

Application of Fredholm integral equations inverse theory to the radial basis function approximation problem[☆]

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Abstract

This paper reveals and examines the relationship between the solution and stability of Fredholm integral equations and radial basis function approximation or interpolation. The underlying system (kernel) matrices are shown to have a smoothing property which is dependent on the choice of kernel. Instead of using the condition number to describe the ill-conditioning, hence only looking at the largest and smallest singular values of the matrix, techniques from inverse theory, particularly the Picard condition, show that it is understanding the exponential decay of the singular values which is critical for interpreting and mitigating instability. Results on the spectra of certain classes of kernel matrices are reviewed, verifying the exponential decay of the singular values. Numerical results illustrating the application of integral equation inverse theory are also provided and demonstrate that interpolation weights may be regarded as samplings of a weighted solution of an integral equation. This is then relevant for mapping from one set of radial basis function centers to another set. Techniques for the solution of integral equations can be further exploited in future studies to find stable solutions and to reduce the impact of errors in the data.

Keywords: Kernel, Radial basis function, Picard condition, first kind Fredholm integral, Spectral distribution

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1. Introduction

Radial basis functions (RBFs) are increasingly of interest for the solution of practical problems described by partial differential equations. A recent review of Flyer and Fornberg [1] demonstrates the promise of RBFs for numerically approximating sophisticated physical phenomena, particularly with the use of the novel QR-RBF algorithm [2]. Still, questions on how to effectively and efficiently determine stable solutions for RBF approximation remain. Here a new perspective on this interesting problem is presented which reemphasizes the significance of the spectral decomposition of the associated systems described by RBF interpolation/approximation.

Given the complete eigen-system $\{(\lambda_\ell, \mathbf{v}_\ell)\}_{\ell=1}^N$ for a matrix Φ , any $\mathbf{x} \in \mathbb{R}^N$ can be written as $\mathbf{x} = \sum_{\ell=1}^N \alpha_\ell \mathbf{v}_\ell$ yielding $\mathbf{y} = \Phi \mathbf{x} = \sum_{\ell=1}^N \alpha_\ell \lambda_\ell \mathbf{v}_\ell$. It is thus immediate that the contribution of \mathbf{v}_ℓ to \mathbf{y} is less than that to \mathbf{x} when $|\lambda_\ell| < 1$. Suppose now that this eigen-system is ordered with $\{|\lambda_\ell|\}_{\ell=1}^N$ arranged in a descending order, and such that this ordering also corresponds to an ordering of the basis vectors $\{\mathbf{v}_\ell\}_{\ell=1}^N$ from low to high frequency. Here, a vector is said to be a high frequency vector when, regarding the vector entries as the sampling of an underlying function, the function is highly oscillating. Then, the matrix Φ is said to have a *smoothing effect* when the vector $\mathbf{y} = \Phi \mathbf{x}$ is *smoother* than \mathbf{x} in the sense that *high frequency information* in \mathbf{x} is damped and, correspondingly, the low frequency content of \mathbf{y} is more dominant than that of \mathbf{x} . Immediately, observe that Φ has a smoothing effect only if the decay of $\{|\lambda_\ell|\}_{\ell=1}^N$ is fast and the eigenvalues of small magnitude correspond to the high frequency eigenvectors. In the context of this paper, such eigenvalues are referred to as *smoothing eigenvalues*; their distribution determines the smoothing properties of Φ .

For the forward problem the smoothing properties of Φ can be useful; for example it is often beneficial to smooth a given measured data signal \mathbf{x} by applying a *moving average technique* [3, 4], [5, chap 2]. The simplest moving average, in which each data point is replaced by the mean of its k nearest neighbours, generates a non-negative matrix Φ where all the non-zero entries are $1/k$ and

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25 each row sum is equal to 1. According to the Perron-Frobenius theorem [6, chap 2], the lowest
 26 frequency eigenvector of Φ , with all elements identical, corresponds to the largest eigenvalue
 27 (Perron root) 1. All other higher frequency eigenvectors, have at least one entry with opposite
 28 sign to the others, and correspond to smoothing eigenvalues $|\lambda| < 1$. This smoothing property is,
 29 however, a hindrance for the inverse, or data recovery, problem. In particular, suppose that Φ^{-1}
 30 exists, then it has a complete eigen-system $\{(\frac{1}{\lambda_\ell}, \mathbf{v}_\ell)\}_{\ell=1}^N$ and $\mathbf{z} = \Phi^{-1}\mathbf{x} = \sum_{\ell=1}^N \frac{\alpha_\ell}{\lambda_\ell} \mathbf{v}_\ell$. Whereas a high
 31 frequency eigenvector \mathbf{v}_ℓ is smoothed by application of Φ , it will be amplified by application of
 32 Φ^{-1} particularly when $|\lambda_\ell| \ll 1$, and any perturbation in the \mathbf{v}_ℓ component will be correspondingly
 33 amplified rather than damped when forming \mathbf{z} rather than \mathbf{y} .

34 The so-called *smoothing kernels* are typically discussed in the context of the solution of ill-
 35 posed inverse problems, and particularly in the context of the solution of first kind Fredholm
 36 integral equations [7, 8]. The focus in this paper is on the kernel matrices that arise in using RBFs
 37 for interpolation/approximation and extends results concerning the stability of the solutions of
 38 discrete first kind Fredholm integral equations to RBF approximation. For this problem, Micchelli
 39 has proved that the RBF interpolation (kernel) matrices are always invertible provided that the
 40 interpolation points are distinct [9]. But it is well recognized that the difficulty of solving systems
 41 determined by such kernel matrices, Φ , depends significantly on the condition of Φ [10–13]. Here
 42 it is emphasized how the smoothing properties of Φ , which depend on the smoothing properties of
 43 the RBFs and the location of the interpolation points, are directly related to the difficulty of solving
 44 the RBF interpolation/approximation problem. The stability of the solution found by solving the
 45 inverse problem, finding \mathbf{z} given Φ and \mathbf{x} , is investigated using the discrete Picard condition [7, 14,
 46 p.37], which is an extension of the continuous Picard condition that is pertinent for examining the
 47 stability of the solution of first kind Fredholm integral equations [8]. To show how these results are
 48 relevant, the connection between the RBF interpolation problem and the solution of an ill-posed
 49 integral equation of the first kind is now established.

50 *1.1. The connection between the radial basis function interpolation problem and integral equa-*
 51 *tions*

52 Given observational data $\{f_k : k = 1, \dots, M\}$ of a function $f(\mathbf{x})$ at scattered data points $\{\mathbf{x}_k \in$
 53 $\mathbb{R}^d : k = 1, \dots, M\}$, an approximating function $s(\mathbf{x}) \approx f(\mathbf{x})$ is defined by

$$54 \quad s(\mathbf{x}) = \sum_{j=1}^N \alpha_j \phi(\mathbf{x}, \mathbf{y}_j). \quad (1)$$

55 For RBFs $\phi(\mathbf{x}, \mathbf{y}) = \phi(\|\mathbf{x} - \mathbf{y}\|_2)$. Selecting weights $\{\alpha_j\}$ such that $s(\mathbf{x}_k) \approx f(\mathbf{x}_k)$, $k = 1, \dots, M$
 56 defines the *approximation* problem for finding $\{\alpha_j\}$. The *interpolation* problem, $M = N$, yields the
 57 defining system of linear equations

$$58 \quad \mathbf{f} = \Phi \boldsymbol{\alpha}, \quad \Phi_{kj} = \phi(\|\mathbf{x}_k - \mathbf{y}_j\|), \quad 1 \leq k, j \leq N. \quad (2)$$

59 In matrix form

$$60 \quad \begin{pmatrix} \phi(\|\mathbf{x}_1 - \mathbf{y}_1\|_2) & \phi(\|\mathbf{x}_1 - \mathbf{y}_2\|_2) & \cdots & \phi(\|\mathbf{x}_1 - \mathbf{y}_N\|_2) \\ \phi(\|\mathbf{x}_2 - \mathbf{y}_1\|_2) & \phi(\|\mathbf{x}_2 - \mathbf{y}_2\|_2) & \cdots & \phi(\|\mathbf{x}_2 - \mathbf{y}_N\|_2) \\ \vdots & \vdots & \ddots & \vdots \\ \phi(\|\mathbf{x}_N - \mathbf{y}_1\|_2) & \phi(\|\mathbf{x}_N - \mathbf{y}_2\|_2) & \cdots & \phi(\|\mathbf{x}_N - \mathbf{y}_N\|_2) \end{pmatrix} \begin{pmatrix} \alpha_1 \\ \alpha_2 \\ \vdots \\ \alpha_N \end{pmatrix} = \begin{pmatrix} f_1 \\ f_2 \\ \vdots \\ f_N \end{pmatrix}. \quad (3)$$

61 Under the assumption $\mathbf{y}_k = \mathbf{x}_k$, $k = 1, \dots, N$, $\Phi_{kj} = \phi(\|\mathbf{x}_k - \mathbf{x}_j\|)$ and matrix Φ is symmetric.

62 Suppose that $\hat{\phi}(\boldsymbol{\omega})$ is the Fourier transformation of $\phi(\|\mathbf{x}\|)$ of dimension d , then

$$63 \quad \phi(\|\mathbf{x}\|) = \frac{1}{(2\pi)^{d/2}} \int_{\mathbb{R}^d} \hat{\phi}(\boldsymbol{\omega}) e^{i\boldsymbol{\omega}^T \mathbf{x}} d\boldsymbol{\omega}. \quad (4)$$

64 For any $\boldsymbol{\alpha} \in \mathbb{C}^N$, using the standard notation, $\bar{\alpha}$, for the complex conjugate of α , the following
 65 relation holds [15, p.67]:

$$66 \quad \boldsymbol{\alpha}^H \Phi \boldsymbol{\alpha} = \sum_{j,k=1}^N \bar{\alpha}_k \alpha_j \phi(\|\mathbf{x}_k - \mathbf{x}_j\|) = \frac{1}{(2\pi)^{d/2}} \int_{\mathbb{R}^d} \hat{\phi}(\boldsymbol{\omega}) \left| \sum_{j=1}^N \alpha_j e^{i\boldsymbol{\omega}^T \mathbf{x}_j} \right|^2 d\boldsymbol{\omega}. \quad (5)$$

68 Therefore if the Fourier transform $\hat{\phi}$ is positive, Φ is symmetric and positive definite (SPD)
 69 and thus *non-singular*. Examples of RBFs which have a positive Fourier transform include the
 70 Gaussian basis functions $\phi(\|\mathbf{x}\|_2) = e^{-\epsilon^2 \|\mathbf{x}\|_2^2}$, $\epsilon > 0$ which are globally supported [15, p.74] and

71 Wendland's functions [15, p.129] which are compactly supported. Notice the dependence of the
72 RBF on the *shape* parameter $\epsilon > 0$, explicitly denoted throughout this paper where relevant by
73 $\phi = \phi_\epsilon(\|\mathbf{x}\|_2)$. Non-singular interpolation matrices are also guaranteed under other mild con-
74 ditions, see for example Micchelli [9]. Indeed it has been demonstrated over recent years that
75 RBFs are of considerable use for the approximation of high dimensional scattered data because
76 the resulting interpolation matrices can often be shown to be non-singular [9, 15]. Despite such
77 theoretical results the resulting linear systems may be highly ill-conditioned, dependent on the
78 shape parameter defining ϕ_ϵ . Discussions and numerical investigations of the ill-conditioning can
79 be found in significant references, for example, [10–13]. In particular, such results demonstrate
80 that Φ may exhibit ill-conditioning of the kind often seen when solving inverse problems; hence
81 theory on the solution of ill-posed integral equations of the first kind, [7, 8, 16–18], can be relevant
82 in the context of the RBF interpolation problem.

