

# 1 Martian dust storm impact on atmospheric water and D/H observed by

## 2 ExoMars Trace Gas Orbiter

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39 Global dust storms on Mars are rare events<sup>1,2</sup>, but can affect the Martian atmosphere for several  
40 months. They can cause changes in atmospheric dynamics and an inflation of the atmosphere<sup>3</sup>,  
41 primarily due to solar heating of the dust<sup>3</sup>. In turn, changes in atmospheric dynamics can affect the  
42 distribution of atmospheric water vapour with potential implications for atmospheric  
43 photochemistry and climate on Mars<sup>4</sup>. Recent observations of the Martian atmospheric water  
44 vapour abundance during dust storm conditions revealed a high-altitude increase in atmospheric  
45 water vapour that was more pronounced at high northern latitudes<sup>5,6</sup>, and a decrease in the water  
46 column at low latitudes<sup>7,8</sup>. Here we present concurrent, high-resolution measurements of dust, H<sub>2</sub>O  
47 and HDO at the onset of a global dust storm from the ExoMars Trace Gas orbiter NOMAD and ACS  
48 instruments. We report on the vertical distribution of the HDO/H<sub>2</sub>O ratio from the planetary  
49 boundary layer up to 80 km altitude. Our findings suggest that before the onset of the dust storm,  
50 HDO abundances are reduced to below detectability above 40 km altitude. The decrease in HDO

51 coincides with the observed presence of water-ice clouds. During the storm, an increase in the  
52 abundance of H<sub>2</sub>O and HDO is observed above 40 km and up to 80 km altitude. We propose that the  
53 increased abundances may be the result of warmer temperatures during the dust storm, causing a  
54 stronger atmospheric circulation and preventing ice clouds formation, which may confine water  
55 vapour to lower altitudes as a result of gravitational fall and subsequent sublimation of ice crystals<sup>3</sup>.  
56 The observed changes in H<sub>2</sub>O and HDO abundance occurred within a few days during dust storm  
57 development, suggesting a fast impact of dust storms on the atmospheric environment.

58

59

60 Although dust is ubiquitous in Mars' atmosphere, global-scale dust storms (GDS) are relatively rare  
61 events<sup>1,2</sup> which only occurred twice in the last 17 years (in 2001 and 2007). The physical processes  
62 responsible for these phenomena are not yet fully understood, although several mechanisms have  
63 been proposed<sup>3</sup>. The ExoMars Trace Gas Orbiter (TGO) arrived at Mars in October 2016 and started  
64 its first science observations in April 2018, just before the beginning of the 2018 global dust storm.  
65 The NOMAD and ACS instruments on board TGO witnessed the onset and development of this GDS  
66 and its impact on water vapour abundance in the Martian atmosphere.

67

68 The 2018 GDS started on 30 May near the northern autumn equinox ( $L_s \sim 185^\circ$ ) and within a few  
69 weeks, the planet was covered with atmospheric dust. Instruments on other Mars orbiting and  
70 landed spacecraft also witnessed the storm's evolution (e.g. PFS and VMC<sup>9</sup> on board Mars Express,  
71 MARCI and MCS<sup>10</sup> on Mars Reconnaissance Orbiter and THEMIS<sup>11</sup> on Mars Odyssey). Observations by  
72 Curiosity<sup>12</sup> in Gale Crater indicated that the dust opacity rose from 0.65 on 7 June to 6.7 on 24 June,  
73 consistent with the values found by NOMAD and ACS which observed dust opacity to increase by a  
74 factor larger than 10 (see Methods).

75

76 TGO has a 2-hour orbit and can perform atmospheric measurements during two solar occultation  
77 events per orbit when the geometry is favourable. NOMAD and ACS measure the solar radiation  
78 spectrum that is filtered by the atmosphere and from which the vertical distribution of atmospheric  
79 compounds, in particular, water vapour (both isotopologues, H<sub>2</sub>O and HDO), can be retrieved.  
80 Atmospheric opacity variation with altitude can also be obtained directly from the decrease in the  
81 continuum part of the transmitted solar intensity, thus allowing the instruments to monitor the  
82 onset and further evolution of the GDS (Figure 1).

83

84 In solar occultation mode, while the TGO-to-Sun line of sight sweeps tangent altitudes above the top  
85 of the atmosphere, the sampled line-of-sight optical depth is zero (*i.e.*, no attenuation of the solar  
86 signal). When the line of sight to the Sun transects the atmosphere, the line-of-sight optical depth  
87 gradually increases, owing to the presence of dust and ice particles, until the atmosphere becomes  
88 completely opaque at some tangent altitude. Here the transmittance drops to zero, which usually  
89 occurs due to enhanced dust presence in the lowermost part of the atmosphere or, in rarer cases, by  
90 the planetary surface. Dust and/or cloud layers in the atmosphere cause local increases in optical  
91 depth, with the effect being most pronounced in the equatorial region (Figure 1.D-F). The  
92 characteristics of the individual vertical profiles of optical depth vary with latitude, before, during  
93 and after the dust storm.

94

95 The observations in Figure 1.A-C, made north of 60° latitude, indicate that the continuum line-of-  
96 sight optical depth remains low down to 10-20 km tangent altitude throughout the dust storm. The  
97 apparent increase with time of the tangent altitude at which the atmosphere becomes opaque is  
98 mainly a latitude effect, indicating that the GDS does not impact much the northern latitudes. Some  
99 features at 25-40 km altitude were observed from June onwards that were not present before the  
100 dust storm: these could be layers of dust that are transported from lower latitudes.

101

102 In the mid-latitudes (Figure 1.D-F), before the dust storm, many layers were observed around 40 km.  
103 Detached dust layers were previously identified on many occasions<sup>13-15</sup>, and their existence has been  
104 explained by uplifting during strong convection processes<sup>16-19</sup>. Water ice clouds may be responsible  
105 for some of the observed layers, as indicated by observations at other wavelengths and by previous  
106 investigations<sup>20-22</sup>. The layers disappear during the dust storm, when the atmosphere is utterly  
107 opaque below 40 km because of high dust abundances, and water ice clouds are expected to  
108 disappear due to the atmospheric warming in the dust storm<sup>13</sup>.

109

110 Furthermore, Figure 1.G-I shows the impact of dust/ice clouds in the high southern latitudes, from  
111 the beginning of southern spring to the onset of the dusty southern summer season. During the GDS,  
112 dust ascended to higher altitudes, comparable to the situation in the mid-latitudes but with more  
113 local variability.

114

115 On Mars, water vapour has a wide variety of effects on atmospheric photochemistry and climate. Its  
116 dissociation by sunlight into hydroxyl radicals controls the overall stability cycle of CO<sub>2</sub>. As frost on  
117 the surface or as ice clouds in the atmosphere, water exerts a strong influence, leading to large  
118 departures from the otherwise dust-controlled radiative balance<sup>3</sup>.

