between NPP and SST of Henson et al. (2011) and Laws, Falkowski, Smith, Ducklow, and McCarthy (2000), aiming to better match dissolved O_2 :Ar-based estimates of net community production, and improve predictive skill over earlier empirical formulations.

Figure S7 presents a validation of these three satellite-based BICEP products estimating POC export flux against our *in situ* dataset from sediment traps and radionuclides. The BICEP-Li product outperforms both the BICEP-Henson and BICEP-Dunne algorithms, as indicated by its higher Spearman rank correlation coefficient (ρ) and lower bias (δ) and root mean square error (RMSE). In contrast, BICEP-Henson and BICEP-Dunne tend to systematically underestimate *in situ* measurements (negative δ), particularly at high flux values, whereas BICEP-Li more accurately captures the magnitude of observed high export peaks.

Text S4. Particle State Equations

We present the Lagrangian superparticle state equations governing particle attributes (fractal scaling, stickiness, and sinking velocity). Parameters are given in **Tables A1–A3** and prescribed variables in **Texts S1–S2**. While many equations are inherited from SLAMS-1.0, some have been extended or modified to reflect the updated process representations in SLAMS-2.0.

Modelled aggregates form through the coagulation of idealised spherical primary particles (sphericity of primaries is a necessary simplification for hydrodynamic calculations, but in reality are irregular in shape). The resulting aggregates are fractal structures: porous, irregular, and statistically self-similar across scales. This self-similarity implies a scaling relationship between size-dependent properties (mass, area, length) and a characteristic scale (radius, diameter). For marine aggregates, the relevant property that scales with size is mass, m_p —or, equivalently the number of primary particles, n_{pp} — following a power law, rather than the cubic relation of a compact sphere:

$$m_p^{\text{fractal}} = m_{pp} \cdot n_{pp} = m_{pp} \cdot \left(\frac{r_p}{r_{pp}} \right)^{D_3}, \quad (10a)$$

$$m_p^{\text{compact}} = \rho_p \cdot V_p = \rho_p \cdot \frac{4\pi}{3} \cdot r_p^3, \quad (10b)$$

where m_{pp} and r_{pp} are the mass and radius of the primary particles, V_p is the bulk (envelope) volume of the aggregate, ρ_p its bulk density, and D_3 the three-dimensional (3-D) fractal dimension. The fractal dimension quantifies how efficiently an object fills space: $D_3 = 1$ for filamentary (one-dimensional) structures, $D_3 = 2$ for surface-like (two-dimensional) forms, and $D_3 = 3$ for compact, volumetric (three-dimensional) bodies. For a compact, non-porous sphere, $D_3 = 3$; in contrast, marine aggregates typically exhibit $D_3 \approx 1.3\text{--}2.6$, reflecting their porous structure, in which solid material occupies only

a *fraction* of the embedding 3-D space, and the object is therefore *fractal* (Logan & Wilkinson, 1990; Kilps, 1993; X. Li & Logan, 1995; Burd & Jackson, 1997; Jackson et al., 1997). Lower D_3 values correspond to open, highly porous structures and higher values to denser, more compact structures.

In SLAMS-2.0, primary particles are assigned $D_3 = 3$, while newly formed aggregates sample D_3 uniformly within 1.3–1.9 (a range adjusted after Kilps, 1993). From **Eq. 10**, since $m_p \propto r_p^{D_3}$ but the volume scales as $V_p \propto r_p^3$, aggregates with $D_3 < 3$ gain volume faster than mass as they grow, producing lower bulk density and higher porosity. This behaviour —i.e., large aggregates that are relatively porous and therefore sink more slowly than compact particles of the same size— is a key observation that the fractal scaling in SLAMS-2.0 is designed to capture. Particle attributes are derived following this scaling below.

Particle Attributes Derived from Fractal Scaling

Mass. The dry mass of a particle, m_p (g), is obtained from its material composition:

$$m_p = \sum_{x=1} q_x \cdot M_x, \quad x = \text{organic matter, CaCO}_3, \text{ bSi, clay}, \quad (11)$$

where q_x is the quantity (mol) of material x and M_x its molar mass. The organic matter component combines organic carbon ($m_{C_{\text{org}}} = q_{C_{\text{org}}} \cdot M_C$) and TEP carbon ($m_{\text{TEPC}} = q_{\text{TEPC}} \cdot M_C$), converted to their respective organic matter (OM) equivalents via carbon-to-matter factors, i.e., $m_{\text{OM}} = m_{C_{\text{org}}}/f_{C:\text{OM}} + m_{\text{TEPC}}/f_{\text{TEPC}:\text{TEP}}$.

Volume. The solid (non-porous) volume of the aggregate particle, V_p^S (μm^3), is:

$$V_p^S = \sum_{x=1} \frac{m_x}{\rho_x}, \quad (12)$$

where ρ_x is the density of each material x .

Radius. The number of constituent primary particles, n_{pp} scales with aggregate radius as r_{p} (μm), as $n_{\text{pp}} = (r_{\text{p}}/r_{\text{pp}})^{D_3}$ (**Eq. 10a**), which can be rearranged to give:

$$r_{\text{p}} = r_{\text{pp}} \cdot n_{\text{pp}}^{1/D_3}. \quad (13)$$

For $D_3 < 3$, aggregates have a more open structure and therefore larger radii for a given n_{pp} . The fractal scaling assumes identical, spherical primary particles of radius:

$$r_{\text{pp}} = \left(\frac{3 \cdot V_{\text{pp}}}{4\pi} \right)^{1/3}, \quad (14)$$

where $V_{\text{pp}} = V_{\text{p}}^{\text{S}}/n_{\text{pp}}$ is the volume of a single primary particle.

Porosity. Let V_{p}^{B} denote the bulk (envelope) volume of a particle. Porosity, P_{p} , is defined as the fraction of the particle volume occupied by seawater:

$$P_{\text{p}} = 1 - \frac{V_{\text{p}}^{\text{S}}}{V_{\text{p}}^{\text{B}}} = 1 - \frac{n_{\text{pp}} \cdot V_{\text{pp}}}{V_{\text{p}}^{\text{B}}} = 1 - n_{\text{pp}} \cdot \left(\frac{r_{\text{pp}}}{r_{\text{p}}} \right)^3. \quad (15)$$

Substituting the fractal relation for n_{pp} (**Eq. 10a**) yields:

$$P_{\text{p}} = 1 - \left(\frac{r_{\text{p}}}{r_{\text{pp}}} \right)^{D_3-3}. \quad (16)$$

This expression implies a size-dependent porosity: smaller aggregates are relatively compact, whereas larger aggregates are increasingly porous as the spacing between primary particles grows, increasing the fraction of seawater within the particle structure.

From the porosity, we define the effective radius, $r_{\text{p}}^{\text{eff}}$, as the radius of an equivalent non-porous sphere containing the same volume of solid material as the aggregate:

$$r_{\text{p}}^{\text{eff}} = r_{\text{p}} \cdot (1 - P_{\text{p}})^{1/3}. \quad (17)$$

Thus, while the bulk radius r_{p} characterises the full particle envelope, including porous space, the effective radius $r_{\text{p}}^{\text{eff}}$ represents only the volume of solid material, and is therefore smaller than the bulk radius.

Density. The bulk (effective) density of a particle, ρ_p (g cm^{-3}), combines the solid and pore water fractions:

$$\rho_p = (1 - P_p) \cdot \rho_p^S + P_p \cdot \rho_w, \quad (18)$$

where ρ_w is the prescribed seawater density, and $\rho_p^S = m_p/V_p^S$ is the density of the solid fraction. Finally, the particle's excess density, which is used to calculate its sinking velocity, is defined as:

$$\Delta\rho_p = \rho_p - \rho_w. \quad (19)$$

Particle Stickiness

Stickiness, α_p (ranging from 0 to 1), represents the likelihood that two particles will attach upon collision. In the ocean, particles do not always stick upon contact: some collisions result in attachment, while others do not. This tendency to stick is enhanced in the presence of organic matter, particularly transparent exopolymer particles (TEPs), which act as natural 'glues' that promote aggregation (Alldredge et al., 1993).