83 Consider the solution of the Fredholm integral equation of the first kind

$$84 \quad \int_{\Omega} K(\mathbf{x}, \mathbf{y})\beta(\mathbf{y})d\mathbf{y} = f(\mathbf{x}), \quad (6)$$

85 from which we wish to determine the underlying source function $\beta(\mathbf{y})$ given the measurements
86 $f(\mathbf{x})$ and a kernel function $K(\mathbf{x}, \mathbf{y})$. Given the observed data vectors \mathbf{x} , \mathbf{f} , and applying numerical
87 quadrature at the nodes \mathbf{y}_j yields the discrete interpolation conditions

$$88 \quad \sum_{j=1}^N K(\mathbf{x}_k, \mathbf{y}_j)\beta(\mathbf{y}_j)w_j = f(\mathbf{x}_k), \quad k = 1, \dots, N \quad (7)$$

89 for quadrature weights w_j . Defining $\alpha(\mathbf{y}) = w(\mathbf{y})\beta(\mathbf{y})$ leads to a system of equations (2) where
90 $\Phi_{kj} = K(\mathbf{x}_k, \mathbf{y}_j) = K(\mathbf{x}_k, \mathbf{x}_j)$ when the nodes and interpolation points are the same. *The connec-*
91 *tion between the solution of the RBF approximation problem which defines the weights $\{\alpha_j\}$ and*
92 *the solution of (6) which defines the source function $\alpha(\mathbf{y})$ is now established for kernel function*
93 $K(\mathbf{x}, \mathbf{y}) = \phi(\mathbf{x}, \mathbf{y})$. Explicitly for ϕ_ϵ the kernel in (6) also depends on ϵ , and is denoted by $K(\epsilon\mathbf{x}, \epsilon\mathbf{y})$,
94 with resulting kernel matrix $\Phi(\epsilon)$.

95 The paper is organised as follows. The Picard condition, which is commonly used to inves-
96 tigate properties of the solution to (2) for discrete ill-posed problems, is introduced in Section 2.

97 This analysis depends on the use of the singular value expansion of the kernel $K(\mathbf{x}, \mathbf{y})$ in the con-
98 tinuous case and the singular value decomposition (SVD) for the operator Φ in the discrete case.
99 For a square integrable kernel the stability of the solution $\alpha(\mathbf{y})$ depends on the decay properties
100 of the kernel singular values and on the projection of the given data to the singular value space.
101 Error analysis of the approximation is facilitated also through the use of the SVD in Section 2.3.
102 The singular value spectrum for the underlying operator is thus significant. Section 3 explores
103 the dependence of the spectrum of Φ on the smoothness of the underlying RBFs, demonstrating
104 that the decay rate of the discrete singular values impacts the smoothing properties of Φ . Such
105 results indicate why well-recognised difficulties with solving the RBF interpolation problem are
106 closely related to the underlying shape parameters. Results presented in Section 4 illustrate these
107 properties in the context of solving (2), with respect to both the underlying spectra for some pre-
108 sented examples, but also in terms of the underlying basis for the solution. Moreover, these results
109 suggest that techniques which are standard in regularizing the solutions for (6) can be applied to
110 obtain feasible solutions for (2). This is demonstrated in Section 5 in the novel context of examin-
111 ing the underlying α as a function which thus permits mapping the centers of the RBFs from one
112 set to another whilst maintaining a good approximation of the data. Equivalently, this amounts to
113 using the functional approximation for α to resample for a different set of quadrature weights in
114 (1). Conclusions of the paper and topics for future areas of investigation are given in Section 6.

115 **2. Stability and Error Analysis of the RBF Interpolation Problem**

116 Given the association between the solutions of (2) and (6), we now appeal to standard tech-
117 niques from the analysis of discrete ill-posed linear problems to study the properties of the solution
118 α and the resulting impact on the the quality of the approximant $s(\mathbf{x})$. To make the paper more
119 self-contained a brief review of the Picard condition now follows; for more details, the reader is
120 directed to [7].

121 *2.1. The Continuous Picard Condition*

122 The Picard condition, proved first by Picard in 1910, [19, p.160], is a necessary and sufficient
123 condition for the existence of a square integrable solution $\beta(\mathbf{y})$ to the Fredholm integral equation of

124 the first kind (6). It reveals how the solution of an ill-posed problem depends both on the spectrum
 125 of the kernel function and on the spectral content of the provided functional data, and is commonly
 126 used to evaluate the degree of ill-posedness of an ill-posed problem [7, 8, 16–18]. In the context of
 127 the RBF approximation problem, the Picard condition is particularly pertinent in illustrating how
 128 the smoothness of a RBF, which is directly related to its shape parameter, immediately impacts the
 129 smoothing properties of Φ and thus determines the ill-conditioning of (2).

130 The well-posedness of solutions to (6) is analyzed using the *singular value expansion* (SVE)
 131 [20], [7, p.10]. Consider the one dimensional case: given two functions u, v defined on a domain
 132 $\Omega \in \mathbb{R}$, their inner product is defined by

$$133 \quad \langle u, v \rangle = \int_{\Omega} u(x)v(x)dx. \quad (8)$$

134 Then a singular value system $\{\sigma_{\ell}, v_{\ell}, u_{\ell}\}$ of the *square integrable* kernel $K(x, y)$ is defined by

$$135 \quad \begin{cases} \int_{\Omega} K(x, y)v_{\ell}(y)dy = \sigma_{\ell}u_{\ell}(x), \\ \langle u_{\ell}, u_j \rangle = \langle v_{\ell}, v_j \rangle = \delta_{\ell j}, \end{cases} \quad (9)$$

136 and yields the SVE for K

$$137 \quad K(x, y) = \sum_{\ell=1}^{\infty} \sigma_{\ell}u_{\ell}(x)v_{\ell}(y). \quad (10)$$

138 Now suppose (10) is an absolutely and uniformly convergent expansion for $K(x, y)$, then

$$139 \quad \int_{\Omega} K(x, y)\beta(y)dy = \int_{\Omega} \sum_{\ell=1}^{\infty} \sigma_{\ell}u_{\ell}(x)v_{\ell}(y)\beta(y)dy = \sum_{\ell=1}^{\infty} \sigma_{\ell}\langle v_{\ell}(y), \beta(y) \rangle u_{\ell}(x) = f(x). \quad (11)$$

140 On the other hand $\{u_{\ell}(x)\}_{\ell=1}^{\infty}$ is an orthogonal basis in $L_2(\mathbb{R})$ and thus

$$141 \quad f(x) = \sum_{\ell=1}^{\infty} \langle u_{\ell}(x), f(x) \rangle u_{\ell}(x). \quad (12)$$

142 Comparing coefficients of $u_{\ell}(x)$ in (11) and (12) gives $\sigma_{\ell}\langle v_{\ell}(y), \beta(y) \rangle = \langle u_{\ell}(x), f(x) \rangle$ which, with
 143 the requirement $\langle u_{\ell}(x), f(x) \rangle = 0$ when $\sigma_{\ell} = 0$, gives the solution to (6) as

$$144 \quad \beta(y) = \sum_{\ell=1}^{\infty} \langle v_{\ell}(y), \beta(y) \rangle v_{\ell}(y) = \sum_{\ell=1}^{\infty} \frac{\langle u_{\ell}(x), f(x) \rangle}{\sigma_{\ell}} v_{\ell}(y). \quad (13)$$

145 It is immediate that a necessary condition for the convergence of (13) is the faster decay of the term
 146 $|\langle u_{\ell}(x), f(x) \rangle|$ in the numerator than the term σ_{ℓ} in the denominator. This leads to the well-known
 147 *Picard condition* (PC):

148 **Condition 1** (Picard Condition [19, III.§11][7, p.12],[8, p.279]). *An integral equation of the first*
 149 *kind, (6), has an unique square integrable solution $\beta(\mathbf{y})$ if and only if the series (13) satisfies*

$$150 \quad \|\beta(\mathbf{y})\|^2 = \sum_{\ell=1}^{\infty} \frac{\langle u_{\ell}(x), f(x) \rangle^2}{\sigma_{\ell}^2} < \infty. \quad (14)$$

151 When the kernel is symmetric, $u_{\ell} = v_{\ell}$, and $\langle u_{\ell}(x), f(x) \rangle$ can be viewed as the generalized
 152 Fourier expansion coefficients of f with respect to the orthogonal system $\{u_{\ell}\}$ [19, p.160]. More-
 153 over, the absolute and uniform convergence of (10) can be guaranteed by the well-known Mer-
 154 cer theorem [21]. From a historical perspective the SVE is a generalization of the eigenfunc-
 155 tion/Fourier series expansion of a symmetric kernel [19, 20]. The singular value system $\{\sigma_{\ell}, u_{\ell}, v_{\ell}\}$
 156 is the analogue of the Fourier series expansion in the sense that low frequency singular functions
 157 correspond to the large singular values and higher frequency singular functions correspond to
 158 smaller singular values [7, p.17]. For a given f , the faster decay of $|\langle u_{\ell}(x), f(x) \rangle|$, on average for
 159 some $\ell > 0$, as compared to σ_{ℓ} is indicative that the high frequency content of $f(x)$ is less than that
 160 of $K(x, y)$. Equivalently, the magnitudes of $|\langle u_{\ell}(x), f(x) \rangle|$ convey information on the *smoothness* of
 161 f with respect to the basis $\{u_{\ell}(x)\}$, while the decay of the singular values is closely related to the
 162 *smoothness* of the kernel function $K(x, y)$. In the equivalence of the expansions for $f(x)$ in (11)
 163 and (12) it is the mapping of u_{ℓ} to $\sigma_{\ell}v_{\ell}$ which helps to explain why higher frequencies (larger ℓ)
 164 in β are damped more in f than the lower frequencies, as observed from [7, p.11].