119

120 Here we present the first water vapour profiles at high vertical resolution (~1 km) that reach down  
121 to the planetary boundary layer and extend up to ~80 km (Figure 2 and Figure 3.A). Besides, for the  
122 first time, the vertical profile of HDO could be measured (Figure 3.B). The first observations from  
123 TGO were carried out prior to the 2018A GDS, and the impact of the global dust storm on the  
124 vertical distribution of water vapour and HDO could be monitored. The ACS observations shown in  
125 the figures were performed at high southern and northern latitudes, while the NOMAD profiles were  
126 obtained in northern mid-latitudes. During the northern autumnal season, when these  
127 measurements were carried out, previous column-integrated measurements<sup>7,8</sup> indicated a dry

128 atmosphere at high latitudes caused by the developing seasonal polar cap in the North and its  
129 receding counterpart in the South. Temperatures are cold enough to condense out CO<sub>2</sub> and as a  
130 consequence also water vapour will be removed from the lower atmosphere by condensation into  
131 clouds and subsequent sedimentation<sup>3</sup>. The seasonal cap development therefore explains the very  
132 low water abundances in the lowest 20 km for the sub-polar profiles (Figure 2). The profiles  
133 observed before the dust storm indicate low abundances of water vapour above 60 km, with values  
134 below 10 ppm, and with large error bars. Profiles from the southern hemisphere are shown in Figure  
135 2.B and Figure 3.A; they correspond to the southern summer season with a lot of dust present in the  
136 atmosphere already before the GDS, explaining the cut-off at ~15 km. Northern hemisphere profiles  
137 were taken in more dust free conditions and reach down to ~4 km.

138

139 Water profiles, both H<sub>2</sub>O and HDO, show a large enhancement in the middle atmosphere after the  
140 onset of the dust storm. The increase in water abundance is observed above 20 km, with water  
141 vapour being lifted upwards up to at least 80 km. Previous studies have reported a sharp decrease of  
142 the total water column in the equatorial region<sup>7,8</sup>, indicative of redistribution of water vapour in a  
143 dust storm. Previous measurements<sup>5</sup> of water vapour profiles already exhibited an increase of the  
144 atmospheric water content at high altitudes and latitudes, as is confirmed by these new data. This  
145 phenomenon was also linked to an increase in the escape of hydrogen from Mars' atmosphere<sup>6,23</sup>.  
146 What is remarkable in the observations presented here (Figure 2), is that this enhancement is  
147 happening very fast, in the course of just a few days during the onset of the dust storm (around 7-8  
148 June, Ls~188-190°).

149

150 The observed changes in the distribution of atmospheric water reported here can be understood as  
151 resulting from a variety of processes. The higher abundance of dust heats large parts of the  
152 atmosphere because of the absorption of solar radiation by the dust particles. Dust absorption and  
153 the subsequent warming of the surrounding gas causes an expansion of the atmosphere, which

154 leads to a redistribution of water vapour to a wider vertical range. The higher atmospheric  
155 temperatures at low and middle latitudes and the resulting higher thermal contrast between the  
156 equatorial and polar regions also strengthen the mean meridional circulation, this leads to an  
157 additional redistribution of water vapour across latitude. Also because of the higher temperatures,  
158 fewer water ice clouds are expected to be present during a dust storm. Under non-dust storm  
159 conditions, the formation of clouds acts to confine water vapour to lower altitudes due to the  
160 gravitational fall and subsequent sublimation of ice crystals. In addition, numerical modelling has  
161 also demonstrated that solar heating of atmospheric dust can drive localized deep convection<sup>17,18</sup>  
162 and larger scale ascent of dust layers<sup>19</sup> that would, along with the dust, also transport water vapour  
163 to higher altitudes. All these processes that contribute to explaining the observed changes in the  
164 water vapour profiles have been quantitatively demonstrated with global circulation models and by  
165 data assimilation of water vapour in previous years on Mars<sup>3,24-26</sup>. For a more quantitative  
166 understanding of the 2018 GDS, more detailed modelling and assimilation studies that simulate the  
167 transition from normal to global dust storm conditions<sup>3,24,25,27</sup> will have to be performed, using dust  
168 constraints derived from instruments that monitored the GDS, including TGO instruments.

169

170 The fractionation between H<sub>2</sub>O and HDO is an important process in planetary atmospheres. The D/H  
171 ratio is a marker of the evolution of the water inventory on Mars<sup>28</sup>. On Mars, the D/H budget is  
172 dominated by H<sub>2</sub>O and HDO which are the unique precursors of the escaping D and H atoms above  
173 the exobase. HDO was previously measured as column-integrated abundances from Earth<sup>29-31</sup> and *in*  
174 *situ*<sup>32</sup> by the Mars Science Laboratory. NOMAD and ACS provide for the first time the capability to  
175 observe the vertical distribution of HDO simultaneously with water vapour, thereby providing key  
176 information on the fractionation processes that are expected to control the amount of hydrogen and  
177 deuterium atoms escaping to space<sup>3</sup>. H<sub>2</sub>O and HDO are fractionated during photolysis and ice  
178 formation<sup>33</sup>. NOMAD observations (Figure 3) reveal that the HDO density profiles during the pre-  
179 storm period exhibit a sudden decline at 40-45 km altitude, just below a layer of water ice clouds,

180 consistent with this view (see Methods). ACS observations show this decrease to occur at 50 km but  
181 were taken at a different latitude, where the hygropause may be located at a different altitude. The  
182 HDO/H<sub>2</sub>O ratio is similar in both profiles below 45 km: 4-6 VSMOW (Vienna Standard Mean Ocean  
183 Water). HDO is distinctly more abundant at high altitudes during the dust storm than before the  
184 storm. We suggest that this might be the result of the strong atmospheric warming during the GDS,  
185 which causes the hygropause to ascend to higher altitudes. The HDO/H<sub>2</sub>O ratio is relatively similar  
186 before and during the GDS, which demonstrates that HDO is advected along with H<sub>2</sub>O to higher  
187 altitudes and latitudes during the onset of the GDS.

188

189 Fractionation of HDO and H<sub>2</sub>O is expected to occur during condensation of both isotopologues<sup>33</sup>.  
190 Model studies predict that by condensation and cloud sedimentation, the abundances of H<sub>2</sub>O and  
191 HDO become very small above the hygropause<sup>4,33</sup>. This is confirmed by our first results (Figure 3).  
192 The coincidence of a strong decrease in HDO, below detectability, with an observed water ice cloud,  
193 may provide the first direct indication of the fractionation process. Because of the fractionation  
194 effect, HDO is expected to be relatively more concentrated into the ice cloud and thus more  
195 scavenged from the condensation level, such that the HDO/H<sub>2</sub>O ratio is expected to decrease with  
196 altitude. The effect of the dust storm is to expand the atmosphere and to lift up the hygropause.  
197 Continued measurements by TGO shall permit us to unveil both the spatial and the seasonal trends  
198 of D/H.