TEPs are carbohydrate-based gels cross-linked by divalent cations and sulfate bridges (Passow, 2002; Meng & Liu, 2016) and are the driver of stickiness in the SLAMS formulation. Empirical studies typically infer stickiness from aggregation experiments involving the presence of TEPs (Kiørboe & Hansen, 1993; Dam & Drapeau, 1995; Engel, 2000). Although no exact formula exists for predicting stickiness directly from TEP concentration, Jokulsdottir and Archer (2016) introduced a mechanistic representation in SLAMS-1.0 linking higher TEP content to an increased probability of particle attachment. In their framework, stickiness increases linearly with the fraction of an aggregate's solid volume

occupied by TEP material:

$$\alpha_p = \frac{m_{\text{TEP}}/\rho_{\text{TEP}}}{V_p^S} = \frac{V_{\text{TEP}}}{V_p^S}. \quad (20)$$

where $V_{\text{TEP}} = m_{\text{TEP}}/\rho_{\text{TEP}}$ is the TEP volume within the aggregate.

While **Eq. 20** provides a simple and physically intuitive formulation, it may not capture all aspects of particle adhesion. We hypothesise that stickiness increases with TEP content up to a saturation point, beyond which additional polymer yields diminishing gains in adhesive contact area, analogous to saturating responses observed in other biogeochemical processes. In addition, mineral content may reduce effective adhesion by occupying surface area and introducing rigid contacts that oppose gel-mediated bridging. Based on these two considerations, SLAMS-2.0 includes an alternative formulation in which stickiness depends on (i) a saturating (Michaelis–Menten) response to the TEP volume fraction and (ii) a multiplicative mineral dilution factor that reduces stickiness in proportion to the mineral fraction:

$$\alpha_p = \underbrace{\frac{\frac{V_{\text{TEP}}}{V_p^S}}{\frac{V_{\text{TEP}}}{V_p^S} + k_{\text{TEP}}}}_{\text{TEP saturation}} \cdot \underbrace{\left(1 - \frac{V_{\text{mineral}}}{V_p^S}\right)}_{\text{mineral dilution}}, \quad (21)$$

where V_{TEP}/V_p^S is the dimensionless TEP volume fraction (0–1), k_{TEP} is the half-saturation constant (dimensionless; equals the TEP fraction at which the TEP term reaches 0.5), V_{mineral}/V_p^S is the mineral volume fraction (0–1).

Particle Sinking Velocity

The particle sinking velocity, v_p (m d^{-1}), follows from a balance of three forces on a sphere falling in viscous fluid (e.g., Maggi, 2013): downward gravitational force, F_g , upward buoyancy force that emerges from the weight of the fluid displaced by the particle,

F_b , and drag force that pulls the particle upward, F_d :

$$F_g = m_p \cdot g = \rho_p \cdot V_p^B \cdot g, \quad (22a)$$

$$F_b = m_w \cdot g = \rho_w \cdot V_p^B \cdot g, \quad (22b)$$

$$F_d = \frac{1}{2} \cdot \rho_w \cdot v_p^2 \cdot C_D \cdot A_p, \quad (22c)$$

where g is the gravitational constant, $V_p^B = (4/3) \cdot \pi \cdot r_p^3$ is the particle (bulk) volume, $A_p = \pi \cdot r_p^2$ is the projected area of the particle, and C_D is the drag coefficient.

When the two retarding forces balance the gravitational force, the net force acting on the particle vanishes ($F_{\text{net}} = F_g - (F_b + F_d) = 0$) and, consequently, the net acceleration of the particle goes to zero, so it keeps falling at a constant velocity, its terminal velocity (v_p). Its generic expression can be isolated from the force balance F_{net} :

$$v_p = \sqrt{\frac{8}{3} \frac{r_p \cdot g \cdot \Delta\rho_p}{\rho_w \cdot C_D}}. \quad (23)$$

The drag coefficient, C_D , is a function of the flow regime around the particle, and is parameterised through the Reynolds number, Re (dimensionless):

$$Re = \frac{\text{inertial force}}{\text{viscous force}} = \frac{d_p \cdot v_p \cdot \rho_w}{\mu}, \quad (24)$$

where v_p is in m s^{-1} , $d_p = 2 \cdot r_p$ is the particle diameter in metres, ρ_w is the prescribed seawater density in kg m^{-3} , and μ is the prescribed seawater dynamic viscosity in $\text{kg m}^{-1} \text{s}^{-1}$.

Re compares inertial to viscous forces: low values ($Re < 0.1$) indicates viscous-dominated (laminar) flow, whereas larger values indicate transitional or turbulent flow (e.g., caused by increasing particle size or settling velocity). Although no exact analytical relationship exists between Re and C_D , several empirical correlations for impermeable spheres are widely used (e.g., Maggi, 2013; Laurenceau-Cornec et al., 2019). SLAMS-2.0 adopts the empirical formulation of White (1974) for solid spheres, which uses the Stokes (1851) drag

in laminar flow and an extended fit for higher Re :

$$C_D = \frac{24}{Re}, \quad Re < 0.1 \quad (\text{laminar}) \quad (25a)$$

$$C_D = \frac{24}{Re} + \frac{6}{1 + \sqrt{Re}} + 0.4, \quad Re \geq 0.1 \quad (\text{transitional/turbulent}) \quad (25b)$$

These drag relations yield two corresponding expressions for v_p (in m s^{-1}): one for laminar flow (known as the ‘Stokes’ velocity’, after Stokes, 1851) and another for transitional/turbulent flow:

$$v_p = \frac{2}{9} \frac{g \cdot \Delta\rho_p \cdot r_p^2}{\mu}, \quad (\text{laminar}) \quad (26a)$$

$$v_p = \sqrt{\frac{4}{3} \frac{g \cdot \Delta\rho_p \cdot d_p}{\rho_w \cdot C_D}}, \quad (\text{transitional/turbulent}) \quad (26b)$$

According to the expression for laminar flow, v_p increases linearly with excess density and quadratically with particle radius, implying a parabolic rise in velocity with particle size. This form is commonly used to compute sinking velocities, but it relies on three simplifying assumptions: (i) laminar flow ($Re < 0.1$); (ii) impermeable, non-fractal particles; and (iii) perfectly spherical geometry. By contrast, for transitional/turbulent flow, C_D and Re both depend on v_p , so the velocity must be solved iteratively. In SLAMS-2.0, this is achieved using a Newton iteration method until convergence.

Text S5. Mesozooplankton Size-Dependent Attributes

In SLAMS-2.0, several mesozooplankton traits relevant to encounter dynamics are parameterised explicitly as functions of body size. Mesozooplankton are represented by discrete size classes indexed by j , each characterised by a representative equivalent spherical diameter (ESD). The characteristic diameter of each size class is defined as the geometric mean of the bin bounds (i.e., the middle point),

$$\text{ESD}_j^{\text{geom}} = \sqrt{\text{ESD}_j^{\text{lower}} \cdot \text{ESD}_j^{\text{upper}}}, \quad (27)$$

where ‘lower’ and ‘upper’ denote the lower and upper bounds of size class j . As explained in Appendix A23 (after Kwong & Pakhomov, 2021, Table S1), there are 26 size classes, with bounds ranging 185–10,950 μm , increasing by a constant factor of 1.17.

Radius. The characteristic radius of zooplankton size class j is:

$$r_{z,j} = \frac{\text{ESD}_j^{\text{geom}}}{2}. \quad (28)$$

Prosome length. Prosome length is estimated from ESD assuming a prolate spheroidal body shape with aspect ratio 0.38, following Kiørboe (2011):

$$\text{PL}_{z,j} = \frac{\text{ESD}_j^{\text{geom}}}{0.38^{2/3}}. \quad (29)$$

Wet weight quota. Each size class is assigned a characteristic wet mass, $m_{z,j}$ (mg WW), computed from the radius assuming spherical geometry and a constant body density of $\rho_z = 1 \text{ g cm}^{-3}$ (after Kiørboe, 2013):

$$m_{z,j} = \rho_z \cdot \frac{4\pi}{3} \cdot r_{z,j}^3, \quad (30)$$

Carbon weight quota. Wet mass, $m_{z,j}$ (mg WW), is converted to carbon mass, $m_{C,j}$ (mg C), using the empirical copepod relationship of Kiørboe (2013, Table 2):

$$\log_{10}(m_{z,j}) = -0.93 + 0.95 \cdot \log_{10}(m_{C,j}). \quad (31)$$

Reaction distance. The effective reaction distance governing particle–zooplankton encounters is given by:

$$r_{z,j}^{\text{detect}} = R_z \cdot r_{z,j}, \quad (32)$$

where R_z is a dimensionless parameter describing mechanoreceptive or chemical sensing efficiency (Table ??).