165 **Remark 1.** *For a given $K(x, y)$ the solution $\beta(\mathbf{y})$ is only square integrable when $f(x)$ is sufficiently*
 166 *smooth with respect to the singular system for $K(x, y)$. In contrast if $f(x)$ is given but $K(x, y)$ can*
 167 *be chosen, then an appropriate kernel is one with a singular system such that coefficients of $f(x)$*
 168 *in the singular basis $\{u_{\ell}\}$ decay faster, on average for a given $\ell > 0$, than the singular values σ_{ℓ} .*

169 This interpretation is not standard when treating the integral equation case, in which $K(x, y)$
 170 and $f(x)$ are prescribed by an underlying model and associated measurements, respectively. For
 171 the RBF problem, measurements $f(x)$ are provided, but the choice of kernel $K(x, y)$ depends on a
 172 number of factors including the context in which the approximation (1) is to be used. Specifically,
 173 the choice for $K(x, y)$ may be impacted by the underlying smoothness of $f(x)$ with respect to the
 174 basis for $K(x, y)$.

175 *2.2. The Discrete Picard Condition*

176 The SVD, [22], is employed to analyse the discrete problem (2). For a general matrix $\Phi \in$
 177 $\mathbb{R}^{M \times N}$ the SVD is given by $\Phi = U\Sigma V^T$, [22], where, $U \in \mathbb{R}^{M \times M}$, $V \in \mathbb{R}^{N \times N}$ are orthogonal
 178 matrices, and $\Sigma \in \mathbb{R}^{M \times N}$ is the diagonal matrix of singular values, ordered from largest to smallest
 179 along the principal diagonal, with $\sigma_1 \geq \sigma_2 \geq \dots \geq \sigma_N \geq 0$. Given the SVD, and explicitly
 180 assuming that Φ has full column rank, $\sigma_N > 0$, the solution of (2) is represented by the dyadic sum

$$181 \quad \alpha = \sum_{\ell=1}^N \frac{\mathbf{u}_\ell^T \mathbf{f}}{\sigma_\ell} \mathbf{v}_\ell = V \Sigma^{-1} U^T \mathbf{f}, \quad (15)$$

182 where \mathbf{u}_ℓ and \mathbf{v}_ℓ are the ℓ^{th} columns of matrices U and V , respectively. This relation is the im-
 183 mediate analogue of the continuous expression (13) using the SVD as the discrete analogue of
 184 the continuous SVE. The solution is a weighted linear combination of the basis vectors \mathbf{v}_ℓ with
 185 weights $\mathbf{u}_\ell^T \mathbf{f} / \sigma_\ell$. Moreover

$$186 \quad \|\alpha\|_2^2 = \sum_{\ell=1}^N \frac{(\mathbf{u}_\ell^T \mathbf{f})^2}{\sigma_\ell^2}, \quad (16)$$

187 which is the equivalent of (14) for the square integrability of the discrete solution α and leads
 188 immediately to the extension of the Picard condition for the discrete case.

189 **Condition 2** (Discrete Picard Condition [23, p.9] [7, p.37]). *Suppose that there exist r and τ such*
 190 *that numerically $\sigma_\ell = O(\tau), \forall \ell > r$. Then the discrete Picard condition (DPC) is said to be*
 191 *satisfied if, for all singular values which are larger than τ , the corresponding coefficients $|\mathbf{u}_\ell^T \mathbf{f}|$, on*
 192 *average, decay faster than σ_ℓ .*

193 Observe first that while it cannot be assumed that satisfaction of the DPC implies satisfaction
 194 of the PC for the underlying continuous problem [23], it does have an equivalent interpretation.
 195 In particular, the DPC depends not only on the smoothness of the data \mathbf{f} but also on the discrete
 196 singular system for the underlying kernel Φ which depends on the choice of RBF, ϕ . It is well-
 197 recognized that the solvability of a given interpolation problem depends on the condition of Φ ,
 198 while the DPC makes it clear that the quality of the solution is also directly related to the frequency
 199 content of the data \mathbf{f} , measured with respect to the singular basis for Φ , even when no noise is
 200 present in \mathbf{f} . The decay rate of $|\mathbf{u}_\ell^T \mathbf{f}|$ should be no slower than the decay rate of σ_ℓ . In practice,

201 the DPC is examined through the Picard plot, which is the plot of $|\mathbf{u}_\ell^T \mathbf{f}|$, σ_ℓ and the ratios $|\mathbf{u}_\ell^T \mathbf{f}|/\sigma_\ell$
 202 against index ℓ . If the DPC is not satisfied, the two norm of the solution α will be large. Moreover,
 203 it is immediate from (16) that

$$204 \quad \|\alpha\|_2^2 = \sum_{\ell=1}^N \frac{|\mathbf{u}_\ell^T \mathbf{f}|^2}{\sigma_\ell^2} \leq \left(\frac{1}{\sigma_N^2} \right) \sum_{\ell=1}^N |\mathbf{u}_\ell^T \mathbf{f}|^2 = \|\Phi^{-1}\|_2^2 \|\mathbf{f}\|_2^2. \quad (17)$$

205 In particular, $\|\alpha\|_2$ large relative to $\|f(\mathbf{x})\|_2$ is indicative of ill-conditioning of Φ and sensitivity of
 206 the solution α to errors in both $f(\mathbf{x})$ and to numerical errors accrued in evaluating α .

207 2.3. Error in the approximation due to ill-conditioning of Φ

208 As already detailed in Section 2.2, the boundedness of the solution α is directly related to the
 209 spectrum of the operator Φ . Error in α , which depends on both the ill-conditioning of Φ and the
 210 accuracy of the measurements f , also impacts the quality of the approximation $s(\mathbf{y})$. Suppose \tilde{f}_k
 211 are the actual measurements approximating the exact data f_k , such that $f_k = \tilde{f}_k + \eta_k$ for some noise
 212 vector $\boldsymbol{\eta}$ with components η_k and that $\tilde{s}(\mathbf{y})$ is the calculated approximation given $\tilde{\mathbf{f}}$. Practically,
 213 one may assume that calculating the solution of (2) in finite precision arithmetic yields $\tilde{\alpha}$ which
 214 exactly solves the perturbed system

$$215 \quad \tilde{\Phi} \tilde{\alpha} = \tilde{\mathbf{f}}, \quad (18)$$

216 where $\tilde{\Phi}$ is a matrix near to Φ in the sense that $\tilde{\Phi} = U \tilde{\Sigma} V^T$, for U and V the orthogonal matrices
 217 from the SVD of Φ , and $\tilde{\Sigma}$ the matrix of perturbed singular values $\tilde{\sigma}_\ell$, with the standard ordering.
 218 Consequently, the obtained \tilde{s} depends on $\tilde{\alpha}$. The quality of the approximation may be measured
 219 through $E = \|\mathbf{f} - s(\mathbf{y})\|_2$ which is bounded as follows:

$$220 \quad E = \|\mathbf{f} - s(\mathbf{y})\|_2 \leq \underbrace{\|\mathbf{f} - \tilde{\mathbf{f}}\|_2}_{E_0} + \underbrace{\|\tilde{\mathbf{f}} - \tilde{s}(\mathbf{y})\|_2}_{E_1} + \underbrace{\|\tilde{s}(\mathbf{y}) - s(\mathbf{y})\|_2}_{E_2}. \quad (19)$$

Here $E_0 = \|\boldsymbol{\eta}\|^2$ is the measurement error, E_1 is the conventional approximation error under the
 assumption that \tilde{s} is obtained to approximate the measured data $\tilde{\mathbf{f}}$, while E_2 is the error in the
 approximation. Observe first that the error in approximating f at vector \mathbf{y} is given by

$$s(\mathbf{y}) - \tilde{s}(\mathbf{y}) = \sum_{j=1}^N (\alpha_j - \tilde{\alpha}_j) \Phi(\mathbf{y}, \mathbf{x}_j) = \Phi(\mathbf{y}, \mathbf{x})(\alpha - \tilde{\alpha}) \quad \text{which yields} \quad (20)$$

$$\|s(\mathbf{y}) - \tilde{s}(\mathbf{y})\|_2 = \|\Phi(\mathbf{y}, \mathbf{x})(\alpha - \tilde{\alpha})\|_2 \leq \|\Phi(\mathbf{y}, \mathbf{x})\|_2 \|\alpha - \tilde{\alpha}\|_2. \quad (21)$$

Now

$$\|\alpha - \tilde{\alpha}\|_2^2 = \|(\Sigma^{-1}\tilde{\Sigma} - I)\tilde{\alpha} + \Sigma^{-1}U^T\boldsymbol{\eta}\|_2^2 \leq \|(\Sigma^{-1}\tilde{\Sigma} - I)\|_2^2\|\tilde{\alpha}\|_2^2 + \|\Phi^{-1}(\mathbf{y}, \mathbf{x})\|_2^2\|\boldsymbol{\eta}\|_2^2,$$

which then yields

$$\begin{aligned} \|s(\mathbf{y}) - \tilde{s}(\mathbf{y})\|_2^2 &\leq \text{cond}^2(\Phi(\mathbf{y}, \mathbf{x}))\|\boldsymbol{\eta}\|_2^2 + \|\Phi(\mathbf{y}, \mathbf{x})\|_2^2\|(\Sigma^{-1}\tilde{\Sigma} - I)\|_2^2\|\tilde{\alpha}\|_2^2 \\ &\leq \text{cond}^2(\Phi(\mathbf{y}, \mathbf{x}))\left(\|\boldsymbol{\eta}\|_2^2 + \sigma_1^2(\Delta)\|\tilde{\alpha}\|_2^2\right), \quad \Delta = \Sigma - \tilde{\Sigma}, \end{aligned} \quad (22)$$