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283

284 **Competing interest**

285 The authors declare no competing financial interests.

286

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306

307 **Author contributions**

308 ACV and OK conceived the study, collected inputs and wrote the paper. SA, GV, GL retrieved trace  
309 gas abundances, in particular those of H<sub>2</sub>O and HDO from the NOMAD instrument. IRT analysed the

310 SO solar occultation data. LT provided transmittances from the NOMAD SO v0.3a. JTE and SR  
311 provided and analysed the a priori knowledge and initial GCM fields. FD and LN provided the GCM  
312 fields. SV, FGG, FL, SL, and JK provided the global circulation model background and discussion. FA,  
313 OzK and VW are coordinating the dust observations between IR and UVIS, and nadir and occultation.  
314 MLV, JCG, MGC, MLP, BF analysed the NOMAD limb data. MLP provided the dust profiles from  
315 NOMAD IR data. MS, TC, and MW provided contextual information from Themis/MO instrument.  
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319 the calibration of the IR channels. CD, DB and YW were involved in UVIS calibration and data  
320 pipeline. BR and EN designed the NOMAD observations helped by JM for the UVIS channel. AAF  
321 calibrated the ACS NIR data and analysed the water profiles assisted with FM, AT, DBet, and JLB. CO<sub>2</sub>  
322 data was analysed by DBel The ACS datasets were prepared for NIR by AT and AP and for TIRVIM by  
323 NI, AS, IM, AT, and AG have designed the ACS observations. ML and DP have analysed the TIRVIM  
324 occultation profiles. KO, JAP, and LB provided support for the water retrieval. YI helped in the MIR  
325 calibration. MRP, GB and JLM provided support in the selection of the NOMAD observations based  
326 on scientific interest. FF, CW, DR, JLV and HS Coordinated the observations of the various  
327 instruments on the TGO. All authors assisted ACV and OK with the preparation of the manuscript.

328

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331

### 332 **Additional information**

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334

335 Figure captions

336

337 Figure 1: Evolution of the dust/cloud extinction obtained by the NOMAD SO channel during the onset of the global dust  
338 storm: from the first observations in April and May (left panels) to the August-September 2018 timeframe (right panels),  
339 spanning  $L_s = 163^\circ$  to  $246^\circ$  (late northern summer to autumn). The data, consisting of 536 individual measurements, is  
340 split into 3 latitude bins, with the colour of the line indicating the latitude within each bin. The latitudinal coverage is  
341 dependent on the orbit and solar position, and so the latitude ranges were selected based on the data available:  
342 northern profiles for latitude  $> 60^\circ\text{N}$  (upper panels); mid-latitude profiles for latitudes between  $-30^\circ\text{S}$  and  $30^\circ\text{N}$  (middle  
343 panels); and southern profiles for latitudes between  $-70^\circ\text{S}$  and  $-50^\circ\text{S}$  (lower panels). In the early phase of the TGO  
344 mission, more solar occultations occurred near the northern pole, as is evident in the figures. Plotted here is the  
345 continuum line-of-sight optical depth versus tangent altitude of the centre of the line of sight above the Mars reference  
346 areoid. The line-of-sight optical depth is inferred from the transmittance after the removal of atmospheric absorption  
347 lines. Diffraction order 121 was used for this study, covering the  $2720\text{-}2740\text{ cm}^{-1}$  spectral range. Horizontal error bars are  
348 not shown here, as they are very small: for an SNR of 1000, the 1-sigma error is 0.003 units for an optical depth of 1, and  
349 0.06 units for an optical depth of 4.

350

351 Figure 2:  $\text{H}_2\text{O}$  volume mixing ratio (vmr) profiles observed by ACS NIR, during the onset of the global dust storm. Left:  
352 northern latitudes; black:  $L_s = 188.28^\circ$  – Lat =  $77.5^\circ\text{N}$ ; blue:  $L_s = 188.75^\circ$  – Lat =  $76.4^\circ\text{N}$ ; green:  $L_s = 189.41^\circ$  – Lat =  $74.8^\circ$   
353 N; yellow:  $L_s = 189.90^\circ$  – Lat =  $73.8^\circ\text{N}$ . Right: southern latitudes; blue:  $L_s = 188.62^\circ$  – Lat =  $68.2^\circ\text{S}$ ; cyan:  $L_s = 189.19^\circ$  –  
354 Lat =  $70.0^\circ\text{S}$ ; yellow:  $L_s = 189.67^\circ$  – Lat =  $71.3^\circ\text{S}$ ; orange:  $L_s = 190.05^\circ$  – Lat =  $72.4^\circ\text{S}$ ; red:  $L_s = 190.50^\circ$  – Lat =  $73.8^\circ\text{S}$ .  
355 Water abundances were deduced from ACS NIR observations (order 56 covering the  $1.38\text{ }\mu\text{m}$  band,  $7225\text{-}7300\text{ cm}^{-1}$ ); the  
356  $\text{CO}_2$  density was measured in order 49,  $6320\text{-}6390\text{ cm}^{-1}$ ). The uncertainty in the local number densities is given by the  
357 covariance matrix of the solution error (as described in the text).

358

359 Figure 3:  $\text{H}_2\text{O}$ , HDO and D/H detections before and during the storm. Panel A: NOMAD  $\text{H}_2\text{O}$  observations before the  
360 storm (blue:  $L_s = 171.45^\circ$ , Lat:  $43^\circ\text{N}$  to  $68^\circ\text{N}$ ), and during the storm (red:  $L_s = 196.64^\circ$ , Lat:  $51^\circ\text{N}$  to  $59^\circ\text{N}$ ), and ACS MIR  
361 observations before the storm (cyan:  $L_s = 168.75^\circ$ , Lat:  $39^\circ\text{S}$  to  $43^\circ\text{S}$ ), and during the storm (yellow:  $L_s = 196.64^\circ$ , Lat:  $80^\circ\text{S}$   
362 to  $83^\circ\text{S}$ ). The corresponding HDO VMR profiles are shown in Panel B. Panel C shows the D/H ratio obtained for each of  
363 the  $\text{H}_2\text{O}$ -HDO observations. All errors in Panel A–C are  $1\sigma$ . VSMOW is the Vienna Standard Mean Ocean Water reference  
364 value, 312 ppm HDO/ $\text{H}_2\text{O}$ .