Swimming speed. For cruising mesozooplankton, swimming speed, $v_{z,j}$ (m s^{-1}), follows the empirical size-based formulation of Kiørboe (2011, Fig. 3):

$$\log_{10}(v_{z,j}) = 0.39 + 0.79 \cdot \log_{10}(\text{ESD}_j^{\text{geom}}), \quad (33)$$

where $v_{z,j}$ is expressed in cm s^{-1} and $\text{ESD}_j^{\text{geom}}$ in cm.

Faecal pellet volume. Faecal pellet volume depends on the PL of the mesozooplankton (**Eq. 29**), with a relationship given by the empirical formula of Stamieszkin et al. (2015):

$$\log_{10}(V_{\text{fp}}^{\text{B}}) = 5.4 + 2.58 \cdot \log_{10}(\text{PL}_{z,j}), \quad (34)$$

where V_{fp}^{B} is in μm^3 and $\text{PL}_{z,j}$ is in mm.

Gut volume. The mesozooplankter gut volume is estimated from the bulk volume of faecal pellets that could be egested from the ingested food particle. As per Jansen (2001, Table 1, footnote), a good approximation is:

$$V_{\text{gut}} = 5 \cdot V_{\text{fp}}^{\text{B}}. \quad (35)$$

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Table S1. Pearson correlation coefficient (r^2) between monthly means of depth-integrated NPP from five satellite-based NPP models (VGPM, CbPM, CAFE, BICEP, and Carr (2002)) and their *in situ* ^{14}C -based observations counterparts at stations ALOHA ($N = 12$), BATS ($N = 12$), EqPac ($N = 5$), and OSP ($N = 10$) (all data are plotted in Figure S5). For each location, the best-performing model statistics are highlighted in green, and the second-best in yellow.

| | ALOHA | BATS | EqPac | OSP |
|------------------------|-------|------|-------|------|
| VGPM (MODIS) | 0.28 | 0.51 | 0.50 | 0.44 |
| CbPM (MODIS) | 0.31 | 0.07 | 0.00 | 0.49 |
| CAFE (MODIS) | 0.88 | 0.67 | 0.02 | 0.46 |
| BICEP (merged sensors) | 0.66 | 0.81 | 0.14 | 0.49 |
| Carr 2002 (MODIS, C02) | 0.07 | 0.79 | 0.00 | 0.60 |

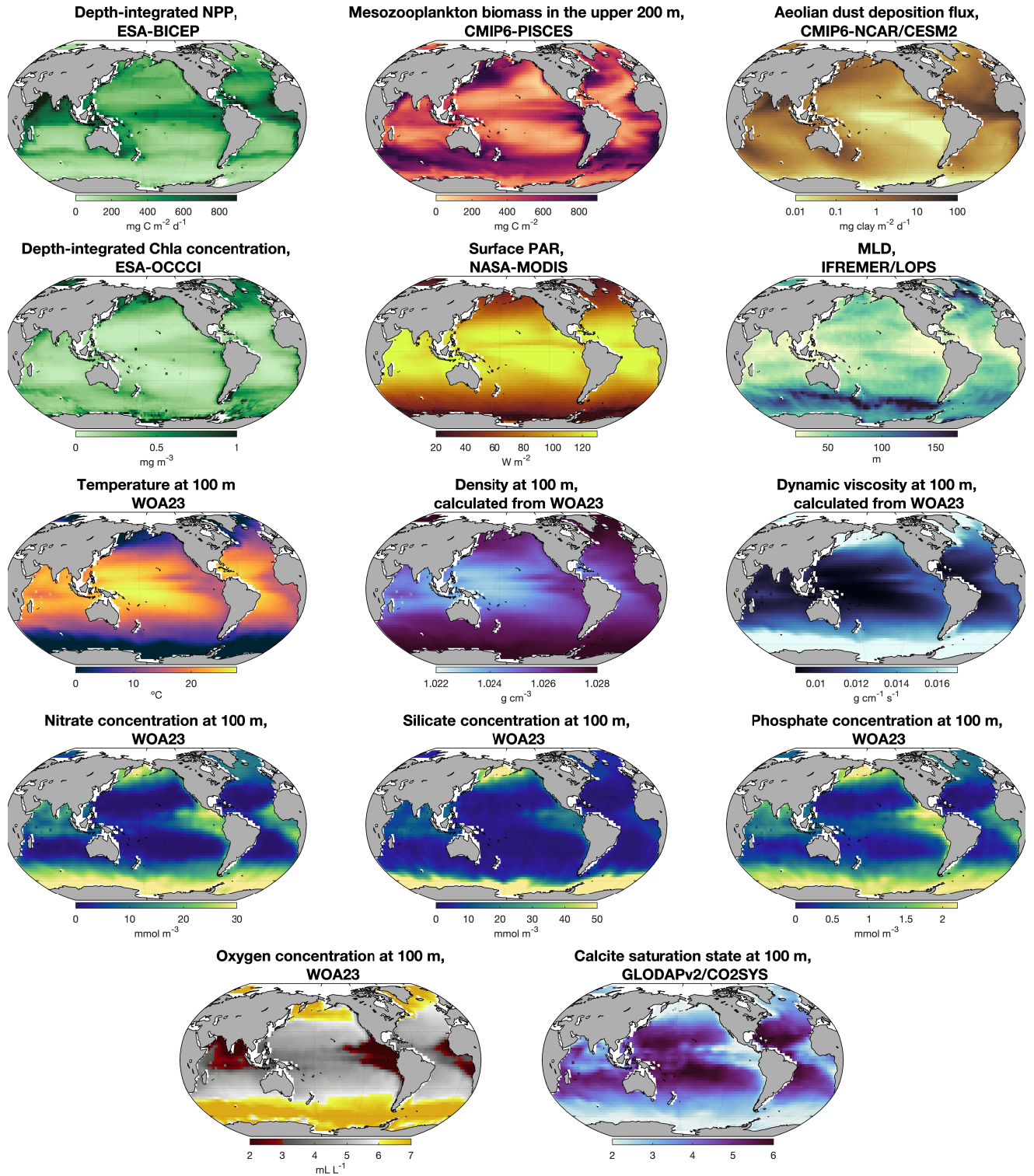


Figure S1. Global ocean distribution of the climatological annual mean of the variables used to force SLAMS-2.0.