221 and where the last inequality follows by, for example, [22, Corollary 8.6.2]. The matrix Δ measures
 222 the error in the approximation of $\Phi(\mathbf{y}, \mathbf{x})$ by the calculated $\tilde{\Phi}(\mathbf{y}, \mathbf{x})$, which is, in general, unknown,
 223 but can certainly be assumed to be bounded below by the machine epsilon. Of course the bound
 224 in (22) can be replaced by a bound in terms of $\|\tilde{\mathbf{f}}\|_2^2$, but this is less useful and requires further
 225 information on the maximum ratio in the singular values $\{\sigma_\ell/\tilde{\sigma}_\ell\}$. On the other hand, from (22)
 226 one sees immediately the contamination of the approximation $s(\mathbf{y})$ due to the ill-conditioning
 227 of $\Phi(\mathbf{y}, \mathbf{x})$ and the relation to the size of the calculated coefficients $\tilde{\alpha}$. Moreover, ignoring the
 228 contribution due to the measurement error, the error in the approximation is bounded by a scaling
 229 of the size of the calculated weight vector, where the scaling is controlled if $\sigma_1(\Delta)$ is such that the
 230 product $\sigma_1(\Delta)\text{cond}(\Phi(\mathbf{y}, \mathbf{x}))$ is not too large. To reiterate:

231 **Remark 2.** *The size of the calculated $\|\alpha\|_2$ relative to $\|\mathbf{f}\|_2$ reflects the smoothing effects of the*
 232 *underlying kernel matrix and can itself be viewed as an indication of the ill-conditioning of the*
 233 *underlying problem.*

234 3. Spectral Distribution of Kernel Matrices

235 From the results presented in Section 2 it is clear that an exponentially decaying spectrum for
 236 the interpolation matrix Φ negatively impacts the accuracy of the estimation of α and, through
 237 this, also the quality of the resulting approximation $s(\mathbf{y})$. Indeed the stability of the solution of
 238 the RBF interpolation problem is closely related to the spectral distribution of the underlying
 239 kernel $K(\mathbf{x}, \mathbf{y}) = \phi(\mathbf{x}, \mathbf{y})$ which is explicitly assumed to be symmetric for the the interpolation
 240 problem described by (2). In this case, with the standard ordering for the singular values $\sigma_\ell =$

241 $|\lambda_\ell|$ where λ_ℓ is an eigenvalue of Φ , it is therefore sufficient to analyse the eigenvalues of Φ .
 242 Moreover, when the continuous kernel $K(\mathbf{x}, \mathbf{y})$ is square integrable, the discrete singular values
 243 from the SVD for the finite operator are increasingly better approximations, with increasing N ,
 244 to the continuous singular values of the SVE when calculated using a standard quadrature, [24].
 245 Given these connections we therefore study the continuous eigenvalue problem

$$246 \quad \int_{\Omega} K(\mathbf{x}, \mathbf{y})\alpha(\mathbf{y})d\mathbf{y} = \lambda\alpha(\mathbf{x}), \quad (23)$$

247 which has been studied by many authors including in Bôcher [25], Weyl [26], Hille and Tamarkin
 248 [27], Little and Reade [28], as well as in the texts of Smithies [20] and Cochran [29], in which
 249 also important early results of Fredholm and Hilbert are detailed.

250 Consistent with the standard ordering of the singular values it is assumed throughout that the
 251 eigenvalues are ordered by decreasing magnitude with increasing ℓ . Moreover, a kernel is symmetric
 252 if $K(\mathbf{x}, \mathbf{y}) = K(\mathbf{y}, \mathbf{x})$, while it is translation invariant if $K(\mathbf{x}, \mathbf{y}) = k(\mathbf{x} - \mathbf{y})$ for some function $k(\mathbf{x})$.
 253 If the underlying function $k(\mathbf{x})$ is also periodic, then the kernel is said to be translation invariant and
 254 periodic. These definitions are relevant in the context of the translation invariant symmetric RBF
 255 $\phi_\epsilon(\|\mathbf{x} - \mathbf{y}\|_2)$. Generally, it is assumed here that the kernel is real and square integrable, and thus,
 256 asymptotically, the results may be viewed as providing results which are relevant in the discrete
 257 case.

258 3.1. Scalar kernels in \mathbb{R}

259 **Theorem 1** (Weyl [26, p.449][16, p.123] [20]). *Suppose $K(x, y)$ is a symmetric kernel and its
 260 partial derivatives $\partial^q K / \partial^q x$, $q = 0, \dots, \nu$ exist and are continuous, then $|\lambda_\ell| = O(\ell^{-\nu-1/2})$.*

261 While it can be shown that Weyl's result is sharp for the general case [30], there are cases in
 262 which sharper results can be obtained.

263 **Theorem 2** ([31]). *Suppose $K(x, y)$ is a symmetric positive definite kernel which is periodic in
 264 both variables with the same period, and its partial derivatives $\partial^q K / \partial^q x$, $q = 0, \dots, \nu$ exist and
 265 are continuous, then $\lambda_\ell = o(\ell^{-\nu-1})$.*

266 **Theorem 3.** [32, Theorem 4] *Suppose $K(x, y) = k(x - y)$ is a translation invariant kernel on
 267 $[-1, 1]$.*

268 1. If $k(x) \in C^{v-1}$ and its v^{th} derivative is of bounded variation, then $|\lambda_\ell| = O(\ell^{-v})$ as $\ell \rightarrow \infty$.

269 2. If $k(x) \in C^\infty$, then $|\lambda_\ell|$ decay faster than any $O(\ell^{-v})$ as $\ell \rightarrow \infty$ for every $v > 0$.

270 When the kernel $K(x, y)$ is both analytic and symmetric, its eigenvalues decay much faster.

271 **Theorem 4** (Little-Read [28]). Suppose \mathcal{E}_R is an ellipse with foci at ± 1 and semi-axis sum $R > 1$.

272 If $K(x, y) = K(y, x) \in C[-1, 1]^2$, and for each $y \in [-1, 1]$ there is an analytic continuation to

273 $K(z, y)$ for z inside \mathcal{E}_R which is uniformly bounded in z, y in this range, then $|\lambda_\ell| = O(R^{-\ell})$.

274 **Remark 3.** For a translation invariant kernel $k(x)$, the smoother $k(x)$ the more quickly its eigenval-

275 ues decay to zero. This is a negative result with respect to finding stable solutions of the underlying

276 integral equation (6).

277 The shape parameter, $\epsilon > 0$, of the kernel is also relevant. The explicit dependence of the

278 eigenvalues on the shape parameter is noted by introducing the notation $\lambda_\ell(\epsilon)$.

279 **Theorem 5** ([32]). Suppose that $K(x, y) = k(x - y)$ is a translation invariant kernel.

280 1. If $k(x) \in C^{v-1}$ and its v^{th} derivative is of bounded variation, then $|\lambda_\ell(\epsilon)| = O((\epsilon/\ell)^v)$ as
281 $\ell \rightarrow \infty$.

282 2. If $k(x) \in C^\infty$, then $|\lambda_\ell(\epsilon)|$ decay faster than $O((\epsilon/\ell)^v)$ as $\ell \rightarrow \infty$ for every $v > 0$.

283 It is immediate from Theorem 5 that the eigenvalues decay faster for smaller values of ϵ and

284 thus the problem becomes more ill-posed when the underlying RBF is referred to as *flat*, [33, 34].

285 **Remark 4.** For RBF ϕ_ϵ , the flatter the basis function, the more significant the smoothing effect, in

286 the limit of large N , of the associated kernel matrix $\Phi(\epsilon)$.

287 Note that a function containing only low frequency content, such as $\sin(x)$ is said to be *smoother*

288 than a function with high frequency content such as $\sin(1000x)$, even though both functions have

289 the same number of continuous derivatives. Figure 1 illustrates this result. The *flatter* the basis

290 function, the faster the spectrum decays to zero.

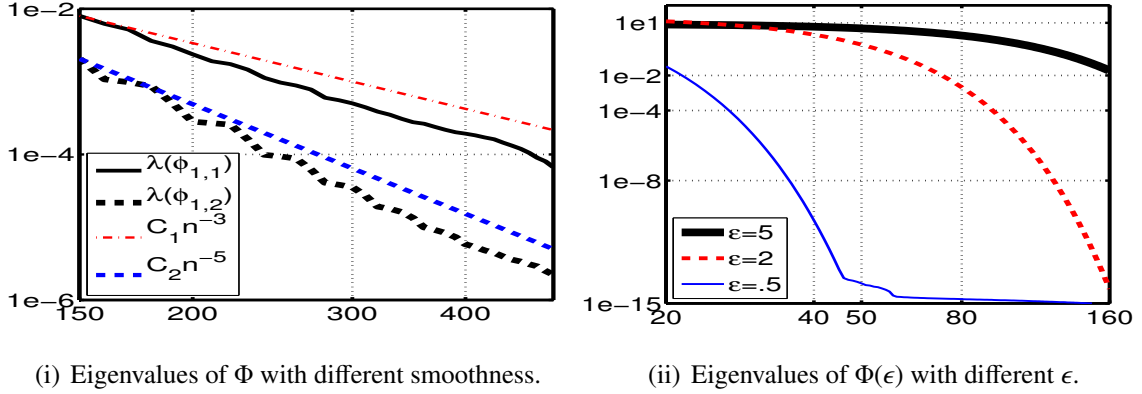


Figure 1: (i) illustrates the eigenvalue distribution of kernel matrices for Wendland functions $\phi_{1,1} = (1-x)_+^3(3x+1) \in C^2$ and $\phi_{1,2} = (1-x)_+^5(8x^2+5x+1) \in C^4$, where $a_+ = \max\{a, 0\}$. The straight lines are the predicted trends for the eigenvalue decay according to Theorem 2, where $C_1 = \lambda_{150}(\phi_{1,1})150^3$, $C_2 = \lambda_{150}(\phi_{1,2})150^5$. (ii) shows the spectral distribution of kernel matrices for Gaussian radial basis functions with different scales, $\phi(x, \varepsilon) = e^{-(\varepsilon x)^2}$, with scale parameters $\varepsilon = 5, 2, 0.5$. The matrices Φ are obtained using an equal sampling on the interval $[-10, 10]$ using 501 points. In (ii) the first 160 eigenvalues are shown.