365 **Methods**

366 **The NOMAD instrument and dataset.** NOMAD, the “Nadir and Occultation for MArS Discovery”  
367 spectrometer suite<sup>34-36</sup>, is part of the payload of the ExoMars 2016 Trace Gas Orbiter mission<sup>37</sup>. The  
368 instrument is conducting a spectroscopic survey of Mars’ atmosphere in ultraviolet (UV), visible and  
369 infrared (IR) wavelengths covering large parts of the 0.2-4.3  $\mu\text{m}$  spectral range. NOMAD is composed  
370 of three spectrometers: a solar occultation only spectrometer (SO – Solar Occultation) operating in  
371 the infrared (2.3-4.3  $\mu\text{m}$ ), a second infrared spectrometer (2.3-3.8  $\mu\text{m}$ ) capable of nadir, but also  
372 solar occultation and limb observations (LNO – Limb Nadir and solar Occultation), and an  
373 ultraviolet/visible spectrometer (UVIS – UV visible, 200-650 nm) that also has all three observation  
374 modes. The spectral resolution of SO ( $0.15 \text{ cm}^{-1}$  at  $3000 \text{ cm}^{-1}$ ) surpasses previous surveys from orbit  
375 in the infrared by at least one order of magnitude. NOMAD offers an integrated instrument  
376 combining a flight-proven concept and innovations based on existing instrumentation: SO is a copy  
377 of the Solar Occultation in the IR (SOIR) instrument<sup>38</sup> on Venus Express (VEx<sup>39</sup>), LNO is a modified  
378 version of SOIR, and UVIS has heritage from the development in the context of the Humboldt lander.  
379 NOMAD provides vertical profiling for atmospheric constituents at unprecedented spatial and  
380 temporal resolution. Indeed, in solar occultation, the vertical resolution is less than 1 km for SO and  
381 UVIS, with a sampling rate of 1 s (one measurement every 1 km), and occultations range from the  
382 surface to 200 km altitude. NOMAD also provides mapping of several constituents in nadir mode  
383 with an instantaneous footprint of  $0.5 \times 17 \text{ km}^2$  (LNO spectrometer) and  $5 \text{ km}^2$  (UVIS spectrometer),  
384 with a repetition rate of 30 Martian days.

385

386 For this work we analysed SO channel data measured between 21 April and 30 September. SO  
387 measures 4 spectra for 5 or 6 different diffraction orders per second in solar occultation mode.

388

389 **The ACS instrument and dataset.** ACS<sup>40</sup> consists of three infrared channels featuring high accuracy,  
390 high resolving power, and a broad spectral coverage (0.7 to 17  $\mu\text{m}$ ). The near-infrared (NIR) channel

391 is based on the principle of echelle-spectrometer with selection of diffraction orders by an acousto-  
392 optical tuneable filter (AOTF). The same principle was employed by SOIR on VEx<sup>4</sup> and by the infrared  
393 channels of NOMAD described above. ACS NIR covers a spectral range of 0.7-1.7  $\mu\text{m}$  in diffraction  
394 orders 101 through 49. The instrument capitalises at the science heritage of SPICAM-IR<sup>41</sup> on board  
395 ESA's Mars Express, benefiting from much higher resolving power of  $\lambda/\Delta\lambda\approx 25,000$ . During an  
396 occultation ACS NIR measures 10 preselected diffraction orders in two seconds, including the  
397 absorption bands of H<sub>2</sub>O at 1.13, 1.38, and 1.40  $\mu\text{m}$ , and CO<sub>2</sub> at 1.27, 1.43, 1.54, and 1.57  $\mu\text{m}$ . The  
398 mid-infrared (MIR) channel is a newly developed crossed dispersion echelle spectrometer dedicated  
399 to solar occultation measurements in the 2.3-4.5  $\mu\text{m}$  range. The spectral resolving power is  
400  $\lambda/\Delta\lambda\approx 50,000$ . For each acquired frame MIR measures up to 20 adjacent diffraction orders, covering  
401 an instantaneous spectral range of 0.15-0.3  $\mu\text{m}$ . To achieve the full spectral coverage a secondary  
402 dispersion grating can be rotated to one out of 11 positions. The H<sub>2</sub>O and HDO profiles can be  
403 measured simultaneously by MIR using the positions 4, 5 and 11.

404

405 The concept of Fourier-transform spectrometer TIRVIM is close to that of Planetary Fourier  
406 Spectrometer (PFS<sup>42</sup>) on board MEx, though TIRVIM features a cryogenic detector and the solar  
407 occultation capability. In occultation, TIRVIM is operated mostly in 'climatology' mode, covering  
408 instantaneously, each 0.4 s, the full spectral range of 1.7-17  $\mu\text{m}$  (effectively 1.7-5  $\mu\text{m}$ ) with spectral  
409 resolution  $\leq 1 \text{ cm}^{-1}$ . These three channels are used to observe in solar occultation; NIR and TIRVIM  
410 are operated also in nadir to measure atmospheric gases and to characterise the atmospheric state:  
411 dust loading and condensation clouds. The atmospheric temperature profile is retrieved from the  
412 15- $\mu\text{m}$  CO<sub>2</sub> band measured by TIRVIM in nadir.

413

414 In this work we used NIR occultation profiles (Figure 2) obtained at high latitudes in the southern  
415 and northern hemispheres (see ED Table 1). MIR simultaneous H<sub>2</sub>O and HDO profiles (Figure 3) were  
416 obtained in the southern hemisphere in order 224 (position 4 of secondary grating). TIRVIM aerosol

417 profiling (ED Figure 7) was done using solar occultation data obtained in the southern hemisphere,  
418 orbit 2556, Ls=197°, latitude 81°S during the egress (local time 9:26), i.e. during the same occultation  
419 of the MIR results shown in Figure 3.

420

421 **Solar occultation technique.** The solar occultation technique is a powerful method to gain  
422 information on the vertical structure of atmospheres. At sunset, the recording of spectra starts well  
423 before the occultation occurs (the solar spectrum outside the atmosphere is used for referencing),  
424 and continues until the line of sight crosses the planet. At sunrise, the recording of spectra continues  
425 well above the atmosphere to provide the corresponding reference. Transmittances are obtained by  
426 dividing the spectra measured through the atmosphere by the reference spectrum recorded outside  
427 the atmosphere<sup>43</sup>. In this way, transmittances become independent of instrumental characteristics,  
428 such as the absolute response or the ageing of the instrument and in particular of the detector. Such  
429 observations provide high vertical resolution (< 1 km for NOMAD SO and ACS NIR and 2.0-2.5 km for  
430 ACS MIR observations) profiles of the structure and composition of the atmosphere. ACS TIRVIM  
431 observes the full Sun disk during an occultation, resulting in a coarser vertical resolution (~9 km).

432

433 **Profiles of dust extinction.** To calculate the extinction due to dust and/or clouds, it is necessary to  
434 remove the absorption lines of atmospheric gas species, leaving the background continuum. For the  
435 analysis here, diffraction order 121 of NOMAD SO was chosen, as 1) this order is measured routinely,  
436 so has high spatial/temporal coverage and 2) it is relatively simple to remove the atmospheric  
437 absorption lines. A 4<sup>th</sup> order polynomial is fitted to the data. The optical depths in Figure 1 are  
438 inferred from the value of the continuum in the centre of the detector (pixel 160). The fitting  
439 algorithm fails at low and high altitudes, where either the absorption lines from molecular species  
440 are saturated or the signal is so low that it is effectively noise. Therefore, any spectra where  
441 transmittance > 99.5% are assumed to have an optical depth of 0, and points where transmittance <  
442 0.5% are not plotted; hence the lines end abruptly at low altitudes when the optical depth becomes

443 high. The observations in Figure 1 are split into North, South and mid-latitudes using the following  
444 criteria: greater than 60° North, between -70° and -50° South, and between -30° and +30° for the  
445 mid-latitudes. The tangent altitude is calculated as the shortest distance between the line of sight of  
446 the centre of the field of view and the MGM1025 Areoid (i.e. the Mars geoid)<sup>44</sup>. The latitude is the  
447 point on the areoid closest to the centre of the field of view, i.e. the tangent point, at the midpoint  
448 of the solar occultation measurement. The characteristics of the individual vertical profiles of optical  
449 depth vary with latitude, as seen when optical depth is plotted vs latitude and Mars longitude (ED  
450 Figure 1).