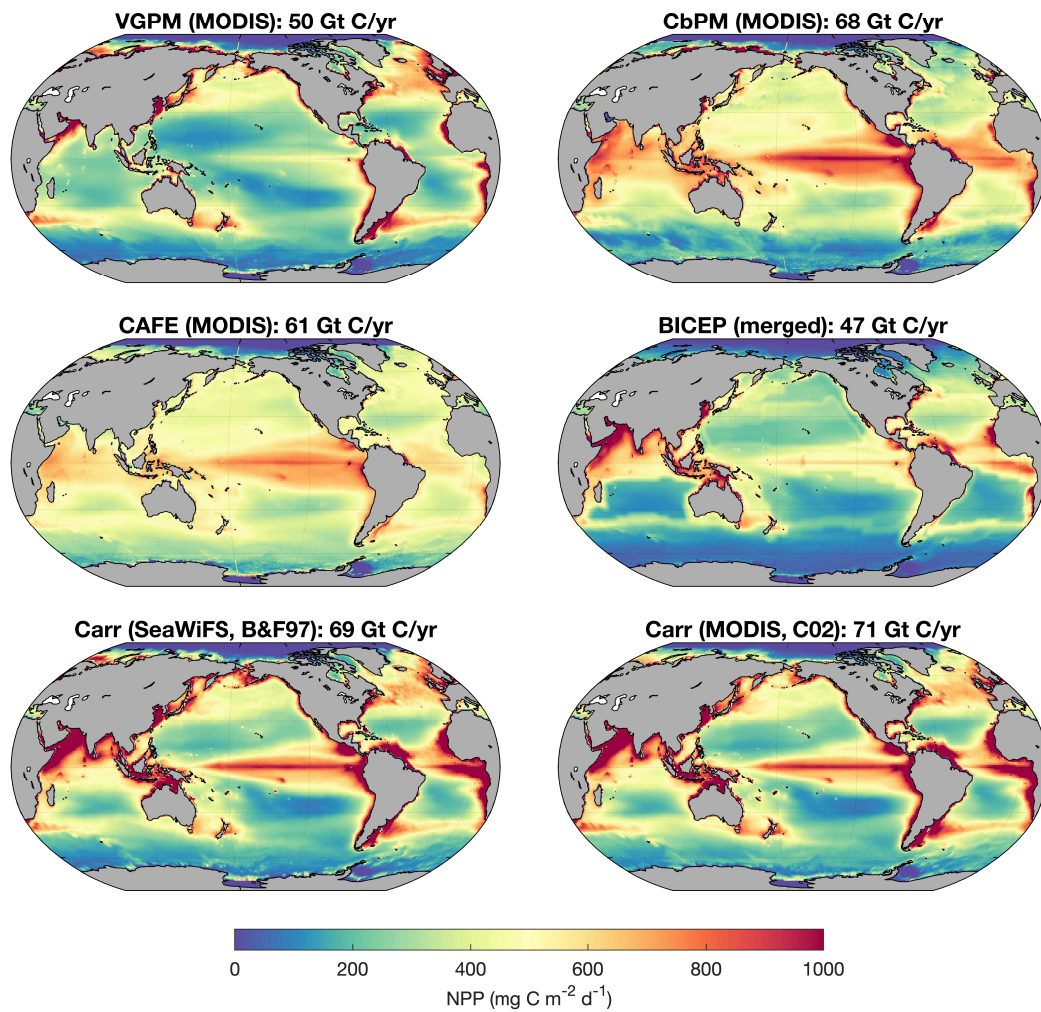


Figure S2. Global distribution of depth-integrated annual net primary production (NPP) from five ocean-colour-based models: (A) VGPM (Vertically Generalized Productivity Model, Behrenfeld & Falkowski, 1997); (B) CbPMv2 (Carbon-based Production Model-2, Westberry et al., 2008); (C) CAFE (Carbon, Absorption, and Fluorescence Euphotic-resolving model, Silsbe et al., 2016); (D) BICEP (Biological Pump and Carbon Exchange Processes project model, Kulk et al., 2020); (E–F) NPP model of Carr (2002), shown in two configurations. VGPM, CbPMv2, and CAFE use MODIS-Aqua inputs (Chl_a , PAR_0 , SST) and are available as finished products at the Ocean Productivity site (<http://sites.science.oregonstate.edu/ocean.productivity/>). The BICEP product is a merged multi-sensor product and is accessible via the ESA Ocean Colour Climate Change Initiative website (<http://www.oceancolour.org>). The Carr model was run using (E) SeaWiFS Chl_a and PAR_0 with the euphotic depth formulation of Behrenfeld and Falkowski (1997) (as applied by Henson et al. (2012)), and (F) MODIS-Aqua Chl_a and PAR_0 with the original euphotic depth of Carr (2002); both configurations use AVHRR Pathfinder v5.0 SST (1985–2001) (<https://www.ncei.noaa.gov/products/avhrr-pathfinder-sst>). MODIS-Aqua and SeaWiFS data were obtained from NASA Ocean Color website (<https://oceandata.sci.gsfc.nasa.gov/13/>).

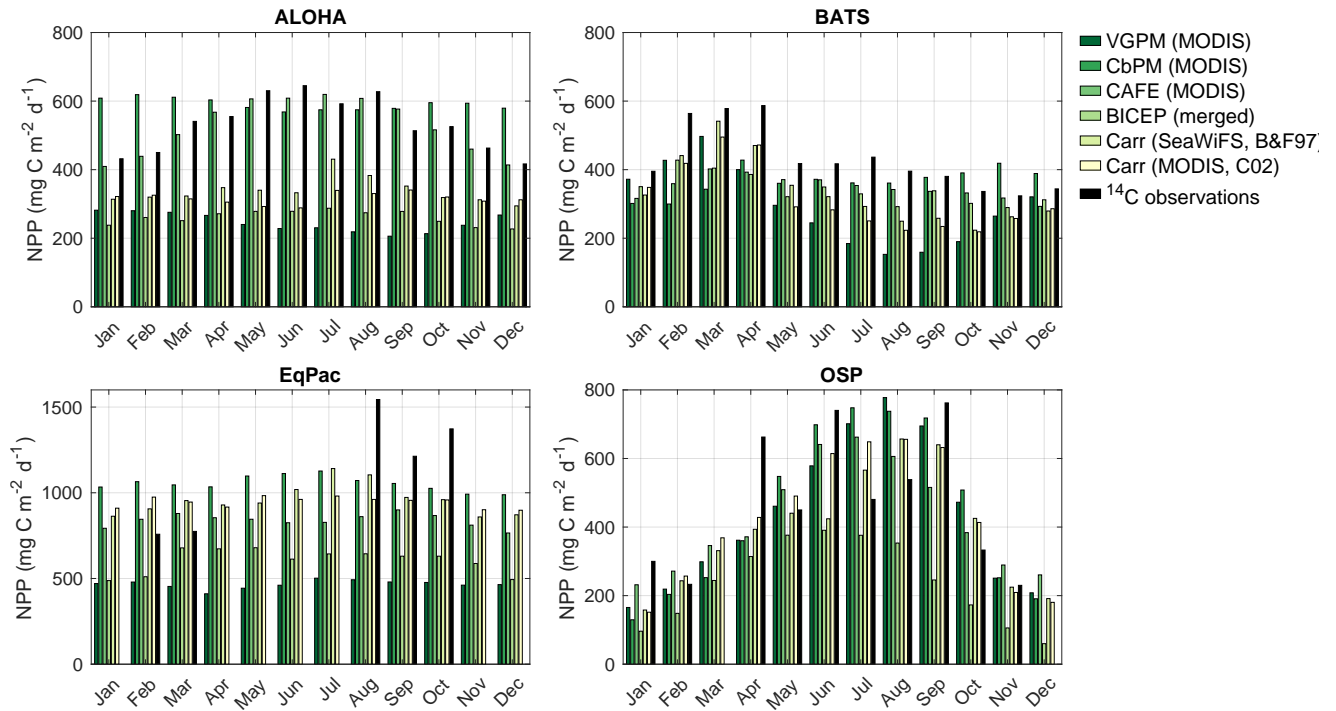


Figure S3. Comparison of monthly mean depth-integrated NPP at four ocean sites (ALOHA, BATS, EqPac, and OSP). Data include satellite-derived estimates from five ocean-colour-based models (VGPM, CbPM, CAFE, BICEP, and Carr (2002); see Figure S2) and *in situ* ^{14}C uptake measurements. ^{14}C data sources are: ALOHA (1989–2018, Karl et al., 2021, Supplementary Table 1); BATS (1988–2022, Bates et al., 2023); EqPac (Murray et al. (1996), Table 5, 2°N–2°S; Le Borgne and Gesbert. (1995), Ch. 11, stations 94 and 100); and OSP (compiled from Long et al. (2021) and sources therein —Booth (1988), Miller et al. (1991), Welschmeyer et al. (1993), Wong et al. (1995), Boyd and Harrison (1999), Kawakami et al. (2010), Giesbrecht et al. (2012), Timmerman and Hamme (2021)— as well as in the synthesis study for the OSP time-series site (1987–1996, Timothy et al., 2013, Fig. 2(i)).