291 3.2. Separable kernels in \mathbb{R}^d

292 For the two dimensional case $\mathbf{x}, \mathbf{y} \in \mathbb{R}^2$, separability means that $K(\mathbf{x}, \mathbf{y}) = K_1(x_1, y_1)K_2(x_2, y_2)$.
 293 While this does limit the kernels that can be considered, an example is given by the Gaussian
 294 kernels: $K(\mathbf{x}, \mathbf{y}) = \exp(-\|\mathbf{x} - \mathbf{y}\|_2^2) = \exp(-(x_1 - y_1)^2) \exp(-(x_2 - y_2)^2)$. One can show that if
 295 λ is an eigenvalue of $K_1(x_1, y_1)$ and μ is an eigenvalue of $K_2(x_2, y_2)$, then $\lambda\mu$ is an eigenvalue of
 296 $K(\mathbf{x}, \mathbf{y})$ [20]. Suppose that each K_j satisfies the conditions of Theorem 4 and that the eigenvalues
 297 of $K_j(x_j, y_j)$ are ordered in magnitude from largest to smallest $O(1), O(R^{-1}), O(R^{-2}), \dots$ relative to
 298 an underlying ellipse with semi axes summing to R . It is not hard to determine that the related
 299 separable kernel $K(\mathbf{x}, \mathbf{y})$ therefore has the following ordering of eigenvalues: 1 of order $O(1)$,
 300 2 of order $O(R^{-1})$, 3 of order $O(R^{-2})$ and $\binom{\ell}{1}$ of order $O(R^{-\ell})$. See for example the ordering of
 301 eigenvalues for a two dimensional tensor product kernel in Figure 2. The result extends in a
 302 similar manner for the tensor product kernel of dimension d .

303 **Theorem 6** ([32]). *Suppose that $K(\mathbf{x}, \mathbf{y})$ is a separable, tensor product-like analytic kernel function*
 304 *in \mathbb{R}^d , and that the ℓ^{th} eigenvalue of each kernel K_j is of order $O(R^{-\ell-1})$, then there are exactly*
 305 *$\binom{\ell+d-1}{d-1} = \frac{(\ell+d-1)!}{\ell!(d-1)!}$, $\ell = 1, \dots$, eigenvalues of the same order, namely $O(R^{-\ell-1})$ for some $R > 1$.*

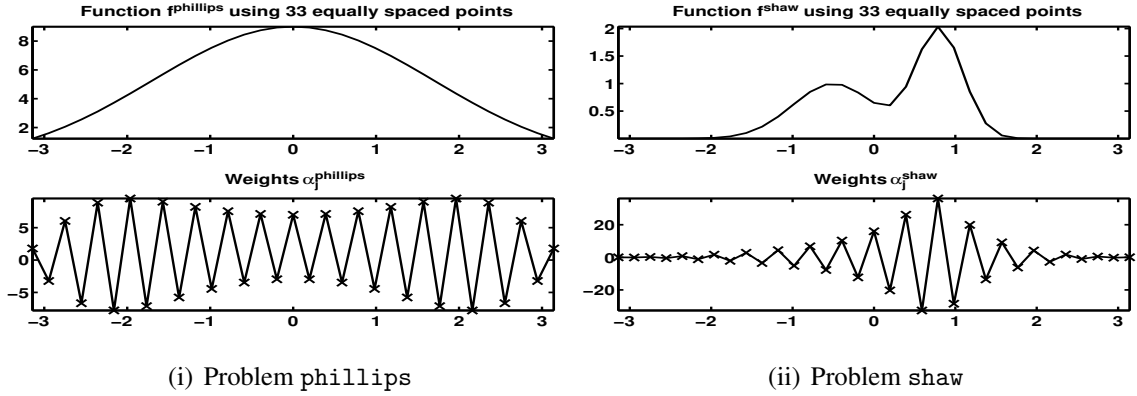


Figure 3: Solving (2) with $N = 33$ for functions f^{phillips} and f^{shaw} without addition of any noise to obtain the weights $\{\alpha_j^{\text{phillips}}\}$ and $\{\alpha_j^{\text{shaw}}\}$ from samples at 33 equally spaced points on $[-\pi, \pi]$.

319 for two one dimensional examples based on the following two functions

$$320 \quad f^{\text{phillips}}(x) = (6 - |x|)(1 + .5 \cos(\pi x/3)) + \frac{4.5}{\pi} \sin(\pi|x|/3) \quad \text{and} \quad (24)$$

$$321 \quad f^{\text{shaw}}(x) = 2 \exp(-6(x - .8)^2) + \exp(-2(x + .5)^2). \quad (25)$$

323 These are the right hand sides for the phillips and shaw examples used in the Regularization
324 toolbox of Hansen [35]. The RBF interpolation equation (2) is solved for the RBF $\phi_\epsilon(x) =$
325 $\exp(-(\epsilon x)^2)$, $\epsilon = 2$, using $N = 33$, with equally spaced $\mathbf{x}_k = -\pi + (k - 1)\Delta\mathbf{x}$, $k = 1, \dots, N$,
326 $\Delta\mathbf{x} = 2\pi/(N - 1)$. The problem is moderately ill-conditioned, $\text{cond}(\Phi) = 2.45e6$. The presented
327 solutions $\{\alpha_j^{\text{phillips}}\}$ and $\{\alpha_j^{\text{shaw}}\}$ found by the naïve inversion of (2), using the dyadic sum (15),
328 clearly illustrate the higher frequency content of the interpolating functions than that of the target
329 functions, i.e. functions f^{phillips} , f^{shaw} are much smoother than the solutions $\alpha^{\text{phillips}}(\mathbf{y})$, $\alpha^{\text{shaw}}(\mathbf{y})$.
330 On the other hand, given that the underlying system matrix $\Phi(\epsilon)$ is not well-conditioned, one can-
331 not immediately assume that the given results are actually reasonable approximations to the true
332 infinite precision solutions for α . Thus the conclusion concerning the frequency content may be
333 an artifact of the sensitivity of the solutions due to the spectral system for $\Phi(\epsilon)$ and, as noted in
334 Section 2.2, it is important to consider the DPC for the system.

335 Figure 4 shows the Picard plot, [23], as discussed in Section 2.2, the plot of $|\mathbf{u}_\ell^T \mathbf{f}|$, σ_ℓ and the
336 ratios $|\mathbf{u}_\ell^T \mathbf{f}|/\sigma_\ell$ against index ℓ , for the problems illustrated in Figure 3. As ℓ increases the singular
337 values σ_ℓ decrease to zero more quickly than the numerators $|\mathbf{u}_\ell^T \mathbf{f}|$. Although the matrix Φ has full

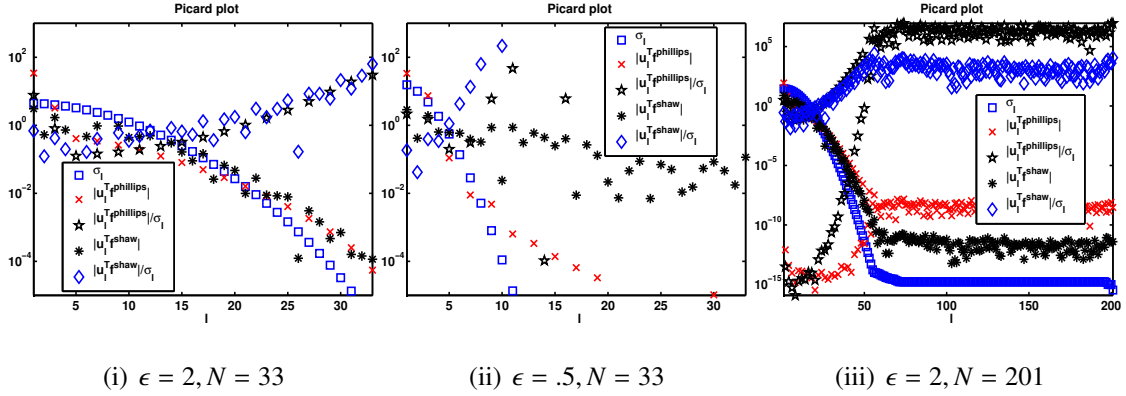


Figure 4: The Picard plot for problems f^{phillips} and f^{shaw} illustrated in Figure 3, (i) for $\epsilon = 2$, in (ii) for $\epsilon = .5$ and in (iii) the same problems but using $N = 201$, $\epsilon = 2$. For (ii) $\text{cond}(\Phi) = 7.3e17$ and for (iii) it is $6.0e18$. Note that in order to emphasize the growth in the coefficients the plots in (i)-(ii) are given for a limited range on the y-axis.

338 column rank the singular values $\sigma_\ell \rightarrow 0$ in such a way that $\Phi(\epsilon)$ is ill-conditioned, particularly
 339 when the size of the problem is increased as seen in Figure 4(iii), or when ϕ_ϵ becomes *flatter* as in
 340 Figure 4(ii). The observations are consistent with the discussion in Section 3. The impact of the
 341 shape parameter is clearly seen in the solutions illustrated in Figure 5 which have $\epsilon = .5$; notice in
 342 particular the scale of the y-axis in each case, the amplification in α for shape parameter $\epsilon = .5$ is
 343 significant.

344 To highlight the problem with these naïve solutions obtained from direct inversion of (2) it is
 345 relevant to calculate *approximate* solutions using the truncated SVD

$$346 \quad \alpha_\ell = \sum_{\ell=1}^{N_\ell} \frac{\mathbf{u}_\ell^T \mathbf{f}}{\sigma_\ell} \mathbf{v}_\ell, \quad (26)$$

347 for truncation parameters $1 \leq N_\ell \leq N$. The dependence of $\|\alpha_\ell\|^2$ on ℓ for the examples used in
 348 Figures 3 and 5(i)-5(ii), respectively, namely for the two different shape parameters, is illustrated
 349 in Figure 5. It is clear that as N_ℓ increases the solutions violate the DPC. Indeed the parameter N_ℓ
 350 is commonly regarded as a *regularization* parameter; in calculating the solution α it is necessary
 351 to either use truncation to obtain a solution without blow-up or to introduce some alternative reg-
 352 ularization. The choice of regularization parameter is extensively studied in the inverse problems
 353 literature, in which the dependence of the solution on the conditioning of the problem and the
 354 noise in the data is well-described, [23]. The focus here is on results using truncation, for which

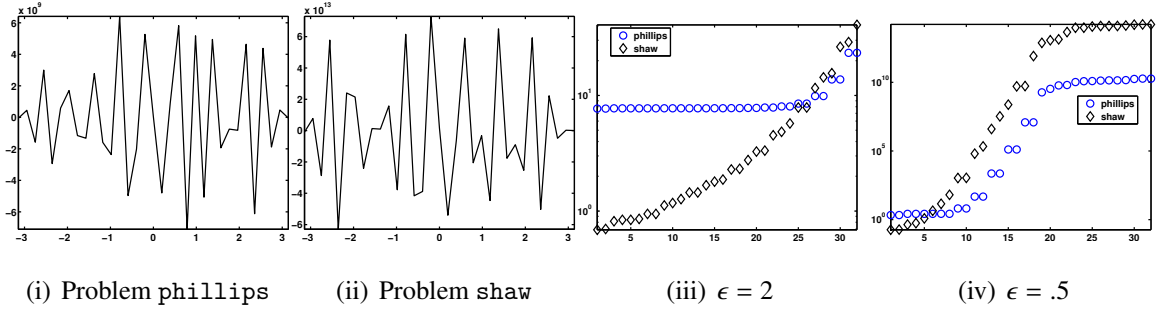


Figure 5: (i)-(ii) show the solution of (2) with $N = 33$ for functions f^{phillips} and f^{shaw} . The weights $\{\alpha_j^{\text{phillips}}\}$ and $\{\alpha_j^{\text{shaw}}\}$ are obtained from samples at 33 equally spaced points on $[-\pi, \pi]$ using a Gaussian kernel with $\epsilon = .5$. The solutions are significantly amplified as compared to the solutions with $\epsilon = 2$ shown in Figure 3. In (iii)-(iv) the discrete Picard condition, Condition 2, for functions f^{phillips} and f^{shaw} for the truncated solutions (26) $N_\ell \leq N$. $\|\alpha_\ell\|^2$ is plotted against index N_ℓ . The data is sampled at 33 equally spaced points on $[-\pi, \pi]$ using a Gaussian kernel with the noted shape parameter ϵ .