451

452 To further investigate the impact of the dust storm, two orbits covering the same footprint and solar  
453 illumination conditions on Mars have been considered; they were acquired by the nadir channel of  
454 NOMAD, respectively, before (April 26th) and during (July 11th) the global dust storm. ED Figure 2  
455 compares the dust radiance signature before and after the storm, but in contrast to Figure 1, now in  
456 a nadir geometry and in a different wavelength, at 2.3 μm. Comparison with radiative transfer  
457 modelling suggests a factor of ~10 increase in opacity at 2.3 μm during the storm. Note also how the  
458 surface albedo features are obscured by the increase of the atmospheric dust load. The radiance  
459 variation with latitude is mainly dominated by the total albedo (surface + atmosphere) and solar  
460 zenith angle, which varies along the track. The radiative transfer model includes multiple scattering  
461 and a layered atmosphere with pressure/temperature profiles from the LMD General Circulation  
462 Model<sup>45</sup>. Further details on the radiative transfer model can be found in Villanueva, et al.<sup>46</sup>.

463

464 **Vertical profiles of H<sub>2</sub>O and HDO volume mixing ratio.** The vertical profiles of H<sub>2</sub>O and HDO volume  
465 mixing ratio are investigated from the NOMAD dataset shown in ED Table 1. These NOMAD spectra  
466 are all taken in the northern hemisphere at the same local time (at 18h). The NOMAD SO channel  
467 can record spectra for multiple diffraction orders during an occultation. The occultation performed  
468 on 7 May includes the measurements of diffraction order 168 (3775.53 – 3805.63 cm<sup>-1</sup>) and order

469 136 (3056.39 – 3080.75  $\text{cm}^{-1}$ ) where strong  $\text{H}_2\text{O}$  lines are present and of order 119 (2674.34 –  
470 2695.65  $\text{cm}^{-1}$ ) with strong HDO lines. The occultation measurement on 20 June contains two  
471 diffraction orders for  $\text{H}_2\text{O}$  - order 168 and 134 (3011.44 – 3035.44  $\text{cm}^{-1}$ ), and diffraction order 121 for  
472 HDO (2719.28 – 2740.96  $\text{cm}^{-1}$ ) (ED Figure 3).

473

474 We retrieved  $\text{H}_2\text{O}$  volume mixing ratio using the whole spectral range of those diffraction orders, in  
475 order to maximize the information content at every occultation tangent altitude. In this study,  $\text{CO}_2$   
476 and  $\text{H}_2\text{O}$  gas absorptions were included. The absorption coefficients of these gases are calculated  
477 based on a line-by-line method using the water vapour line list for a  $\text{CO}_2$ -rich atmosphere for  $\text{H}_2\text{O}$ <sup>47,48</sup>  
478 and HITRAN 2016<sup>49</sup> for  $\text{CO}_2$ . Temperature, pressure, and  $\text{CO}_2$  volume mixing ratio are taken from the  
479 values predicted by GCMs for each altitude. The calculated synthetic spectra are convolved with a  
480 Gaussian function that corresponds to the spectral resolving power of the NOMAD SO channel  
481 ( $R=11000-15000$ ). The final synthetic spectra are then built by considering an instrument model that  
482 comprises the effects of the Acousto-Optic Tunable Filter (AOTF) and the grating (i.e., Blaze  
483 function)<sup>50</sup>. The free parameters in the retrievals are the vertical profiles of volume mixing ratio and  
484 the parameters for the polynomial function to model the continuum of each spectrum. Retrievals  
485 are performed using an Optimal Estimation approach<sup>51</sup> implemented in a Gauss-Newton iterative  
486 scheme. ED Figure 4 shows an example of fit results.

487

488 The water vapour profiles shown in Figure 2 are retrieved from the ACS NIR spectra (see also ED  
489 Table 1). Wavelength drift is corrected using positions of gaseous absorption lines. The spectra  
490 fitting and the profile retrieval follow the method described for SPICAM MEx 1.38- $\mu\text{m}$  band<sup>5,52</sup>. All  
491 the altitudes of the profiles are fitted simultaneously (global fit) using a Levenberg–Marquardt  
492 iterative algorithm<sup>53,54</sup>. Tikhonov regularization is then applied, customary for vertical inversions in  
493 order to smooth the profile. The uncertainty in the local number densities is given by the covariance  
494 matrix of the solution errors. The water vapour abundances were retrieved from spectra acquired in

495 diffraction order 56 (covering the 1.38  $\mu\text{m}$  band, or 7220-7300  $\text{cm}^{-1}$ ). ED Figure 5 shows an example  
496 of fit results. The spectral line parameters for  $\text{H}_2\text{O}$  are taken from HITRAN 2016<sup>49</sup> with a correction  
497 coefficient for the  $\text{CO}_2$  rich atmosphere<sup>5</sup>. Temperature and pressure for the radiative transfer  
498 computations are taken from the MCD<sup>45</sup> global circulation model (GCM). To obtain the VMR profiles  
499 of water vapour, the  $\text{CO}_2$  density was retrieved from ACS NIR spectra in order 49, 6320-6390  $\text{cm}^{-1}$ .

500

501 The vertical profiles of  $\text{H}_2\text{O}$  and HDO volume mixing ratio investigated from the ACS dataset (Figure  
502 3.B and C) were obtained in the southern hemisphere at middle and high latitudes (ED Table 1).  
503 During these observations MIR channels recorded spectra at position 4 (diffraction orders 210-224).  
504 To obtain the  $\text{H}_2\text{O}$  and HDO density the order 224 (3763-3775  $\text{cm}^{-1}$ ) was used for both observations  
505 in 2<sup>th</sup> May and 20<sup>th</sup> June. We retrieved  $\text{H}_2\text{O}$  and HDO volume mixing ratio using several lines present  
506 in this diffraction order. The spectral line parameters for  $\text{H}_2\text{O}$  are taken from HITRAN 2016 with a  
507 correction coefficient for the  $\text{CO}_2$  broadening<sup>5</sup>. Temperature and pressure for radiative transfer  
508 computations are taken from GCM MCD<sup>45</sup>. The calculated synthetic spectra are convolved with a  
509 Gaussian function that corresponds to the spectral resolving power of the ACS MIR channel  
510 ( $R \sim 30000$ -35000).