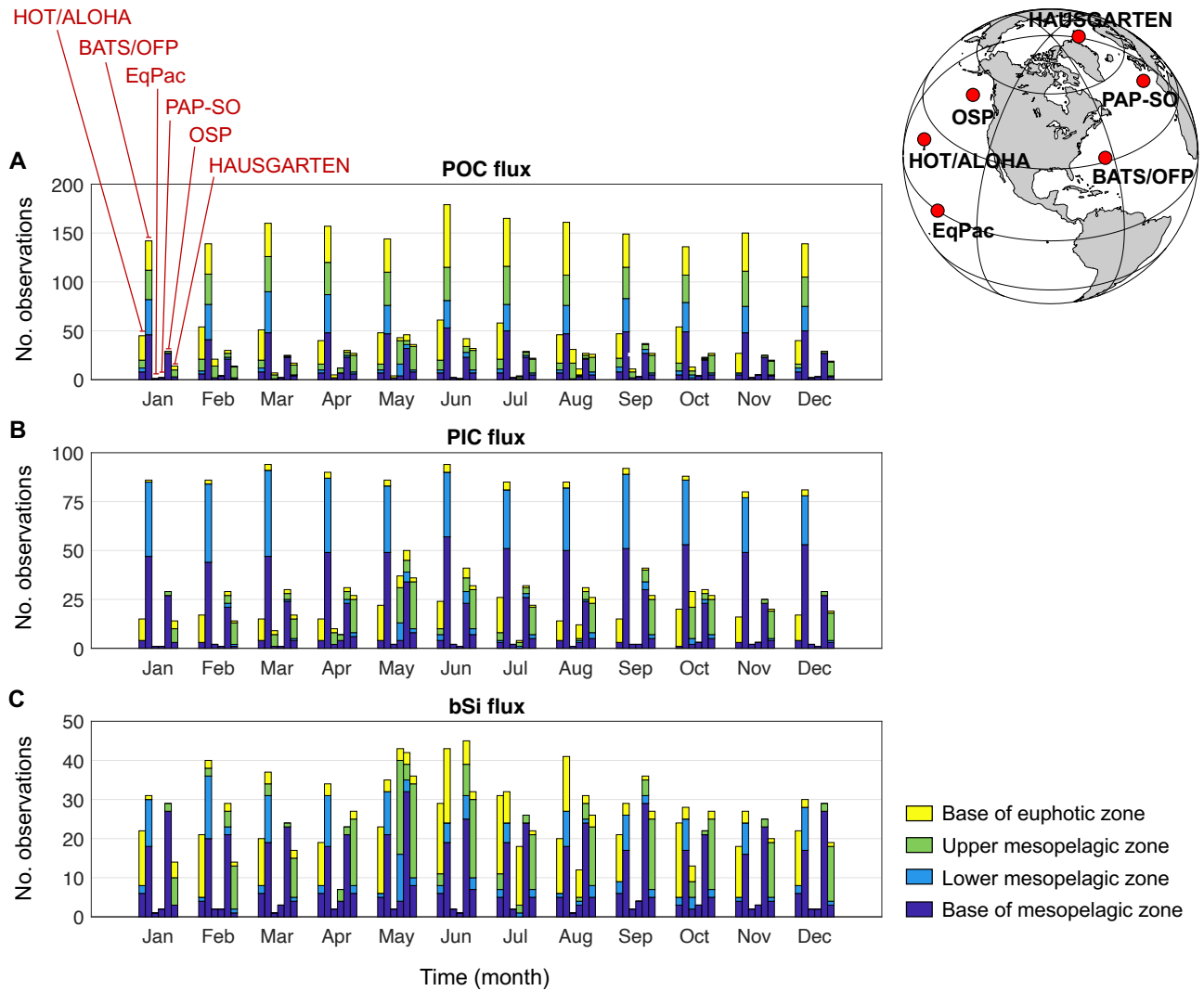


Figure S4. Number of (A) POC, (B) PIC, and (C) bSi flux observations from sediment traps and radionuclide disequilibrium contributing to the monthly-averaged depth values shown in Figure S5. Data are grouped by month for each of our six study sites (HOT/ALOHA, BATS/OFP, EqPac, PAP-SO, OSP, HAUSGARTEN) and by depth zone (base of the euphotic zone, upper, lower, and base mesopelagic; see Table S2). In total, $N = 7162$ POC, $N = 5282$ PIC, and $N = 3856$ bSi flux measurements were compiled, of which $N = 3240$, $N = 2079$, and $N = 1451$ were used for model optimisation after depth selection. The inset map shows the geographical locations of the six sites: HOT/ALOHA (subtropical NE Pacific), BATS/OFP (subtropical NW Atlantic), EqPac (central equatorial Pacific), PAP-SO (subpolar NE Atlantic), OSP (subpolar NE Pacific), and HAUSGARTEN (subarctic Atlantic).

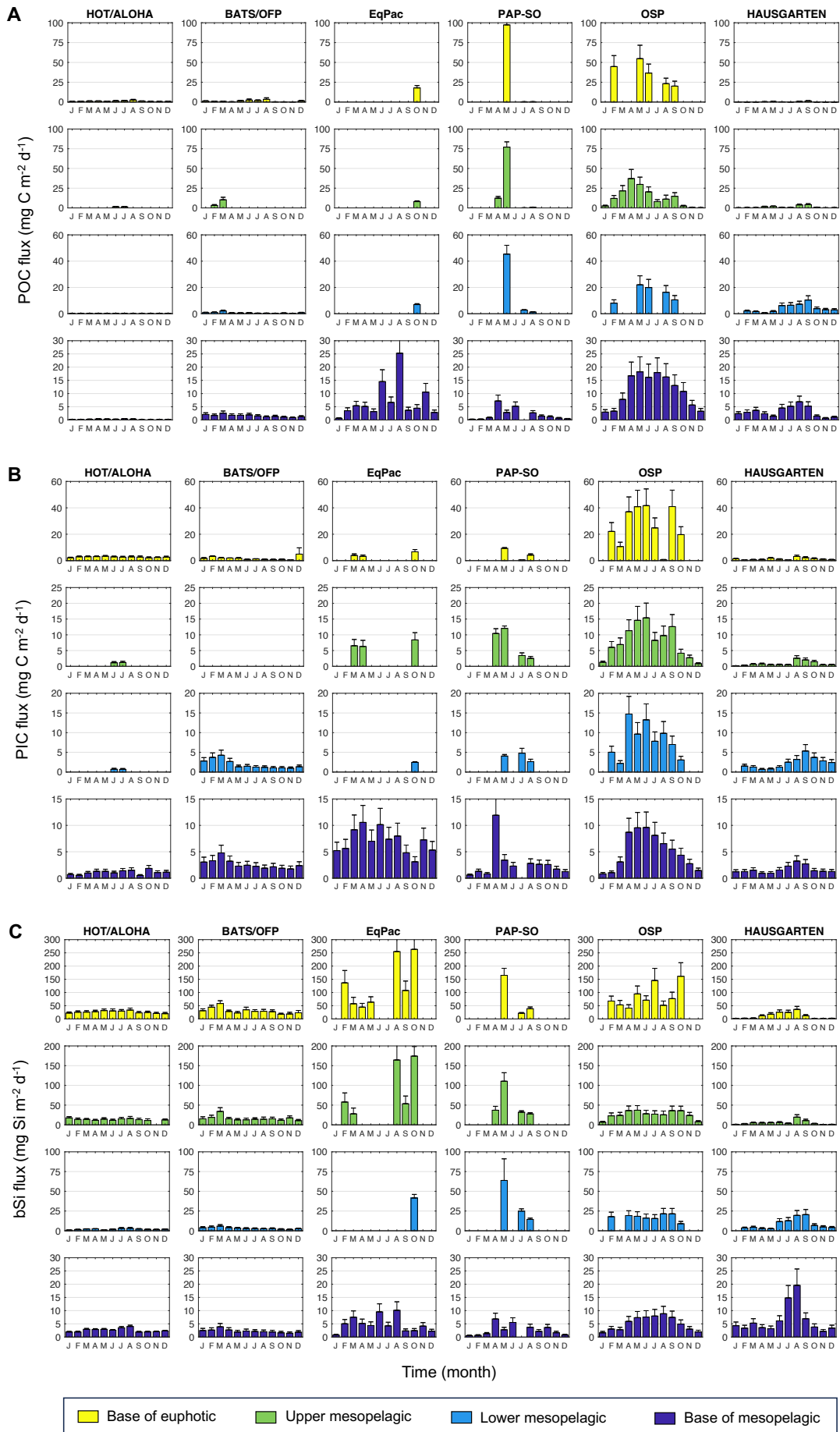


Figure S5. Observed monthly mean fluxes of (A) POC, (B) PIC, and (C) bSi compiled from sediment traps and radionuclides at our six ocean study sites. Vertical error bars indicate net uncertainty and has been calculated as the quadratic sum of systematic and random errors propagated from measured values.