355 Figure 4(i) suggests that when $\epsilon = 2$ one should expect to truncate at about $N_\ell = 16$ in order
 356 to avoid the coefficients which start to increase with index ℓ , although strictly speaking for this
 357 choice the singular values have not reached the machine tolerance, τ , such that $\sigma_\ell = O(\tau)$, rather
 358 $\sigma_\ell > O(\tau)$ for $\ell > N_\ell$. In general, this choice depends not only on these singular values but also
 359 on the noise level in the measured data f .

360 4.2. Basis Vectors for the Solution

361 It is clear that the spectrum of the kernel matrix impacts the determination of feasible solutions
 362 of the RBF interpolation problem. The more significant the smoothing effect of the kernel matrix,
 363 the more difficult it is to find a satisfactory solution of (2). While this discussion has so far
 364 emphasized the effect of the poor conditioning of Φ due to the exponentially decaying spectrum, it
 365 should also be apparent that the existence of the small smoothing singular values will also lead to
 366 finite precision errors in the calculations of the basis vectors. High frequency basis vectors carry
 367 information corresponding to the extremely small smoothing eigen/singular values and are thus
 368 more likely to be contaminated by the effects of finite precision arithmetic. Incorrect estimates of
 369 these basis vectors may, therefore, contaminate the solutions and impact the quality of the RBF
 370 approximation.

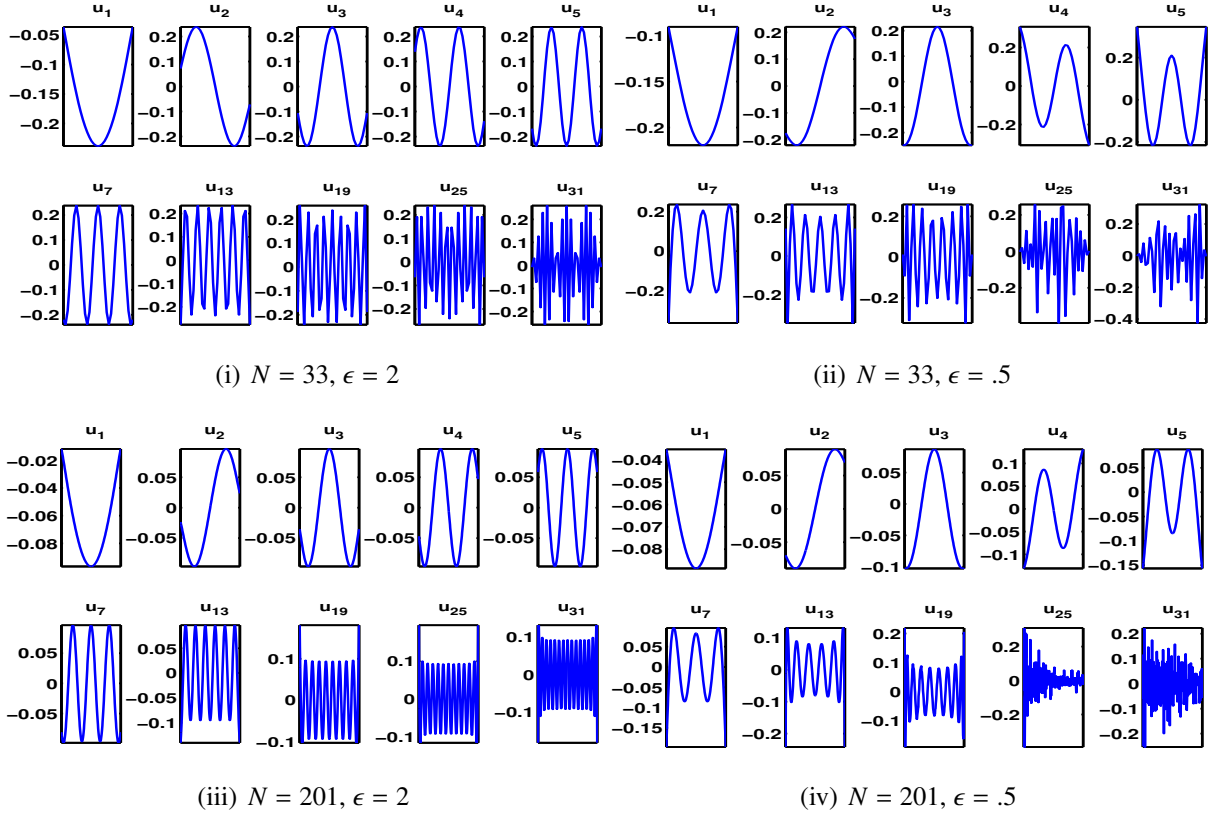


Figure 6: The basis vectors \mathbf{u}_ℓ for the the same system matrices used to generate the Picard plots in Figure 4, for $N = 33$ and $N = 201$. Note that these vectors are for the Gaussian radial basis function with $\epsilon = 2$ and $\epsilon = .5$.

371 Figure 6 illustrates the basis vectors \mathbf{u}_ℓ for a selection of ℓ for both cases $N = 33$ and $N = 201$,
372 for the examples in Figure 4. The y -axis is not shown because the vectors are normalized to
373 norm 1 and thus the scales in each case are comparable, with definitely no blow up in any case.
374 The dependence of the basis vectors on the shape parameter ϵ is shown in Figures 6(ii) and 6(iv).
375 There is an increasing number of oscillations, corresponding to increasing frequency content of
376 the vectors, with increasing ℓ . This confirms that the solutions for α will contain broad frequency
377 content according to (26) dependent on the choice for N_ℓ . Still it is not possible to examine the
378 frequency content of the basis vectors and then immediately determine an appropriate choice for
379 N_ℓ . Indeed there is apparent mixing of frequency content for larger ℓ , when considered for the
380 broader kernel with $\epsilon = .5$. The frequency content of a given vector \mathbf{u}_ℓ can be examined using the
381 normalized cumulative periodogram (NCP), which is a standard time series analysis technique,
382 e.g. [5]. A brief description of the approach follows.

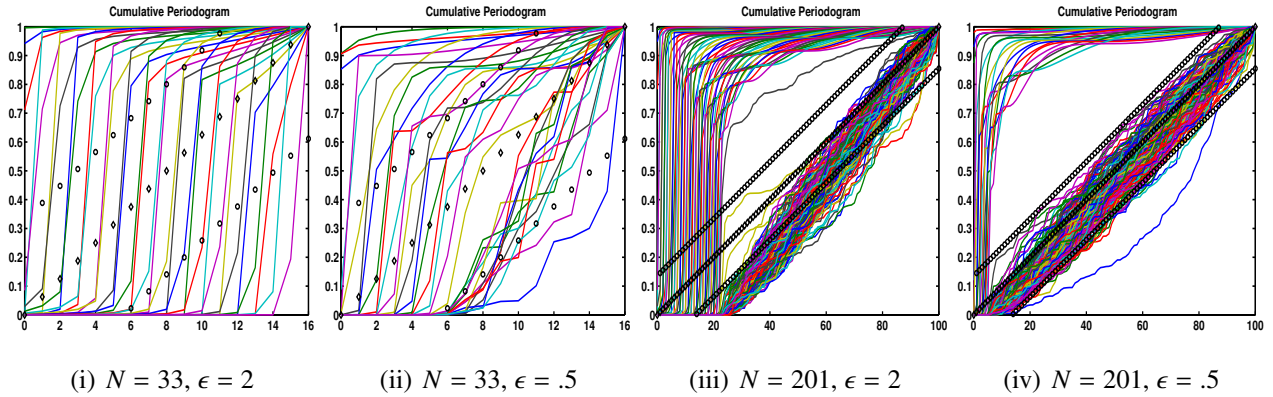


Figure 7: The NCP for vectors \mathbf{u}_ℓ for $N = 33$ and $N = 201$, with $\epsilon = 2$ and $\epsilon = .5$, using equally spaced sampling. The black markers indicate the upper and lower KS-5% bounds, and the diagonal white noise line given by $y = \ell/\tilde{N}$.

383 The power spectrum of a given vector \mathbf{b} of length N is determined via the discrete Fourier
 384 transform (dft) of the vector treated as a time series. For given \mathbf{b} define $\mathbf{c}_\ell = |(\text{dft}(\mathbf{b}))_\ell|^2$, for
 385 $\ell = 1, \dots, \tilde{N} = \lfloor N/2 \rfloor + 1$. Then the NCP is generated as the ratio of the cumulative sum of entries

$$386 \quad \mathbf{z}_j = \frac{\sum_{\ell=1}^j \mathbf{c}_\ell}{\sum_{\ell=1}^{\tilde{N}} \mathbf{c}_\ell}, \quad j = 1, \dots, \tilde{N}. \quad (27)$$

387 This definition includes the bias term (or DC term) \mathbf{c}_0 . It is then well-known that for a white noise
 388 time series, the plot of \mathbf{z}_j against frequency j normalized between 0 and 1 should lie within the
 389 Kolmogorov-Smirnoff (KS) limits of a straight line, which has slope 1 and length $\sqrt{5}/2$ [5, p. 363].
 390 Moreover, vectors with lower frequency content lie primarily above the KS limits, while those with
 391 high frequency lie distinctly below the KS limits. In all cases the normalization reaches 1 for the
 392 last index; see Figures 7 for the basis vectors $\ell = 1 : N$ for the same cases as Figure 6, and observe
 393 not only the dependence on both N and ϵ , but also the finite precision impact of the calculations.
 394 In particular for large N the lack of accuracy in the calculation of the singular values, translates
 395 to inaccurate estimates of the underlying singular vectors, as exhibited by the large number of
 396 vectors within the white noise KS bounds for the case with $N = 201$, i.e. these vectors degenerate
 397 to becoming white noise for large ℓ [36]. On the other hand, the lower frequency content for the
 398 *smoother* kernel with $\epsilon = .5$ translates to a larger number of vectors with low frequency content;
 399 more vectors above the upper KS line.