511

512 For both NOMAD and ACS, the D/H values are obtained by dividing the HDO and  $\text{H}_2\text{O}$  densities  
513 retrieved at the same altitude. The D/H errors are calculated by weighted average:  
514  $D/H_{\text{Error}} = [D/H] * \text{sqrt} ( (HDO_{\text{error}}/[HDO])^2 + (H2O_{\text{error}}/[H2O])^2 )$ , where [] are used to  
515 indicate the density. Because the  $\text{H}_2\text{O}$  and HDO are retrieved independently and coming from  
516 different datasets (here different orders and different absorption lines), they are truly independent.  
517 The error-bars are therefore non-correlated. The error-bars capture the instrumental systematic  
518 effects, since they are based on the final chi-square

519

520 **Water ice clouds.** ED Figure 6 shows the aerosol optical depth derived from NOMAD SO on May 7th,  
521 2018 before the dust storm, during which orders 119, 136, 148, 168 and 189 were measured,  
522 corresponding to the central wavenumbers 2685.0, 3068.0, 3339.0, 3790.0, and 4265.0  $\text{cm}^{-1}$   
523 respectively. The optical depths have been derived by averaging the transmittances with a sampling  
524 of 3 km, and deriving for each tangent height the equivalent optical depth rescaled by the  
525 occultation path. Each optical depth has been determined simultaneously with the abundances of  
526 the other gases detectable in each order; hence a full retrieval is used. Each spectrum has been  
527 processed with the Planetary Spectrum Generator (PSG<sup>46</sup>) forward model and a retrieval scheme  
528 based on Optimal Estimation in a Gauss-Newton iterative scheme. Optical depth is derived for each  
529 single sampled altitude, and is compared to the extinction of water ice with different particle sizes  
530 (Top panel of ED Figure 6). This Figure shows that the detached layer observed by NOMAD at 40-50  
531 km is a water ice cloud with particle sizes between 0.1 and 1  $\mu\text{m}$ .

532

533 Aerosol properties from TIRVIM solar occultation data were retrieved from 20 wavenumbers in the  
534 spectral range of 1500–4500  $\text{cm}^{-1}$  chosen outside of strong gas absorption bands. This channel is  
535 operated continuously, and therefore remains very stable during an occultation. To increase signal-  
536 to-noise ratio, spectra are averaged using simple moving mean within 50  $\text{cm}^{-1}$  spectral window  
537 centred at the chosen wavenumbers, spectra errors were calculated as s.e.m with  $n = 78$ . The  
538 procedure to obtain transmittances from the TIRVIM dataset is straightforward. This channel is  
539 operated continuously, and therefore remains very stable during an occultation. Slant optical depth  
540 is calculated as  $\tau_v(L) = -\ln(T(L))$ , where  $T$  is the transmittance over the line of sight  $L$ . Transmission  
541 and slant optical depth errors are calculated using propagation of uncertainty formulas. Vertical  
542 profiles of extinction and corresponding errors are retrieved using the standard ‘onion peeling’  
543 method in Fedorova, et al. <sup>55</sup>. Further steps involve Mie modelling of the spectral dependence of the  
544 extinction coefficient assuming known optical properties for the aerosols<sup>56,57</sup> fit to the experimental  
545 data to retrieve vertical profiles of the size distribution and number density as described in

546 Fedorova, et al.<sup>58</sup>. A log-normal size distribution<sup>59</sup> of the aerosol particles with a width (the effective  
547 variance) of 0.3 was assumed. To distinguish between water ice and dust particles, we apply the  
548 optimal estimation retrieval scheme independently for both types, and make the decision based on  
549 the fit quality (ED Figure 7). The algorithm is able to retrieve the number density (typically ~1  
550 particles cm<sup>-3</sup>), and the effective radius (1–1.5 μm). Errors on these values are estimated from the  
551 covariance matrix.

552

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621 **Data availability**

622 The datasets generated by the NOMAD and ACS instruments and analysed during the current study  
623 will be available in the ESA PSA repository, <https://archives.esac.esa.int/psa>, after the proprietary  
624 period. The datasets directly used in this study, and especially the data used for the figures, are  
625 available from the corresponding author upon reasonable request.

626

627 **Code availability**

628 The codes used to calculate the dust/aerosols optical depths shown in figure 1 are available upon  
629 request to the corresponding author. The code used to inverse the NOMAD and ACS spectra and  
630 derive density profiles, have been favourably compared to the Planetary Spectrum Generator (PSG)  
631 tool which can be accessed at <https://psg.gsfc.nasa.gov/> and which is part of this study. A version of  
632 the retrieval code is available at <https://psg.gsfc.nasa.gov/helpatm.php#retrieval>

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## 763 Extended Data - Figure captions

764 ED Figure 1: Continuum optical depth vs latitude and solar longitude. The colour denotes the lowest  
765 altitude at which the optical depth is less than 1.0, i.e. the lowest altitude where sunlight can still  
766 penetrate the atmosphere easily. There is a strong latitudinal dependence, where northern and  
767 southern high latitudes are relatively clear until the line of sight drops below 10-15km (blue and dark  
768 blue) – except during the  $L_s = 200^\circ - 240^\circ$  period where the global dust storm appears to have raised  
769 this altitude to 20-25km (light blue and cyan).

770

771 ED Figure 2: Impact of the dust storm on NOMAD LNO nadir observations. The calibrated radiance at  
772  $2.3 \mu\text{m}$  is shown for two orbits before (Panel a) and during (Panel b) the dust event as a function of  
773 the latitude. In red, the comparison to a radiative transfer model is also presented. The dust opacity  
774 before the global dust storm is  $\tau=0.46$  at  $3 \mu\text{m}$ , while during the event, there is an increase by at  
775 least of a factor 10 ( $\tau=4.6$ ). The 1-sigma error of the data is  $8.2 \times 10^{-5} \text{ W/m}^2/\text{sr}/\text{cm}^{-1}$ . Panel c shows  
776 the surface albedo, in black OMEGA albedo at  $2.33 \mu\text{m}$  (order 190), in red TES bond albedo scaled to  
777 the OMEGA one.

778

779 ED Figure 3: Atmospheric transmittances measured by NOMAD during the storm ( $L_s = 196.64^\circ - \text{Lat}$   
780  $= 51^\circ - \text{Lon} = 148^\circ\text{E}$ ) showing HDO absorption features (arrows) appearing at tangent heights up to  
781 50 km; most of the other absorption features originate from  $\text{CO}_2$ .

782

783 ED Figure 4: Example of results of the  $\text{H}_2\text{O}$  retrieval from NOMAD. Top panel: black: transmittance  
784 measured at the tangent height of 22.2 km; blue: best fit; cyan and green: different simulations with  
785 1 ppm and 100 ppm water respectively. Bottom panel: residuals between the observation and the  
786 best fit. The error for transmittance were calculated from 1-sigma noise value.

787

788 ED Figure 5: Example of results of the  $\text{H}_2\text{O}$  retrieval from ACS NIR. Top panel: black: transmittance  
789 measured at the tangent height of 34.1 km; blue: best fit; cyan, red, and green: different simulations  
790 with no water, 1 ppm and 50 ppm water respectively. Bottom panel: residuals between the  
791 observation and the best fit. The error for transmittance were calculated from 1-sigma noise value.