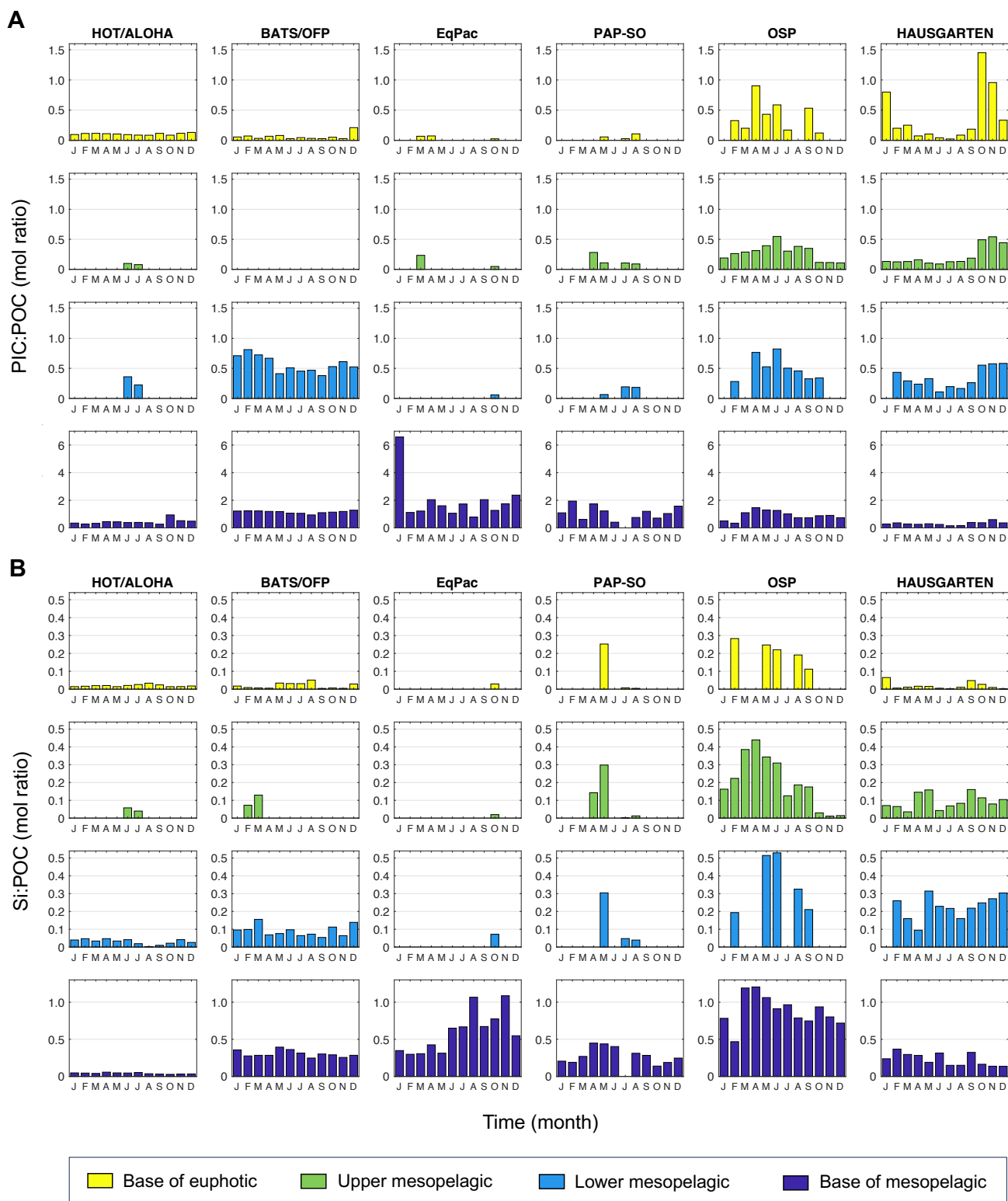


Figure S6. Observed monthly mean flux ratios of two ballast minerals relative to POC across our six ocean study sites. (A) Calcium carbonate (PIC, mainly from coccolithophore carbonated shells) to POC. (B) Biogenic silica (Si, mainly from diatom frustules) to POC. Ratios are derived from sediment trap and radionuclide data shown in Figure S7.

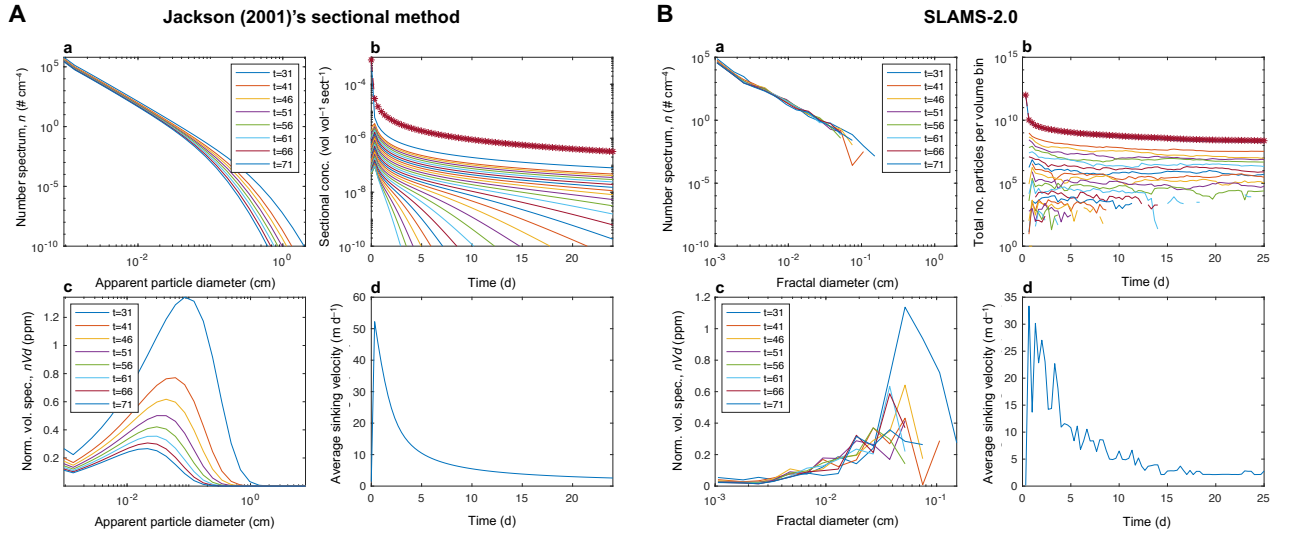


Figure S7. Comparison of two particle coagulation approaches. (A) Solution of the Smoluchowski equation using a sectional population-balance approach with fractal kernels following Jackson (2001) (implementation from Jackson, 2004). (B) SLAMS-2.0's coagulation subroutine based on the Super-Droplet Method (SDM; Shima et al., 2009), which visually converges to the Smoluchowski solution. Both simulations are initialised with a particle concentration of 10^6 particles cm^{-3} in the smallest size class and run for 25 modelled days (with a 1/3-day time step) in a 50 m deep mixed layer. In SLAMS-2.0, particles are represented by 1000 superparticles with fractal dimension 2. Panels (a) and (b) show the evolution of the particle number spectrum and total mass concentration across 27 size classes, respectively. Both simulations exhibit rapid transfer of particles toward larger sizes via coagulation, with SLAMS-2.0 converging to the Smoluchowski solution as coagulation and settling losses balance. For smaller particles, coagulation losses exceed settling, while for larger particles, settling losses become larger than coagulation losses.