400 It should be apparent from the NCP that it is important to include in any given basis for a

401 solution, e.g. in the truncated dyadic sum (26), only those basis vectors which represent *true*
402 frequency content. Thus one needs a test to detect for the contamination of a vector by white
403 noise, with respect to some tolerance, so as to estimate an appropriate choice for $N_\ell \leq N$. Observe
404 that this determination is totally independent of the measured data \mathbf{f} . A number of suggestions for
405 automatically determining whether a residual vector for the data fit term, $\mathbf{r} = \Phi\alpha - \mathbf{f}$, represents
406 white noise have been presented in the literature, [37–39], all of which in some manner measure
407 the deviation of the NCP from the diagonal white noise line $y = \ell/\tilde{N}$. In this case it is not a residual
408 which is of interest, rather given a basis of orthogonal vectors, here the set $\{\mathbf{u}_\ell\}$, what is the largest
409 N_ℓ such that all vectors with $\ell > N_\ell$ are effectively, with respect to some tolerance, contaminated
410 by white noise?

411 Two standard approaches for measuring this deviation from white noise are considered, [37–
412 39]; the first is the total least squares deviation from the diagonal white noise line summed over
413 all entries for each vector, and the second is the percentage count of vector entries outside the KS
414 bounds for each vector. Given these measures it remains to determine an appropriate tolerance for
415 accepting the hypothesis that the vector is effectively a white-noise vector, relative to the chosen
416 KS confidence level. Here the approach of [36] is followed: for the Kolmogorov-Smirnov test
417 with a 5% confidence interval set $N_\ell = r$ which is the least index r such that for vector \mathbf{u}_{r+1} more
418 than 50% of the NCP lies within the KS-5% bounds. The 50% choice is in some sense arbitrary,
419 as is the choice of any regularization tolerance, but in this case one can see that there is a distinct
420 cut-off for these basis vectors. This is illustrated in Figure 8(i)-8(ii) for the case with $N = 201$
421 with shape parameters $\epsilon = 2$ and $\epsilon = .5$, respectively. With respect to both measures – the red and
422 blue curves – the sudden drop for a specific r is preserved. It is also clear from Figure 7 that the
423 noise is much more significant for $N = 201$ than for $N = 33$ and the number of reliable vectors
424 when $\epsilon = .5$ is much fewer than when $\epsilon = 2$.

425 Figure 8(iii) illustrates the underlying frequency content of the right hand sides for *phillips*
426 and *shaw* with $N = 33$ and $N = 201$ equally spaced samples, and the Gaussian radial basis
427 functions sampled on the same grids, with $\epsilon = 2$ and $\epsilon = .5$, with the total frequency normalized
428 to 1. The dft of each vector, normalized to 1 is plotted, rather than the cumulative vector which is
429 not useful due to the limited frequency content of these examples, Figure 8(iii). One can see that

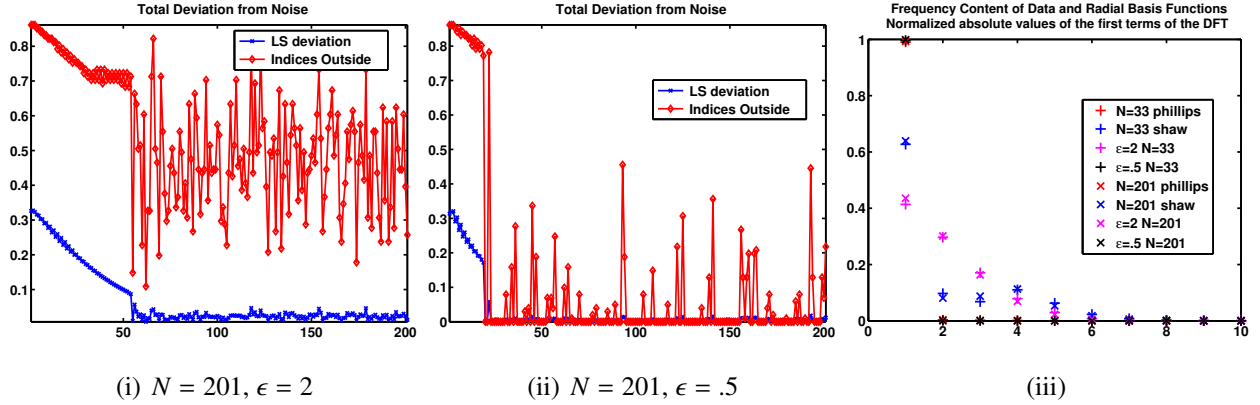


Figure 8: (i)-(ii) the deviation of vectors \mathbf{u}_ℓ from white noise for $N = 201$, with $\epsilon = 2$ and $\epsilon = .5$, and equal sampling. In (iii) the dft for f^{phillips} and f^{shaw} with $N = 33$ and $N = 201$ and for the Gaussian radial basis functions with $\epsilon = 2$ and $\epsilon = .5$.

430 the analysis is virtually independent of the sampling. Moreover, phillips is far smoother than
 431 shaw, and the RBF with shape parameter $\epsilon = .5$ is smoother than that with $\epsilon = 2$.

432 In summary, the DPC provides significant information on the impact of both noise in the
 433 measurements f and the accuracy of the smaller spectral values for estimating the number of
 434 potentially reliable terms in the TSVD estimate of the solution. The NCP, on the other hand, is
 435 measurement independent and may more precisely show the first deterioration in the basis and so
 436 more effectively indicate a good truncation parameter for the TSVD solution.

437 5. Applying Results to find interpolants

438 5.1. Impact on interpolation

439 The relationship between the integral equation (6) and the quadrature (7) leads to the following
 440 observation.

441 **Remark 6.** *The weights $\{\alpha_j\}$ in (1) may be regarded as the sampling of a function $\alpha(\mathbf{y})$.*

442 The impact of this remark is now discussed with respect to the use of both the naïve and the
 443 truncated solutions with $N_\ell = 16$ when $\epsilon = 2$ for interpolating the functions f^{phillips} and f^{shaw} .
 444 Results are shown first for the naïve solutions in Figure 9, and then for the truncated solutions in
 445 Figure 10. As before, note that the choice $N_\ell = 16$ is a regularization parameter choice, here based

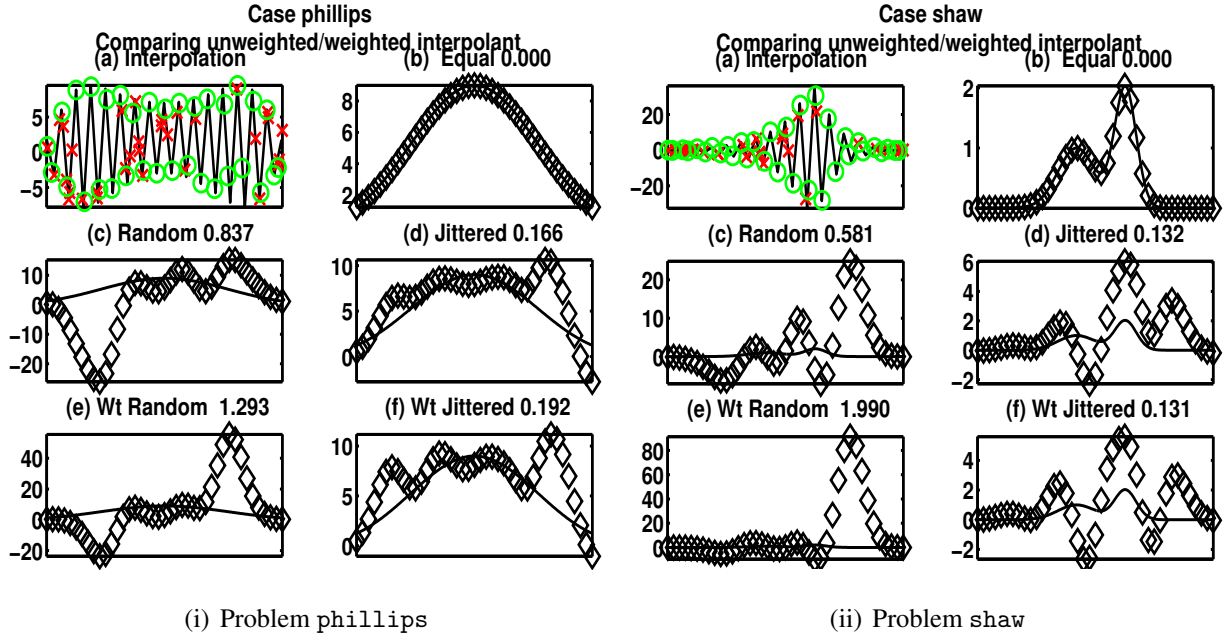


Figure 9: Using the direct solution α from (2) to obtain approximants to f^{phillips} and f^{shaw} , the left and right blocks (i)-(ii) for each problem, respectively. In panels (a) the solutions α are given as the solid black lines. This data is interpolated as described in the text to find estimates for $\alpha(\mathbf{y})$ for random and jittered sampling in \mathbf{y} , giving weights α_R and α_J , red x and green o symbols respectively. The resulting weights are used to generate approximations to f^{phillips} and f^{shaw} at 201 equally spaced points, using the original α in panel (b), α_R in panel (c) and α_J in panel (d). In panels (c)-(d) no weighting to account for the quadrature is performed, whereas in panels (e)-(f) weighted quadrature based on (7), assuming trapezoidal approximation, is used. The titles show the respective root mean square errors calculated using the approximated and exact data. The shape parameter is $\epsilon = 2$ as in Figure 3.