792

793 ED Figure 6: Extinction of water ice with different particle sizes (Panel A) and slant optical depth  
794 (Panel B, in units of  $\text{km}^{-1}$ ) for the solar occultation performed by NOMAD before the dust storm, on 7  
795 May between 05.40 and 05.46 UTC (local time 18h), which covers the latitude range  $44^\circ \text{N}$  to  $57^\circ \text{N}$   
796 and the longitude range  $-122.6^\circ \text{E}$  to  $-121.4^\circ \text{E}$ .

797

798 ED Figure 7: Illustration of the independent retrieval of dust and water ice from the TIRVIM dataset  
799 for a typical southern hemisphere occultation (20.06.2018; Lat  $81^\circ \text{N}$ ; Lon  $-66^\circ \text{E}$ ; egress): a selection  
800 of transmission profiles at 5 wavelengths (left), corresponding to slant opacities (top left) and  
801 extinction profiles (top right), retrieved effective radius ( $r_{\text{eff}}$  in  $\mu\text{m}$ , bottom right) and aerosol number  
802 density (in  $\text{cm}^{-3}$ , bottom right). The occultation has been performed on June 20<sup>th</sup> at orbit 2556 (local  
803 time 21:25), and covers the latitude range  $81^\circ \text{N}$  to  $82^\circ \text{N}$  and the longitude range  $-67^\circ \text{E}$  to  $-60^\circ \text{E}$ . The  
804 observation corresponds to the MIR  $\text{H}_2\text{O}$ -HDO profiles during the dust storm (yellow curves of Fig.3).  
805 The water ice and dust are well distinguished using the  $3 \mu\text{m}$  water ice absorption band  
806 (wavenumber  $3263 \text{ cm}^{-1}$  in the figure). In this case the water ice cloud is detected at 25-30 km. All  
807 reported errors are 1-sigma.