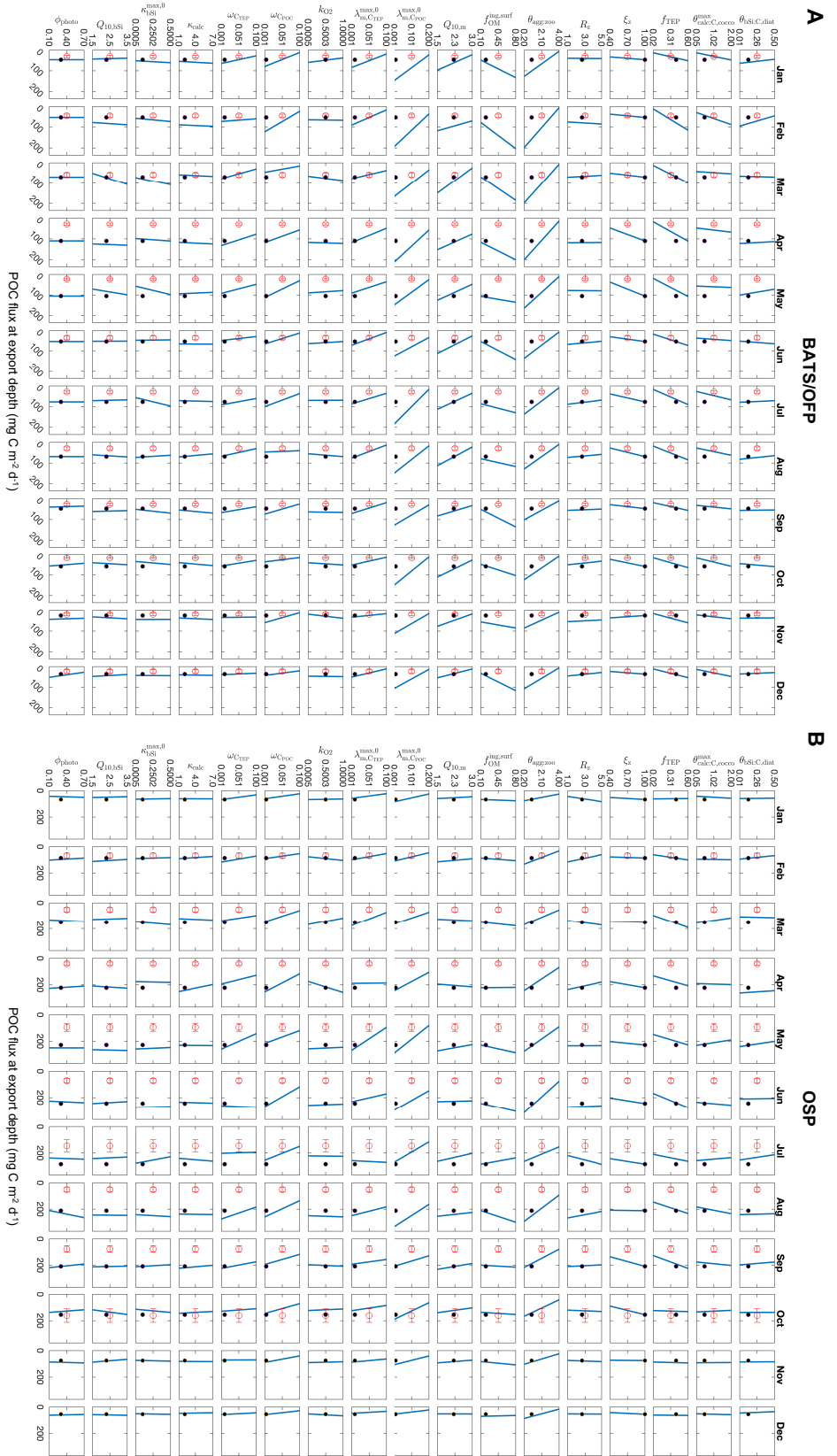


Figure S8. Parameter sensitivity of simulated particulate organic carbon (POC) flux at the export depth for two study sites (A and B). Rows correspond to the 17 tested parameters and columns to months of the year. Blue lines connect model outputs obtained using the minimum and maximum literature-supported parameter values; black dots indicate default model output, and red dots indicate observations.

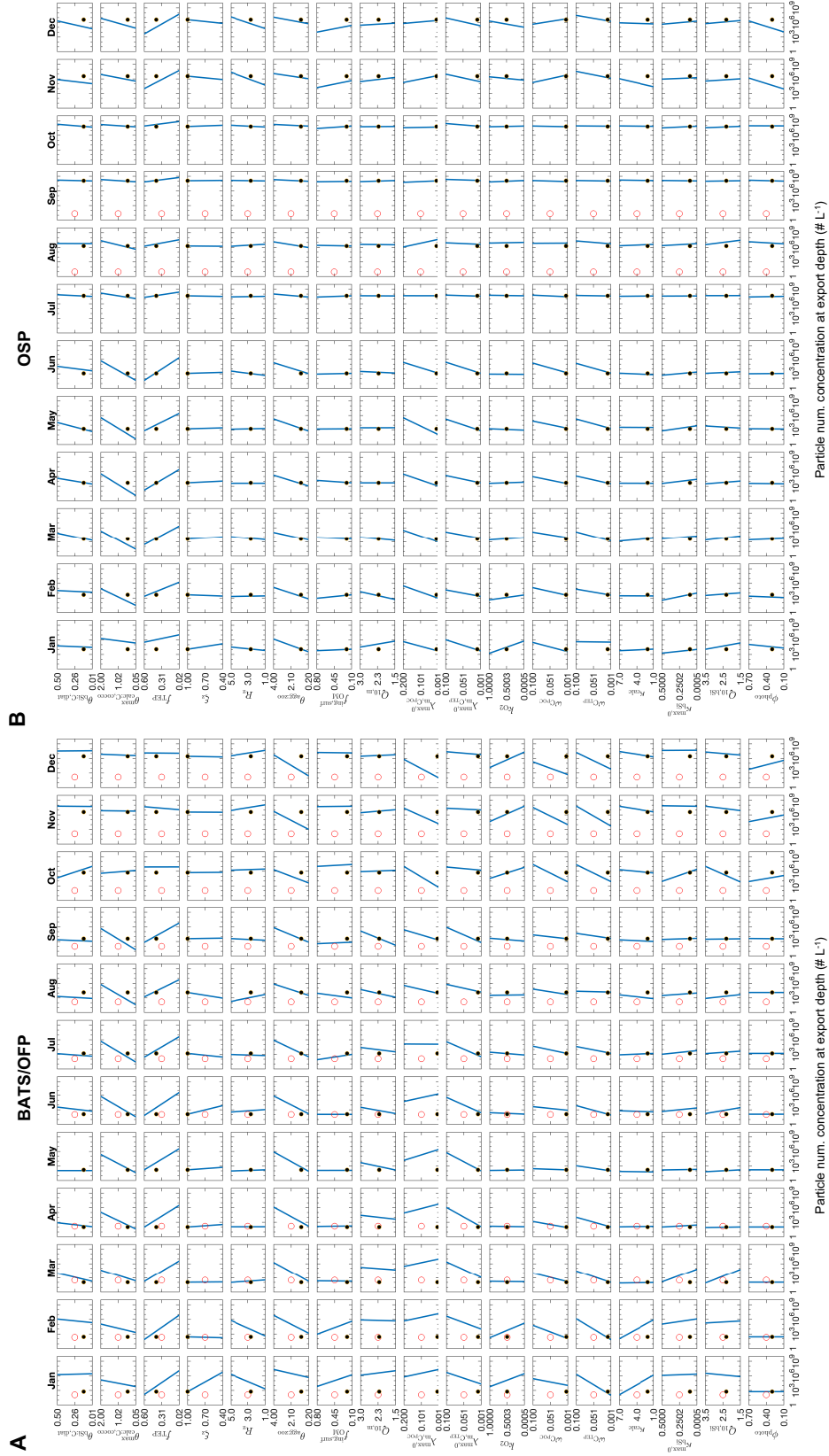


Figure S9. As in Figure S8, but for particle number concentration at the export depth.