446 on the results of the NCP analysis for the case with no noise in the measurements and chosen only
 447 as an illustration of the use of $\alpha(\mathbf{y})$ as an interpolating function to the weights $\{\alpha_j\}$.

448 Based on Remark 6 one should be able to resample α at a set $\{\mathbf{y}_j\}$ to obtain (1), provided
 449 that the resampled points $\{\alpha(\mathbf{y}_j)\}$ are appropriately weighted in accordance with (7). Consider
 450 now resampling for two additional data sets, \mathbf{y}_R and \mathbf{y}_J , each at 33 points on the interval $[-\pi, \pi]$,
 451 where \mathbf{y}_R is randomly sampled on the interval, and \mathbf{y}_J represents a random jitter of the points
 452 away from equally spaced points by $.2\zeta_j(2\pi)/33$ for ζ_j , $j = 1, \dots, 33$ sampled from a normal
 453 distribution with mean 0 and variance 1. The interpolation from equally spaced points to the
 454 random or jittered points uses, here for purposes of illustration only, a cubic spline interpolant and

455 provides two new sets of weights $\{\alpha_j^{\text{random}}\}$ and $\{\alpha_j^{\text{jitter}}\}$ illustrated as the red x and green o symbols
 456 in panels (a) in Figures 9. Given a set of weights one can assess their accuracy in approximating
 457 the underlying target function f for a different set of sampled points by using (2). Taking 201
 458 equally spaced points on the given interval and forming the appropriate matrices Φ of size 201×33
 459 yields then approximations for f^{phillips} , f^{shaw} at these 201 equally spaced points. The *root mean*
 460 *square error* calculated using the exact values for f^{phillips} and f^{shaw} given by (24) and (25) can
 461 then be measured. The results using $\{\alpha_j^{\text{equispaced}}\}$, $\{\alpha_j^{\text{random}}\}$ and $\{\alpha_j^{\text{jitter}}\}$ are illustrated in Figure 9
 462 panels (b)-(d), respectively, for each problem. In these plots the interpolants are shown as the
 463 black diamonds as compared to the underlying exact solutions given by the continuous black
 464 lines. The diamonds are plotted for indices $1 : 4 : 201$ so as to permit visualization of the
 465 differences between exact and interpolated solutions. The title in each case gives the root mean
 466 square of the approximation. While the interpolation works well for $\{\alpha_j^{\text{equispaced}}\}$, it is immediate
 467 that the errors are large for the other two samplings, apparently contradicting Remark 6. But these
 468 results are obtained without appropriately reweighting the sampled $\{\alpha_j\}$ to be consistent with the
 469 integral equation (7). Results using the appropriately weighted samples using weights w_j in (7)
 470 chosen for a trapezoidal approximation to the integrals are illustrated in panels (e)-(f), yielding no
 471 improvement and further suggesting a problem with Remark 6, unless the solutions $\{\alpha_j^{\text{equispaced}}\}$ are
 472 not good approximations to the underlying infinite precision solutions.

473 Figure 10 illustrates the same experiments as in Figure 9 but using the truncated solutions. First
 474 from the solutions in panels (a) of Figure 10 it is clear that $\alpha^{\text{phillips}}(\mathbf{y})$, $\alpha^{\text{shaw}}(\mathbf{y})$ are far smoother than
 475 those illustrated in Figure 3. Carrying out the resampling process for the same sets of jittered and
 476 random samplings, yields the results in Figure 10. Now the obtained results are much improved
 477 and support the hypothesis of Remark 6, while also indicating that the frequency content of the
 478 underlying functions $\alpha(\mathbf{y})$ need not be as severe as originally suggested. These results emphasize
 479 that weights $\{\alpha_j\}$ may indeed be considered as sampling of a function $\alpha(\mathbf{y}) = w(\mathbf{y})\beta(\mathbf{y})$. Moreover,
 480 it is relevant to use techniques for the solution of the ill-posed integral equation (6), when finding
 481 *good* approximations to the interpolating function $\alpha(\mathbf{y})$.

482 To summarize: the presented results illustrate that the solution of the RBF approximation
 483 problem is closely connected to the solution of ill-posed integral equations and may therefore be

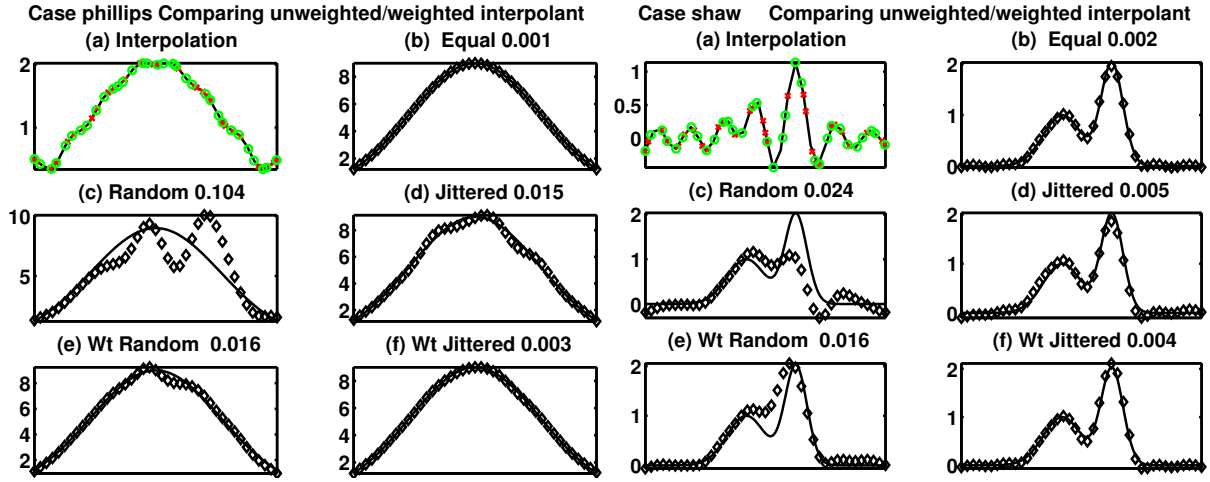


Figure 10: Using the truncated solution (26) with $N_\ell = 16$ to obtain approximants to f^{phillips} and f^{shaw} , the left and right blocks for each problem, respectively. Illustrated are then the equivalent results as shown in Figure 9 but obtained from the truncated solutions. The shape parameter is $\epsilon = 2$ as in Figure 3.

484 analysed using relevant techniques from the inverse problem literature.

485 6. Conclusions

486 Previous investigations of the properties of RBF interpolation/approximation seldom consid-
 487 ered the close connection between the solution of the interpolation problem and the inverse prob-
 488 lem for linear integral equations. Carefully examining this relationship shows that the system
 489 matrices of RBF interpolation may exhibit the same exponentially decaying spectrum that is com-
 490 monly seen when solving ill-posed inverse problems. In this case the spectral properties depend
 491 directly on the shape parameters of the RBFs, which determine whether a given RBF is *flat*. The
 492 *flatter* (more smooth) the RBF ϕ_ϵ , the faster the exponential decay of the spectrum, due to the lack
 493 of high frequency content in ϕ_ϵ , and the consequently greater difficulty in solving the underlying
 494 system of equations. Many prior results have focused on the solvability and stability of the RBF
 495 system, stressing the importance of the condition number of the system, [10–13, 40–44]. This
 496 paper, in contrast, demonstrates that it is the exponential decay of the spectrum, related directly
 497 to the smoothing effects of the operator, which better explains the numerical difficulties of solving
 498 the underlying problems.

499 The Picard condition is used in standard investigations of ill-posed problems to determine a
500 meaningful solution of a given set of equations, by revealing the number of reliable spectral coeffi-
501 cients in the SVD expansion of the kernel. For the RBF problem, the Picard condition reveals that
502 the solution depends on both the smoothness of the underlying function to be approximated (the
503 target function) and the *smoothness / flatness* of ϕ_ϵ . In contrast to the situation when seeking the
504 solution of a given integral equation, in which the kernel function is prescribed by an underlying
505 model of the problem, in the RBF case ϕ_ϵ can be chosen based on required criteria for the given
506 problem. For example, one may deduce that for a function which is itself relatively *flat* that ϕ_ϵ
507 should also be relatively *flat*. On the other hand, the Picard condition analysis demonstrates that it
508 is necessary to balance the smoothness of the target function and the SVE of the kernel function,
509 in order to find a solution which is meaningful.

510 In analyzing the stability of the RBF interpolation problem it is important to look at the com-
511 plete spectral decomposition, namely the singular vectors which provide a basis for the solution,
512 as well as the singular values. An approach for determining the number of reliable, *noise free*
513 basis vectors has been presented, which goes beyond the standard Picard condition analysis, and
514 is independent of the functional data. When the norm of the resulting weight vector is relatively
515 large in contrast to the norm of the data, one may deduce that this is indicative of ill-conditioning
516 of the problem, and that one may need to look for a more stable kernel, or to apply some kind
517 of regularization. Here regularization based on using the truncated singular value expansion of
518 the solution has been shown to provide a meaningful solution for the case of noise-free data. The
519 truncation parameter can be estimated by examining the noise properties of the basis vectors.

520 The presented numerical results demonstrate that the obtained solutions are reliable in the
521 context of quadrature applied to the integral equation; the weights for the RBF interpolant are
522 samplings of a function $\alpha(\mathbf{y})$ weighted by relevant quadrature weights dependent on the sampling.
523 As such, the functional form of the weights is amenable to resampling so as to obtain weights
524 for an alternative set of centers. Sampling from the functional form which is obtained from the
525 truncated expansion yields reliable approximations, while resampling from the solution obtained
526 without truncation is not meaningful. The idea of treating the weights as samplings from an
527 underlying continuous function promises to provide an approach to adaptively move the centers of

528 the RBF kernels so as to either yield more efficient approximation evaluation, or, for example, to
529 calculate derivative approximations at a different set of points which may be useful in improving
530 the conditioning of the obtained operators for partial differential equation solvers.

531 In general there is significant scope to apply analysis techniques from the field of ill-posed
532 inverse problems to the solution of the RBF approximation problem, with respect to areas such
533 as investigation of the impact of error in measured data, determining feasible regularization tech-
534 niques, and understanding the impact of iterative techniques in solving the resulting systems of
535 equations. Likewise, developments from the RBF literature, most notably the QR-RBF algorithm
536 [2], may provide useful insights for finding solutions of ill-posed integral equations. These are
537 topics, amongst others, for future research relating to the use of the inverse theory for RBFs.

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