808

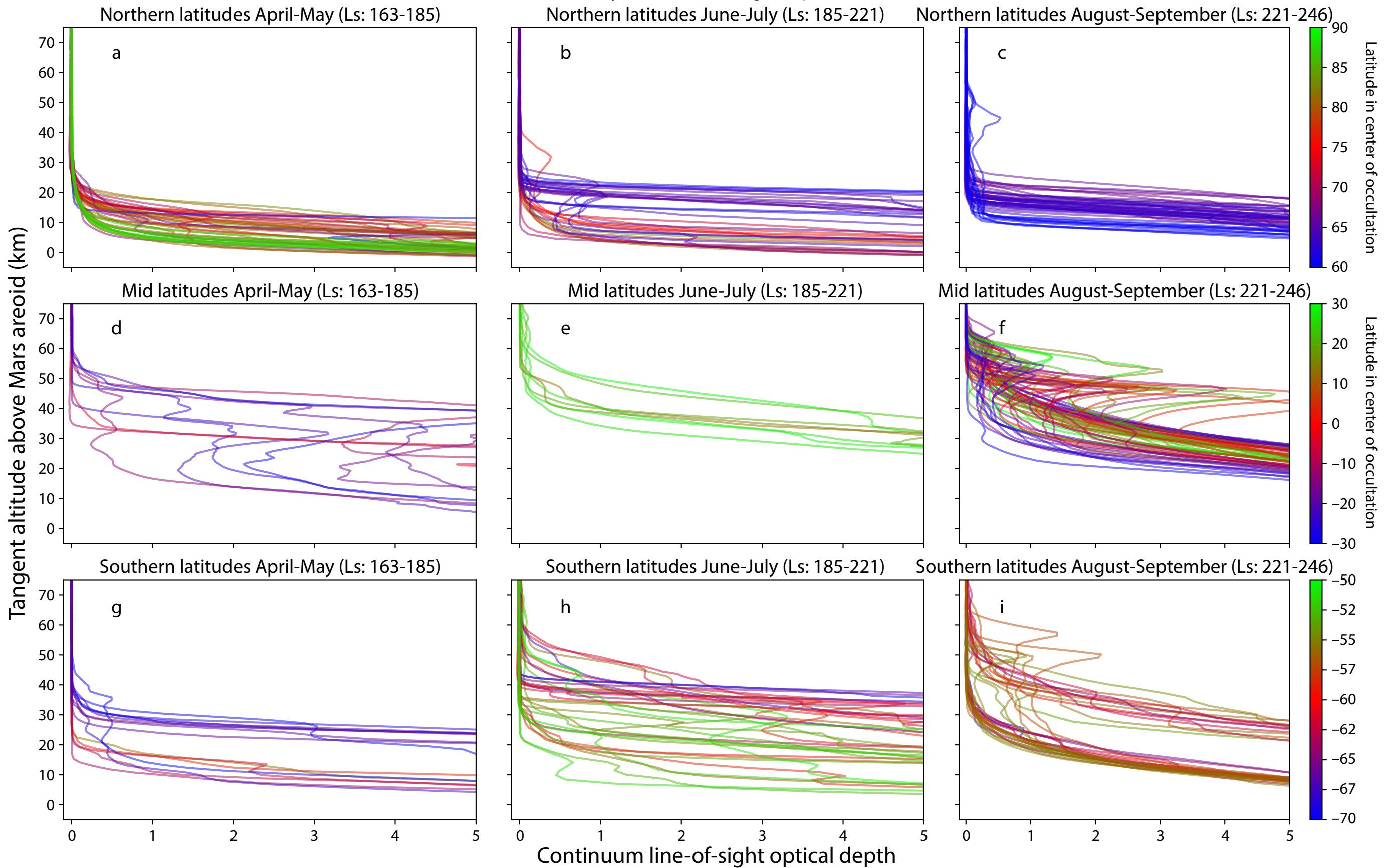
### 809 Table caption

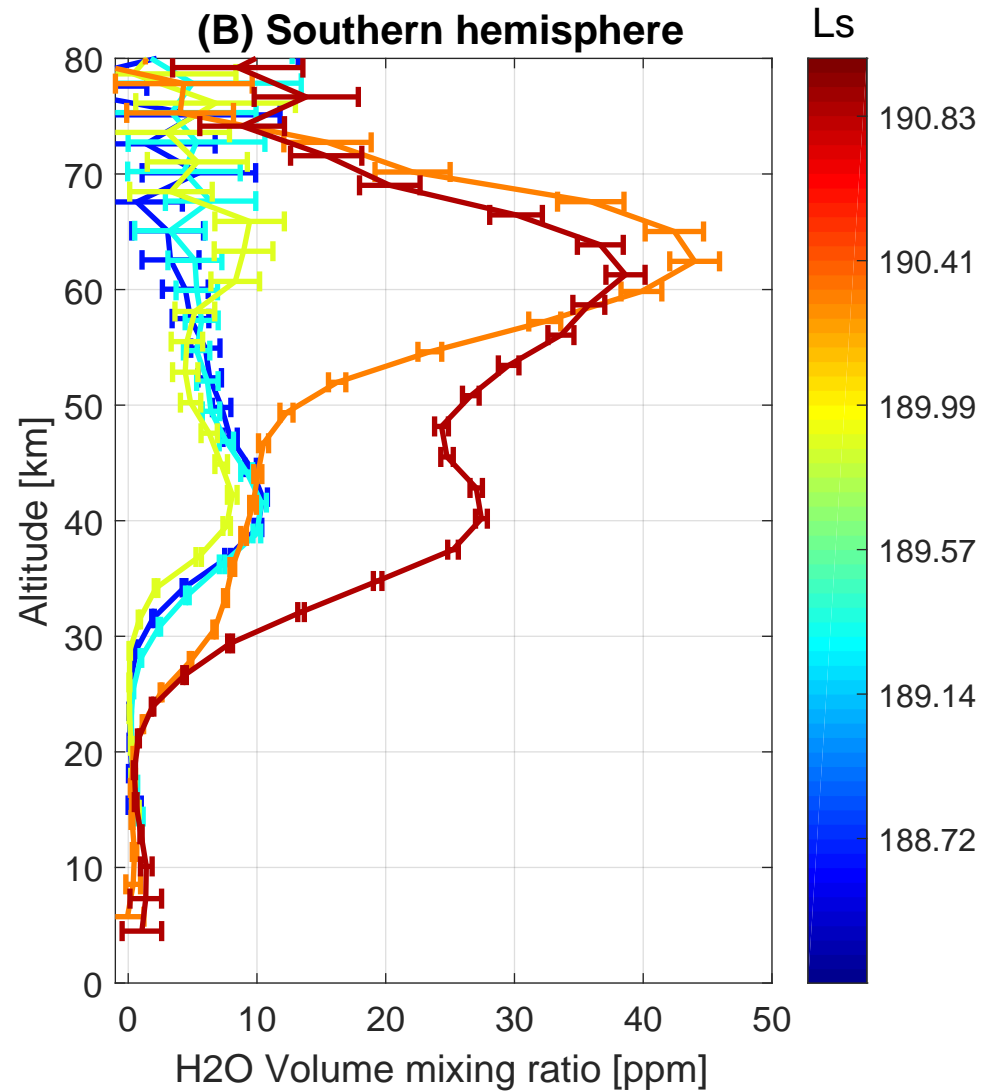
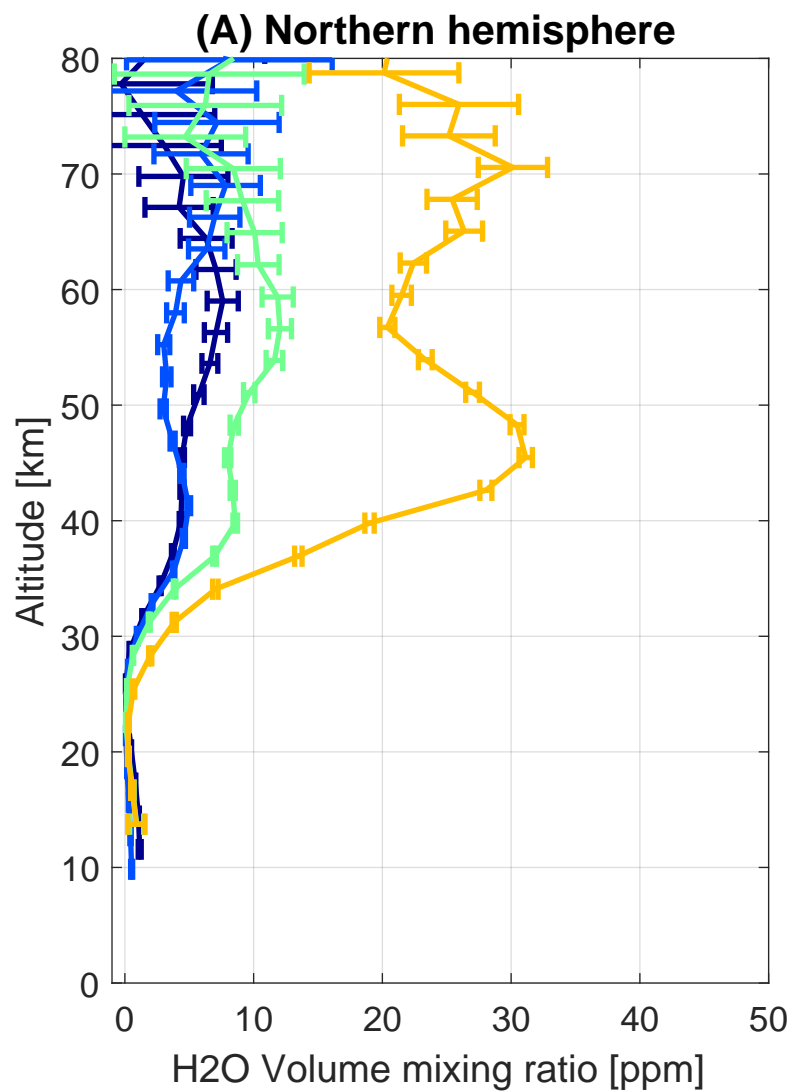
810 ED Table 1: Overview of the NOMAD and ACS observations of  $\text{H}_2\text{O}$  and HDO used and shown in this  
811 study.

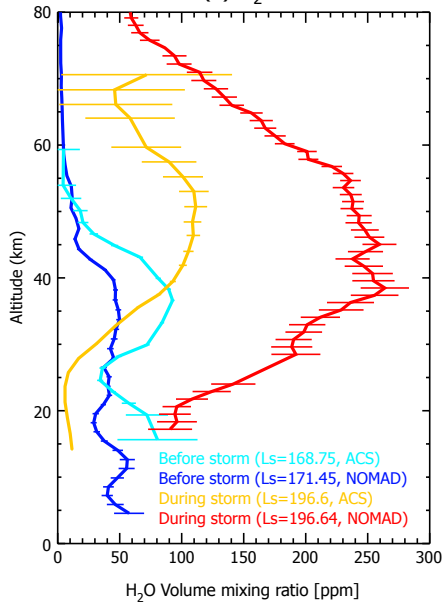
812



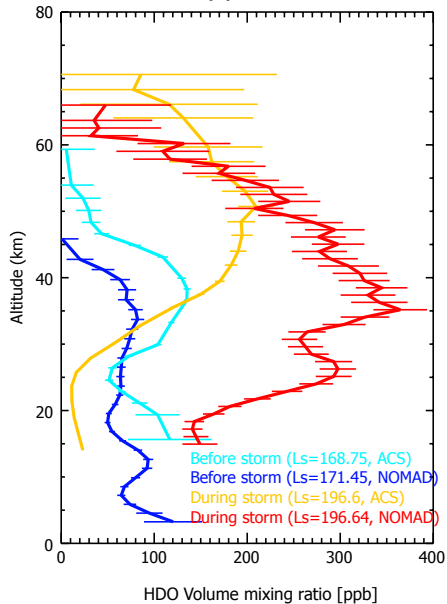
Continuum line-of-sight optical depths for SO diffraction order 121:  
Variations by time and tangent point latitude





(a) H<sub>2</sub>O

(b) HDO

(c) HDO/H<sub>2</sub>O