Table S2. Depth zone ranges (in metres) used to extract both observed and modelled fluxes of POC, PIC, and bSI at our six study sites: HOT/ALOHA, BATS/OFP, EqPac, PAP-SO, OSP, and HAUSGARTEN. For the “Base of euphotic” category, monthly depth ranges are shown, and blank cells indicate months with no available observational data. For the mesopelagic layers, depth ranges are shown as fixed values across all months, even though observational data were not always available in every month.

| | HOT/ALOHA | | | | | BATS/OFP | | | | | EqPac | | | | | |
|---------------------|------------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|
| | z_{eu} | POC | PIC | bSI | z_{eu} | POC | PIC | bSI | z_{eu} | POC | PIC | bSI | z_{eu} | POC | PIC | bSI |
| Base of euphotic | 115-142 | 135-165 | 135-165 | 135-165 | 96-118 | 135-165 | 135-165 | 135-165 | 83-103 | – | – | – | – | – | – | – |
| | 119-147 | 135-165 | 135-165 | 135-165 | 92-114 | 135-165 | 135-165 | 135-165 | 83-102 | 90-110 | – | – | – | – | – | – |
| | 95-117 | 135-165 | 135-165 | 135-165 | 93-115 | 135-165 | 135-165 | 135-165 | 68-84 | 67-83 | 72-88 | – | – | – | – | – |
| | 97-120 | 135-165 | 135-165 | 135-165 | 85-105 | 135-165 | 135-165 | 135-165 | 68-84 | 90-110 | 72-88 | – | – | – | – | – |
| | 101-125 | 135-165 | 135-165 | 135-165 | 105-130 | 135-165 | 135-165 | 135-165 | 80-98 | 90-110 | – | – | – | – | – | – |
| | 102-126 | 135-165 | 135-165 | 135-165 | 111-137 | 135-165 | 135-165 | 135-165 | 77-95 | – | – | – | – | – | – | – |
| | 98-121 | 135-165 | 135-165 | 135-165 | 115-142 | 135-165 | 135-165 | 135-165 | 75-93 | – | – | – | – | – | – | – |
| | 97-120 | 135-165 | 135-165 | 135-165 | 117-144 | 135-165 | 135-165 | 135-165 | 77-95 | 75-100 | – | – | – | – | – | – |
| | 125-154 | 135-165 | 135-165 | 135-165 | 110-136 | 135-165 | 135-165 | 135-165 | 81-100 | 90-110 | – | – | – | – | – | – |
| | 120-148 | 135-165 | 135-165 | 135-165 | 136-167 | 135-165 | 135-165 | 135-165 | 81-100 | 94-116 | 90-110 | 94-116 | – | – | – | – |
| | 115-141 | 135-165 | 135-165 | 135-165 | 116-143 | 135-165 | 135-165 | 135-165 | 83-103 | – | – | – | – | – | – | – |
| | 110-136 | 135-165 | 135-165 | 135-165 | 101-125 | 135-165 | 135-165 | 135-165 | 83-103 | – | – | – | – | – | – | – |
| Upper mesopelagic | 290-310 | 290-310 | 290-310 | 290-310 | 290-310 | 290-310 | 290-310 | 150-200 | 150-200 | 150-200 | 150-200 | 150-200 | 150-200 | 150-200 | 150-200 | 150-200 |
| Lower mesopelagic | 790-810 | 490-510 | 490-510 | 790-810 | 490-510 | 490-510 | 490-510 | 310-330 | 310-330 | 310-330 | 310-330 | 310-330 | 310-330 | 310-330 | 310-330 | 310-330 |
| Base of mesopelagic | 1490-1510 | 1490-1510 | 1490-1510 | 1490-1510 | 1490-1510 | 1490-1510 | 1490-1510 | 870-890 | 870-890 | 870-890 | 870-890 | 870-890 | 870-890 | 870-890 | 870-890 | 870-890 |
| Upper mesopelagic | PAP-SO | | | | | | | | | | | | | | | |
| | OSP | | | | | | | | | | | | | | | |
| | HAUSGARTEN | | | | | | | | | | | | | | | |
| | z_{eu} | POC | PIC | bSI | z_{eu} | POC | PIC | bSI | z_{eu} | POC | PIC | bSI | z_{eu} | POC | PIC | bSI |
| | 82-101 | – | – | – | 71-88 | – | – | – | 60-75 | 161-197 | 161-197 | 161-197 | 161-197 | 161-197 | 161-197 | 161-197 |
| | 82-101 | – | – | – | 73-90 | 90-110 | 90-110 | 90-110 | 63-78 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 |
| | 77-95 | – | – | – | 72-89 | 90-110 | 90-110 | – | 65-81 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 |
| | 61-76 | – | – | – | 73-90 | 90-110 | 90-110 | – | 68-84 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 |
| | 57-71 | 54-66 | 67-83 | 54-66 | 69-85 | 90-110 | 90-110 | 90-110 | 37-46 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 |
| | 53-66 | – | – | – | 71-88 | 90-110 | 90-110 | 90-110 | 35-43 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 |
| | 56-70 | 61-75 | 61-75 | 48-64 | 61-75 | 45-55 | 45-55 | – | 44-55 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 |
| | 60-75 | 50-68 | 50-68 | 50-68 | 65-81 | 65-80 | 65-80 | 65-80 | 55-68 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 |
| 63-78 | – | – | – | 64-79 | 45-55 | 45-55 | 45-55 | 50-62 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | |
| 62-77 | – | – | – | 63-78 | 45-55 | 45-55 | – | 53-65 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | |
| 68-84 | – | – | – | 62-77 | – | – | – | 55-69 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | |
| 77-95 | – | – | – | 64-79 | – | – | – | 58-72 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | 72-88 | |
| Upper mesopelagic | 171-200 | 171-200 | 171-200 | 171-200 | 250-300 | 250-300 | 250-300 | 270-320 | 270-320 | 270-320 | 270-320 | 270-320 | 270-320 | 270-320 | 270-320 | 270-320 |
| Lower mesopelagic | 446-500 | 446-500 | 446-500 | 446-500 | 590-610 | 590-610 | 590-610 | 790-810 | 790-810 | 790-810 | 790-810 | 790-810 | 790-810 | 790-810 | 790-810 | 790-810 |
| Base of mesopelagic | 990-1010 | 990-1010 | 990-1010 | 990-1010 | 990-1010 | 990-1010 | 990-1010 | 1220-1260 | 1220-1260 | 1220-1260 | 1220-1260 | 1220-1260 | 1220-1260 | 1220-1260 | 1220-1260 | 1220-1260 |

Table S3. Comparison of modelled and observed particle attributes of six primary particle types seeded in SLAMS-2.0: diatoms, coccolithophores, large flagellated phytoplankton, picophytoplankton, transparent exopolymer particles (TEPs), and aeolian-borne clays. Observational data are compiled from 32 references in Rufas (2021). Where observed diameters were unavailable, they were estimated from volume assuming spherical shape. Modelled sinking velocities were calculated using a global mean seawater density (1.0262 g cm^{-3}) and dynamic viscosity ($0.0134 \text{ g cm}^{-1} \text{ s}^{-1}$) for the upper 200 m, based on World Ocean Atlas 2023 data.

| Primary particle | Density (g cm^{-3}) | Volume (μm^3) | | Diameter (μm) | | Sinking velocity (m d^{-1}) | |
|------------------|--------------------------------|----------------------------|--------------|----------------------------|---------------|--|---------------|
| | | Model | Obs. | Model | Obs. | Model | Obs. |
| Diatom | 1.21 | 256 – 127, 940 | $100 - 10^7$ | 7.9 – 63 | 4.0 – 200 | 0.40 – 25 | $0.30 - 35^c$ |
| Cocco. | 1.84 | 32 – 197 | $25 - 10^4$ | 3.9 – 7.2 | 3.0 – 14 | 0.44 – 1.5 | $0.15 - 10^d$ |
| Flagel. | 1.06 | 8, 393 – 209, 840 | $200 - 10^7$ | 25 – 74 | 10 – 200 | 0.76 – 6.5 | |
| Pico. | 1.06 | 0.042 – 2.1 | 0.10 – 25 | 0.43 – 1.6 | 0.60 – 4.0 | 0.0002 – 0.0030 | |
| TEP | 0.80 | 158 – 15, 804 | | 6.7 – 31 | $1 - 100^a$ | –(0.36 – 7.7) | |
| Clay | 2.70 | 14 | | 3.0 | $0.10 - 20^b$ | 0.54 | |

^a (Passow, 2002). ^b (Schulz et al., 2012). ^c (Miklasz & Denny, 2010). ^d (Jansen et al., 2002).